

# Anomalous transport: Theory and Applications - III

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**Yesterday:** stare at these equations until you think they are beautiful ...

$$\vec{J}_A = d_{ABC} \frac{\mu_B}{4\pi^2} \vec{B}_C + \left( d_{ABC} \frac{\mu_B \mu_C}{4\pi^2} + b_A \frac{T^2}{12} \right) \vec{\Omega}$$

$$\vec{J}_\epsilon = \left( d_{ABC} \frac{\mu_B \mu_C}{8\pi^2} + b_A \frac{T^2}{24} \right) \vec{B}_A + \left( d_{ABC} \frac{\mu_A \mu_B \mu_C}{6\pi^2} + b_A \frac{\mu_A T^2}{6} \right) \vec{\Omega}$$

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**Today:** Rendezvous with the **DEVIL** !!

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- **Experiment**

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**Today:** Rendezvous with the **DEVIL** !!

- **Experiment**
- **Condensed Matter Physics**



*Please, allow me  
to introduce myself*

...

*Hot Stuff!*

# Outline:

- Weyl semi-metals
- Physics at the Edge
- A prediction from holography?
- Summary

# Weyl semi-metals

Bloch wave functions  $\Psi(\vec{x}) = e^{i\vec{k}\cdot\vec{x}} u_{\mathbf{k}}(\vec{x})$

$$u_{\mathbf{k}}(\vec{x} + \vec{a}) = u_{\mathbf{k}}(\vec{x})$$

Physically inequivalent momenta: Brillouin zone

$$\vec{k} \equiv \vec{k} + \vec{K} \quad \vec{K} \cdot \vec{a} = 2\pi n$$

Berry connection on Brillouin zone

$$A_i = \langle u_{\mathbf{k}} | \frac{\partial}{\partial k_i} | u_{\mathbf{k}} \rangle$$

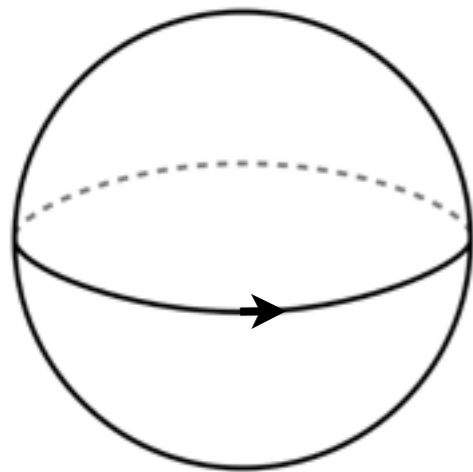
$$u_{\mathbf{k}} \rightarrow e^{\phi(\mathbf{k})} u_{\mathbf{k}}$$

$$A_i \rightarrow A_i + \frac{\partial}{\partial k_i} \phi(\mathbf{k})$$



# Weyl semi-metals

Berry curvature  $\mathcal{F} = d\mathcal{A}$



$$\Delta\phi = \oint \mathcal{A} = \int_{\text{upper}} \mathcal{F} = \int_{\text{lower}} \mathcal{F} + 2\pi n$$

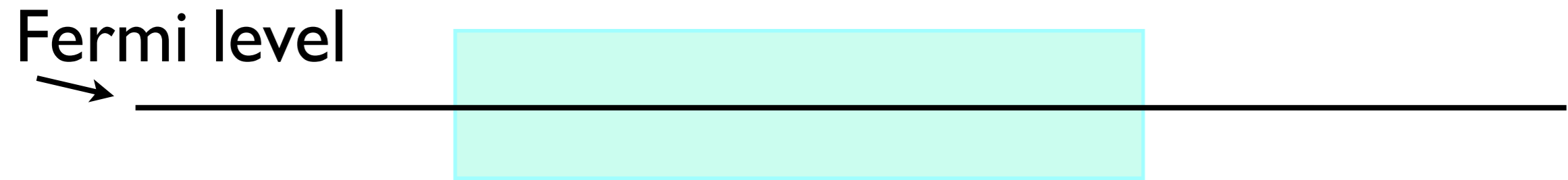
Flux of Berry curvature is quantized  $\oint \mathcal{F} = 2\pi n$

Topological Material = non trivial Berry connection

(anti) Chiral fermion = (anti) monopole in Berry connection

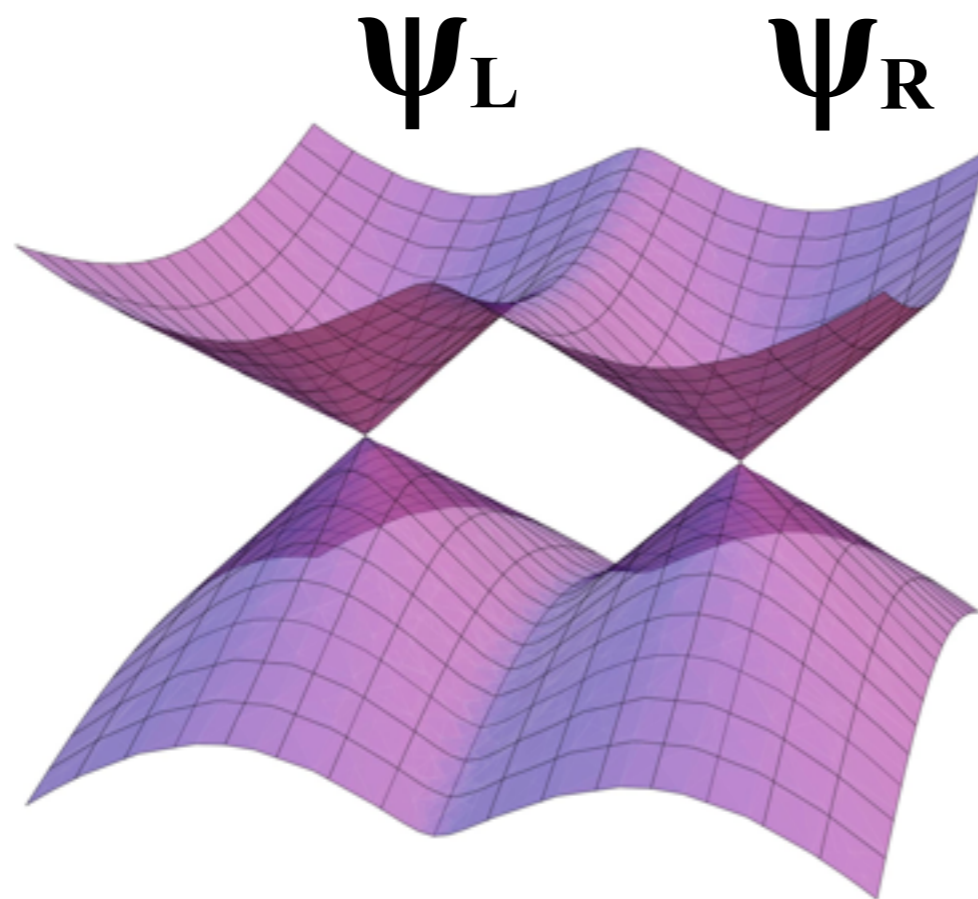
# Map to WSMs

## Band structure of WSM



# Weyl semi-metal

Close up: Linear band touching



$$\gamma^\mu (i\partial_\mu + A_\mu^5 \gamma_5) \Psi = 0$$

Nielsen-Ninomiya: left-right come in pairs

Stable: Monopole charge can not vanish  
(no mass for chiral fermions)

# Weyl semi-metal

Topological constraint

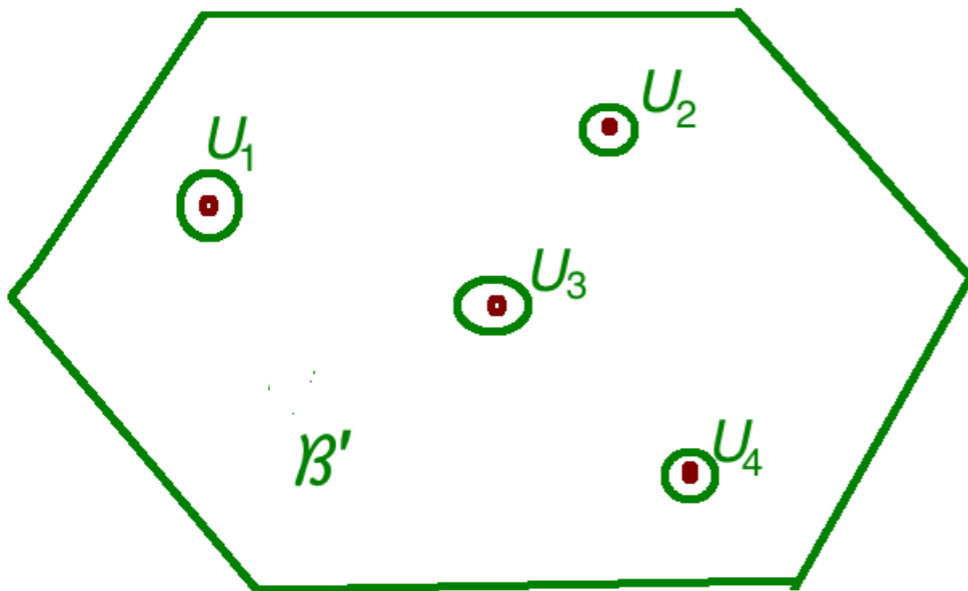
Berry connection

$$\mathcal{A} = \langle \psi(k) | \frac{\partial}{\partial k_i} | \psi(k) \rangle dk_i$$

$$\mathcal{F}_B = d\mathcal{A}$$

$$d\mathcal{F}_B = 0$$

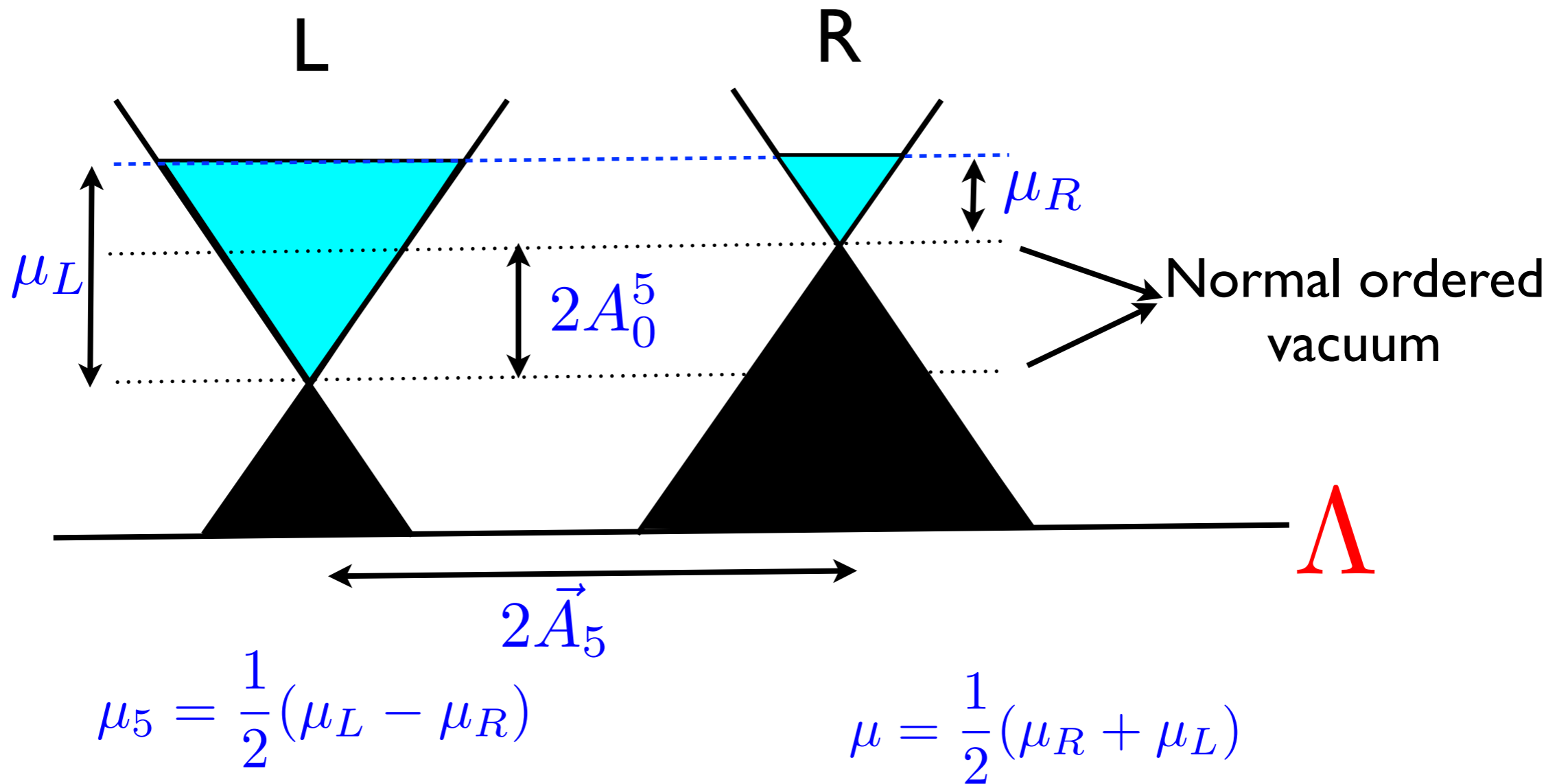
$$\int \frac{d\mathcal{F}_B}{2\pi} = \sum_i \oint_{U_i} \frac{\mathcal{F}_B}{2\pi} = 0$$



[Kiritsis]

BZ has no boundary !

# CME in WSMs



**CME:** 
$$\vec{J} = \frac{1}{2\pi^2} (\mu_5 - A_0^5) \vec{B} = 0$$

CME vanishes in equilibrium in WSM

“Bloch theorem”

# Negative Magnetoresitivity

**CME**  $\vec{J} = \sigma \vec{E} + \frac{\delta\mu_5}{2\pi^2} \vec{B}$

**Anomaly**  $\partial_t \delta\rho_5 = \frac{1}{2\pi^2} \vec{E} \vec{B}$

**Susceptibility**  $\delta\rho_5 = \chi_5 \delta\mu_5$

$$\vec{J} = \left( \sigma + \frac{iB^2}{\omega\chi_5} \right) \vec{E}$$

$$\Re(\sigma_{\text{tot}}) = \sigma + \frac{B^2}{\chi_5} \delta(\omega)$$

**Anomaly induces infinite DC conductivity !!**

# Negative Magnetoresitivity

CME  $\vec{J} = \sigma \vec{E} + \frac{\delta\mu_5}{2\pi^2} \vec{B}$

Anomaly  $\partial_t \delta\rho_5 = \frac{1}{2\pi^2} \vec{E} \vec{B} - \tau_c \delta\rho_5$

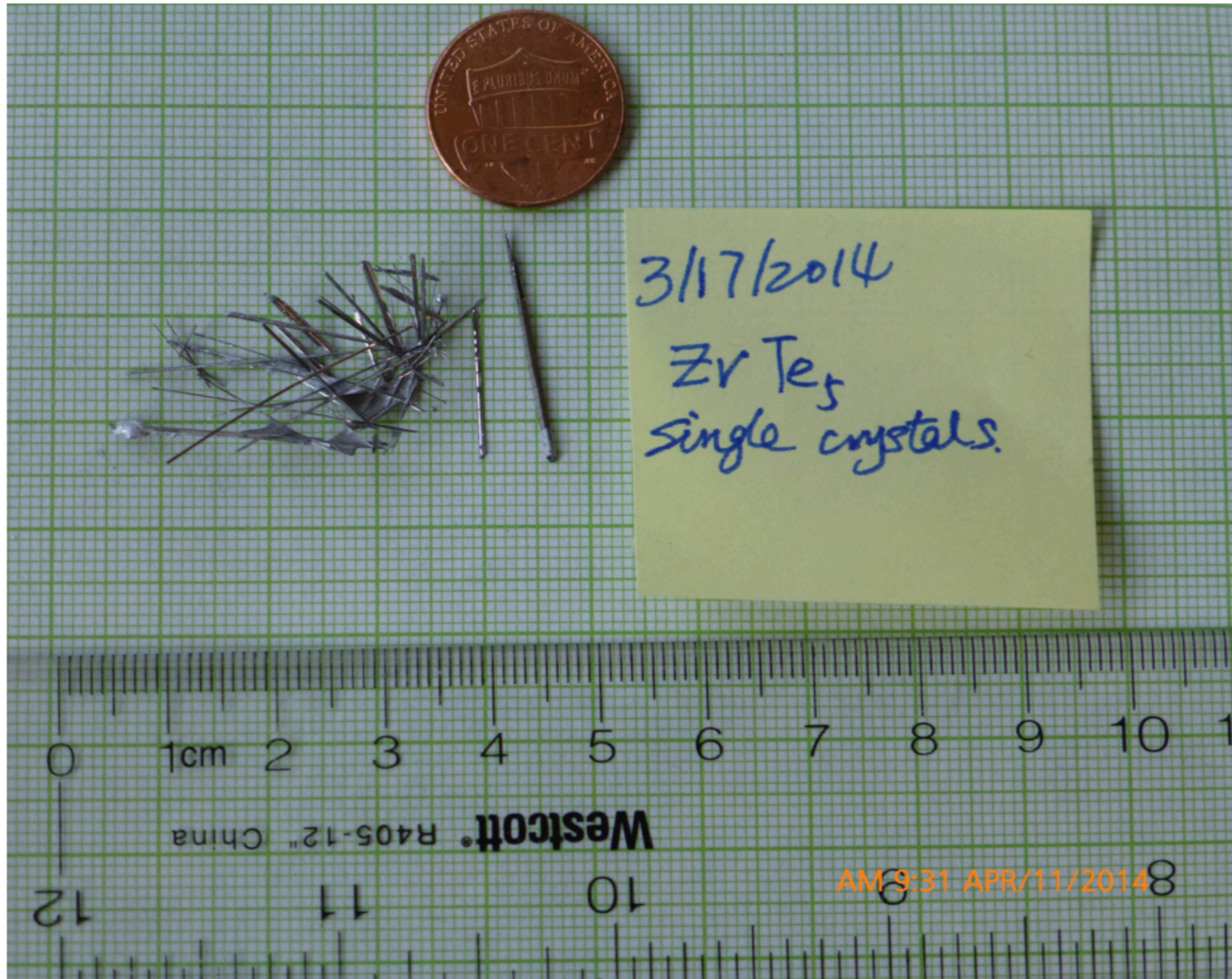
Susceptibility  $\delta\rho_5 = \chi_5 \delta\mu_5$

$$\vec{J} = \left( \sigma + \frac{\tau_c B^2}{\chi_5} \right) \vec{E}$$

Anomaly induces finite DC conductivity !!



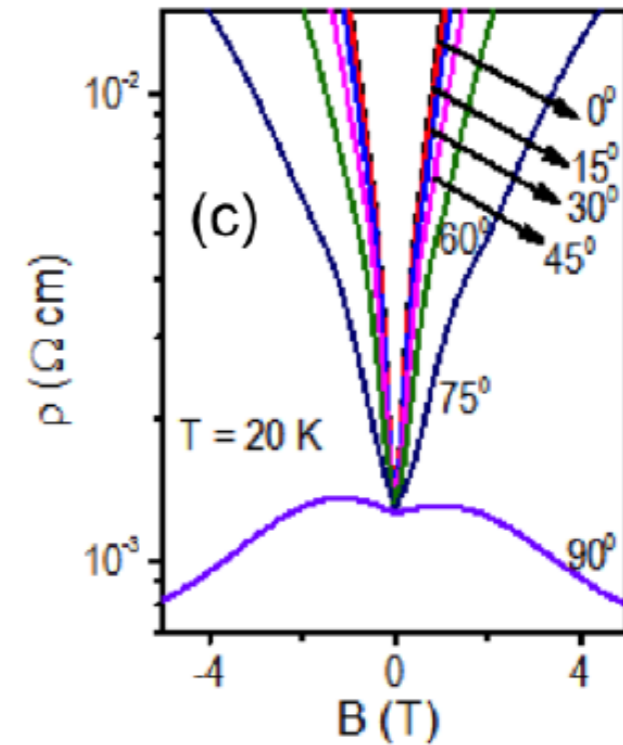
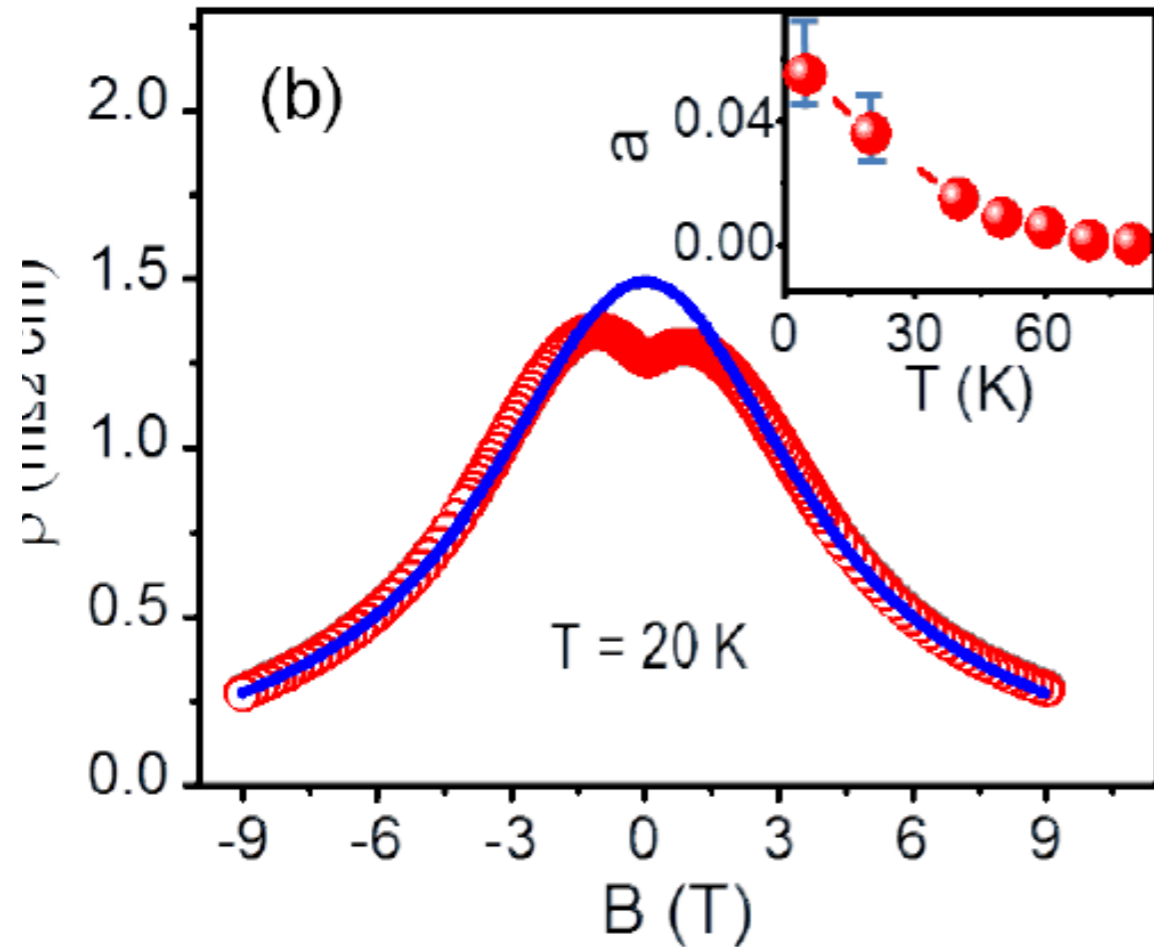
# Recent Experiments



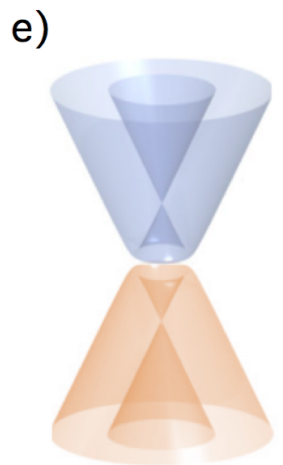
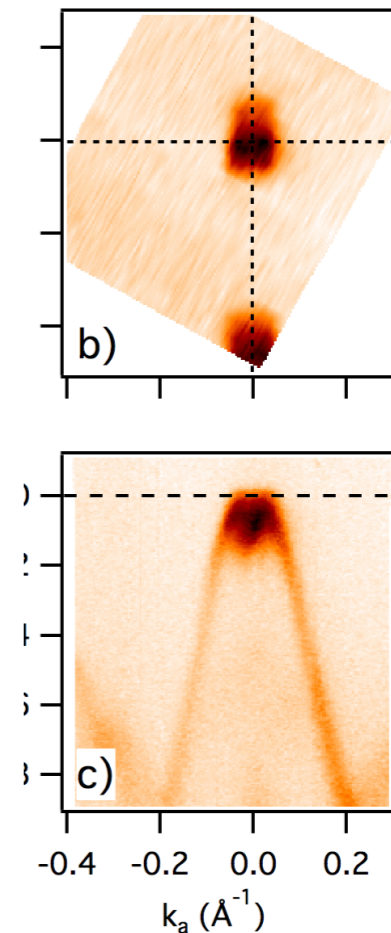
[D. Kharzeev: Lectures at Schladming Winterschool 2015]



# Recent Experiments

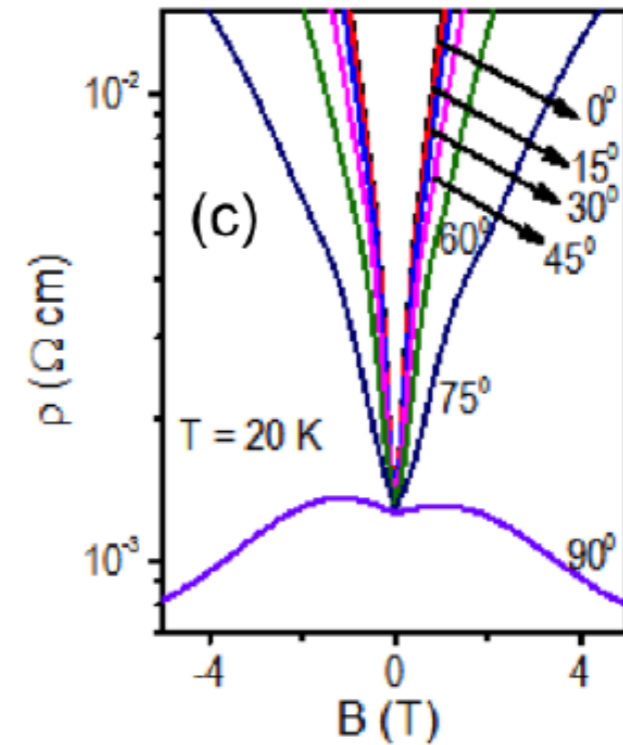
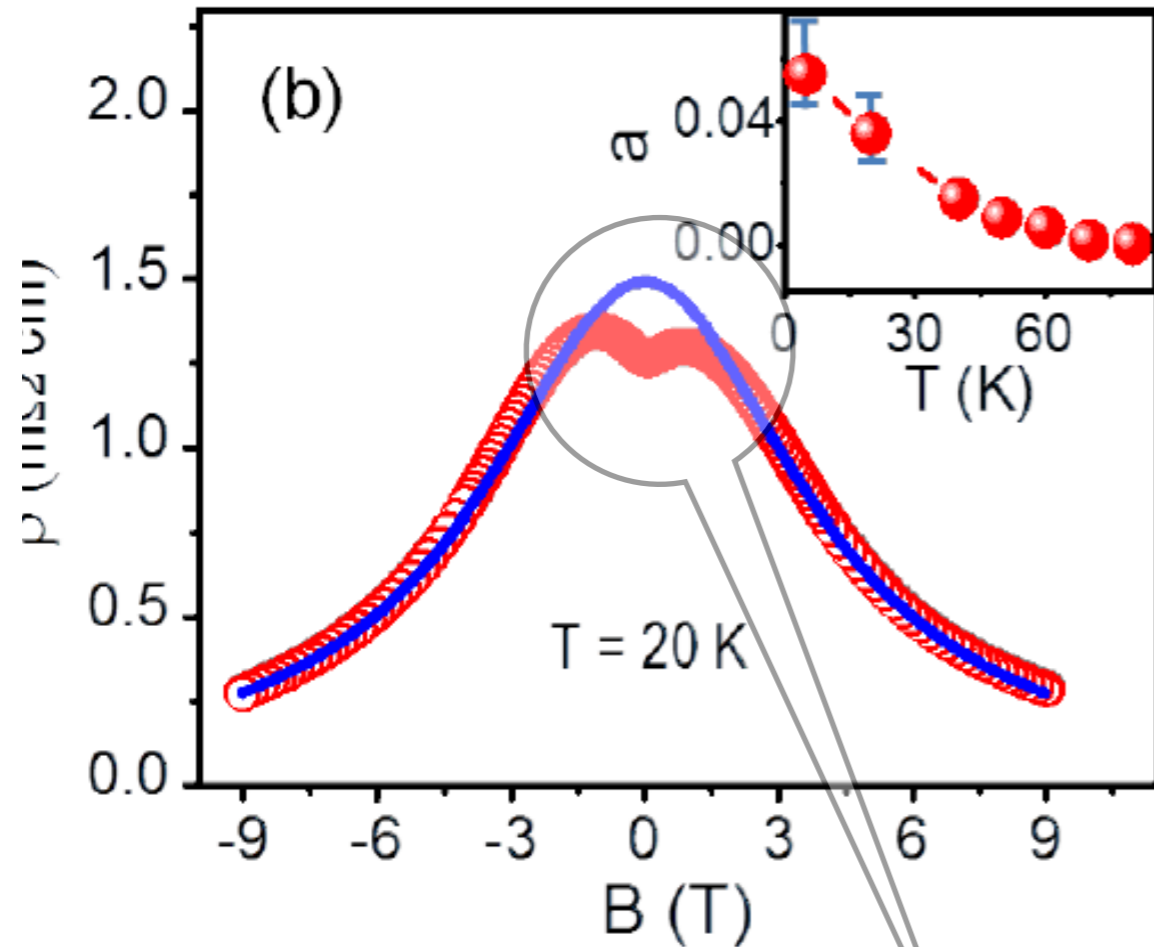


fits to  $1/(a+b B^2)$   
 **$B^2$  behaviour!**



[Li, Kharzeev et al. NaturePhysics 3648 (2016)]

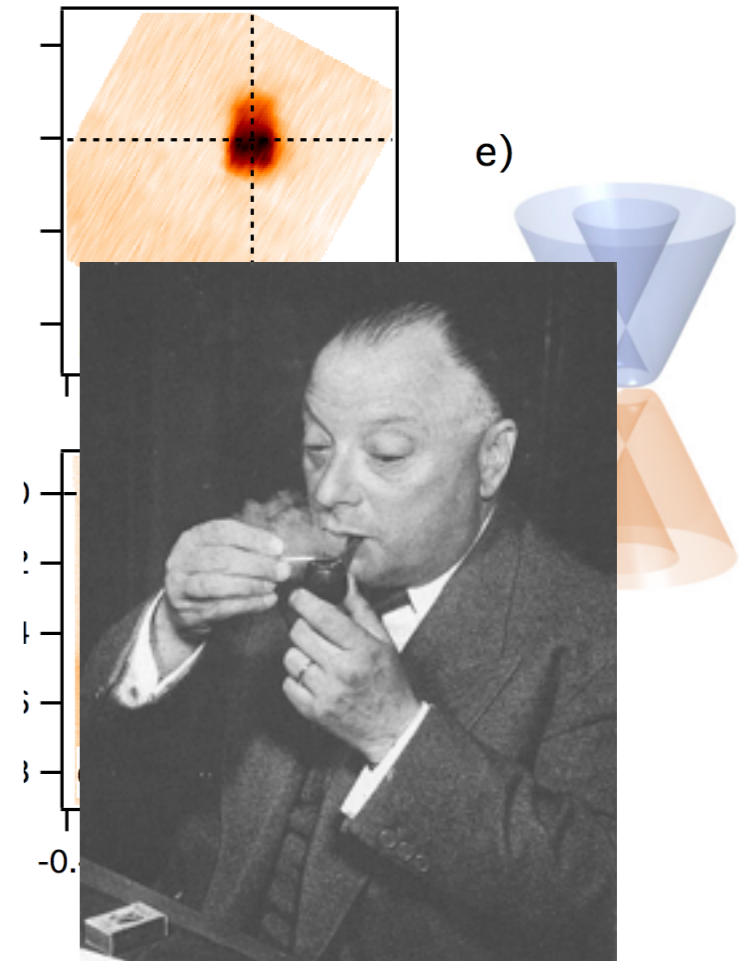
# Recent Experiments



fits to  $1/(a+b B^2)$   
 $B^2$  behaviour!

“Schmutzphysik”  
(dirt physics)

[Li, Kharzeev et al. NaturePhysics 3648 (2016)]



# AHE in WSMs

$$J_{cons}^{\mu} = J_{cov}^{\mu} + \frac{1}{4\pi^2} \epsilon^{\mu\nu\rho\lambda} A_{\nu}^5 F_{\rho\lambda}$$

$\vec{A}^5$  Separation of Weyl nodes in Brillouin zone

$$\vec{J} = \frac{1}{2\pi^2} \vec{A}_5 \times \vec{E}$$

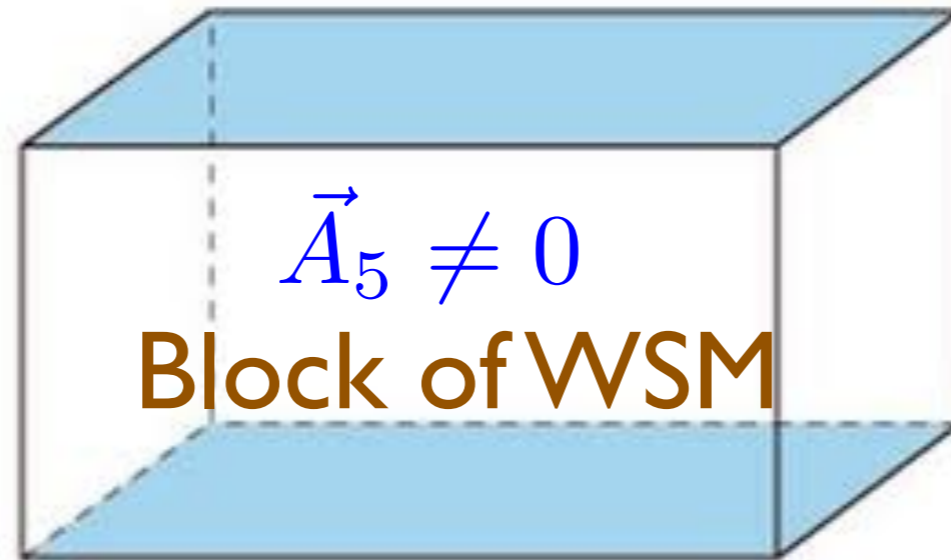
No contribution from covariant current  
(no analog of chemical potential in momentum)

# Edge Physics

Vacuum

Vacuum

$$\vec{A}_5 = 0$$



Vacuum

Vacuum

- Vacuum: (infinitely) massive Dirac fermions
- Need gradient in axial vector = axial magnetic field

$$\vec{J} = \frac{\mu}{2\pi^2} \vec{B}_5 \quad \vec{B}_5 = \nabla \times \vec{A}_5$$

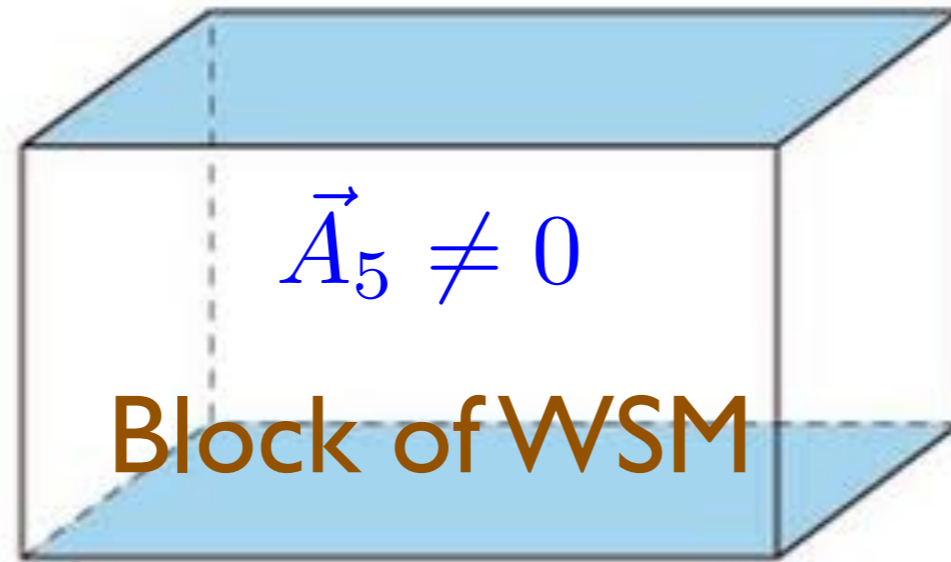
- Edge current
- States: “Fermi arcs” (degeneracy =  $|A_5|$ )

# Edge Physics

Vacuum

Vacuum

$$\vec{A}_5 = 0$$



Vacuum

Vacuum

- Thermal Hall Effect



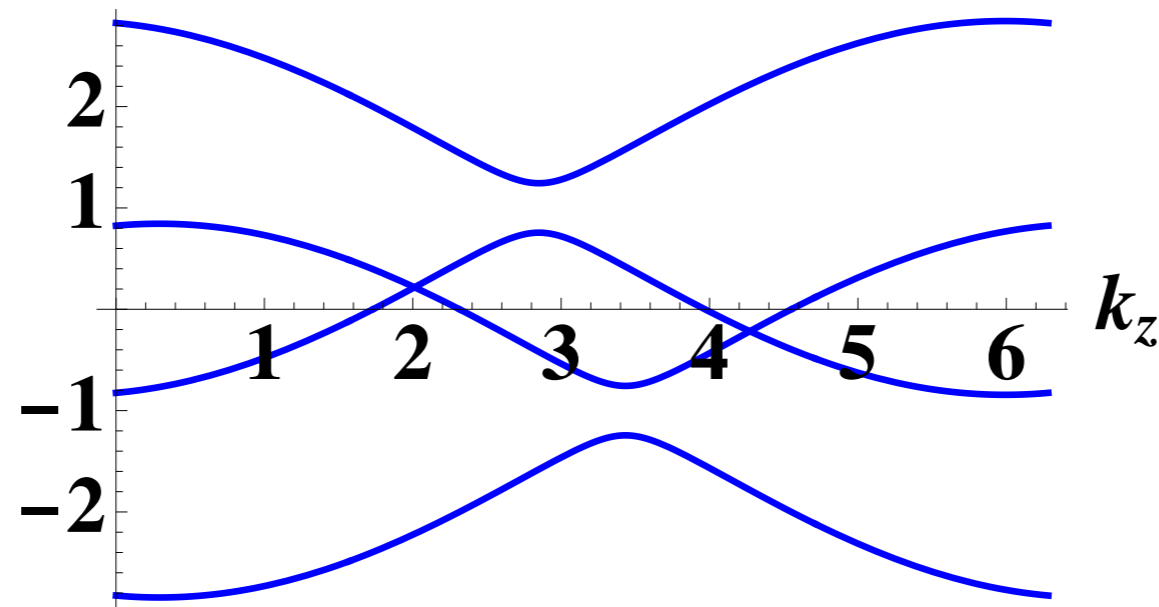
$$\vec{J}_\epsilon = \left( \dots + \frac{T^2}{12} \right) \vec{B}_5 \quad \vec{B}_5 = \nabla \times \vec{A}_5$$

- Net energy current at edge
- Effect of *gravitational anomaly!!*

$$J_{\text{thermal}} = \frac{T}{6} |A_5| \Delta T$$

# WSM

In crystal left- and right-handed chiralities are connected along the band dispersions



Can we find a QFT and a holographic model with this property?

**YES!!**

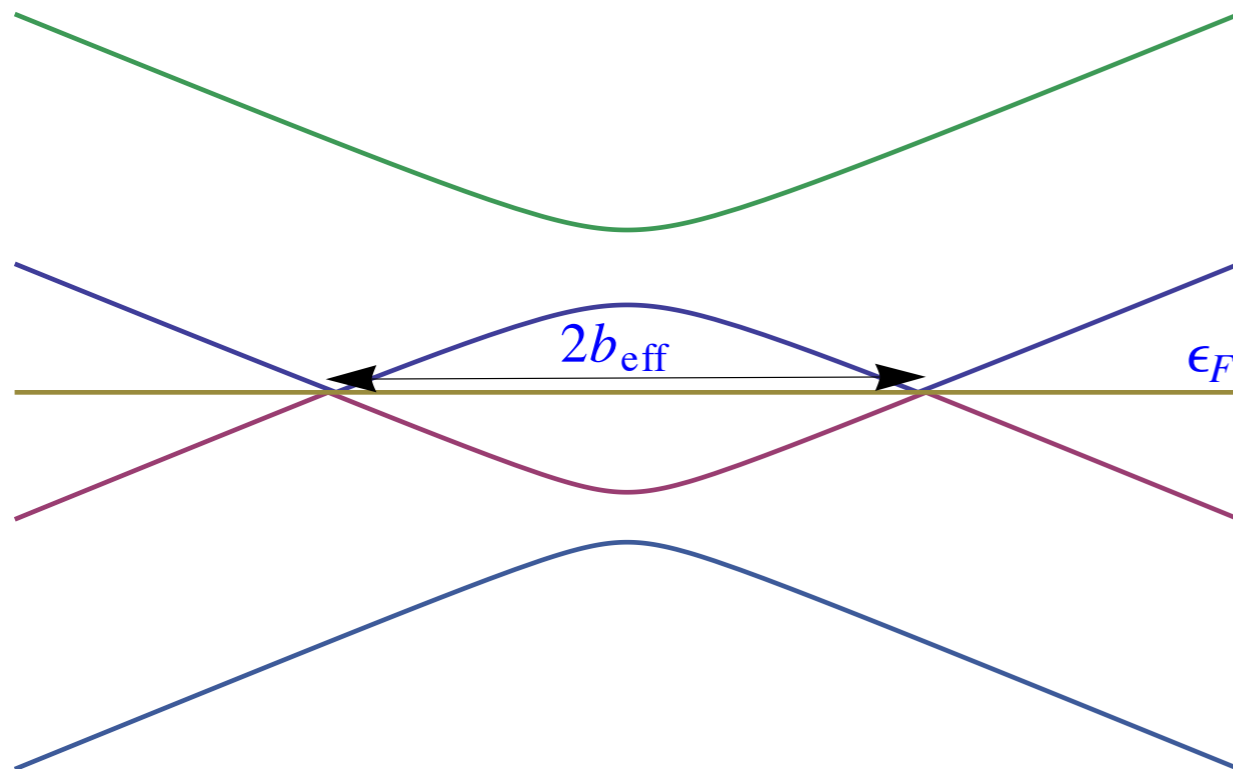
# QFT model

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu \partial_\mu + M + \gamma_5 \gamma_z b)\psi$$

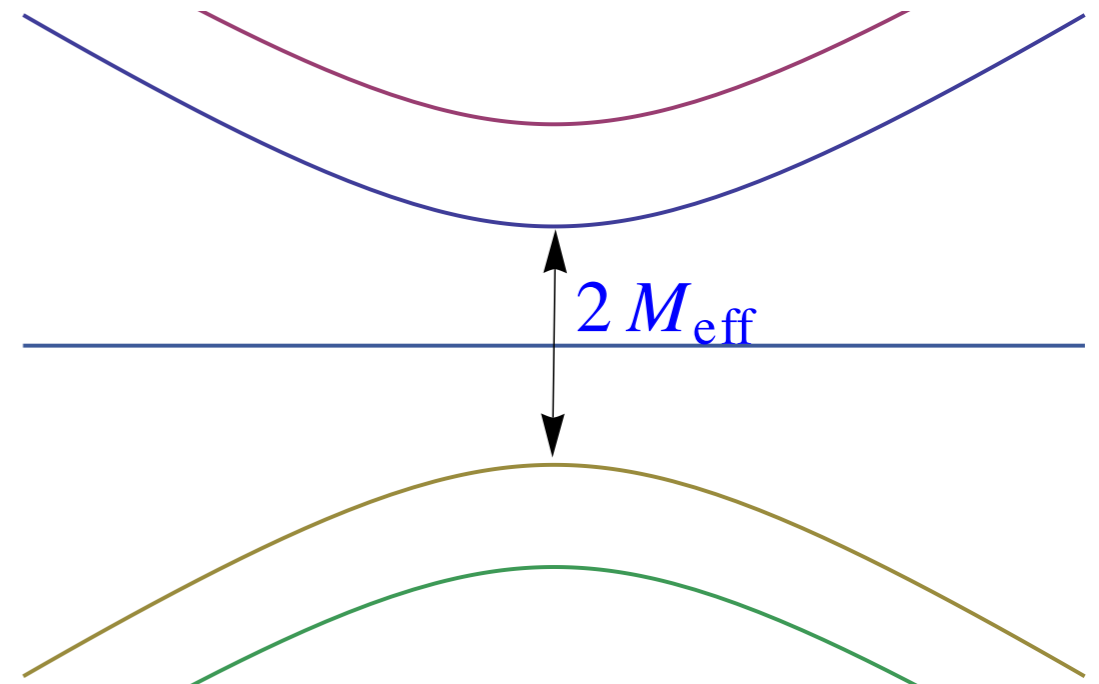
[Kostolecky et al. ] [Jackiw]  
[Burkov, Balents] [Grushin]

spectrum:

$$M < b : \quad b_{\text{eff}} = \sqrt{b^2 - M^2}$$



$$M > b : \quad M_{\text{eff}} = \sqrt{M^2 - b^2}$$



# Action of HoloWSM

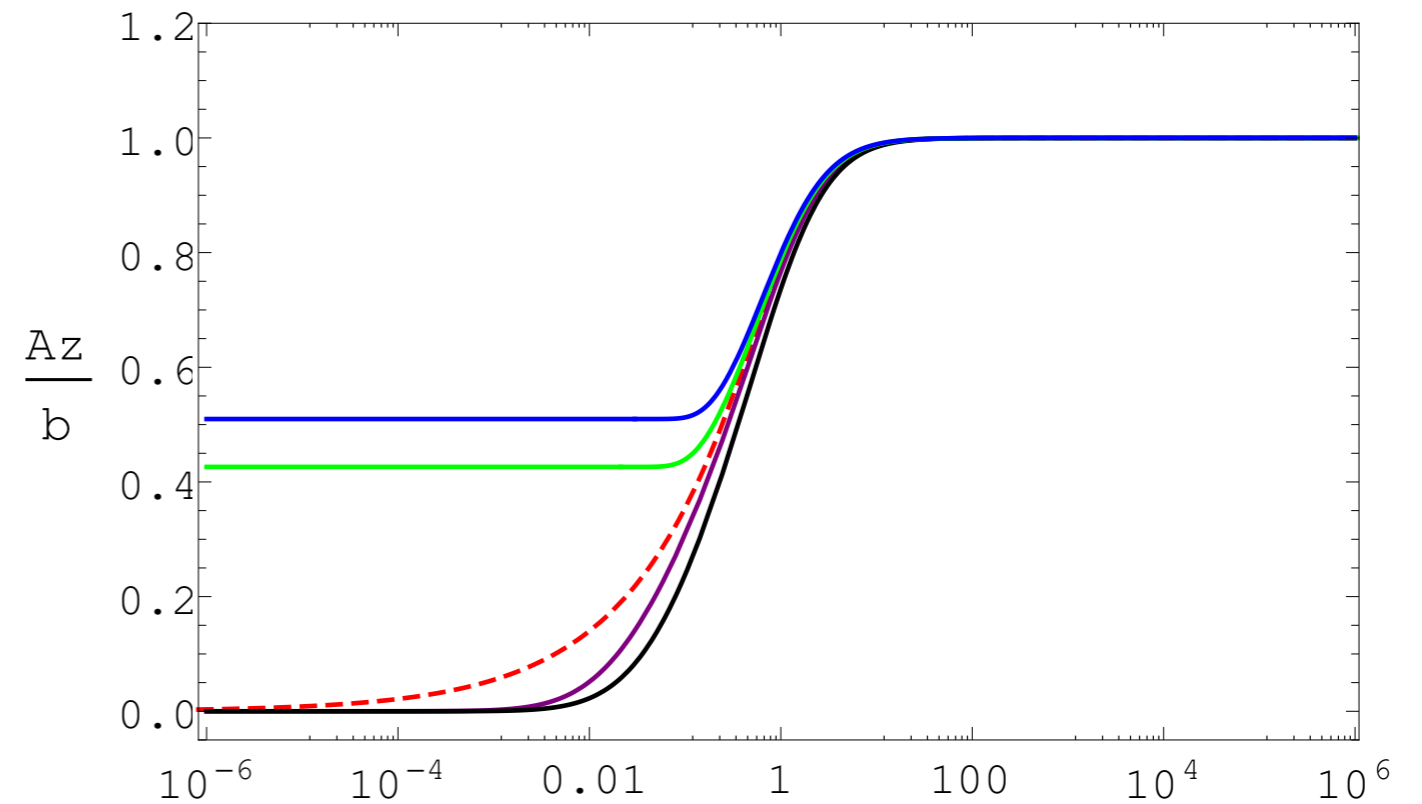
$$\begin{aligned}\mathcal{L} = & \frac{1}{2\kappa^2}(\mathcal{R} + 12) - \frac{1}{4}\mathcal{F}^2 - \frac{1}{4}F_5^2 + \\ & |(\partial_\mu + iqA_\mu^5)\Phi|^2 - V(|\Phi|) + \\ & + A_5 \wedge \left( \frac{\alpha}{3}F_5 \wedge F_5 + \alpha\mathcal{F} \wedge \mathcal{F} + \zeta R \wedge R \right)\end{aligned}$$

- Cosmological constant = AdS
- Very specific CS term = form of Anomaly
- Scalar potential determines dimension of dual scalar operator (we chose dim=3) i.e. mass deformation

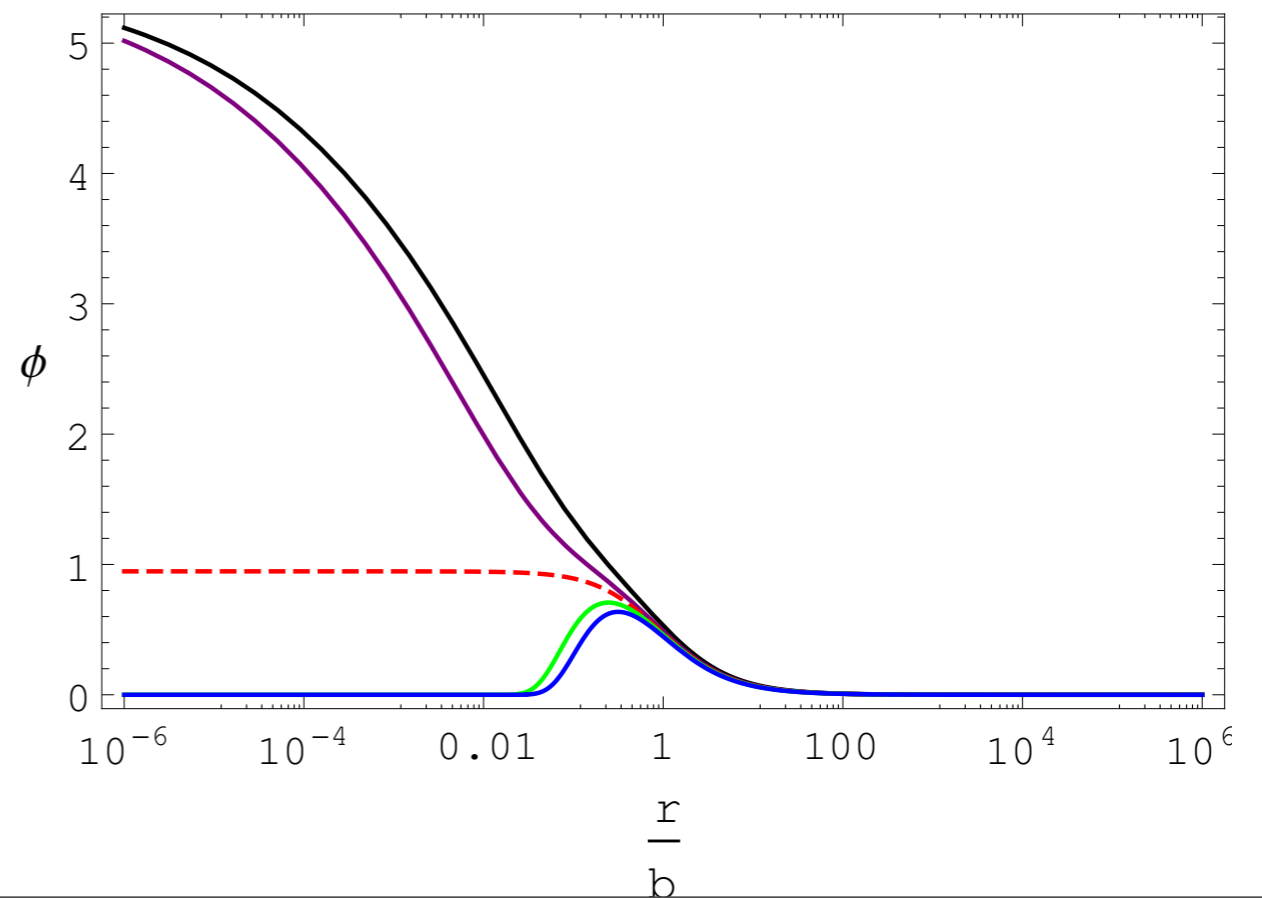


# A holographic topological quantum phase transition

Running of axial gauge field:



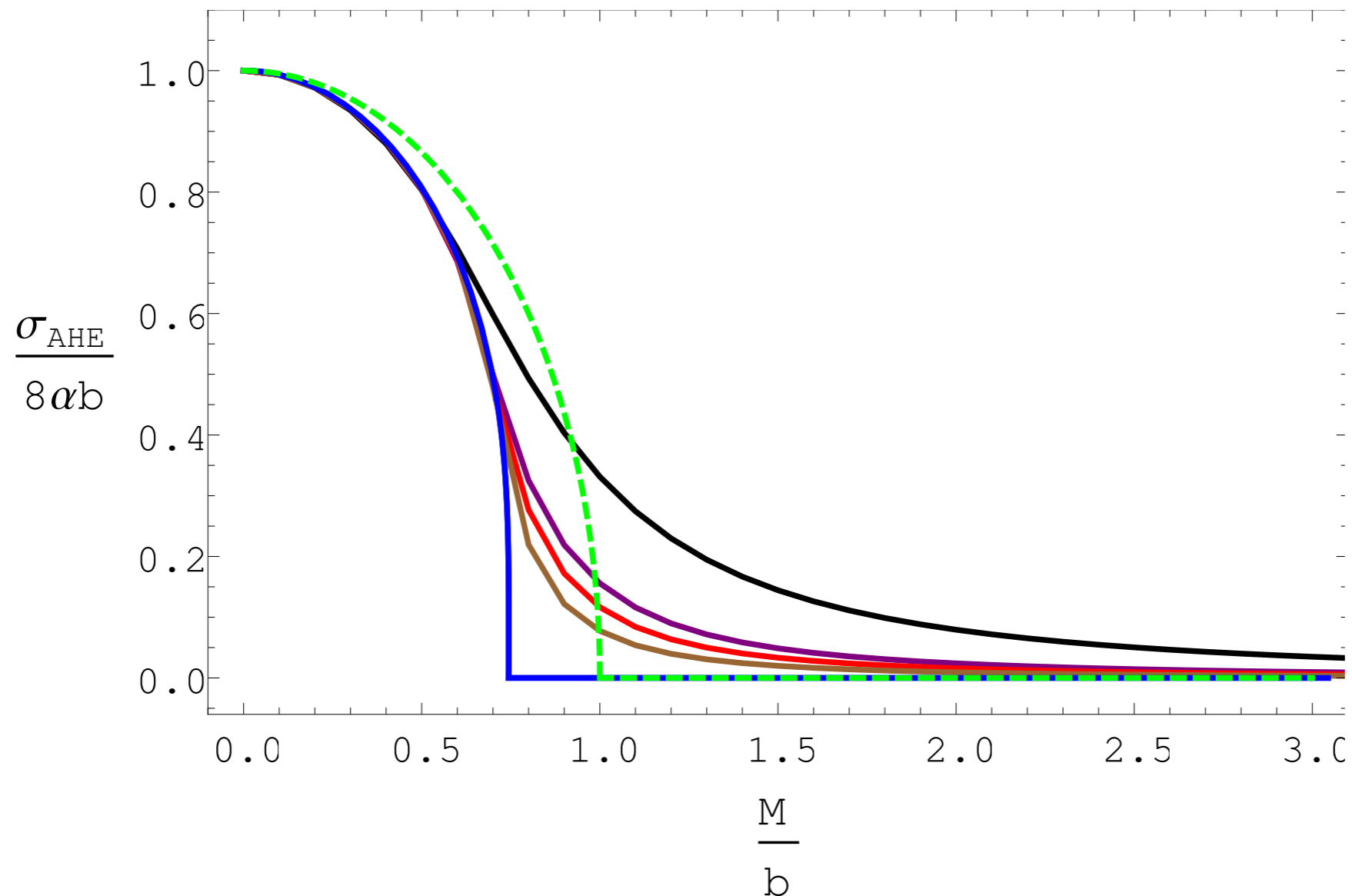
Running of scalar field:



# Holographic WSM

Smoking gun of topological state of matter : AHE

$$\sigma_{xy} = 8\alpha A_z^5(0)$$



# Odd viscosity

- Hall viscosity in 2D Quantum Hall states [Avron, Seiler, Zograf]
- Time reversal breaking necessary
- 2D : invariant  $\epsilon$  tensor
- 3D: need some additional vector
- Hydrodynamics of magnetized plasmas [Landau, Lifshytz Vol. 10]

$$\tau_{xy} = \eta_{\perp} V_{xy} - \eta_{\perp}^H (V_{xx} - V_{yy})$$

$$\tau_{xz} = \eta_{\parallel} V_{xz} + \eta_{\parallel}^H V_{yz}$$

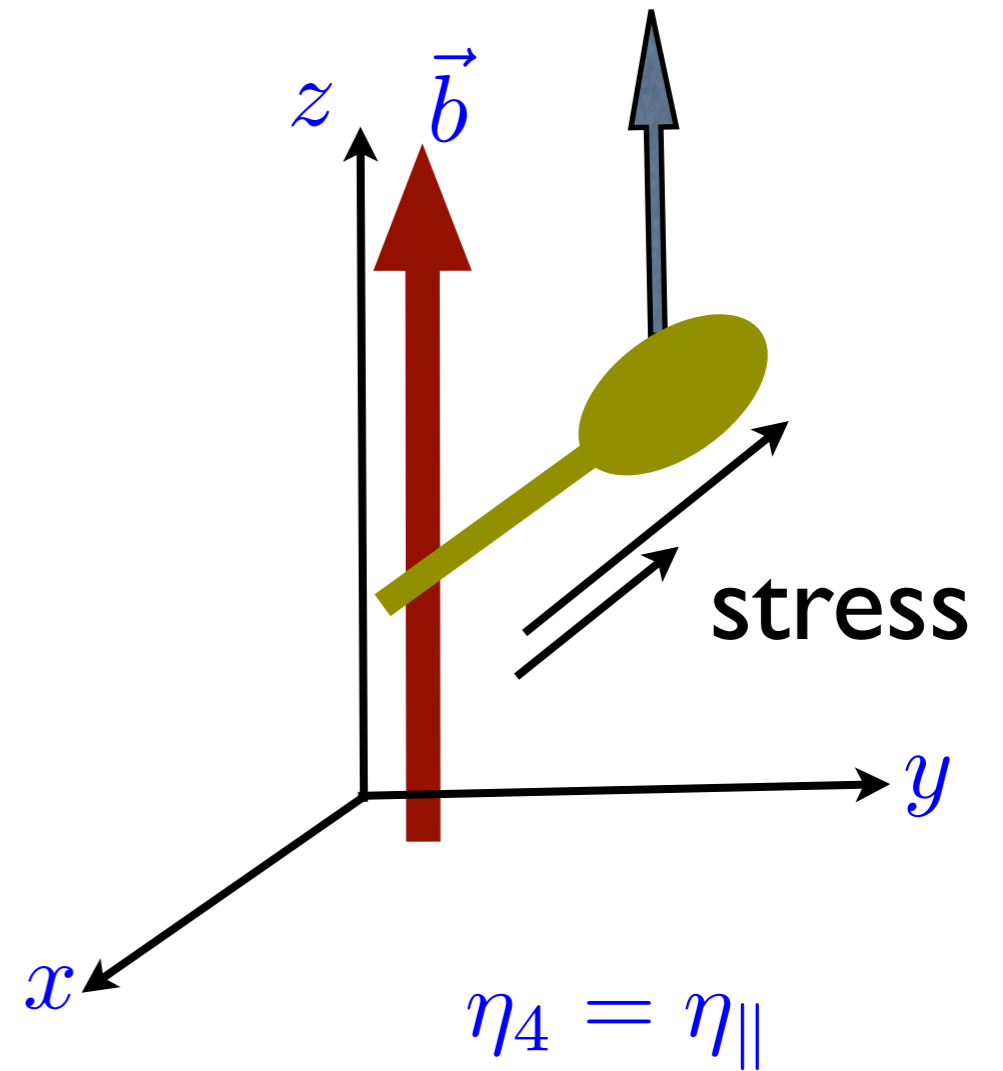
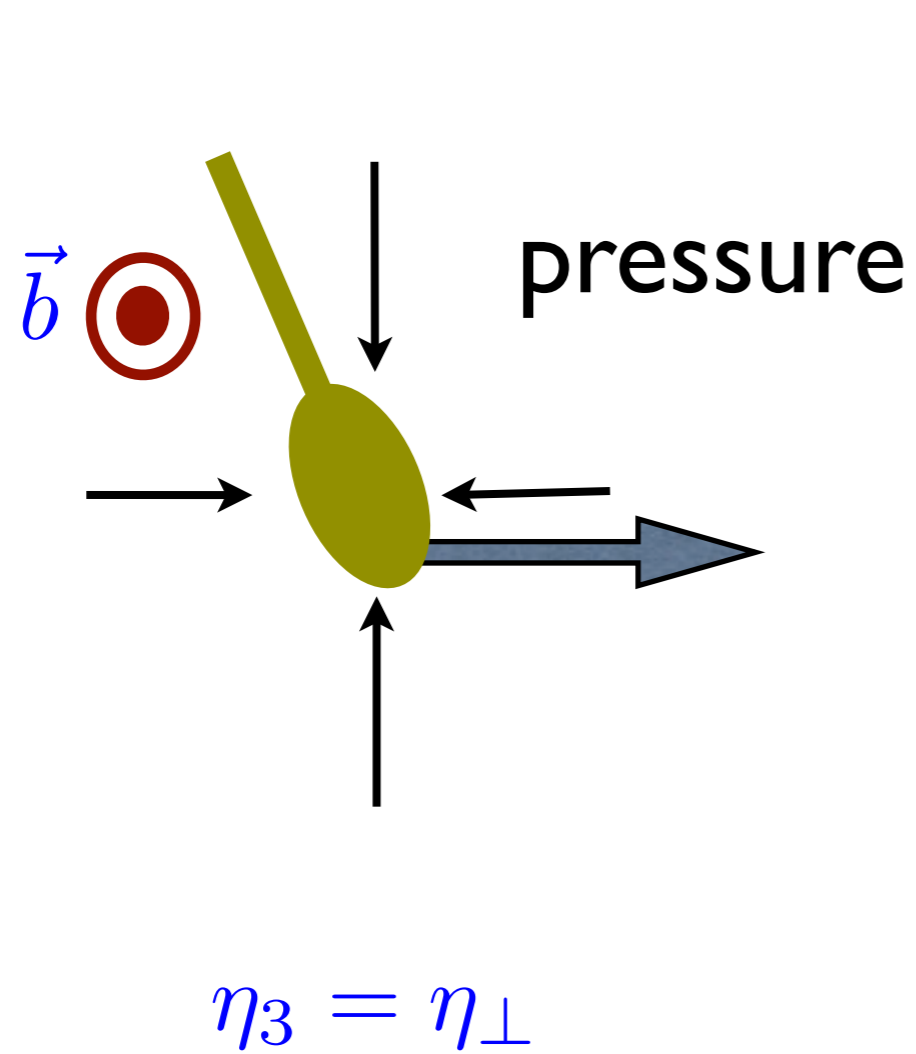
$$\tau_{yz} = \eta_{\parallel} V_{yz} - \eta_{\parallel}^H V_{xz}$$

$$V_{ij} = \frac{1}{2} (\partial_i v_j + \partial_j v_i)$$

odd viscosities

- In total: 3 shear, 2 “bulk” and 2 odd viscosities
- $\eta_{odd}(-B) = -\eta_{odd}(B)$

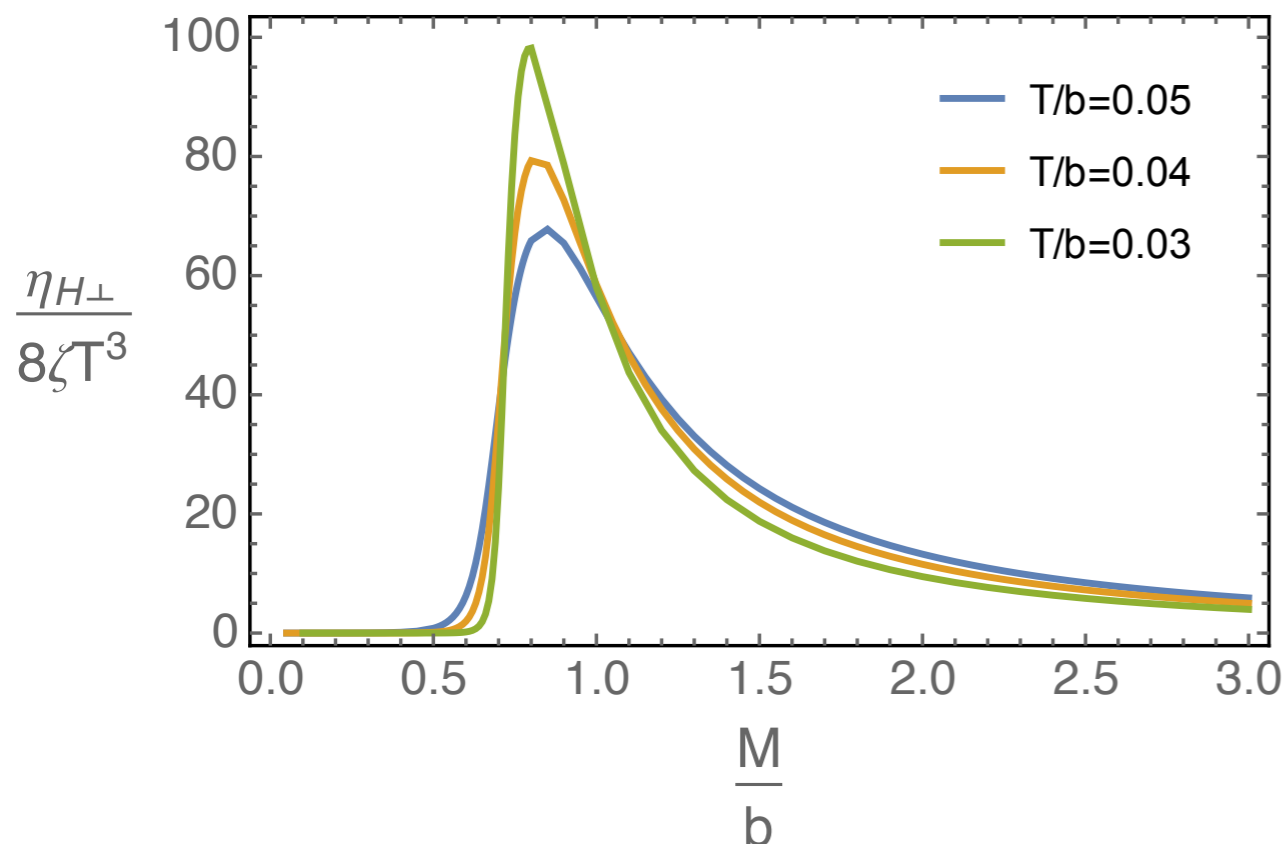
# Odd viscosity



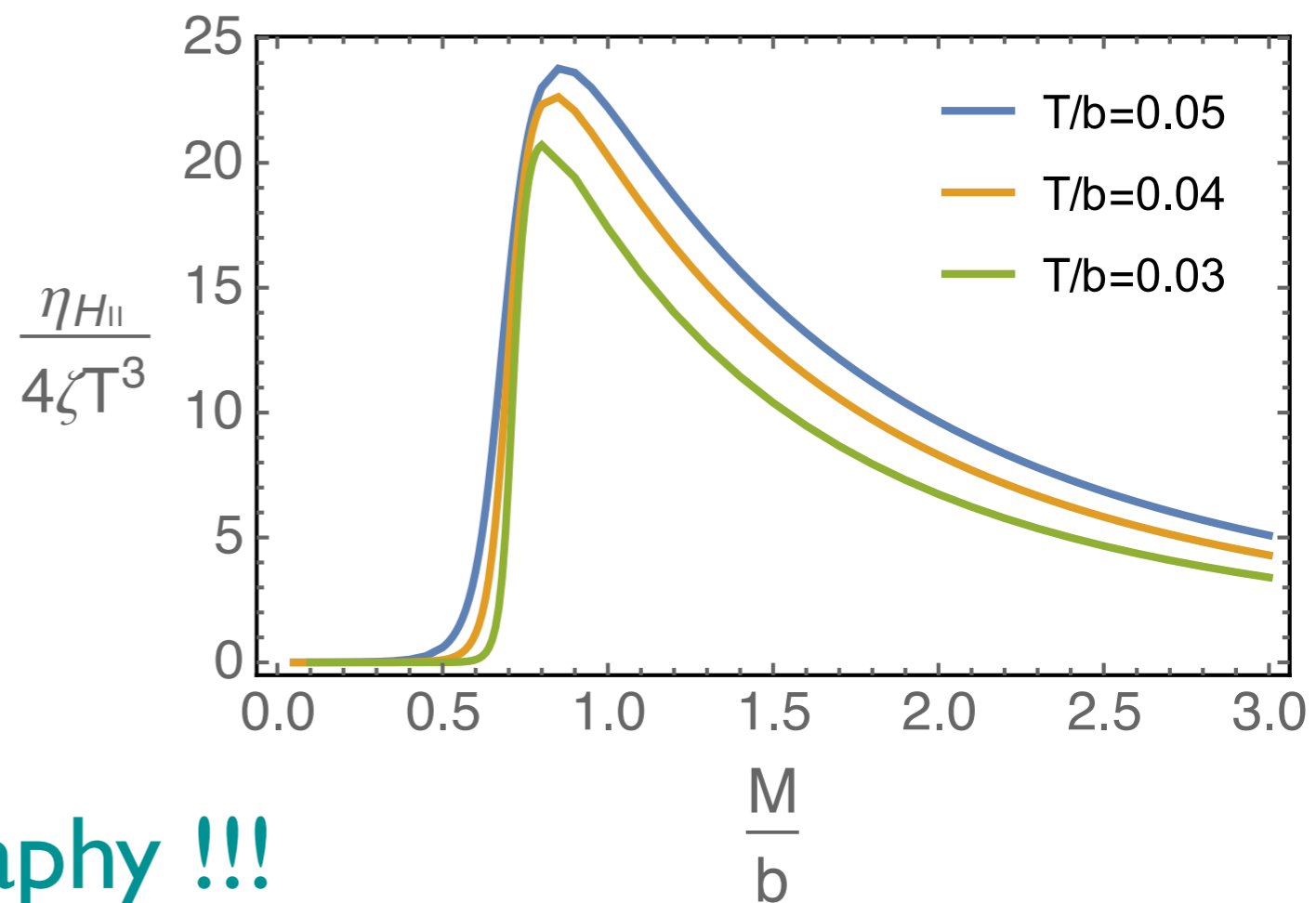
# Odd viscosity

- Odd viscosities
- Probe IR region of geometry: Low T

*transverse*



*parallel*



- Prediction from Holography !!!
- Again: gravitational Anomaly at first order !

# Summary

# Summary

