Holography of non-conformal field theory and entanglement entropy

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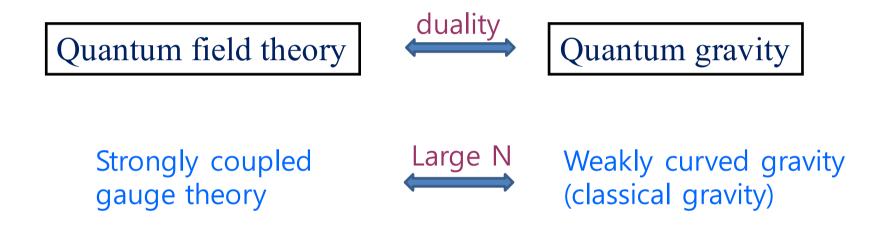
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Outline

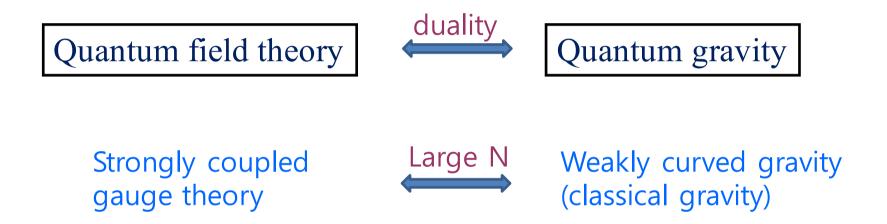
- Gauge/gravity duality for finite N
- Exact holography for finite N
- Holographic entanglement entropy for the LLM
- Gauge/gravity duality and entanglement entropy
- Summary

Duality properties of field theory and gravity theory



Very useful but difficult to check the duality!
 For some BPS objects which have no quantum corrections, it was possible to check the duality using supersymmetry and conformal symmetry in the large N limit.

Duality properties of field theory and gravity theory



To test Maldacena's conjecture, large N is necessary!

AdS/CFT Conjecture by Maldacena:

We started with a quantum theory and saw that it includes gravity, so it is natural to think that this correspondence goes beyond the supergravity approximation. We are led to the conjecture that <u>type IIB string theory on</u> $(AdS_5 \times S^5)_N$ plus some appropriate boundary conditions (and possibly also some boundary degrees of freedom) is dual to N = 4, d = 3 + 1, U(N) super-Yang-Mills.

- \rightarrow No restriction on N, g_{YM} Strong conjecture!
- → to support this conjecture he considered the supergravity limit

• finite N?

$$g_{YM}^2 = g_s$$
 $\lambda = g_{YM}^2 N$ $L^4 = 4\pi g_{YM}^2 N \alpha'^2 = 4\pi \lambda \alpha'^2$

for limiting cases

$$\lambda\gg 1\iff g_s\gg 1 \& \alpha'\ll 1 \Rightarrow$$
 String loop effects $\lambda\ll 1\iff g_s\ll 1 \& \alpha'\gg 1 \Rightarrow$ Stringy effects

→ difficult to test the gauge/gravity duality!

No example of exact gauge/gravity duality for finite N up to now!



One example of exact gauge/gravity duality:

Conformal symmetry → non-conformal symmetry

Large N → finite N

ABJM (CFT) + mass deformation = mABJM (non-CFT)

$$AdS_4 \times S^7/\mathbb{Z}_k$$

$$\mathcal{M}_{1}^{(n)} = \begin{pmatrix} \sqrt{n} & 0 & & \\ & \sqrt{n-1} & 0 & & \\ & & \sqrt{2} & 0 & \\ & & & 1 & 0 \end{pmatrix}$$
 Discrete Higgs vacua denoted by numerical values \rightarrow Quantitative comparison

$$\mathcal{M}_{2}^{(n)} = \begin{pmatrix} 0 & 1 & & \\ & 0 & \sqrt{2} & & \\ & & \ddots & \ddots & \\ & & & 0 & \sqrt{n-1} & \\ & & & 0 & \sqrt{n} \end{pmatrix}$$
 [Gomis et al 08] [Cheon-Kim-Kim 2011]

LLM geometry with \mathbb{Z}_k orbifold

[Lin-Lunin-Maldacena, 2004]

Exact gauge/gravity for finite N

- To show the dual relation of two theories, correspondence between vacua of QFT and BPS solutions in gravity is not enough.
 - mABJM theory: 3-dim. gauge theory
 - LLM geometry: BPS solution in 11-dim. SUGRA
 KK-reduction
 4-dim. Gravity theory

Exact gauge/gravity for finite N

 To show the dual relation of two theories, correspondence between vacua of QFT and BPS solutions in gravity is not enough.

mABJM theory: 3-dim. gauge theory
 LLM geometry: BPS solution in 11-dim. SUGRA
 duality
 4-dim. Gravity theory

Exact gauge/gravity for finite N

• Vev of CPO with $\Delta = 1$ in field theory side

$$\langle \mathcal{O}^{(1)} \rangle_0 = \frac{k\mu}{4\sqrt{2}\pi} \sum_{n=0}^{2N_{\rm B}+1} n(n+1)(N_n - N_n')$$

$$\mathcal{O}^{(1)} = \frac{1}{2\sqrt{2}} \text{Tr} \left(Z^a Z_a^{\dagger} - W^{\dagger a} W_a \right)$$

Gauge/gravity mapping (GKP-W relation)

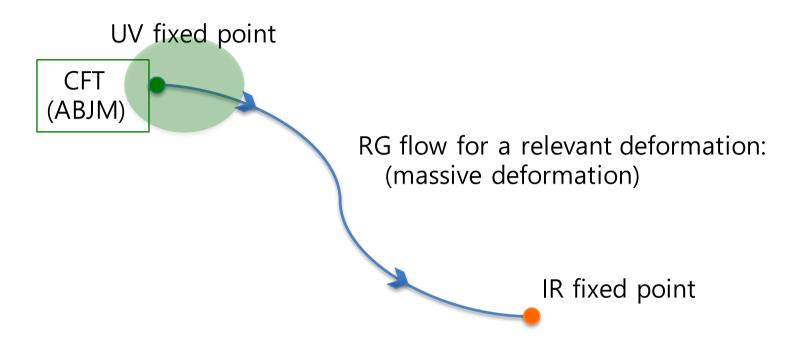
$$\langle \mathcal{O}^{(1)} \rangle_0 = \mathbb{N} \psi_{(1)} = -6\sqrt{2} \mathbb{N} \beta_1 \mu_0$$

$$\beta_1 = 2C_1^3 - 3C_1C_2 + C_3$$

$$C_p = \sum_{i=1}^{2N_{\rm B}+1} (-1)^{i+1} \left(\frac{\tilde{x}_i}{2\pi l_{\rm P}^3 \mu_0 \sqrt{A}} \right)^p$$

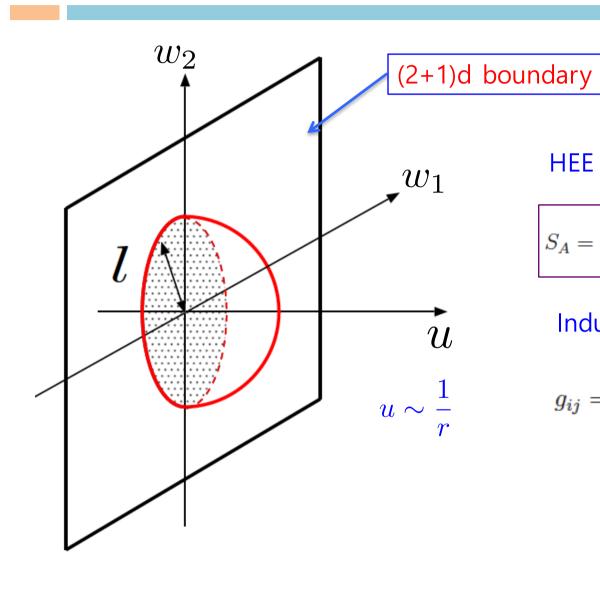
Fix the normalization factor (k=1):

$$\langle \mathcal{O}^{(1)} \rangle_0 = \frac{N^{\frac{3}{2}} \mu_0}{3\sqrt{2}\pi} \beta_1$$
 $N = -\frac{N^{\frac{3}{2}}}{36\pi}$ $N \ge 2$



- Exact holographic renormalization near the UV fixed point
- Scalar fields deformation induces the metric deformation
 - deformation of the entanglement entropy
- How can we relate vevs of CPOs in QFT with the holographic entanglement entropy?

Holographic Entanglement Entropy



HEE proposal

$$S_A = \frac{\operatorname{Min}(\gamma_A)}{4G_N}, \quad \gamma_A = \int d^9 \sigma \sqrt{\det g_{ij}}.$$

Induced metric $i = 1, 2, \dots, 9$

$$g_{ij} = \frac{\partial X^M \partial X^N}{\partial \sigma^i \partial \sigma^j} G_{MN}$$

Half-BPS solutions with SO(2,1)XSO(4)XSO(4) isometry in 11-dimensional supergravity [04, Lin-Lunin-Maldacena]

$$ds^{2} = -G_{tt} \left(-dt^{2} + dw_{1}^{2} + dw_{2}^{2} \right) + G_{xx} (dx^{2} + dy^{2}) + G_{\theta\theta} ds_{S^{3}/\mathbb{Z}_{k}}^{2} + G_{\tilde{\theta}\tilde{\theta}} ds_{\tilde{S}^{3}/\mathbb{Z}_{k}}^{2}$$

$$-G_{tt} = \left(\frac{4\mu_0^2 y \sqrt{\frac{1}{4} - z^2}}{f^2}\right)^{2/3},$$

$$G_{xx} = \left(\frac{f \sqrt{\frac{1}{4} - z^2}}{2\mu_0 y^2}\right)^{2/3},$$

$$Z(x, y) = \sum_{i=1}^{2m+1} \frac{(-1)^{i+1} (x - x_i)}{2\sqrt{(x - x_i)^2 + y^2}},$$

$$\mathcal{X}$$

$$G_{\theta\theta} = \left(\frac{f y \sqrt{\frac{1}{2} + z}}{2\mu_0 (\frac{1}{2} - z)}\right)^{2/3},$$

$$\mathcal{X}$$

$$C(x, y) = \sum_{i=1}^{2m+1} \frac{(-1)^{i+1} (x - x_i)}{2\sqrt{(x - x_i)^2 + y^2}},$$

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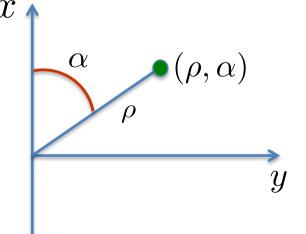
$$C(x, y) = \sum_{i=1}^{2m+1} \frac{(-1)^{i+1}}{2\sqrt{(x - x_i)^$$

$$G_{\tilde{\theta}\tilde{\theta}} = \left(\frac{fy\sqrt{\frac{1}{2}-z}}{2\mu_0\left(\frac{1}{2}+z\right)}\right)^{2/3}.$$

$$f(x,y) = \sqrt{1 - 4z^2 - 4y^2V^2}$$

$$z(x,y) = \sum_{i=1}^{2m+1} \frac{(-1)^{i+1}(x-x_i)}{2\sqrt{(x-x_i)^2 + y^2}},$$

$$V(x,y) = \sum_{i=1}^{2m+1} \frac{(-1)^{i+1}}{2\sqrt{(x-x_i)^2 + y^2}}$$



To calculate the HEE in the LLM geometry we have to solve

PDE equations

HEE proposal [Ryu-Takayanagi 06]

$$S_A = \frac{\operatorname{Min}(\gamma_A)}{4G_N}, \quad \gamma_A = \int d^9 \sigma \sqrt{\det g_{ij}}, \qquad i = 1, 2, \cdots, 9$$

Induced metric

$$g_{ij} = \frac{\partial X^M \partial X^N}{\partial \sigma^i \partial \sigma^j} G_{MN}$$

Mapping for the Disk case:

$$w_1 = \sigma_2 \cos \sigma_1, \quad w_2 = \sigma_2 \sin \sigma_1, \quad u = u(\sigma_2, \sigma_3), \quad \alpha = \alpha(\sigma_2, \sigma_3),$$

 $\sigma_4 \sim \sigma_9 = \theta_1 \sim \theta_6.$

$$\sigma_2 = \rho, \ \sigma_3 = \phi$$

$$w_1 = \rho \cos \sigma^1$$
, $w_2 = \rho \sin \sigma^1$, $u = u(\rho, \phi)$, $\alpha = \alpha(\rho, \phi)$

[C.Kim, K.Kim, OK work in progress]

$$u = u_0(\tilde{\rho}) + u_1(\tilde{\rho}, \phi)\mu_0 l + u_2(\tilde{\rho}, \phi)(\mu_0 l)^2 + \cdots,$$

$$\tilde{g}(\tilde{u}, \phi) = 1 + g_1(\phi)u_0(\mu_0 l) + (g_1(\phi)u_1 + g_2(\phi)u_0^2)(\mu_0 l)^2 + \cdots$$

$$\begin{split} g_1(\phi) &= D_1 \cos \phi, \\ g_2(\phi) &= D_2 + D_3 \cos(2\phi) \\ D_0 &= \sqrt{2} \sqrt{C_2 - C_1^2} \;,\; D_1 = -\frac{(C_1 C_2 - C_3)}{\sqrt{2}} \qquad \qquad C_k = \sum_{i=1}^{2N_b + 1} (-1)^{i+1} \left(\frac{\hat{x}_i}{\sqrt{Nk}}\right)^k \\ D_2 &= \frac{1}{16} \left(-5C_2^2 - 2\left(C_3 - C_1C_2\right)^2 - 4C_1C_3 + 9C_4\right) \;, \qquad x_i = 2\pi l_p^3 \mu_0 \hat{x}_i. \\ D_3 &= \frac{1}{16} \left(-3C_2^2 - 2\left(C_3 - C_1C_2\right)^2 - 12C_1C_3 + 15C_4\right) \end{split}$$

$$z(x,y) = \sum_{i=1}^{2m+1} \frac{(-1)^{i+1}(x-x_i)}{2\sqrt{(x-x_i)^2 + y^2}}, \qquad V(x,y) = \sum_{i=1}^{2m+1} \frac{(-1)^{i+1}}{2\sqrt{(x-x_i)^2 + y^2}}.$$

Solution satisfying the boundary condition

$$\dot{u}(\rho = 0) = 0 \text{ and } u(\rho = 1) = 0$$

$$u_0(\rho) = \sqrt{1 - \rho^2}$$

$$u_1(\phi,\phi) = -\frac{D_1}{2} \left(1 - \rho^2\right) \cos \phi$$

$$u_2(\rho,\phi) = -\frac{1}{6\sqrt{1-\rho^2}} \left(D_1^2 + 20D_2 - 12D_3 \right) \log \left(1 + \sqrt{1-\rho^2} \right)$$

$$+ \frac{1}{48} \left[\left(8 + (9 - 13\rho^2)\sqrt{1-\rho^2} \right) D_1^2 + 16 \left(10 - (6 - \rho^2)\sqrt{1-\rho^2} \right) D_2 \right]$$

$$- 48 \left(2 - \sqrt{1-\rho^2} \right) D_3 + \frac{1}{48} (11D_1^2 - 16D_3) (1 - \rho^2)^{3/2} \cos(2\phi).$$

Entanglement entropy

$$S_{\text{disk}} = \frac{\pi^5 R^9}{24G_N k} \left\{ \frac{l}{\epsilon} - 1 - \mu_0^2 l^2 \left[\frac{4}{3} + \frac{1}{24} (C_3 - 3C_1 C_2 + 2C_1^3)^2 \right] \right\} + \mathcal{O}(\mu_0^3)$$

Renormalized Entanglement entropy

$$\begin{split} \mathcal{F}_{\text{disk}} &\equiv \left(l\frac{\partial}{\partial l} - 1\right) S_{\text{disk}} \\ &= \frac{\pi^5 R^9}{24 G_N k} \left\{1 - \mu_0^2 l^2 \left[\frac{4}{3} + \frac{1}{24} (C_3 - 3C_1 C_2 + 2C_1^3)^2\right]\right\} + \mathcal{O}(\mu_0^3) \end{split}$$

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Positive for all LLM solutions

→ Monotonically decreasing behavior of c-function

 Entanglement entropy of the LLM geometry in 11d SUGRA (RT formula)

$$S_{\text{disk}} = \frac{\pi^5 R^9}{24G_N k} \left\{ \frac{l}{\epsilon} - 1 - \mu_0^2 l^2 \left[\frac{4}{3} + \frac{1}{24} (C_3 - 3C_1 C_2 + 2C_1^3)^2 \right] \right\} + \mathcal{O}(\mu_0^3)$$

Exact holography for the vev of CPO

$$\langle \mathcal{O}^{(1)} \rangle_0 = \frac{N^{\frac{3}{2}} \mu_0}{3\sqrt{2}\pi} \left(C_3 - 3C_1 C_2 + 2C_1^3 \right)$$

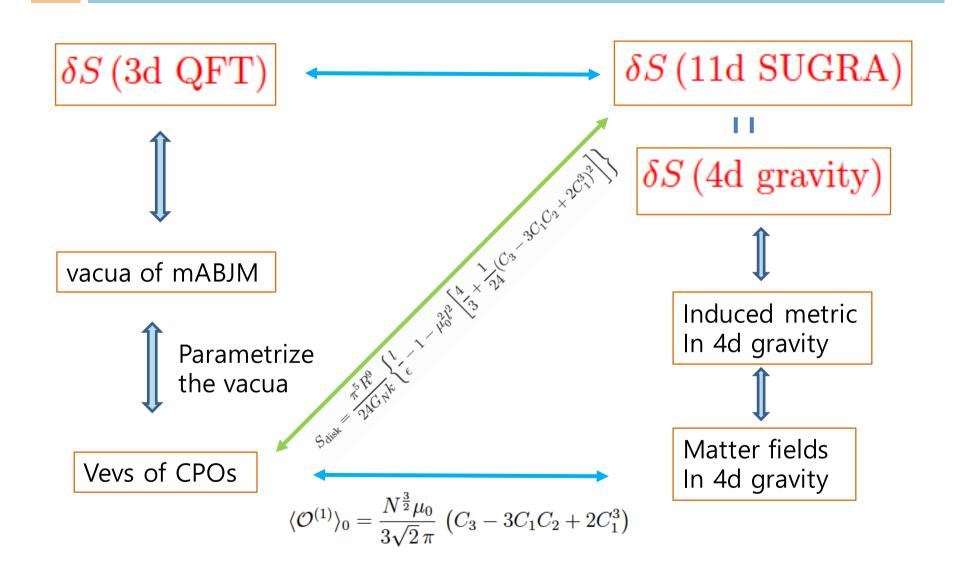
Entanglement entropy of the vacua in mABJM theory(RT formula)

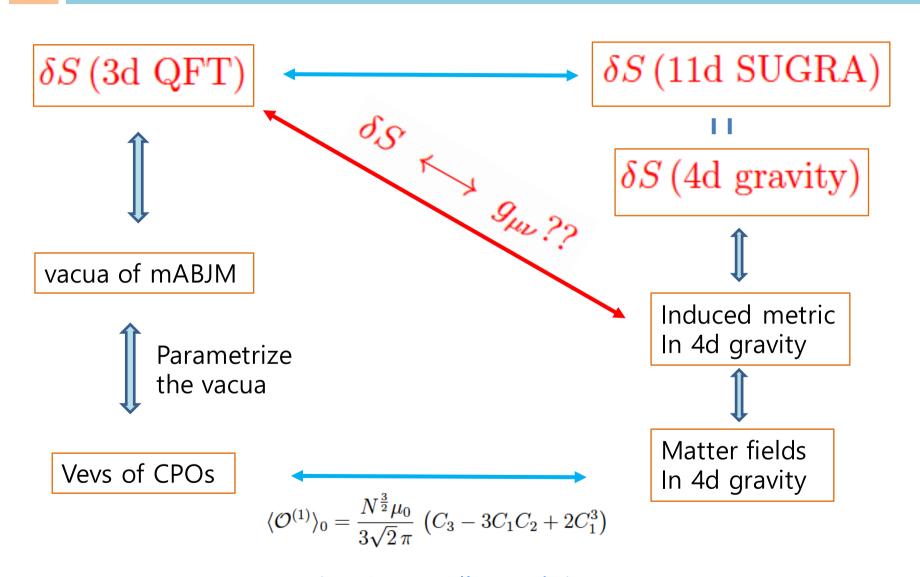
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Exact holography for the vev of CPO

$$\langle \mathcal{O}^{(1)} \rangle_0 = \frac{N^{\frac{3}{2}}\mu_0}{3\sqrt{2}\pi} \left(C_3 - 3C_1C_2 + 2C_1^3 \right)$$

How can we consider this correspondence?





[D. Jang, Y. Kim, OK, D.Tolla, work in progress]