

Role of Dense Matter in Supernova

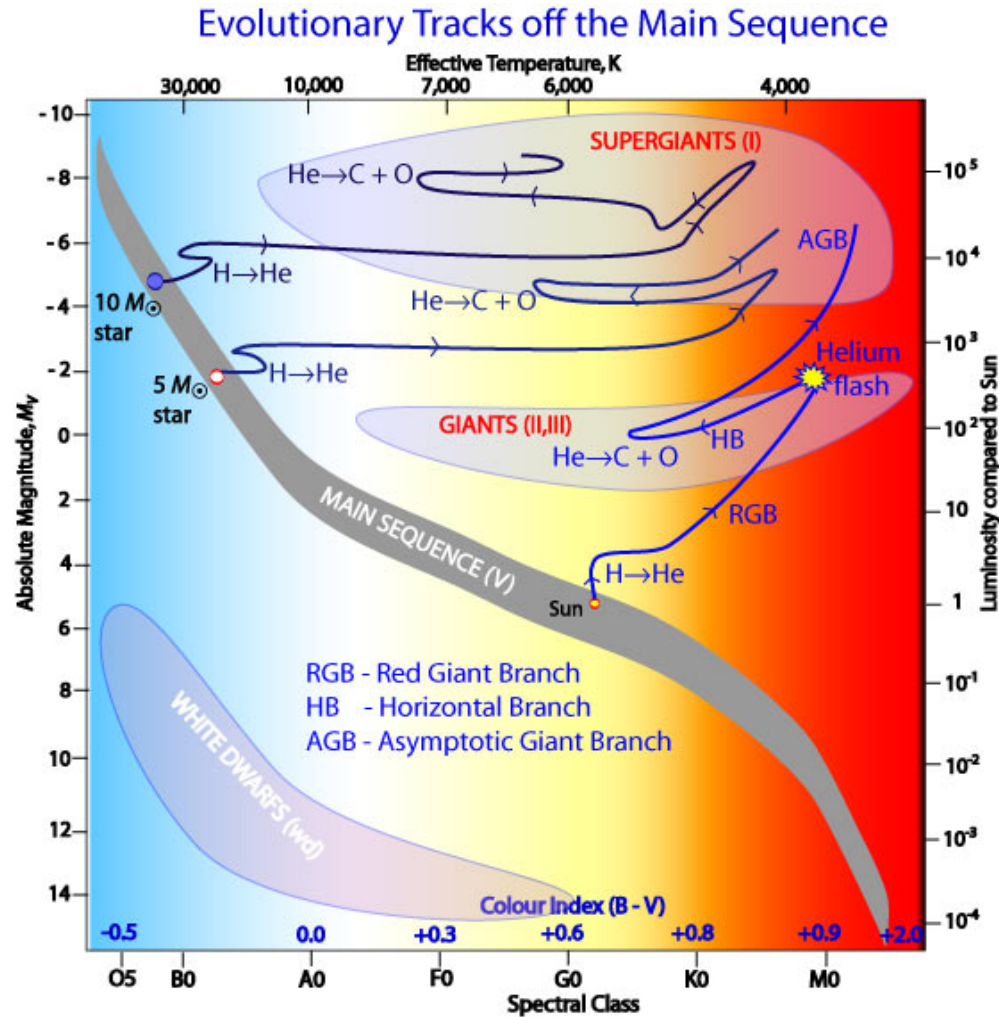
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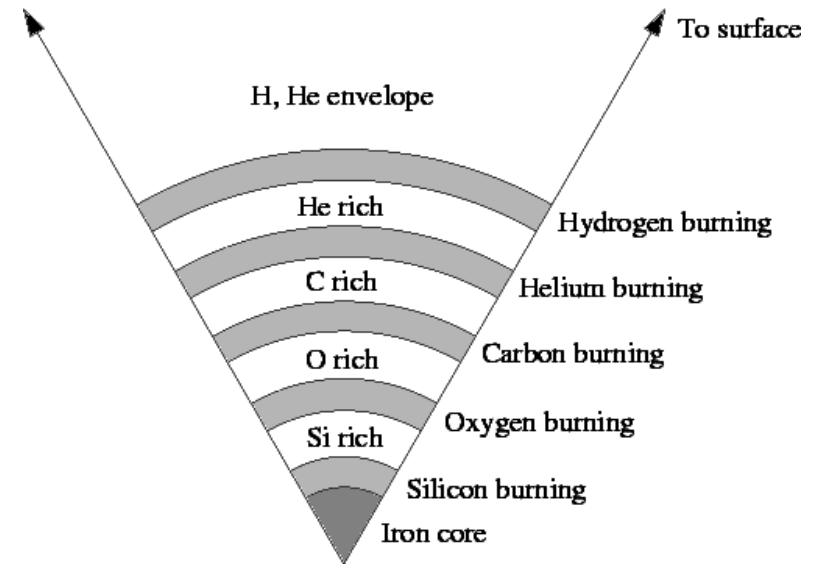
December 1, 2018

APCTP HaPhy Meeting @Seoul

Stellar Evolution



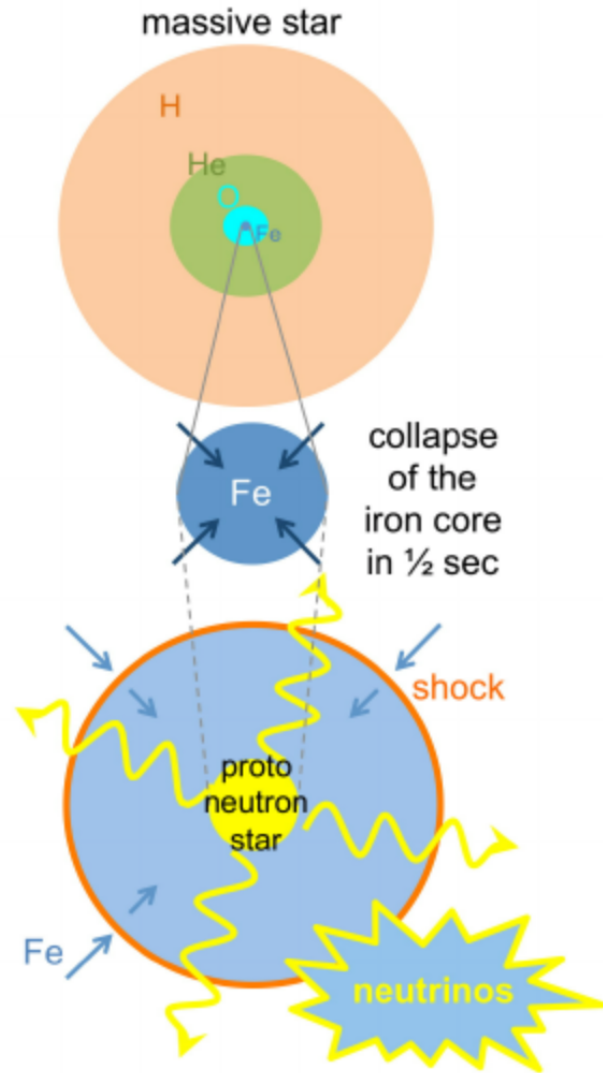
www.atnf.csiro.au



Pre-supernova at the end of the massive star evolution

<http://burro.astr.cwru.edu>

On the Way to Supernova Explosion

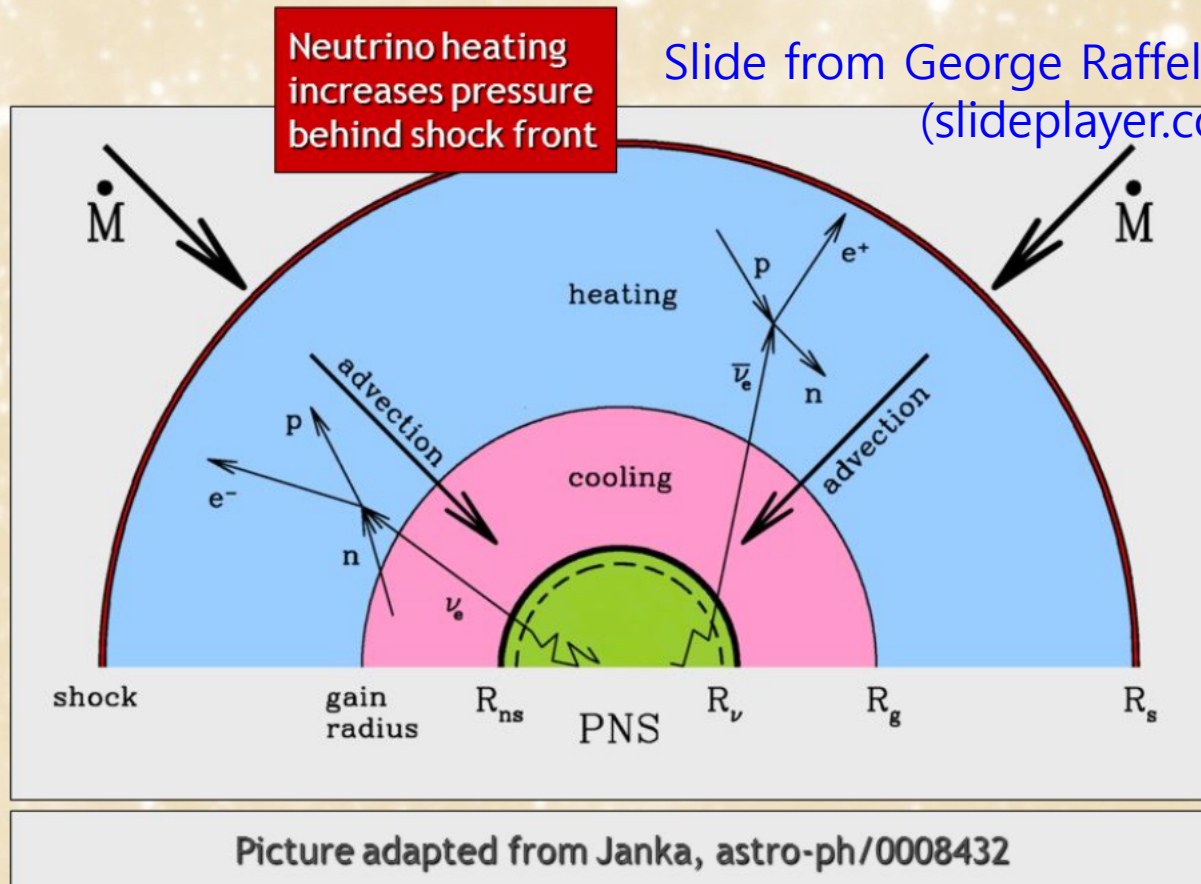


- ✓ Inner core (about $0.5 M_{\odot}$) contracts homologously.
- ✓ Size of inner core weakly depends on the pre-SN structure.
- ✓ Outer core falls supersonically.
- ✓ Central region exceeds nuclear saturation density, which leads to bounce depending of equation of state. Alternatively, it collapses into a black hole.
- ✓ Bouncing results in shock wave that forms near the edge of inner core.

(From Introduction of Pejcha & Thomson 2015. See the references therein.)

Figure from www.researchgate.net.

Neutrino-Driven Delayed Explosion



- ✓ Shock wave is stalled by losing energy to dissociate iron.
- ✓ Standing accretion shock forms through the balance between neutrino emission from proto-neutron star (PNS) and infalling matter from outer core. (From Introduction of Pejcha & Thomson 2015. See the references therein.)

Uncertainties in Supernova Simulations

- Uncertain physics
 - Neutrino self-interactions including oscillation at high luminosity
 - Equation of state for dense (nuclear/quark) matter that affects the emissivity and opacity of neutrino transport
 - Other mechanisms (e.g., rotational energy)
- Technical (numerical) challenges
 - Proper implementation of general relativity
 - 3D simulations with neutrino radiative transfer
 - Other effects (e.g., inclusion of magnetic field)

EoS for SN Simulations

- EoS for (cold) neutron star
 - At beta equilibrium and $T \approx 0$
- A wider range (general purpose EoS)
 - $0 \leq T \leq 100 \text{ MeV}$
 - $10^4 \text{ g/cm}^3 \leq \rho \leq 10^{15} \text{ g/cm}^3$
 - $0 \leq Y_e \leq 0.6$
- Needs to cover gaseous nuclei to uniform nuclear matter
- Mixture of nuclei and nucleon
 - Phase transition

Handling the Mixture

- Thomas–Fermi approximation
 - Nucleus at the center (body-centered cubic) to minimize the Coulomb lattice energy
 - Wigner–Seitz cell: approximation with sphere

$$n_i(r) = \begin{cases} (n_i^{\text{in}} - n_i^{\text{out}}) \left[1 - \left(\frac{r}{R_i}\right)^{t_i}\right]^3 + n_i^{\text{out}}, & 0 \leq r \leq R_i, \\ n_i^{\text{out}}, & R_i \leq r \leq R_{\text{cell}}, \end{cases}$$

$$n_\alpha(r) = \begin{cases} -n_\alpha^{\text{out}} \left[1 - \left(\frac{r}{R_p}\right)^{t_p}\right]^3 + n_\alpha^{\text{out}}, & 0 \leq r \leq R_p, \\ n_\alpha^{\text{out}}, & R_p \leq r \leq R_{\text{cell}}, \end{cases}$$