

Parallel tempering algorithm for integration over Lefschetz thimbles

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@APCTP, 22 Sep 2017

“Discrete approaches to the Dynamics of Fields and Space-Time”

based on : [PTEP 2017 (2017) no.7, 073B01 (arXiv:1703.00861)]

and work in progress

Motivation: sign problem

Monte Carlo calculation (real action)

$$\begin{aligned}\langle \mathcal{O}(x) \rangle &\equiv \frac{1}{Z} \int dx \mathcal{O}(x) e^{-S(x)} && \left(Z \equiv \int dx e^{-S(x)} \right) \\ &= \int dx p(x) \mathcal{O}(x) && \left(p(x) \equiv \frac{1}{Z} e^{-S(x)} \right) \\ &\simeq \frac{1}{N} \sum_{i=1}^N \mathcal{O}(x_i)\end{aligned}$$

But, if $S(x)$ is complex,
 $p(x)$ can no longer be regarded as probability distribution.

Reweighting

$$\langle \mathcal{O}(x) \rangle = \frac{\langle e^{-i\text{Im}S(x)} \mathcal{O}(x) \rangle_{\text{Re}S(x)}}{\langle e^{-i\text{Im}S(x)} \rangle_{\text{Re}S(x)}}$$

$$\langle \mathcal{O}(x) \rangle_{S_R} \equiv \frac{1}{Z_R} \int dx \mathcal{O}(x) e^{-\text{Re}S(x)}$$

(local) sign problem:
 $e^{-i\text{Im}S(x)}$ oscillates if $|\text{Im}S(x)|$ become large.

⇒ complex Langevin method, Lefschetz thimble method, ...

Lefschetz thimble method (1/4)

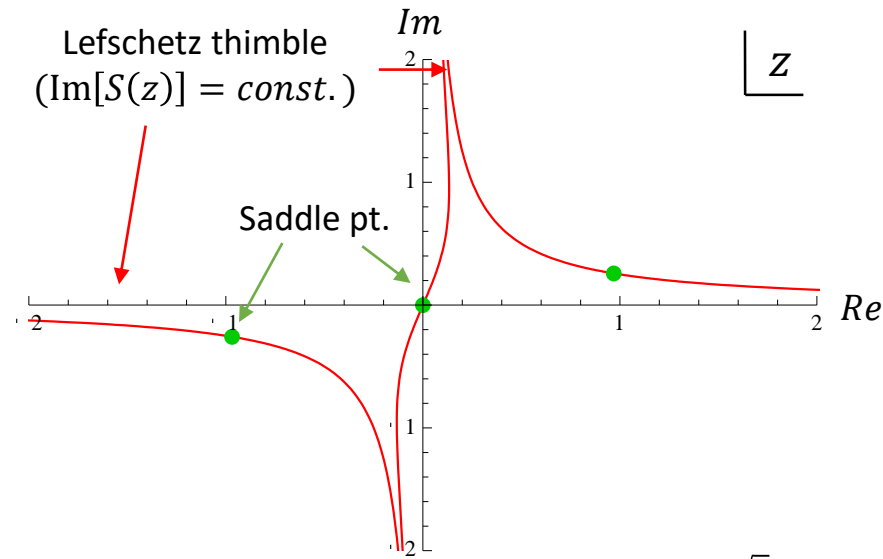
Idea: take better integration contours in complex plane

$$x = (x_i) \in \mathbb{R}^N \rightarrow z = (z_i) \in \mathbb{C}^N$$

Lefschetz thimble [Cristoforetti *et. al.* 2012]

$$\left\{ z(t) \left| \frac{\partial z(t)}{\partial t} = \frac{\overline{\partial S(z(t))}}{\partial z}, z(t = -\infty) = z^{(0)} \right. \right\}$$

- N -dimensional submanifold in \mathbb{C}^N
- Defined by anti-holomorphic gradient flow
- $\text{Im}S(z) = \text{const.}$ on each thimble



Ex.) $S(z) = -\frac{c}{2}z^2 + \frac{1}{4}z^4$ ($c = \frac{\sqrt{3}}{2} + \frac{1}{2}i$)

Transform the integration contour:

$$Z = \int_{\mathbb{R}^N} dx e^{-S(x)}$$

$$= \sum_{\sigma} n_{\sigma} \int_{J_{\sigma}} dz e^{-S(z)}$$

(J_{σ} : Lefschetz thimble, n_{σ} : intersection number)

$$= \sum_{\sigma} n_{\sigma} e^{-i\text{Im}S(z_{\sigma})} \int_{J_{\sigma}} dz e^{-\text{Re}S(z)}$$

\Rightarrow the sign problem can be avoided.

Lefschetz thimble method (2/4)

[Alexandru et. al. (2016)]

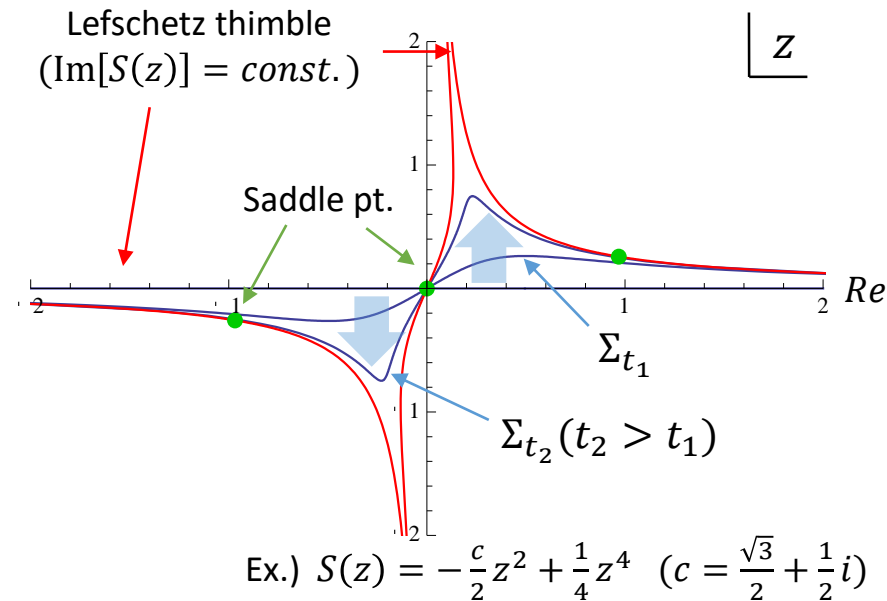
We can use the gradient flow to change the variables:

$$\frac{\partial z_i(x; t)}{\partial t} = \frac{\overline{\partial S(z(x; t))}}{\partial z_i}$$

$$\frac{\partial J_{ij}(x; t)}{\partial t} = \frac{\partial^2 S(z(x; t))}{\partial z_i \partial z_k} J_{kj}(x; t)$$

$$z_i(x; t = 0) = x_i, \quad J_{ij}(x; t = 0) = \delta_{ij}$$

$$(J_{ij} = \frac{\partial z_i}{\partial x_j} : \text{Jacobian})$$



$$Z = \int_{\mathbb{R}^N} dx e^{-S(x)}$$

$$= \int_{\Sigma_t} dz e^{-S(z)}$$

$$= \int_{\mathbb{R}^N} dx \det J(x; t) e^{-S(z(x; t))}$$

$$(S_{\text{eff}}(x; t) \equiv \text{Re}S(z(x; t)) - \log|\det J|)$$

$$= \int_{\mathbb{R}^N} dx e^{-S_{\text{eff}}(x; t)} e^{i[\arg \det J - \text{Im}S(z(x; t))]} \quad : \text{effective action + residual phase}$$

← No need to compute intersection numbers

Lefschetz thimble method (3/4)

$$Z = \int_{\mathbb{R}^N} dx e^{-S_{\text{eff}}(x;t)} e^{i[\arg \det J - \text{Im} S(z(x;t))]}$$

Effective action

Incorporate into operators (reweighting)



Correlation functions can be estimated by using S_{eff} :

$$\langle \mathcal{O}(x) \rangle = \frac{1}{Z} \int dx e^{-S(x)} = \frac{\langle e^{i[\arg \det J - \text{Im} S(z(x;t))]} \mathcal{O}(z(x;t)) \rangle_{\text{eff}}}{\langle e^{i[\arg \det J - \text{Im} S(z(x;t))]} \rangle_{\text{eff}}}$$

This is t -independent.

⇒ The sign problem can be avoided for sufficiently large t .

((empirically) the residual phase do not cause additional problems.)

Lefschetz thimble method (3/4)

However, there are some problems.

- For large t , configurations will be trapped on one thimble.

➔ We can avoid this by introducing tempering algorithm.

[Fukuma-NU (2017)], [Alexandru *et. al.* (2017)]

- If $\frac{d^2 S}{dx^2} = 0$ (Hesse matrix have 0 eigenvalues) on saddle points, Lefschetz thimble is not well-defined.

(i.e. the thimbles can not be regarded as an N -dimensional manifold.)

However, the antiholomorphic gradient flow is well-defined even in this case.

➔ We can use LTM without any change.

[Fukuma-NU (work in progress)]

Plan of talk

- Introduction
- Parallel tempering algorithm for LTM
 - Simulated tempering
 - Parallel tempering
 - Parallel tempering for LTM
- Examples
 - (0+1)D massive Thirring model
 - One matrix model with quartic potential
- Summary

Multimodal problem

(Before reviewing the parallel tempering, I review the simulated tempering method.)

Multimodal problem

If the action have potential barriers, the configuration is trapped at a local minimum.

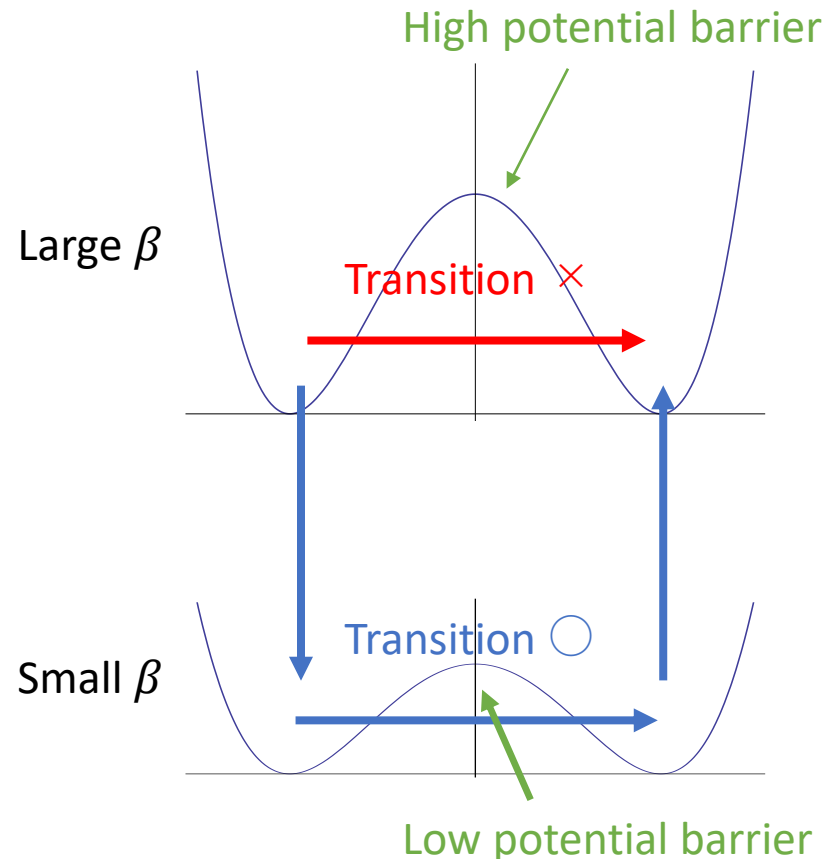
The idea of the tempering algorithm

Expand configuration space:

$$x \rightarrow X \equiv (x, \beta)$$

→ we can evade the potential barrier passing through the configs with smaller β .

$$(S(x) = \beta V(x))$$



Simulated tempering (2/2)

Simple approach: simulated tempering

Expand configuration space: $x \rightarrow X \equiv (x, \beta)$

(i.e. regard β as an additional dynamical variable.)

where, β can take discrete variables: $\beta \in \{\beta_0, \beta_1, \dots, \beta_A\}$ ($\beta_0 > \beta_1 > \dots > \beta_A$)

(\Rightarrow total equilibrium distribution: $\mathcal{P}_{eq}(x, a) = \frac{1}{A+1} e^{-S(x; \beta_a)} / Z(\beta_a)$)


Algorithm

Define two transition matrix:

$$\left\{ \begin{array}{l} P_1: X = (x, \beta_a) \rightarrow X' = (x', \beta_a) \text{ (ordinary transition matrix for fixed } \beta = \beta_a) \\ P_2: X = (x, \beta_a) \rightarrow X' = (x, \beta_{a'}) \\ \text{Accept/reject with probability } \min\{1, \mathcal{P}_{eq}(x, a') / \mathcal{P}_{eq}(x, a)\} \end{array} \right.$$

and apply P_1 and P_2 alternately. $\left(\text{Finally we calculate correlation functions by using a subset of configurations } (x, \beta = \beta_0) . \right)$

Problem

We must (at least roughly) estimate $Z(\beta_a)$ a priori.  *Parallel tempering*

Parallel tempering

Parallel tempering

Consider replicas of configuration spaces: $X = \{x_a\}_{a=0,1,\dots,A}$
The configuration of replica a explores with β_a
and be exchanged among replicas at fixed intervals.

Algorithm

Define two transition matrix:

$$\left\{ \begin{array}{l} P_1: X = (x, \beta_a) \rightarrow X' = (x', \beta_a) \text{ (ordinary transition matrix for fixed } \beta = \beta_a) \\ P_2: \text{exchange configurations between replica } a_1 \text{ and } a_2 \\ \quad \text{with probability } \min\{1, e^{-\Delta S}\} \\ \quad \quad (\Delta S \equiv S(x_{a_1}, \beta_{a_1}) + S(x_{a_2}, \beta_{a_2}) - S(x_{a_1}, \beta_{a_2}) - S(x_{a_2}, \beta_{a_1})) \end{array} \right.$$

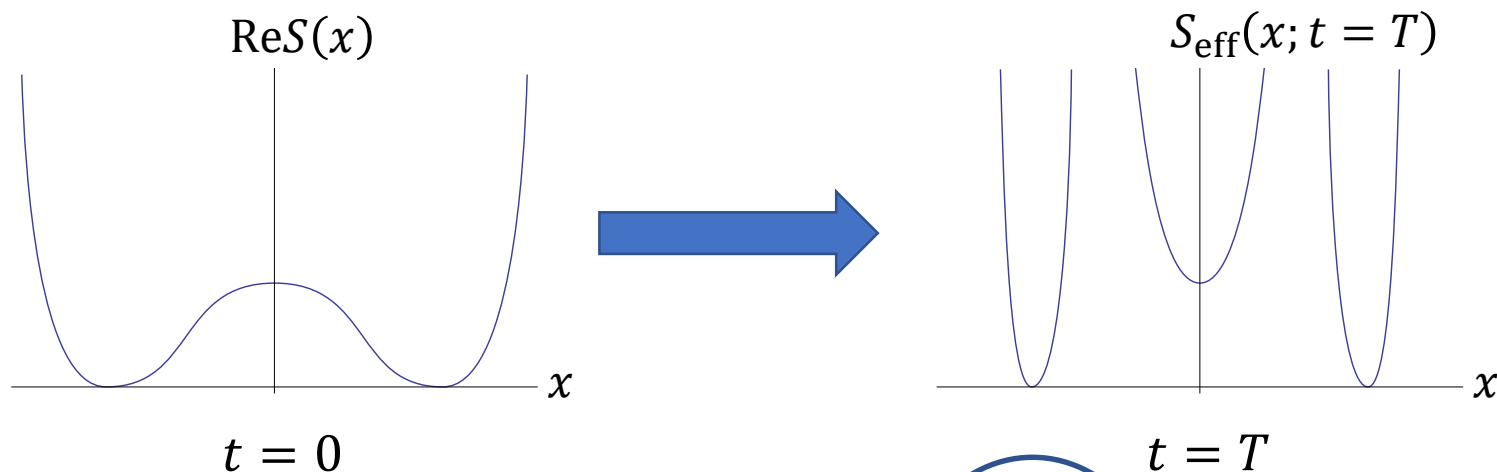
and apply P_1 and P_2 alternately.

We need not to estimate $Z(\beta_a)$.
(\Rightarrow All we need is to choose the set of $\{\beta_a\}$.)

Parallel tempering for LTM (1/2)

Multimodal problem in Lefschetz Thimble Method

If we take the flow time ($= T$) large, there are infinitely (or exponentially) high potential barrier in $S_{\text{eff}}(x; t = T)$:



Flow time ($= T$)	large	small	middle
Sign problem	○	×	△
Multimodal problem	×	○	△

← Previous approach
 [Alexandru *et. al.* (2016)]
Need fine-tuning


Parallel tempering for LTM (2/2)


Our approach:

We introduce parallel tempering algorithm

by regarding the flow time t as a tempering parameter.

Consider several number of flow time $t_a (t_0 = T > t_1 > \dots > t_A = 0)$
 \Rightarrow parallel tempering

 No need to fine-tune. (All we need is to take sufficiently large T (and A).)



Flow time ($= T$)	large	small	middle	Parallel tempering
Sign problem	○	×	△	○
Multimodal problem	×	○	△	○

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(0+1)D massive Thirring model (1/3)

Example 1 : (0+1)-dimensional massive Thirring model at finite density

This model is defined from (1+1)-dimensional massive Thirring model by dimensional reduction:

$$Z = \int_{\text{PBC}} [d\phi(\tau)] \int_{\text{ABC}} [d\bar{\psi}(\tau)d\psi(\tau)] e^{-S[\phi, \bar{\psi}, \psi]}$$

$$S[\phi, \bar{\psi}, \psi] = \int_0^\beta d\tau [-\bar{\psi}(\gamma^0(\partial_0 + i\phi + \mu) + m)\psi + \frac{1}{2g^2} \phi^2] \quad \left(\gamma^0 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right)$$

 Discretize: $\phi(\tau) \rightarrow \phi_n$ ($n = 1, \dots, N$)

$$Z = \int d\phi e^{-S(\phi)}$$

$$S(\phi) = \frac{1}{2g^2} \sum_n \left[1 - \frac{1}{2} (e^{i\phi_n} + e^{-i\phi_n}) \right] - \log \det D$$

$$D_{nl}(\phi) = \frac{1}{2} (e^{i\phi_n + \mu} \delta_{n+1,l} - e^{-i\phi_{n-1} - \mu} \delta_{n-1,l} - e^{i\phi_N + \mu} \delta_{n,N} \delta_{l,1} + e^{-i\phi_N - \mu} \delta_{n,1} \delta_{l,N}) + m\delta_{n,l}$$

- Sign problem occurs for $\mu \neq 0$
- Analytically solvable even for finite N
- More than two thimbles can have non-negligible contributions to Z

(0+1)D massive Thirring model (2/4)

This model can be solved analytically even for finite N:

$$Z = \int d\phi e^{-S(\phi)}$$

$$S(\phi) = \frac{1}{2g^2} \sum_n \left[1 - \frac{1}{2} (e^{i\phi_n} + e^{-i\phi_n}) \right] - \log \det D$$

$$D_{nl}(\phi) = \frac{1}{2} (e^{i\phi_{n+\mu}} \delta_{n+1,l} - e^{-i\phi_{n-1}-\mu} \delta_{n-1,l} - e^{i\phi_{N+\mu}} \delta_{n,N} \delta_{l,1} + e^{-i\phi_{N-\mu}} \delta_{n,1} \delta_{l,N}) + m \delta_{n,l}$$



$$Z = \frac{e^{-N\alpha}}{2^{N-1}} [\cosh(N\mu) I_1^N(\alpha) + \rho_+ I_0^N(\alpha)]$$

$$\left(\alpha \equiv \frac{1}{2g^2}, \quad I_n: \text{modified Bessel function of 1st kind} \right)$$

$$2\rho_+ \equiv (\sqrt{m^2 + 1} + m)^N + (\sqrt{m^2 + 1} - m)^N$$

Chiral condensation

$$\langle \bar{\chi} \chi \rangle = \left\langle \frac{1}{N} \text{tr} D^{-1}(U) \right\rangle$$

$$= \frac{\rho_- I_0^N(\alpha)}{\sqrt{m^2 + 1} [\cosh(N\mu) I_1^N(\alpha) + \rho_+ I_0^N(\alpha)]}$$

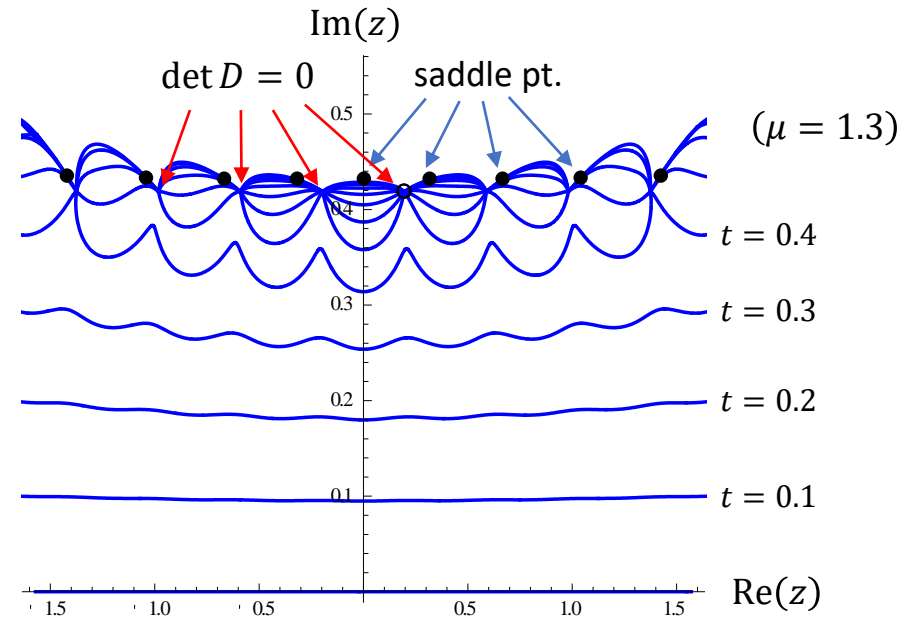
➡ We calculate the chiral condensation numerically by using LTM.

(0+1)D massive Thirring model (3/4)

We calculate the region $0 < \mu < 2$
with $N = 8$, $g^2 = \frac{1}{2}$, and $m = 1$.

The integration contours stick to thimbles
when $T \gtrsim 1$.

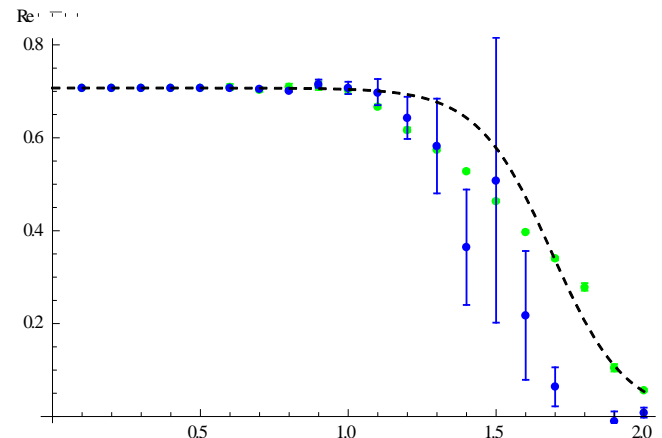
➡ We set $T = 2$.



The main contribution is that from the saddle point $\text{Re}(z) \equiv \text{Re}\left(\frac{1}{N} \sum_n z_n\right) = 0$.
However, the neighboring saddle points have non-negligible contributions.

⇒ ordinary algorithm gives wrong results.

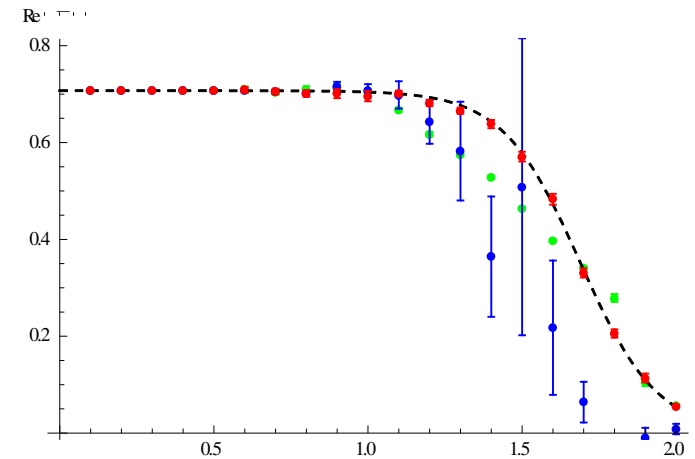
blue: $T=0$ (reweighting)
green: $T=2$ (without PT)
dotted: analytic solution



(0+1)D massive Thirring model (4/4)

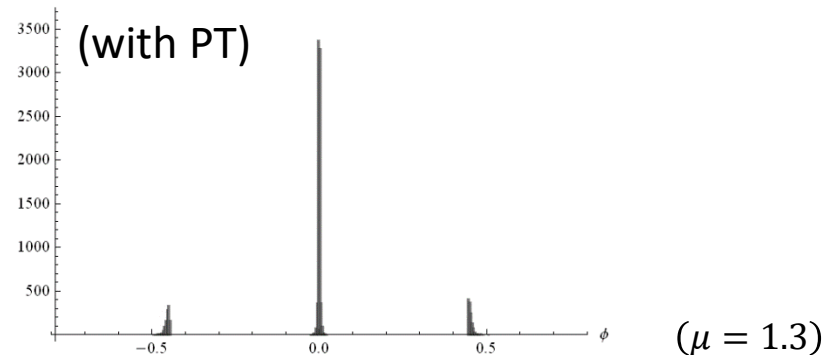
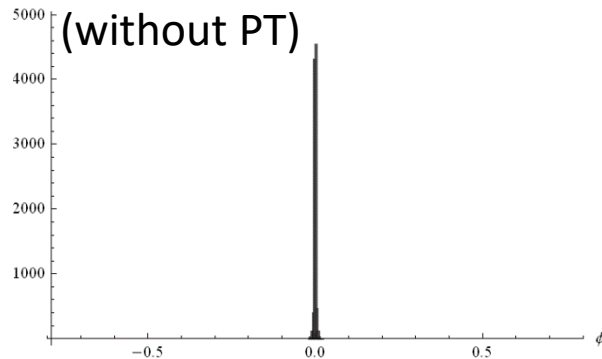
We set $\{t_a\} = \{t_0 = 0, t_1 = 0.1, \dots, t_{19} = 1.9, t_{20} = 2.0\}$ and introduce parallel tempering algorithm.

blue: T=0 (reweighting)
 green: T=2 (without PT)
 red: T=2 (with PT)
 dotted: analytic solution



⇒ The results agree with the analytic solution.

The histogram of $\phi = \frac{1}{N} \sum_n \phi_n$



Several thimbles contribute correctly.

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One matrix model

Example 2 (one-matrix model with quartic potential)

$$S(M) = N\beta \operatorname{tr}V(M)$$

$$V(x) = -\frac{c}{2}x^2 + \frac{1}{4}x^4 \quad (M: N \times N \text{ Hermitian matrix, } c = e^{i\theta})$$

- sign problem occurs if $c \in \mathbb{C}$.
- Analytically solvable in $N \rightarrow \infty$ limit.
- Continuous symmetry ($SU(N)$)  Lefschetz thimble is not well-defined.

However, antiholomorphic gradient flow is well-defined even without Lefschetz thimble.

⇒ LTM can work in this case.

Furthermore, there can be a 3-cut solution for $c \in \mathbb{C}$.



The topic of the rest of my talk

Matrix model (real, 1/2)

First I review the real case of this model.

$$S(M) = N\beta \operatorname{tr} V(M)$$

$$V(x) = -\frac{c}{2}x^2 + \frac{1}{4}x^4 \quad (M: N \times N \text{ Hermitian matrix, } c = \pm 1)$$

$N \rightarrow \infty$ limit

- Eigenvalues distribute continuously $\Rightarrow \rho(x)$
- We can solve this model by using the resolvent $\omega(z) \equiv \left\langle \frac{1}{N} \operatorname{tr} \left(\frac{1}{z-M} \right) \right\rangle = \int dx \frac{\rho(x)}{z-x}$

EOM give quadratic equation of $\omega(z)$:

$$\omega(z)^2 - \beta V'(z)\omega(z) - \beta Q(z) = 0$$

$$\Rightarrow \omega(z) = \frac{\beta}{2} V'(z) - \sqrt{\text{polynomial}}$$

$$\left(Q(z) \equiv \left\langle \frac{1}{N} \operatorname{tr} \left(\frac{V'(z) - V'(M)}{z - M} \right) \right\rangle : \text{polynomial of } z \right)$$

$V(z)$: 4th order \Rightarrow (polynomial): 6th order

Matrix model (real, 2/2)

Properties of $\omega(z)$

$$\omega(z) \rightarrow \frac{1}{z} \quad (|z| \rightarrow \infty)$$

$$\omega(x + i\epsilon) - \omega(x - i\epsilon) = 2\pi i \rho(x)$$

← Eigenvalues distribute in the potential $V(x) = -\frac{c}{2}x^2 + \frac{1}{4}x^4$

Two types of solution

$$\left\{ \begin{array}{l} \text{1-cut solution: } \omega(z) = \frac{\beta}{2} \left(V'(z) - \sqrt{\left(z^2 + \frac{a^2}{2} - c \right)^2 (z^2 - a^2)} \right) \quad \left(a^2 \equiv \frac{2}{3}c \left(1 - \sqrt{1 + \frac{\beta}{12c^2}} \right) \right) \\ \text{2-cut solution: } \omega(z) = \frac{\beta}{2} \left(V'(z) - \sqrt{z^2 \left(z^2 - c + \frac{2}{\sqrt{\beta}} \right) \left(z^2 - c - \frac{2}{\sqrt{\beta}} \right)} \right) \end{array} \right.$$

- $c = -1$: only 1-cut solution is realized.
- $c = +1$: 1-cut solution is realized for $\beta < 4$,
2-cut solution is realized for $\beta > 4$.

Especially, 3rd-order phase transition occurs at $\beta = \beta_c = 4$.

Matrix model (complex, 1/4)

Complex case

$$S(M) = N\beta \operatorname{tr}V(M)$$

$$V(x) = -\frac{c}{2}x^2 + \frac{1}{4}x^4 \quad (M: N \times N \text{ Hermitian matrix, } c = e^{i\theta})$$

- The cut of $\omega(z)$ is no longer on the real axis.
- However, we can find the form of resolvent $\omega(z) = \frac{\beta}{2}(V'(z) - \sqrt{\text{polynomial}})$.

6th order polynomial: in general there should be the 3-cut solution.

We can again solve the model by two steps:

1. Assume the number of cuts
2. Determine the form of $\omega(z)$ by $\omega(z) \rightarrow \frac{1}{z}$ ($|z| \rightarrow \infty$).

Matrix model (complex, 2/4)

Complex case

$$S(M) = N\beta \operatorname{tr}V(M)$$

$$V(x) = -\frac{c}{2}x^2 + \frac{1}{4}x^4 \quad (M: N \times N \text{ Hermitian matrix, } c = e^{i\theta})$$

result

$$\text{1-cut solution: } \omega(z) = \frac{\beta}{2} \left(V'(z) - \sqrt{\left(z^2 + \frac{a^2}{2} - c \right)^2 (z^2 - a^2)} \right) \quad \left(a^2 \equiv \frac{2}{3}c \left(1 - \sqrt{1 + \frac{\beta}{12c^2}} \right) \right)$$

$$\text{2-cut solution: } \omega(z) = \frac{\beta}{2} \left(V'(z) - \sqrt{z^2 \left(z^2 - c + \frac{2}{\sqrt{\beta}} \right) \left(z^2 - c - \frac{2}{\sqrt{\beta}} \right)} \right)$$

$$\text{3-cut solution: } \omega(z) = \frac{\beta}{2} \left(V'(z) - \sqrt{z^6 - 2cz^4 + \left(c^2 - \frac{\beta}{4} \right) z^2 - u(\beta, c)} \right)$$



There remains 1 (complex) parameter.

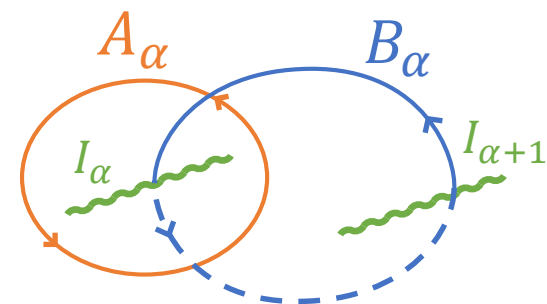
Matrix model (complex, 3/4)

The parameter $u(\beta, c)$ can be determined as following:

David's discussion [David 1992]

The real part of free energy should be minimized with constraint that filling fraction $n_\alpha \in [0,1]$ (especially $n_\alpha \in \mathbb{R}, \sum_\alpha n_\alpha = 1$)

Filling fraction (for cut I_α): $n_\alpha = \int_{A_\alpha} \frac{dz}{2\pi i} \omega(z)$



Minimize $\text{Re}(\text{free energy}) \Leftrightarrow \text{Re}(\text{chemical potential } \mu_\alpha) = \mu$ (α -indep.)

Difference between chemical potentials: $\mu_{\alpha+1} - \mu_\alpha = \int_{B_\alpha} \frac{dz}{2\pi i} \omega(z)$

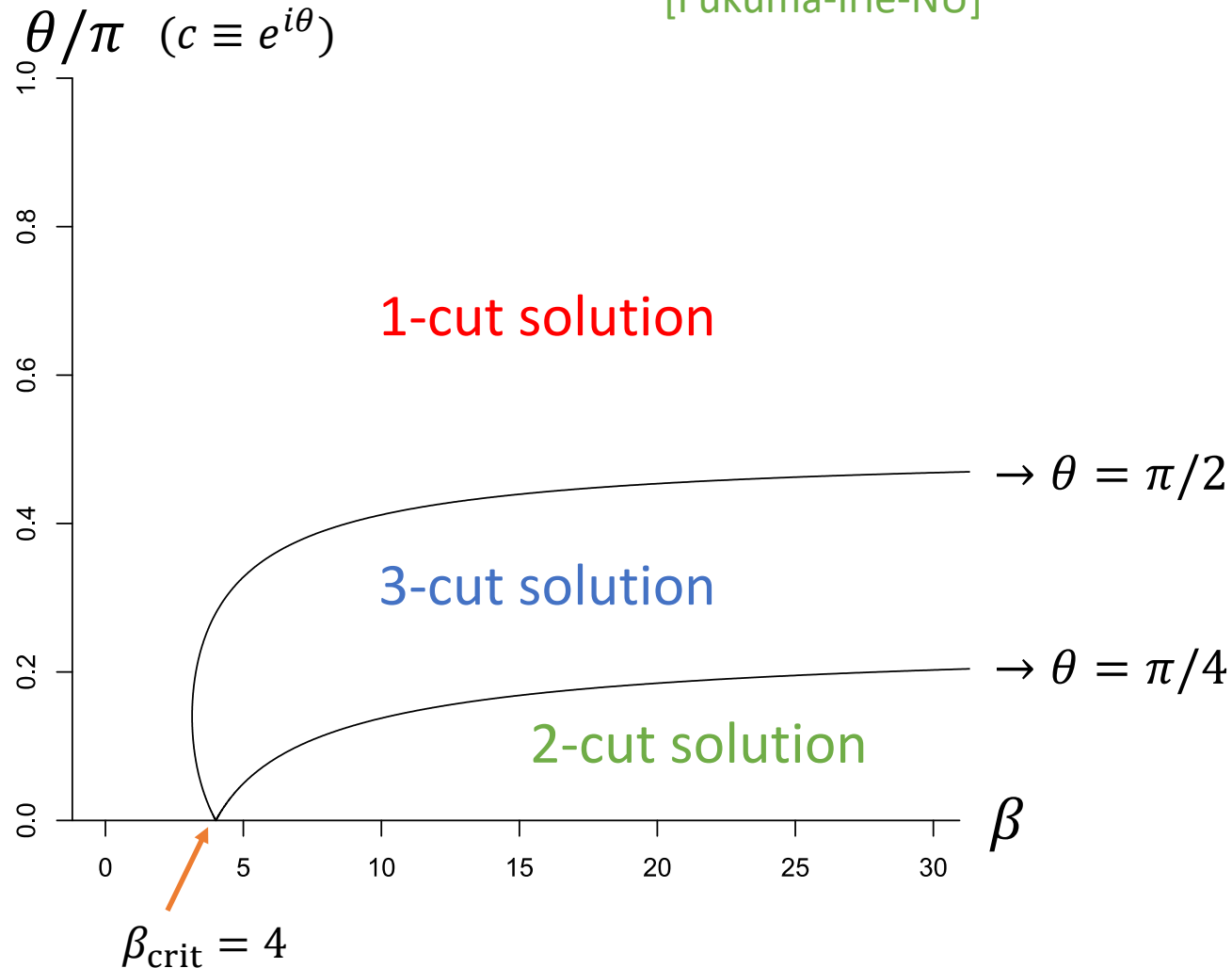
$$\text{Im} \left[\int_{A_\alpha} \frac{dz}{2\pi i} \omega(z) \right] = 0, \quad \text{Re} \left[\int_{B_\alpha} \frac{dz}{2\pi i} \omega(z) \right] = 0$$

➔ $u(\beta, c)$ is determined.

Phase diagram (prospection)

The phase diagram is prospected as follows:

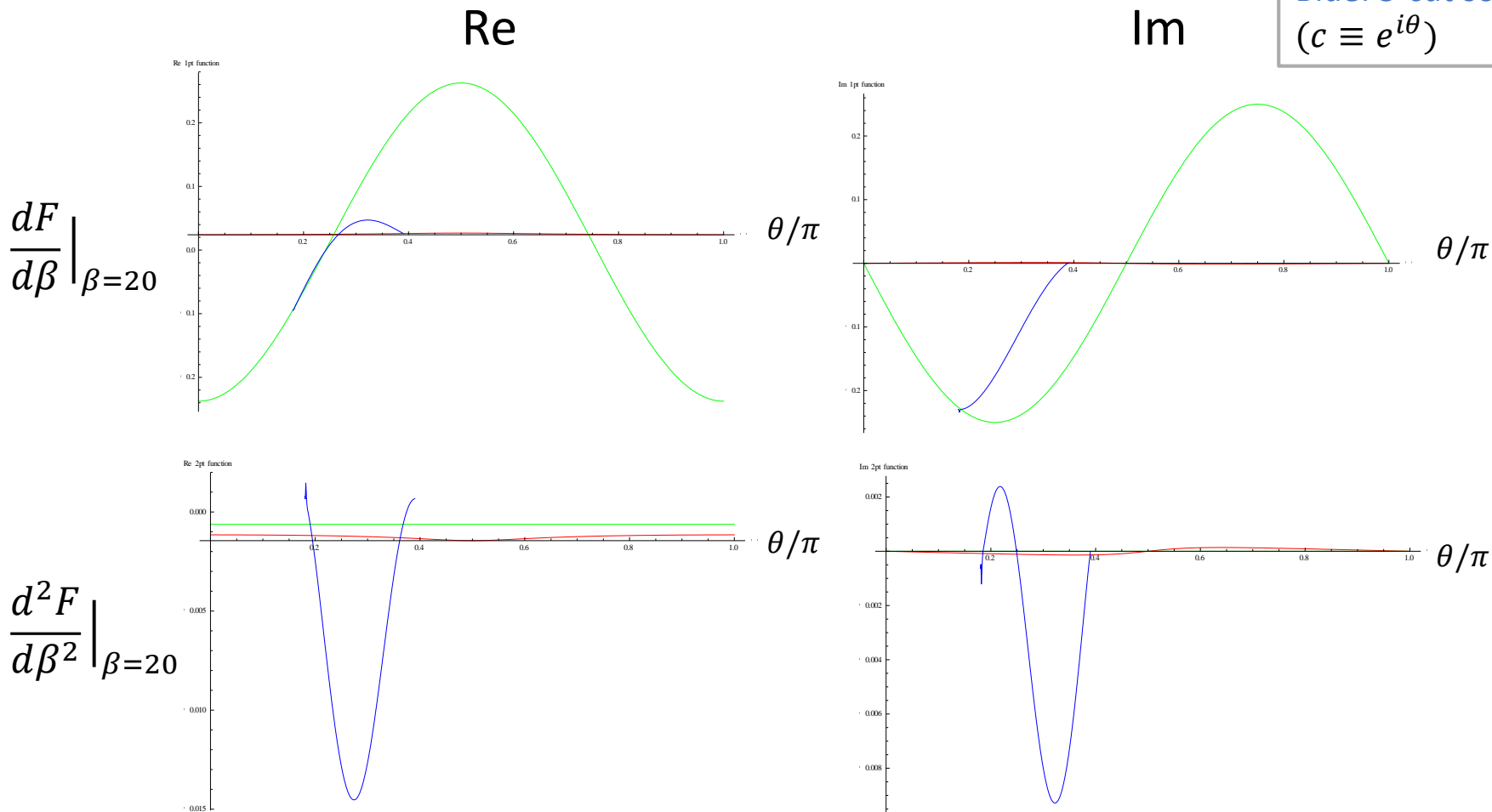
[Fukuma-Irie-NU]



Matrix model (complex, 4/4)

On $\beta = 20$ line, three types of solution should appear.

Red: 1-cut soln.
 Green: 2-cut soln.
 Blue: 3-cut soln.
 ($c \equiv e^{i\theta}$)



In particular, (probably) 3rd-order phase transitions occur at $\theta \sim 0.2\pi, 0.4\pi$.

Numerical results (1/2)

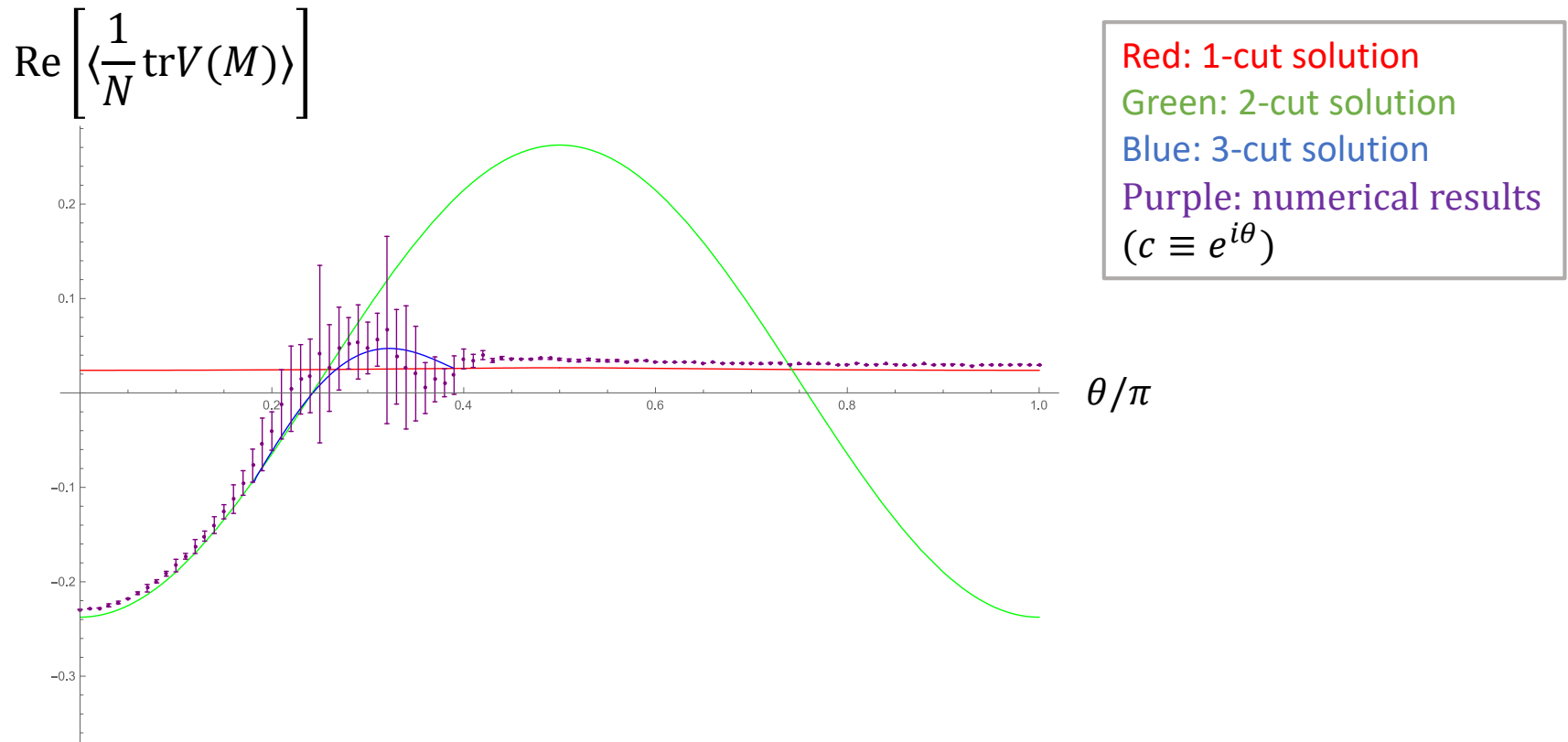
We calculate $\frac{dF}{d\beta} = \langle \frac{1}{N} \text{tr} V(M) \rangle$ of this model by using LTM.

Parameters

- $N = 6$
- $\beta = 20$, $\theta \in [0, \pi]$ ($c \equiv e^{i\theta}$)
➔ Sign problem
- We introduce two types of tempering parameter: (β, t)
set of β : $\{\beta\} = \{5, 10, 15, 20\}$, set of t : $\{t\} = \{0, \frac{1}{4}, \frac{1}{2}, \frac{3}{4}, 1\}$
- Dynamical variables: $N \times N$ Hermitian matrix (not diagonalized)
➔ The models have $SU(N)$ symmetry.

Numerical results (2/2)

Results



Numerical results correctly agree with analytical solutions.

Summary

We introduce the parallel tempering algorithm to LTM by regarding the flow time t as a tempering parameter.



This method is quite versatile:

We can calculate *any* models (in principle) without the fine-tuning of t .

i.e. we can calculate the models {

- which have multi-thimble contributions.
- which have continuous symmetry.

}

Future problems

- LTM is Costly
Computational cost \propto (degrees of freedom)³
- Residual sign problem
- Global sign problem