



Crete Center
for Theoretical Physics

**Geometry and Holography
for Quantum Criticality**

Universal bounds on transport To be or not to be?

Matteo Baggioli

UOC & Crete Center for Theoretical Physics

with B. Gouteraux, E. Kiritsis, W.Li, A.Amoretti

Holography 2013

:Gauge/gravity duality and strongly correlated systems

June 13 (Thu) ~ June 22 (Sat), 2013
APCTP, Pohang, Korea



Main Page

Registration/Participants

Program

Accommodation

Travel Information

■ Main Page

Overview

This focus workshop is to promote research oriented discussions between the world class experts on holography and strongly interacting and strongly correlated system. This year, however, we will also have a school session for students and postdocs.

Organizers

Sang-Jin Sin (Hanyang Univ.,Chair)
Nick Evans (U. of Southampton)
Deog-Ki Hong (Pusan Nat'l Univ.)
Keun-Young Kim (GIST)
Piljin Yi (KIAS)

Geometry and Holography for Quantum Criticality

August 18 (Fri), 2017 ~ August 26 (Sat), 2017

Organizers

Koji Hashimoto (Osaka University)
Deog-Ki Hong (Pusan National University)
Keun-Young Kim (GIST) - *Cochair**
Ki-Seok Kim (Postech)
Nakwoo Kim (Kyunghee University)
Sang-Jin Sin(Hanyang University) - *Chair**
Piljin Yi (KIAS)

***Is there a minimum
(Planckian) timescale in nature ?***

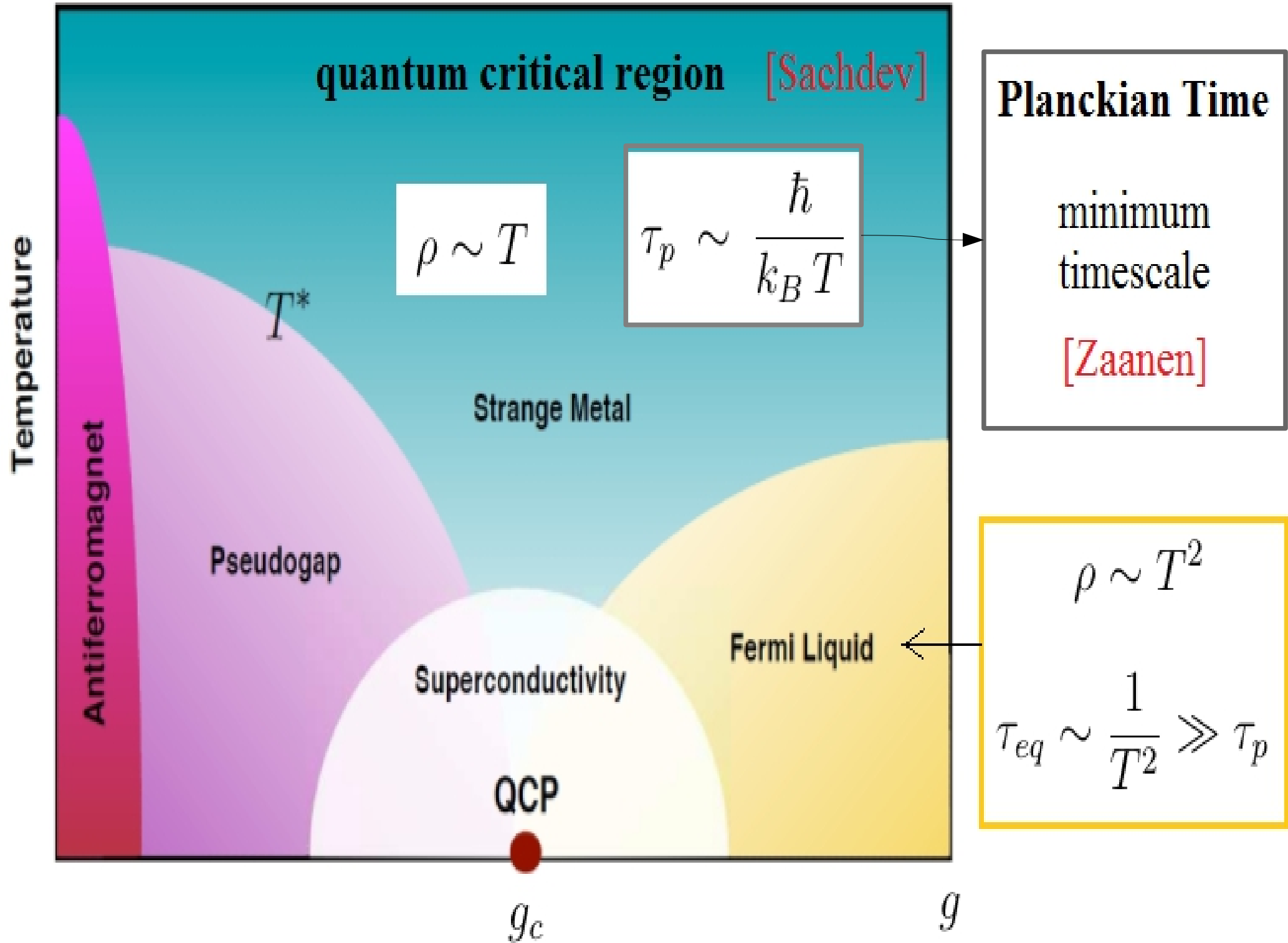


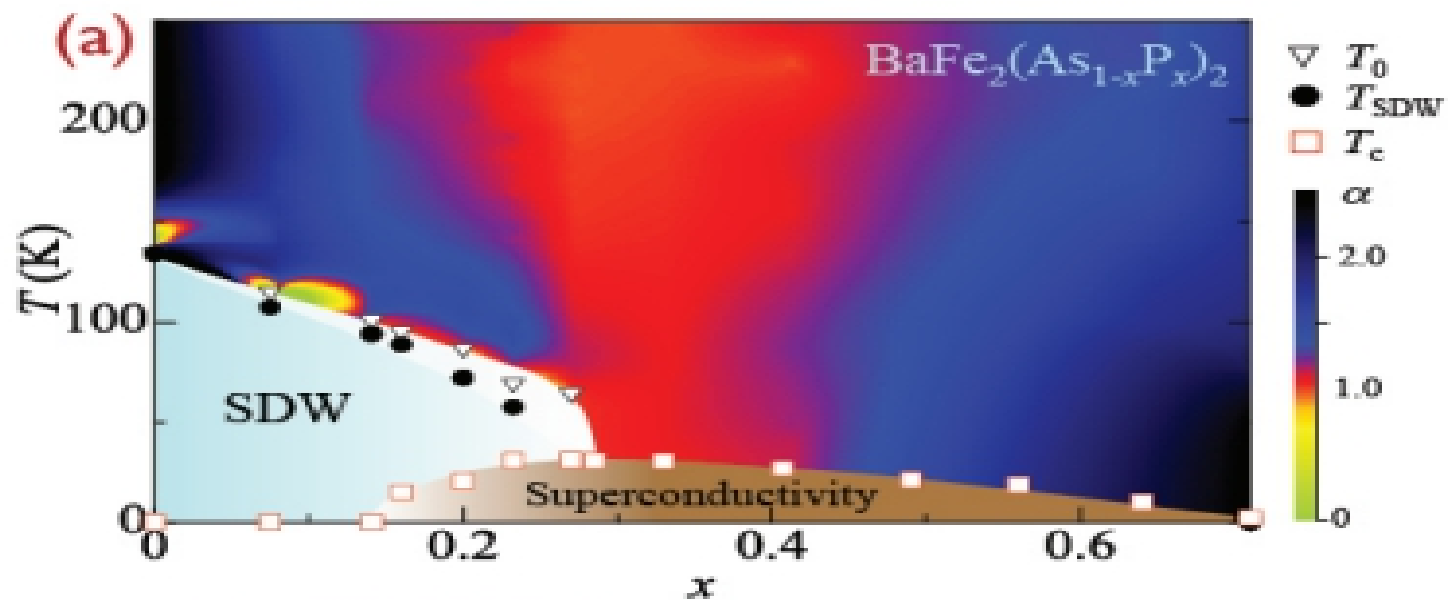
***Imprints
on transport ??***



***Bounds
on transport !??***

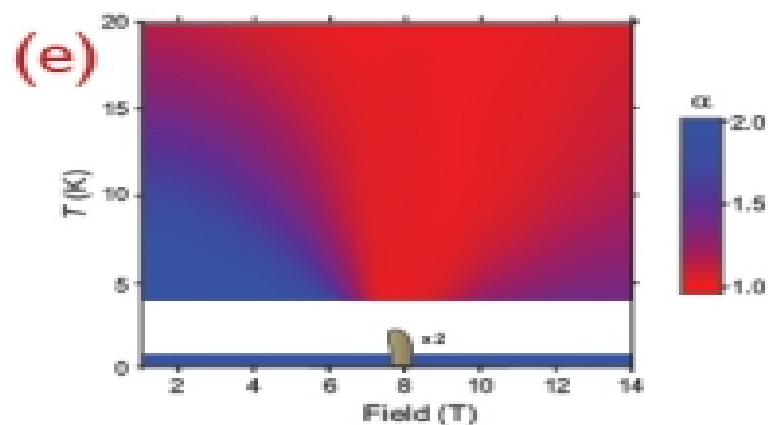
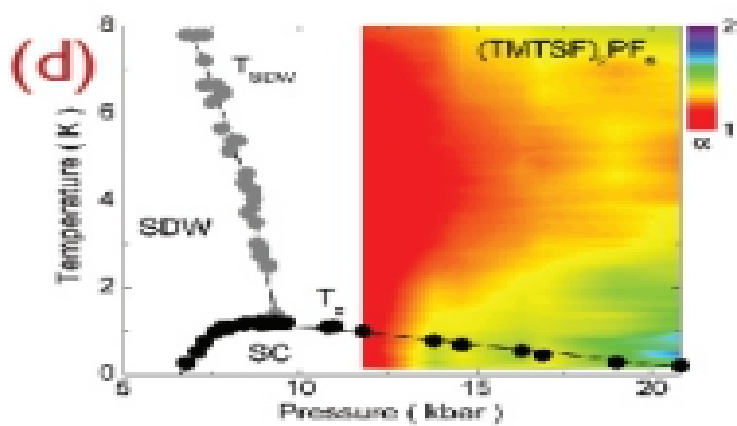
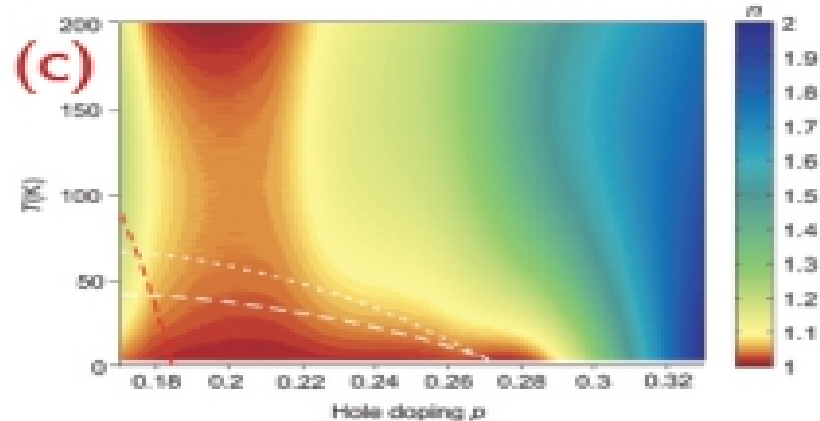
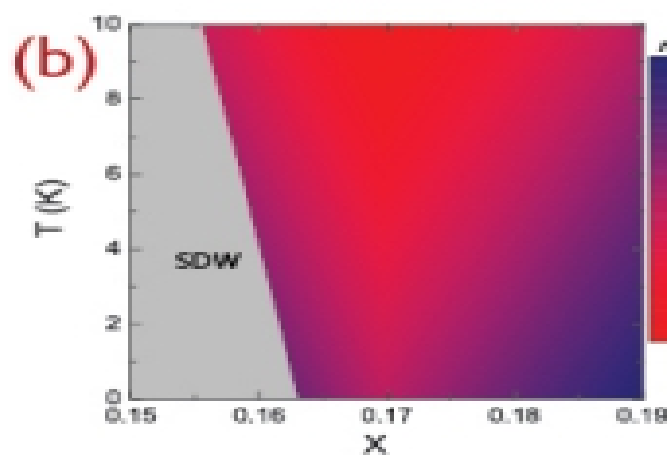






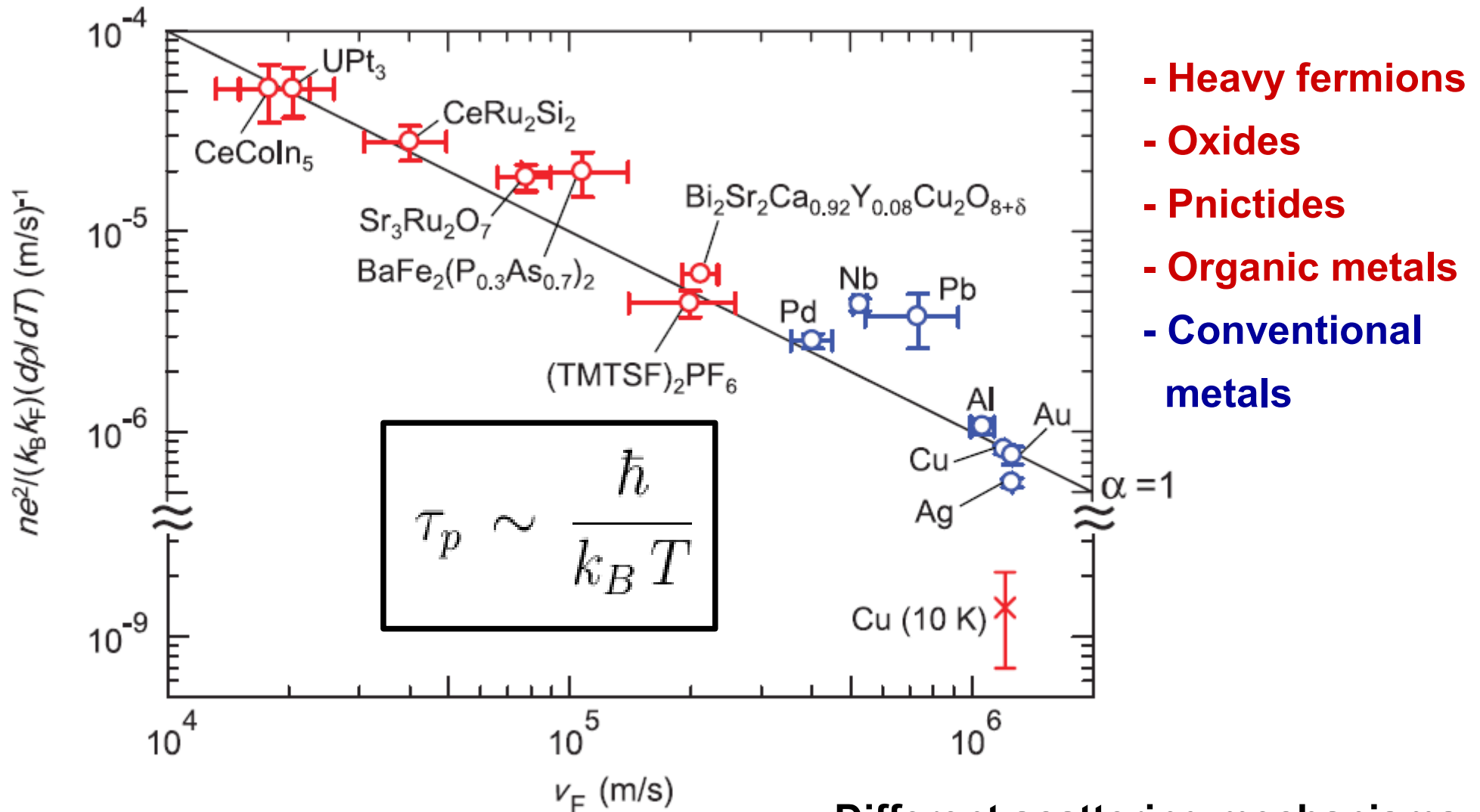
$$\rho \sim T$$

[Sachdev, Keimer]



UNIVERSAL PLANCKIAN SCATTERING TIME

[Bruin JA, Sakai H, Perry RS, Mackenzie AP., Science 2013]



$$\tau_p \sim \frac{\hbar}{k_B T}$$

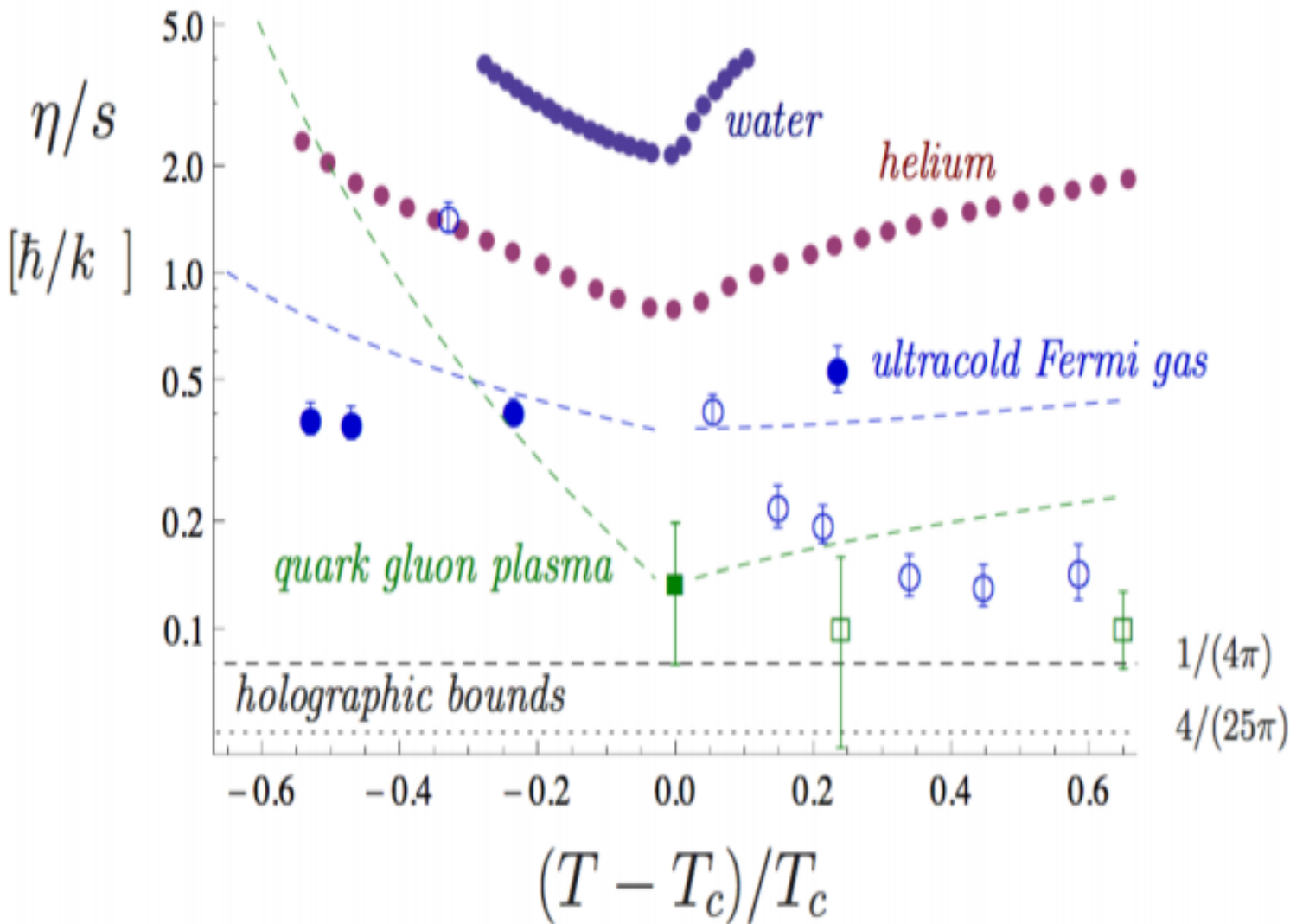
Different scattering mechanisms

2 order of magnitude variations in the Fermi velocity

$$\frac{1}{\tau} = \frac{e^2 \rho}{\hbar d} \sum_i k_{Fi} v_{Fi} \quad (T \tau)^{-1} \sim \alpha \frac{k_B}{\hbar}$$

KSS BOUND

[Kovtun, Son, Starinets, Policastro, 2004]



Minimum
Planckian Time



**UNIVERSAL
BOUND(S)**

$$\frac{\eta}{s} \geq \frac{1}{4\pi} \frac{\hbar}{k_B}$$

$$\tau_p \sim \frac{\hbar}{k_B T}$$

BOUNDS ON DIFFUSIVITIES

[Hartnoll, 2014]

$$[D] \equiv [v^2] * [t] \quad \text{Diffusion can't be arbitrarily fast}$$

$$D \geq v^2 \tau_p = v^2 \frac{\hbar}{k_B T}$$

Generically

$$\tau_{eq} \gg \tau_p$$

Fast(est) dissipation
Strong coupling

KSS story
($v=c$) \rightarrow $D_p \equiv \frac{\eta}{sT} \geq \left(\frac{1}{4\pi}\right) \frac{\hbar}{k_B T}$

Charge and Energy Diffusivities (??) \rightarrow $D_c = \frac{\sigma}{\chi} \quad D_e = \frac{\kappa}{c_v}$

COHERENT METALLIC TRANSPORT

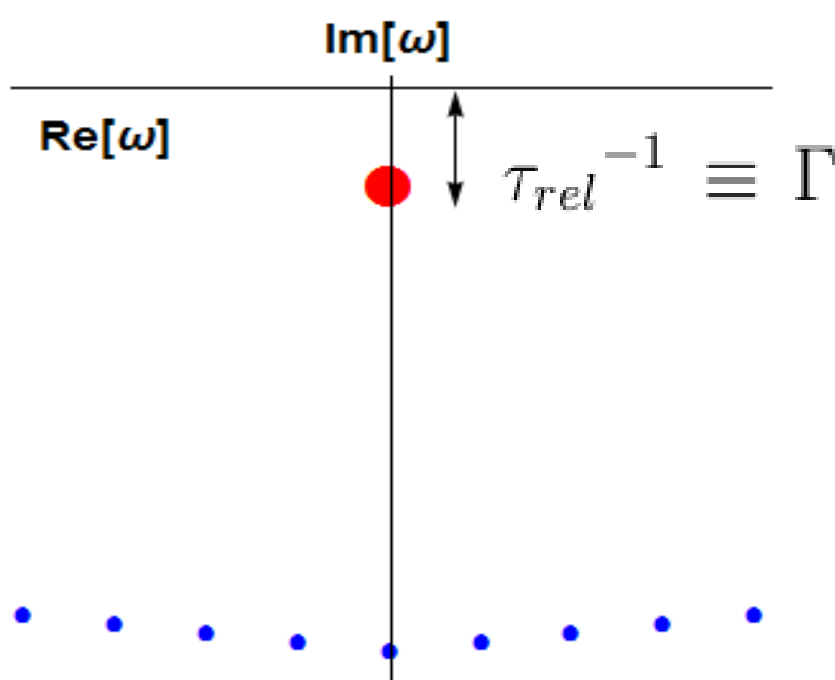
electric current \mathcal{J} + momentum \mathcal{P}

They overlap because a finite charge density !!

Momentum relaxation : $\langle \mathcal{P}(t) \rangle \sim e^{-t/\tau_{rel}}$

Weak momentum relaxation
(long lived momentum)

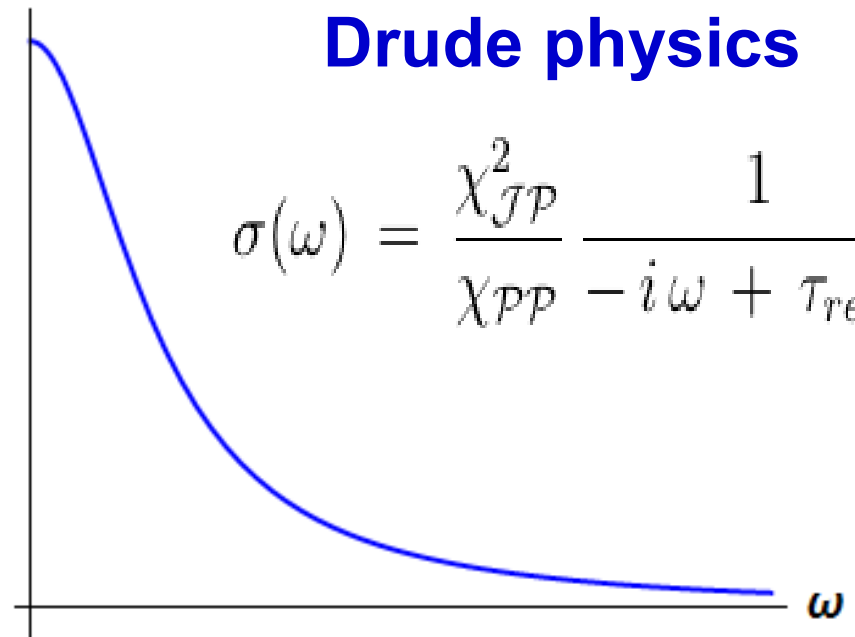
$$\Gamma \ll E \sim k_B T$$



$\text{Re}[\sigma]$

Drude physics

$$\sigma(\omega) = \frac{\chi_{\mathcal{J}\mathcal{P}}^2}{\chi_{\mathcal{P}\mathcal{P}}} \frac{1}{-i\omega + \tau_{rel}^{-1}}$$

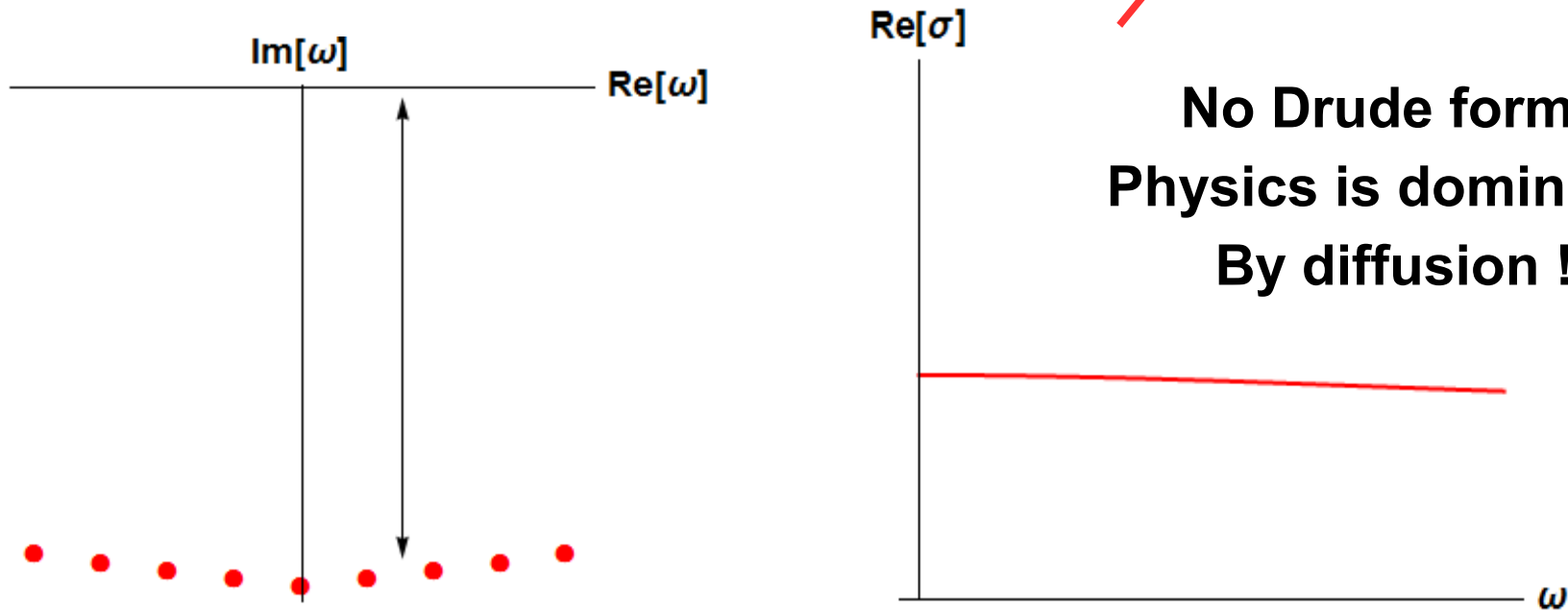


INCOHERENT METALLIC TRANSPORT

no long-lived quantities overlapping with the current operators.

Fast momentum relaxation

$$\Gamma \not\ll E \sim k_B T$$



In this limit we should approach the minimum timescale

$$\frac{\hbar}{k_B T}$$

In this limit $D \geq v^2 \tau_p$ should be saturated !!

TRANSPORT WITHOUT QUASIPARTICLES

$$\ell_{\min} \sim a.$$

**Mott-Ioffe-Regel bound
is brutally violated in
Bad metals**

[Hussey, Takenaka, Takagi]

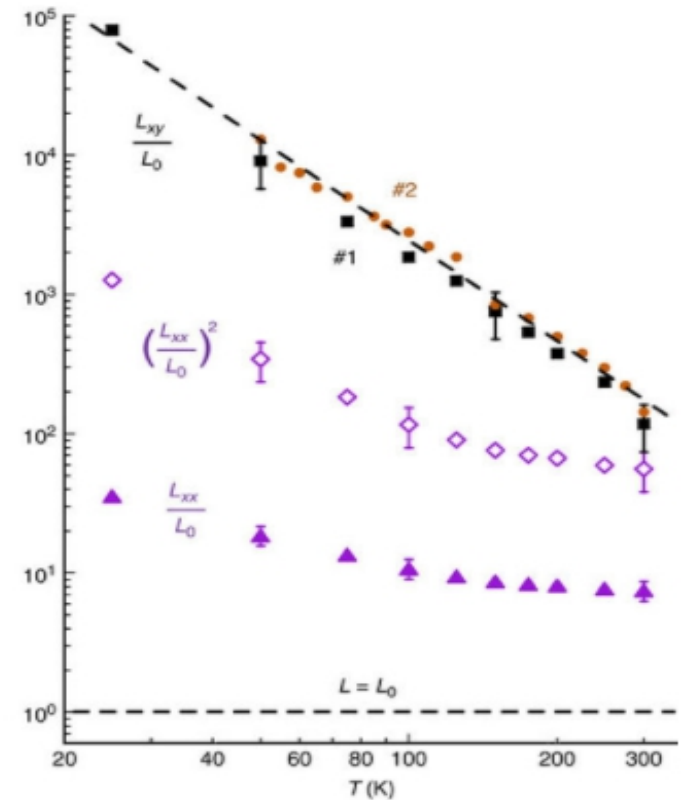
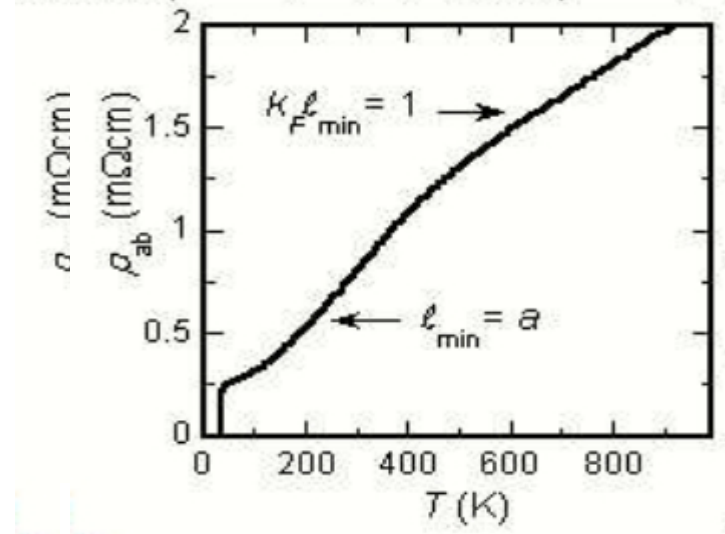
**Wiedemann – Franz law
is violated as well**

[many authors ...]

$$L = \frac{\kappa}{\sigma T} = \frac{\pi^2}{3} \left(\frac{k_B}{e} \right)^2$$

**Concept of Landau quasiparticles
breaks down !!!**

Strong coupling / correlations



The butterfly velocity and quantum chaos

Who is the velocity v ??
It can't be the Fermi velocity !!

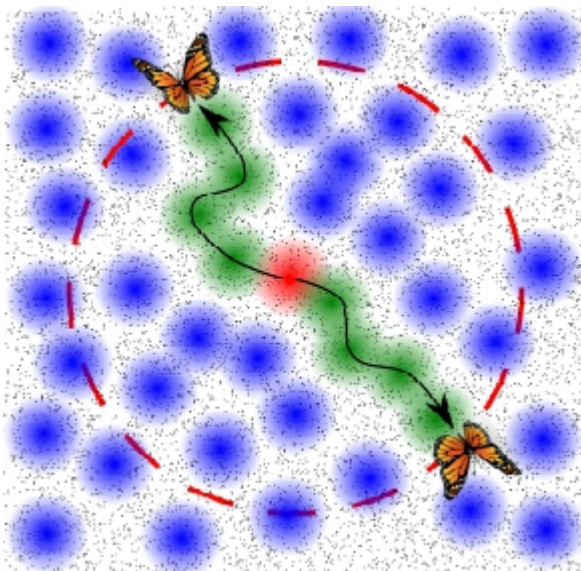
v IS THE BUTTERFLY VELOCITY = SPEED OF INFORMATION PROPAGATION

**QUANTUM CHAOS
SCRAMBLING**

*[Shenker, Stanford, Susskind, Swingle,
Maldacena, Blake, Roberts, Douglas, ...]*

**OUT-OF-TIME
CORRELATOR**

$$\langle [\mathcal{V}(x, t) \mathcal{W}(0, 0)]^2 \rangle_{\beta} \sim e^{\lambda_L(t-t^* - |x|/v_B)},$$



PROPOSAL

$$\frac{D}{v_B^2} \geq c \frac{\hbar}{k_B T}$$

*[Blake,
2016]*

PHYSICS AT STRONG COUPLING

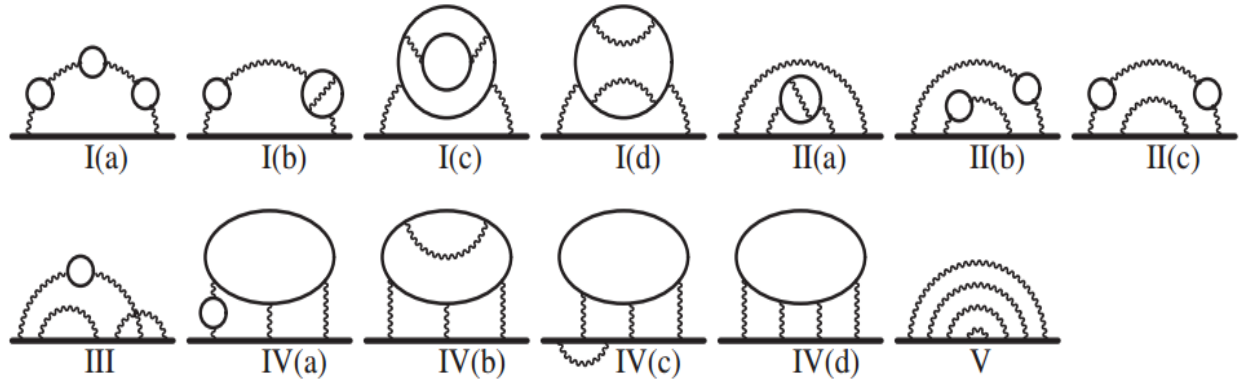
We are *very good*
at computing

$$\blacksquare (g) = \mathcal{O}(g^0) + \mathcal{O}(g) + \mathcal{O}(g^2) + \dots$$

Best example

$$g_e = 2\left(1 + \frac{\alpha}{2\pi} + \dots\right),$$

$$g_e = 2.0023193043617(15).$$



But what if
the coupling
is not small ??

$$\blacksquare \left(\text{strawberry} \right) = \text{strawberry} + \text{strawberry} + \text{strawberry} + \text{strawberry} + \dots$$

$$g \ll 1$$

STRONG COUPLING

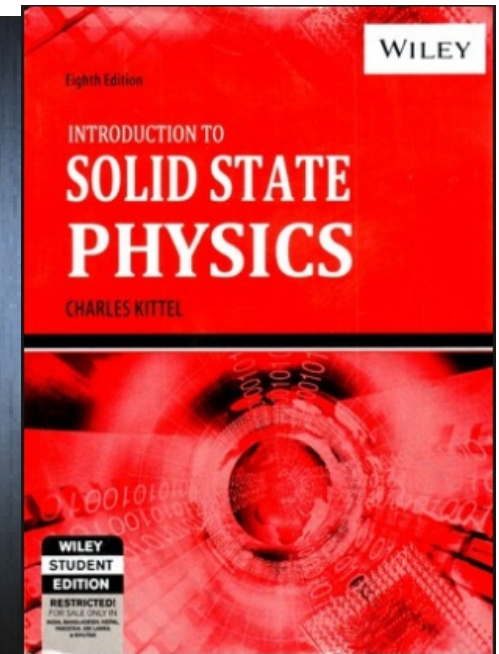
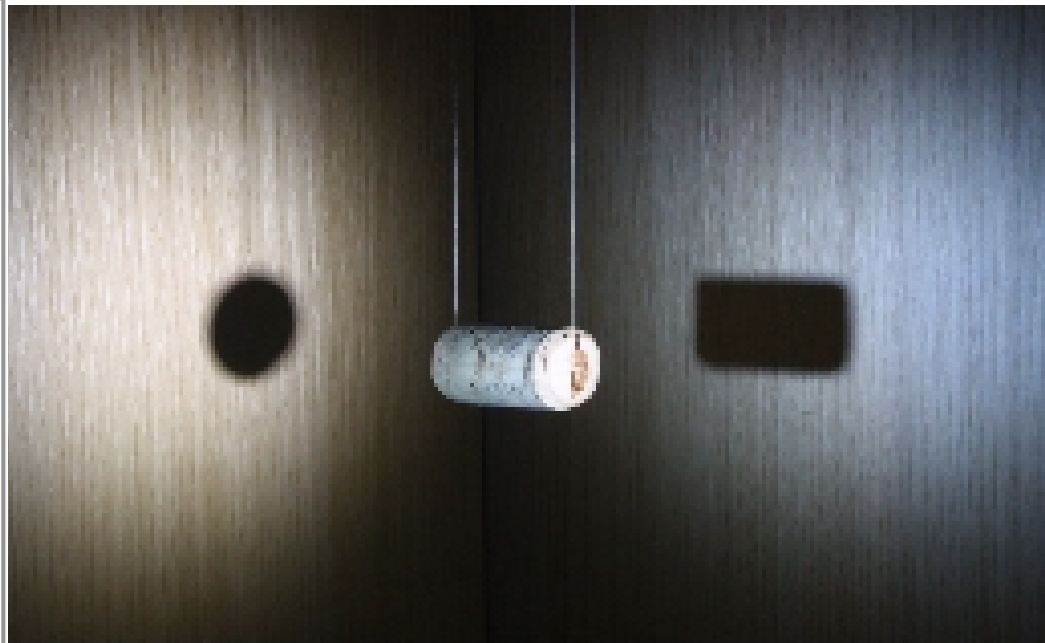
**Standard (usually very efficient) methods
are not useful anymore !!**



IDEA : DUALITY (AdS-CMT)

General Relativity

Robert M. Wald



*Use a dual description in terms of different d.o.f.
where the theory is weakly coupled
and the computations are doable*



BOTTOM – UP HOLOGRAPHY



In some limits :

**STRONGLY COUPLED
(LARGE N) QFT**



**CLASSICAL GR
+ BULK FIELDS**

d

d + 1

$$\mathcal{Z}_{CFT} [\phi_0(\vec{x})] = \langle e^{\int d^d x \phi_0(\vec{x}) \mathcal{O}(\vec{x})} \rangle_{CFT} = e^{-S_{gravity}[\phi(\vec{x}, z=0) = \phi_0(\vec{x})]}$$

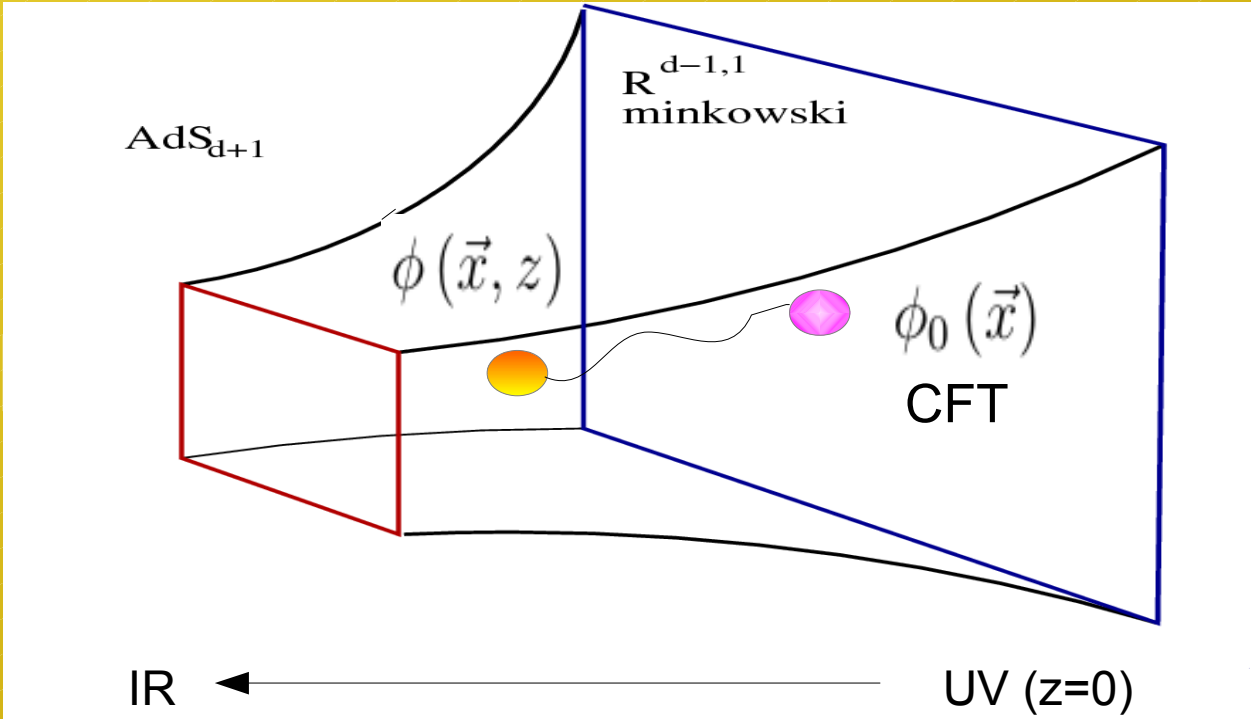
Everything !!

$$J_\mu(\vec{x}) \quad A_\mu(\vec{x}, z)$$

$$T_{\mu\nu}(\vec{x}) \quad g_{\mu\nu}(\vec{x}, z)$$

$$\langle \mathcal{O}(\vec{x}) \mathcal{O}(\vec{y}) \rangle$$

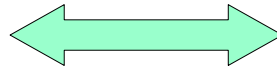
At strong coupling



MOMENTUM DISSIPATION

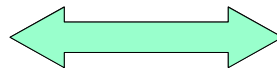
$$T_{\mu\nu}(\vec{x}) \quad \text{---} \quad g_{\mu\nu}(\vec{x}, z)$$

Global symmetry



Gauged symmetry

Translational symmetry
of the CFT

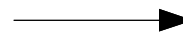


Diffeomorphism invariance
Of the bulk theory

*In order to break translational invariance in the CFT
We need to break (spatial) diffeomorphisms in the bulk*

[Vegh, '13]

Generic effective holographic theory
With momentum dissipation



**MASSIVE
GRAVITY**

[Vegh, Tong, Blake]

Translations broken
Energy conserved



**LORENTZ VIOLATING
MASSIVE GRAVITY**

Massive gravity and phenomenology

e

[Davison]

$$\partial_i T^{ij} = -\frac{1}{\tau_{rel}} T^{tj} \neq 0 \quad \frac{1}{\tau_{rel}} \sim \mathcal{M}_h^2 (T, k, q, g_i, \dots)$$

[cf. MEMORY MATRIX FORMALISM , Andy's lectures]

$$\sigma = \sigma^{\mathcal{I}} + \frac{q^2}{\mathcal{M}_h^2}, \quad \alpha = \bar{\alpha} = \frac{s q}{\mathcal{M}_h^2}, \quad \bar{\kappa} = \frac{s^2 T}{\mathcal{M}_h^2}.$$

[Donos, Gauntlett, ...]



$\mathcal{M}_h^2 \rightarrow 0$ in the case of translational invariance



\mathcal{M}_h^2 is the effective, model dependent, graviton mass !



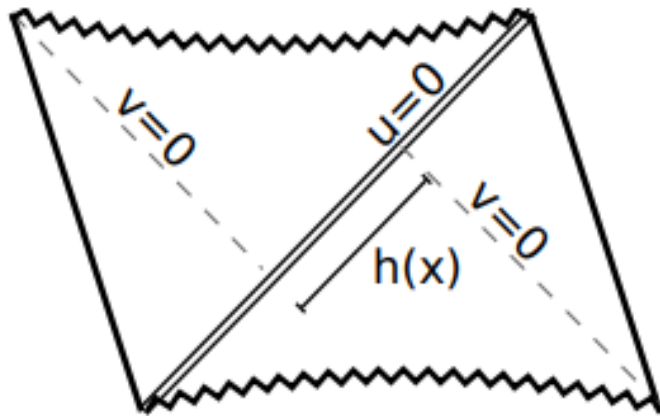
$\mathcal{M}_h^2 \rightarrow \infty$: incoherent limit

Quantum chaos and shockwaves

e

[Susskind, Maldacena,
Stanford, Douglas,
Shenker, Blake, ...]

$$\langle [\mathcal{V}(x, t) \mathcal{W}(0, 0)]^2 \rangle_\beta \sim e^{\lambda_L(t-t^* - |x|/v_B)},$$



Horizon shockwave $\delta T_{uu} \sim E e^{\frac{2\pi}{\beta} t_w} \delta(u) \delta(\vec{x})$

$$(\partial_i \partial_i - m^2) h(x) \sim \frac{16\pi G_N V(0)}{A(0)} E e^{\frac{2\pi}{\beta} t_w} \delta(\vec{x})$$

Time shift :

$$h(x) \sim \frac{E e^{\frac{2\pi}{\beta}(t_w - t_*) - m|x|}}{|x|^{\frac{d-1}{2}}}$$

RESULTS:

$$ds_{d+2}^2 = -U(r)dt^2 + \frac{dr^2}{U(r)} + V(r)d\vec{x}_d^2$$

$$m^2 = d\pi T V'(r_0)$$

$$\lambda_L = \frac{2\pi}{\beta} \quad v_B = \frac{2\pi}{\beta m}$$

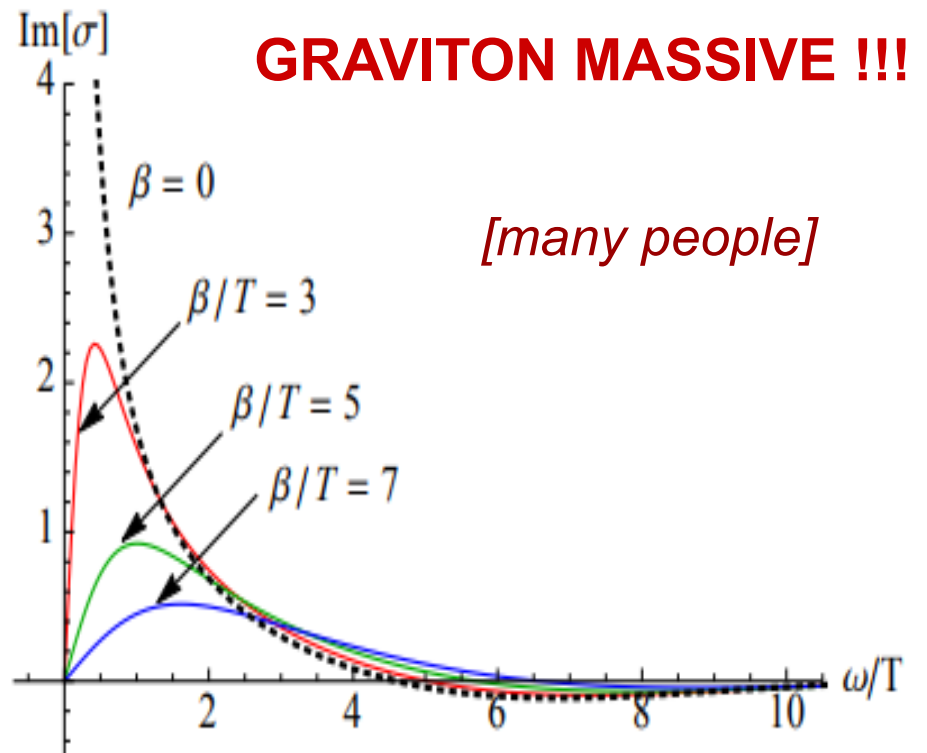
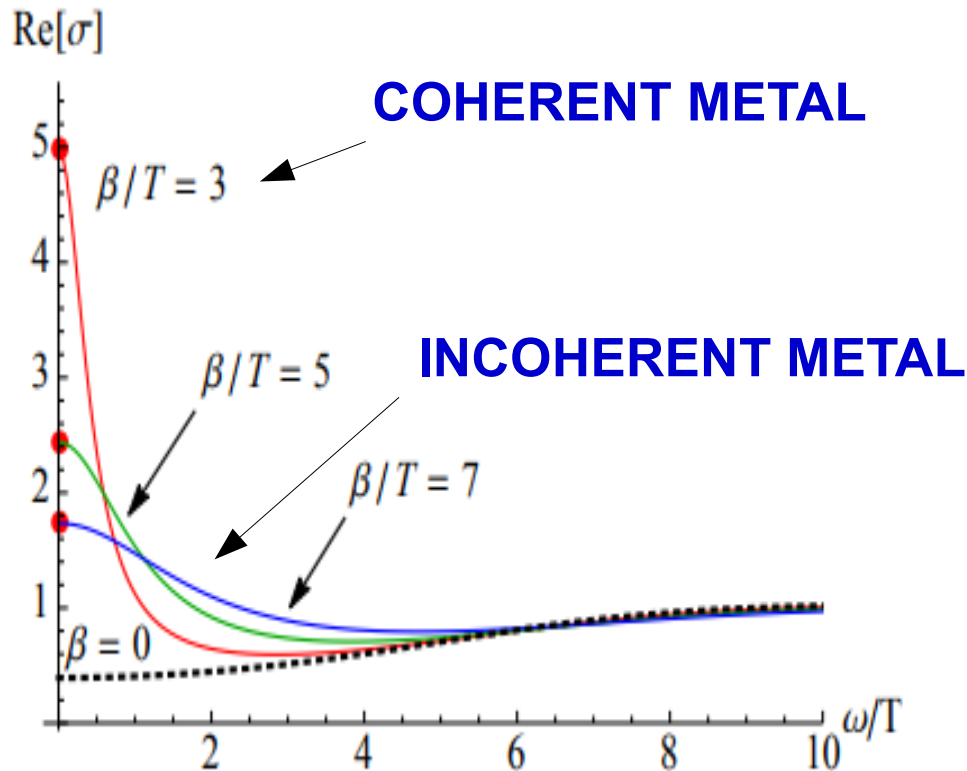
THE SIMPLEST THEORY

[Andrade, Withers, '13]

~~LINEAR (AXIONS) STUECKELBERGS~~

$$S_0 = \int_M \sqrt{-g} \left[R - 2\Lambda - \frac{1}{2} \sum_I^{d-1} (\partial\psi_I)^2 - \frac{1}{4} F^2 \right] d^{d+1}x$$

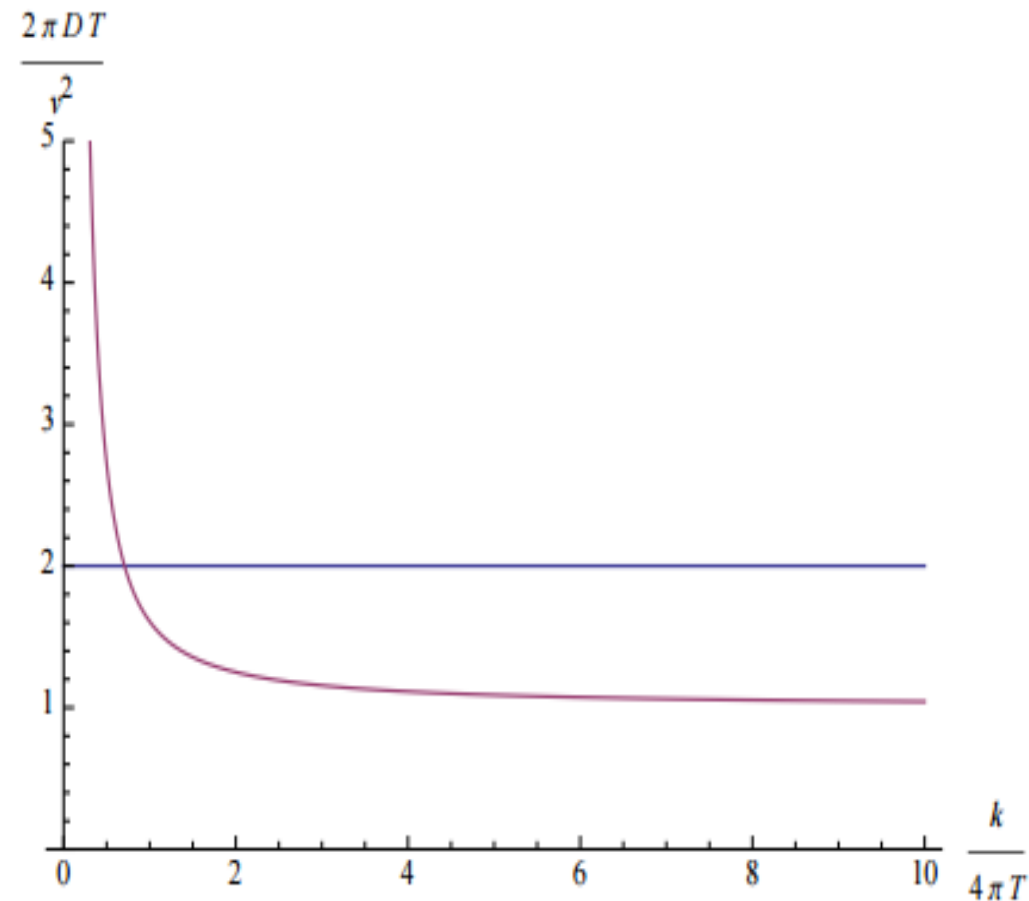
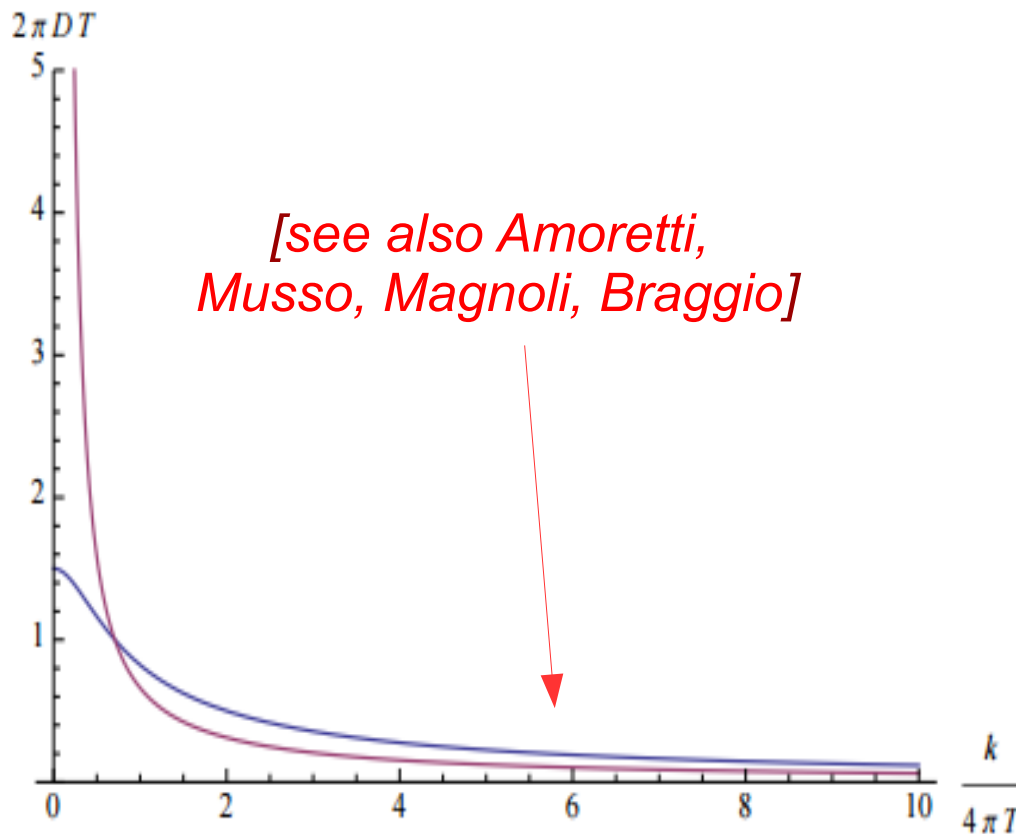
Massless scalars with x-dependent sources : $\psi^I = \beta x^I$



Checking the "bound" : part I

In the simplest holographic theories

[Blake]
$$D_c = \frac{v_B^2}{\pi T} \quad D_e \approx \frac{v_B^2}{2\pi T}$$



HOW GENERAL IS IT ?

[M.B., Kiritsis, Gouteraux, Li]

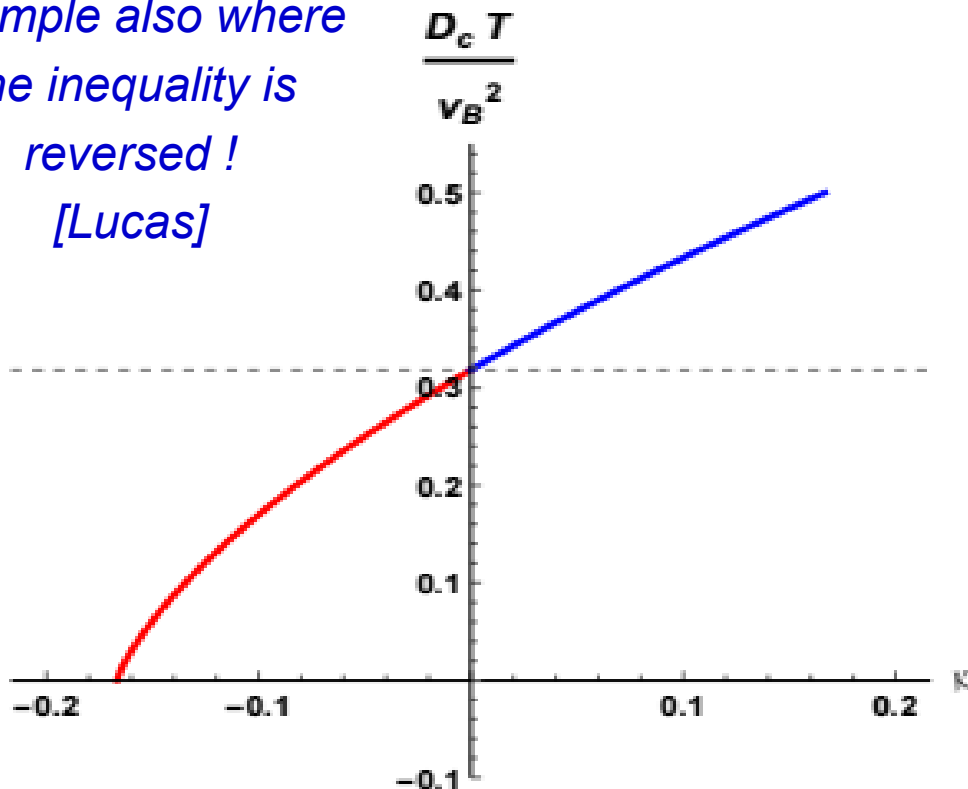
$$\chi^\mu{}_\nu \equiv \frac{1}{2} \sum_{I=x,y} \partial^\mu \phi^I \partial_\nu \phi^I$$

HIGHER DERIVATIVES CHECKS

$$\mathcal{L} = \dots - \frac{\mathcal{J}}{4} \text{Tr} [\chi F^2]$$

$$\mathcal{L} = \dots - \frac{F^2}{4} (1 + \kappa \text{Tr} [\chi])$$

Example also where
The inequality is
reversed!
[Lucas]



CHARGE SECTOR

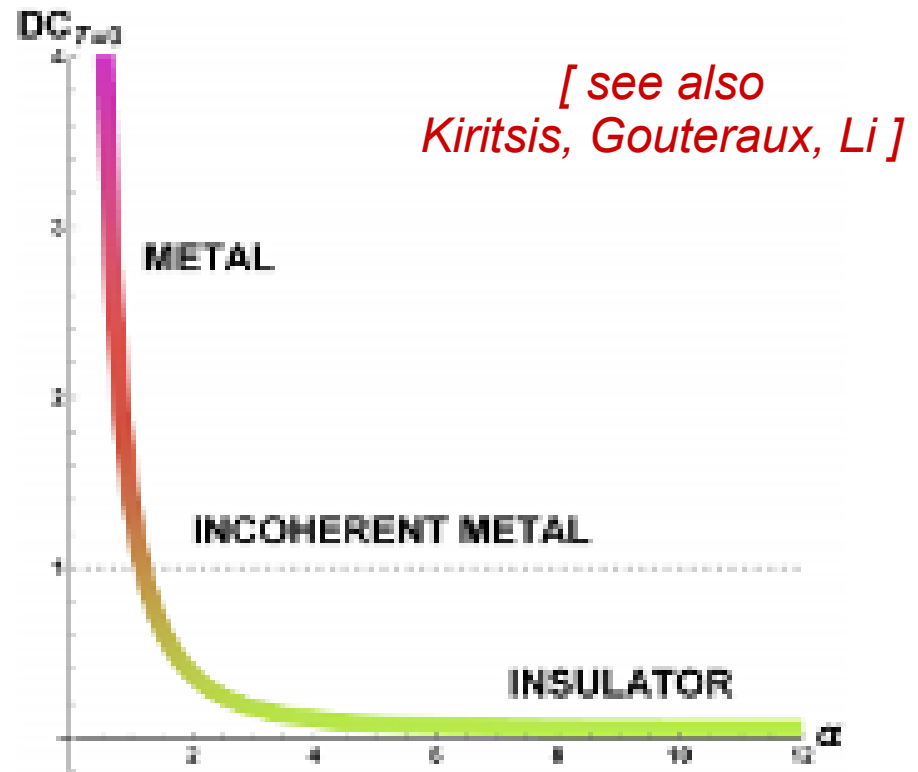
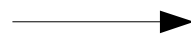
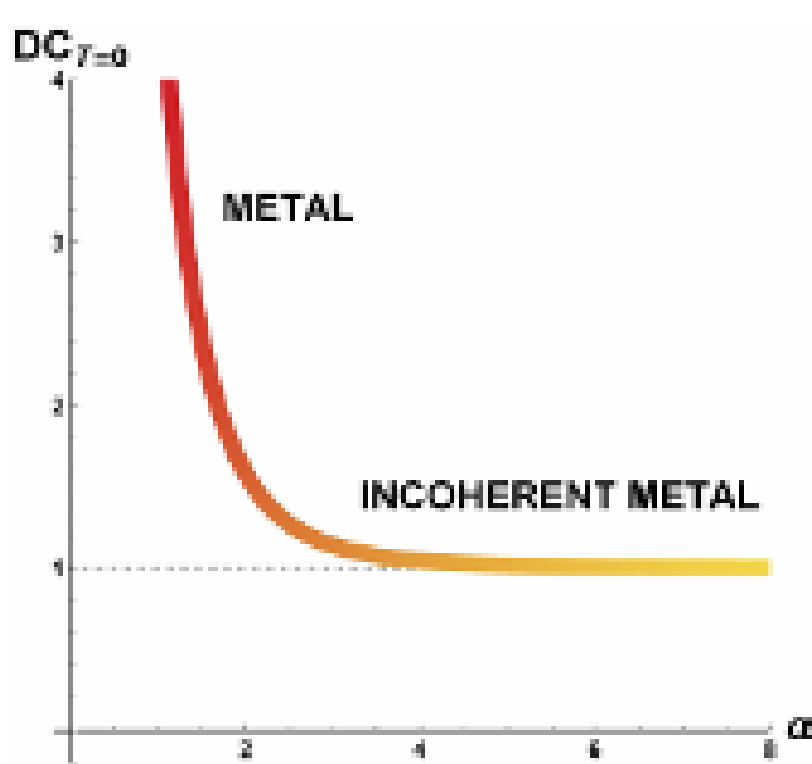


**ENERGY
SECTOR**



ANOTHER BOUND †

$$\mathcal{S} = \int d^4x \sqrt{-g} \left[\frac{R}{2} - \Lambda - \frac{1}{4e^2} Y(X) F^2 - m^2 V(X) \right].$$



[see also
Kiritsis, Gouteraux, Li]

$$\sigma_{DC} \not> \sigma_{min}$$



The electric conductivity
Is not bounded !

STATUS AND (MY) INTUITION

$$D_c \geq v_B^2 \tau_l \quad \sigma_{DC} > \sigma_{min}$$



D_c Maxwell sector  v_B^2 Gravitational sector

$D_c = \frac{\sigma}{\chi}$  **The charge susceptibility is not IR data**
 $\sigma \rightarrow 0$

$$D_e \geq v_B^2 \tau_l \quad \frac{\kappa}{T} \geq C_{min}$$

**DON'T
WANNA
GO HOME**

Higher derivatives check part II

Higher derivative couplings
Charge sector - Momentum
relaxing sector

→ ~~CHARGE BOUND~~

[MB, Gouteraux,
Kiritsis, Li, 2016]

What about Gravity – Momentum Relaxing sector ??

$$S = \int d^4x \sqrt{-g} \left(R - 2\Lambda - \frac{1}{4} F^2 - \frac{1}{2} (g^{\mu\nu} - \gamma G^{\mu\nu}) \sum_{i=1}^2 \partial_\mu \phi^i \partial_\nu \phi^i \right)$$

Holographic Horndeski theories

[MB, Li, 2017]

Results:

$$\frac{D_e}{v_B^2} \geq c \frac{\hbar}{k_B T}$$



More and more checks

**Lifshitz – Hyperscaling FP
(dilaton couplings)**

[Blake, Sachdev, Davison]



AdS2 horizons

[Blake, Donos]



Gauss- Bonnet

[Wu, Wang, Ge, Tian]



→ $L^2 \rightarrow L_{eff}^2 (\lambda_{GB})$

Charge and magnetic field

[Blake, Sachdev, Davison] [Kim, Niu]



SYK

[Davison, Fu, Gu, Georges, Sachdev, Jensen]



→ $D_2 = \frac{v_B^2}{2\pi T}$

Weakly coupled Fermi Liquids

[Aeine, Faoro, Ioffe]



Bose-Hubbard models

[Bohrdt, Endrel, Mendes, Knap]



Diffusive metals

[Swingle, Chowdhury]



Electron-Phonon bad metals

[Werman, Kivelson, Berg]



Critical Fermi Surfaces

[Patel, Sachdev]



→ $D_T \sim v_B^2 \lambda_L^{-1}$

Beyond the incoherent limit : I

ENERGY
DIFFUSION

BUTTERFLY
VELOCITY

[Blake, Donos]

$$ds_{d+2}^2 = -f(r) dt^2 + \frac{dr^2}{f(r)} + h(r) dx_A^2$$

$$h(r) = h_0 + \underline{c_h^0(\rho)} r + \dots$$

$$D = \kappa / c_\rho$$

$$v_B$$

SAME IRRELEVANT DEFORMATIONS OF AdS₂

$$D = E \frac{v_B^2}{2\pi T} \quad 1/2 < E \leq 1$$

*Where E is Related to the
Conformal Dimension of
The (dilaton) Deformation
At the IR fixed point*

No incoherent limit taken !!!

Just a property of the IR fixed point !!

Beyond the incoherent limit : II

[Blake, Davison, Sachdev]

$$ds_{d+2}^2 = -f(r) dt^2 + \frac{dr^2}{f(r)} + h(r) dx_i^2,$$

Generic theories :
Dilaton couplings, matter, magnetic field
(but Einstein gravity)

$$\kappa = 4\pi \frac{f' h^{d-2}}{(f' h^{d/2-1})'} \Big|_{\tau_0},$$

**Thermal conductivity just in
Terms of metric data at the horizon !**

**UNIVERSAL
RELATION**

$$D_T \sim \frac{f' h^{d/2-1}}{(f' h^{d/2-1})'} \frac{h'}{h} \Big|_{\tau_0} v_B^2 \tau_L.$$



**For generic
Hyperscaling - Lifshitz
IR fixed points**



$$D_T = \frac{z}{2z-2} v_B^2 \tau_L,$$

Conclusions

$$\frac{D_e}{v_B^2} \geq c \frac{\hbar}{k_B T}$$



**UNIVERSAL RELATION VERY ROBUST
HOLOGRAPHY + CONDENSED MATTER**

Just one possible violation (inhomogeneous SYK chain) [Lucas, Gu, Qi]

[Gouteraux, Blake, Davison, Sachdev, [Hartnoll, Grozdanov, Lucas, Gentle, Donos, Kiritsis, Patel, Li, Kim, Ling, Wu, Shenker, Liu, Stanford, Phillips, Jensen, Tian, Wang, Swingle] Ge, Niu, Amoretti, Jin-Sin Musso, Magnoli, ...]

**Holography, Condensed Matter, Quantum Chaos,
Hydrodynamics, Quantum Information,
Random matrix theories, Black Holes**

Future

IS THIS BOUND REALLY UNIVERSAL ???

Higher derivative corrections

[IN PROGRESS !! M.B. , A. Amoretti]

**Gravity dual of inhomogeneous SYK chain ??
KSS bound with momentum dissipation ??**

CAN WE PROVE IT ?? UNDERSTAND IT BETTER !!

vB? Strong coupling \leftrightarrow Chaos?

IF UNIVERSAL, CAN WE MEASURE IT ?? CAN WE TEST IT ?

감사합니다
[kamsahamnida]

