

AdS₆/CFT₅ in Type IIB

Part I: Warped AdS_6 solutions and 5-brane webs

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Strings, Branes and Gauge Theories

APCTP, July 2018

arXiv: 1606.01254, 1611.09411, 1703.08186, 1705.01561,

1706.00433, 1802.07274, 1805.11914, 1806.07898, 1806.08374

with E. D'Hoker, M. Gutperle, A. Karch, C. Marasinou, A. Trivella,

O. Varela, O. Bergman, D. Rodríguez-Gómez, M. Fluder

Introduction

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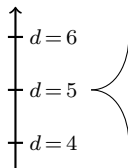
Theories in $d > 4$ crucial part in general understanding of QFT. Often reduce to interesting $d \leq 4$ theories \rightarrow new theories and dualities, geometric realizations of known dualities.

Defining interacting QFTs in $d > 4$ challenging. In particular, Yang-Mills theories non-renormalizable ($\sim \sqrt{-g}R$ in $d = 4$).

Evidence from [Strings, Branes and Gauge theories](#) suggests many interacting $d > 4$ QFTs exist, challenging perturbative arguments.

Introduction

Suitable superconformal algebras in $d \leq 6$; maximal ones with $16_Q + 16_S = 32$ supercharges

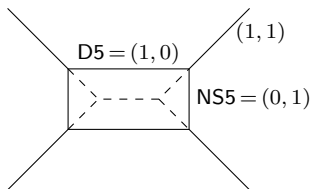
- 
- unique superconf. algebra $F(4)$, 16 supercharges
 - strongly-coupled UV fixed points for large classes of gauge theories w/ 8_Q supercharges [Seiberg '96;...]
 - no standard Lagrangian, existence from Coulomb branch analysis and string theory

Asymptotically safe gauge theories, exceptional global symmetries, dualities, parents to isolated $4d$ theories, relations to $6d$, ...

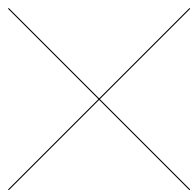
5-brane webs in Type IIB

[Aharony, Hanany, Kol '97]

5-brane web: planar arrangement of (p, q) 5-branes at angles fixed by (p, q) , junctions w/ conserved charges

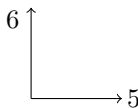


Coulomb branch
finite gauge coupling



UV fixed point CFT

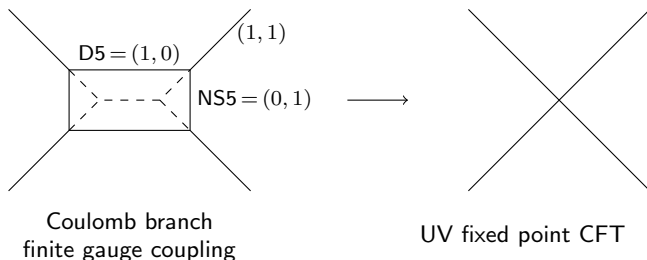
	0	1	2	3	4	5	6	7	8	9
D5	x	x	x	x	x	x	x			
NS5	x	x	x	x	x		x			



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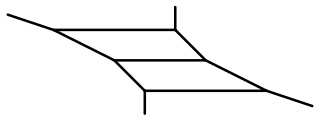
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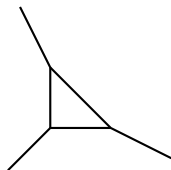
Large classes of 5d SCFTs, with and without gauge theory deformations, quivers, flavors, Chern-Simons terms,...

5-brane webs in Type IIB

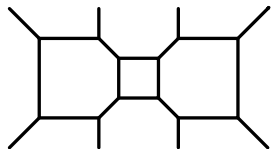
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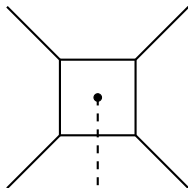
$SU(3)$, $CS=0$



E_0 , no global symmetry



$SU(2) \times SU(2) \times SU(2)$ quiver

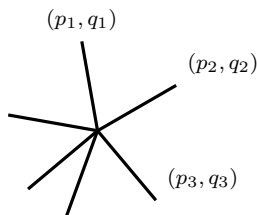


$SU(2) + 1$ flavor

5-brane webs in Type IIB

[Aharony,Hanany,Kol '97]

General picture: any planar 5-brane junction realizes a 5d SCFT on the intersection point



$$p_i, q_i \in \mathbb{Z}$$

$$\sum p_i = \sum q_i = 0$$

Characterized entirely by external 5-brane charges. No standard Lagrangian. May or may not have gauge theory deformations.

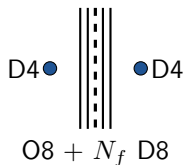
Supergravity duals?

AdS/CFT to access superconformal fixed points? Harder than in $d \neq 5$, no maximally supersymmetric solutions.

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AdS/CFT to access superconformal fixed points? **Harder than in $d \neq 5$, no maximally supersymmetric solutions.**

One well known AdS_6 solution from type I' construction:



N D4 probing $\text{O8} + N_f$ D8: $\text{USp}(N)$
with antisymmetric and N_f flavors

→ warped $\text{AdS}_6 \times S^4$ in massive IIA

[Brandhuber, Oz]

Locally unique [Passias], orbifolds dual to quiver gauge theories
[Rodriguez-Gomez, Bergman], T-duals in IIB [Cvetic et al.; Lozano et al.]

Holographic duals for 5d SCFTs

Holographic duals for SCFTs realized by 5-brane webs in Type IIB?
Not a standard near-horizon limit – fully localized intersections.

Type IIB BPS equations studied by [Apruzzi,Fazzi,Passias,Rosa,Tomasiello;
H.Kim,N.Kim,Suh;H.Kim,N.Kim] . No explicit solutions.

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Part I: Warped AdS_6 solutions in Type IIB (today)

Part II: $\text{AdS}_6/\text{CFT}_5$ – Tests and applications (next week)

Holographic duals for 5d SCFTs

Warped AdS_6 solutions in Type IIB

- Ansatz and general local $AdS_6 \times S^2 \times \Sigma$ solution
- Global solutions on the disc and 5-brane webs
- Solutions with more general Σ ?
- Solutions with $SL(2, \mathbb{R})$ monodromy \rightarrow 7-branes

AdS₆ solutions in Type IIB
– ansatz and local solution –

Symmetries and ansatz

$$\text{AdS}_6 + 16 \text{ susies} \rightarrow \text{F}(4) \supset \text{bosonic } \text{SO}(2,5) \oplus \text{SO}(3)$$

Symmetries and ansatz

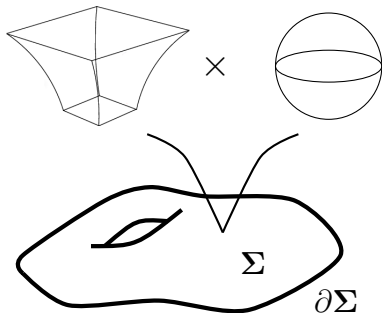
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AdS_6 S^2

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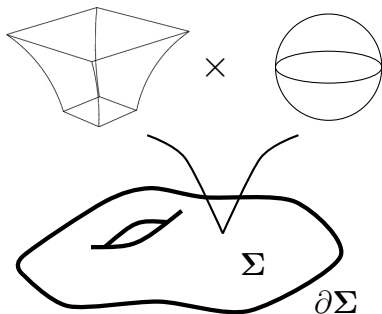


General ansatz: AdS_6 and S^2
warped over Riemann surface Σ

$$\mathcal{M} = (\text{AdS}_6 \times \text{S}^2) \times_{\text{w}} \Sigma$$

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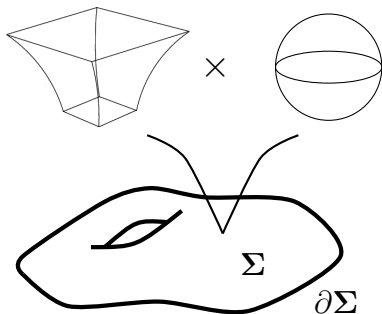
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Remaining bosonic fields: $C_{(4)} = 0$ $C_{(2)} \propto \text{vol}_{S^2}$ $\tau = \chi + ie^{-2\phi}$

General local solution

With complex coordinate w on Σ

$$ds^2 = f_6(w, \bar{w})^2 ds_{\text{AdS}_6}^2 + f_2(w, \bar{w})^2 ds_{\text{S}^2}^2 + 4\rho(w, \bar{w})^2 |dw|^2$$

$$C_{(2)} = \mathcal{C}(w, \bar{w}) \text{vol}_{\text{S}^2} \quad B(w, \bar{w}) = \frac{1 + i\tau(w, \bar{w})}{1 - i\tau(w, \bar{w})}$$

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Preserve 16 supersymmetries \rightarrow BPS eq.

$$\delta\psi_M = D_M\epsilon - \frac{1}{96} (\Gamma_M(\Gamma \cdot G) + 2(\Gamma \cdot G)\Gamma_M) \mathcal{B}^{-1}\epsilon^* \stackrel{!}{=} 0$$

$$\delta\lambda = i(\Gamma \cdot P)\mathcal{B}^{-1}\epsilon^* - \frac{i}{24}(\Gamma \cdot G)\epsilon \stackrel{!}{=} 0$$

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Decomposing Killing spinors, reducing BPS eq. on AdS_6 and S^2
 \rightarrow coupled PDEs on Σ for supergravity fields & Killing spinors

FF ▶▶

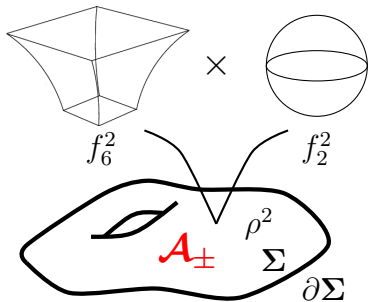
- $\text{AdS}_6 \times \text{S}^2$ gravitino eq. \rightarrow radii in terms of Killing spinors
- dilatino eq. \rightarrow 3-form in terms of B , Killing spinors
- Killing spinors in terms of ρ^2 , B , holomorphic $\partial_w \mathcal{A}_\pm$
- decouple and integrate remaining equations for ρ^2 , B

General local solution

... \rightarrow general local solution to BPS eq., parametrized by two locally holomorphic functions on Σ .

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Arbitrary locally holomorphic

$$\mathcal{A}_\pm : \Sigma \rightarrow \mathbb{C}$$

yield metric functions f_6^2 , f_2^2 , ρ^2 ,
axion-dilaton B , two-form field \mathcal{C}
and Killing spinors solving BPS eq.

$SU(1,1) \otimes \mathbb{C}$ transf. of \mathcal{A}_\pm induce $SL(2, \mathbb{R})$ on supergravity fields

General local solution

Explicit solution for supergravity fields:

$$\begin{aligned}f_6^2 &= \sqrt{6\mathcal{G}} \left[\frac{1+R}{1-R} \right]^{1/2} & f_2^2 &= \frac{1}{9} \sqrt{6\mathcal{G}} \left[\frac{1-R}{1+R} \right]^{3/2} \\ \rho^2 &= \frac{\kappa^2}{\sqrt{6\mathcal{G}}} \left[\frac{1+R}{1-R} \right]^{1/2} & B &= \frac{\partial_w \mathcal{A}_+ \partial_{\bar{w}} \mathcal{G} - R \partial_{\bar{w}} \bar{\mathcal{A}}_- \partial_w \mathcal{G}}{R \partial_{\bar{w}} \bar{\mathcal{A}}_+ \partial_w \mathcal{G} - \partial_w \mathcal{A}_- \partial_{\bar{w}} \mathcal{G}} \\ C &= \frac{4i}{9} \left[\frac{(1+R^2) \partial_w \mathcal{G} \partial_{\bar{w}} \bar{\mathcal{A}}_- - 2R \partial_{\bar{w}} \mathcal{G} \partial_w \mathcal{A}_+}{(1+R)^2 \kappa^2} - \bar{\mathcal{A}}_- - 2\mathcal{A}_+ \right]\end{aligned}$$

with composite quantities

$$\begin{aligned}\kappa^2 &= -|\partial_w \mathcal{A}_+|^2 + |\partial_w \mathcal{A}_-|^2 & \partial_w \mathcal{B} &= \mathcal{A}_+ \partial_w \mathcal{A}_- - \mathcal{A}_- \partial_w \mathcal{A}_+ \\ \mathcal{G} &= |\mathcal{A}_+|^2 - |\mathcal{A}_-|^2 + \mathcal{B} + \bar{\mathcal{B}} & R + \frac{1}{R} &= 2 + 6 \frac{\kappa^2 \mathcal{G}}{|\partial_w \mathcal{G}|^2}\end{aligned}$$

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 \left[\kappa^2 = -|\partial_w \mathcal{A}_+|^2 + |\partial_w \mathcal{A}_-|^2 \right] & \quad \partial_w \mathcal{B} = \mathcal{A}_+ \partial_w \mathcal{A}_- - \mathcal{A}_- \partial_w \mathcal{A}_+ \\
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General local solution

General type IIB supergravity solution with 16 supersymmetries on $\text{AdS}_6 \times S^2$ warped over Σ , in terms of locally holomorphic \mathcal{A}_\pm on Σ .

Solves IIB supergravity equations of motion for arbitrary \mathcal{A}_\pm ✓

[arXiv:1712.04463 Corbino, D'Hoker, CFU]

Generic \mathcal{A}_\pm do not lead to physically regular solutions. ✗

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→ narrow down to globally regular solutions

AdS₆ solutions in Type IIB
– global solutions –

Regularity conditions

Imposing physical regularity conditions imposes constraints. Real geometry with consistent spacetime signature, $\text{Im}(\tau) > 0$:

$$\kappa^2|_{\text{int}(\Sigma)} > 0 \qquad \mathcal{G}|_{\text{int}(\Sigma)} > 0$$

→ Σ must have a boundary ($\partial_w \partial_{\bar{w}} \mathcal{G} = -\kappa^2$ by construction)

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For 10d geometry w/o boundary, collapse S^2 on $\partial\Sigma$ (AdS₆ finite):

$$\kappa^2|_{\partial\Sigma} = 0 \qquad \mathcal{G}|_{\partial\Sigma} = 0$$

Not all independent, $\mathcal{G}|_{\text{int}(\Sigma)} > 0$ implied by the other conditions.

Solving the regularity conditions

Fix topology of Σ , 1) construct locally holomorphic \mathcal{A}_\pm producing regular κ^2 , 2) implement additional constraints for regular \mathcal{G} .

Solving the regularity conditions

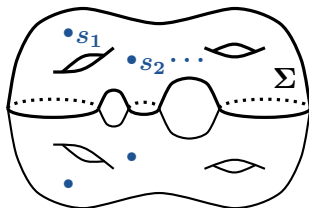
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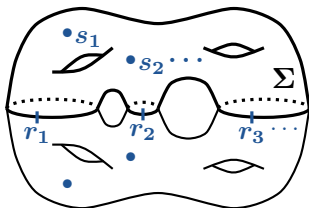
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positive **charges** s_n inside Σ + negative mirror charges

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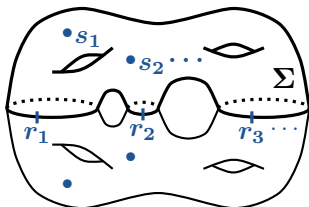


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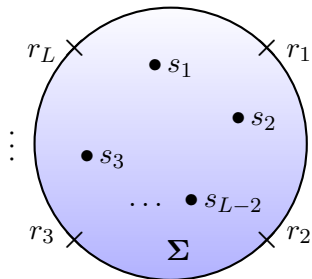
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- 1b) From Φ to $(\partial_w) \mathcal{A}_\pm$: **poles** r_ℓ on $\partial\Sigma$, integ. constants \mathcal{A}_\pm^0
- 2) $\mathcal{G}|_{\partial\Sigma} = 0$: constraints on $\{s_n, r_\ell, \mathcal{A}_\pm^0\}$

Regular solutions on the disc

$\Sigma = \text{disc}/\text{upper half plane}$: L poles $\sim L - 2$ "charges" $\implies L \geq 3$



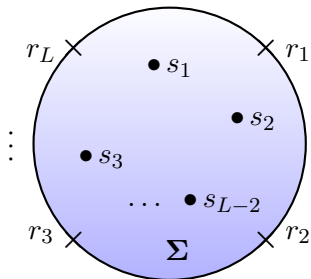
$$\mathcal{A}_{\pm} = \mathcal{A}_{\pm}^0 + \sum_{\ell=1}^L Z_{\pm}^{\ell} \ln(w - r_{\ell})$$

$$Z_{+}^{\ell} = \sigma \prod_{n=1}^{L-2} (r_{\ell} - s_n) \prod_{k \neq \ell}^L \frac{1}{r_{\ell} - r_k}$$

$$\mathcal{A}_{-}(w) = -\overline{\mathcal{A}_{+}(\bar{w})} \quad \sum_{\ell} Z_{+}^{\ell} = 0$$

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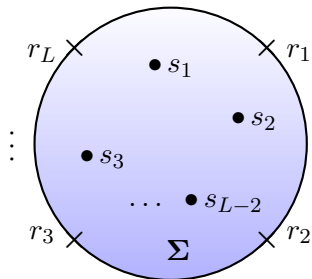
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$\mathcal{G}|_{\partial\Sigma} = 0 \sim$ one local condition per pole $\rightarrow 2L - 2$ free parameters

$$\mathcal{A}_{+}^0 Z_{-}^k - \mathcal{A}_{-}^0 Z_{+}^k + \sum_{\ell \neq k} Z^{[\ell, k]} \ln |p_{\ell} - p_k| = 0$$

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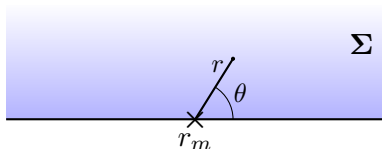
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Solutions regular everywhere, except for possibly the poles...

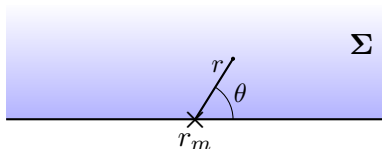
Regular solutions on the disc – behavior near poles



string frame metric, $r \ll 1$:

$$\tilde{d}s^2 \approx \frac{2}{3} |Z_+^m - Z_-^m| \left[3 |\ln r| ds_{\text{AdS}_6}^2 + \overbrace{\frac{dr^2}{r^2} + d\theta^2}^{\Sigma} + \sin^2 \theta ds_{\mathbb{S}^2}^2 \right]$$

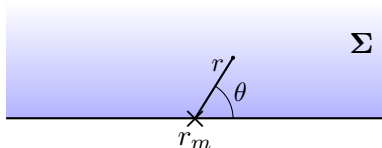
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string frame metric, $r \ll 1$:

$$\tilde{d}s^2 \approx \frac{2}{3} |Z_+^m - Z_-^m| \left[\underbrace{3 |\ln r| ds_{\text{AdS}_6}^2}_{\rightarrow \mathbb{R}^{1,5}} + \overbrace{\frac{dr^2}{r^2} + d\theta^2 + \sin^2 \theta ds_{\mathbb{S}^2}^2}^{\Sigma} \right]$$

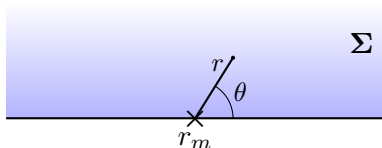
Regular solutions on the disc – behavior near poles



string frame:
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$$dC_{(2)} \approx \frac{8}{3} Z_+^m \text{vol}_{S^3} \quad e^{-2\phi} \approx \frac{\sqrt{3} \kappa_m^2}{4 \text{Re}(Z_+^m)^2} \frac{r}{\sqrt{|\ln r|}} \quad \chi \approx \frac{\text{Im}(Z_+^m)}{\text{Re}(Z_+^m)}$$

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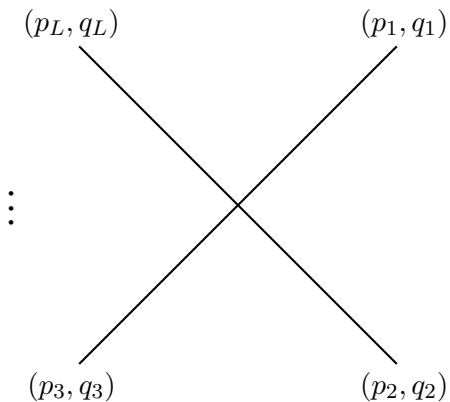
Entire near-pole solution matches (p, q) 5-branes of [Lu,Roy '98]

$$q + ip \longleftrightarrow Z_+^m$$

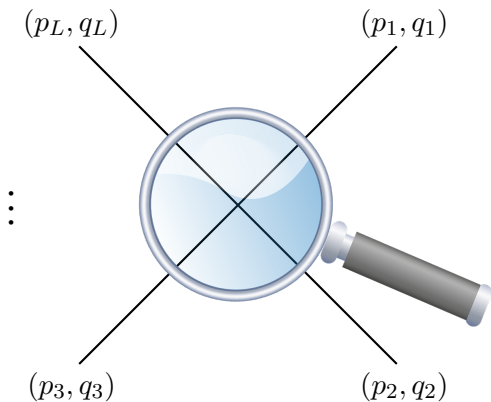
Solutions regular with isolated poles corresponding to 5-branes ✓

AdS₆ solutions in Type IIB
– connection to 5-brane webs –

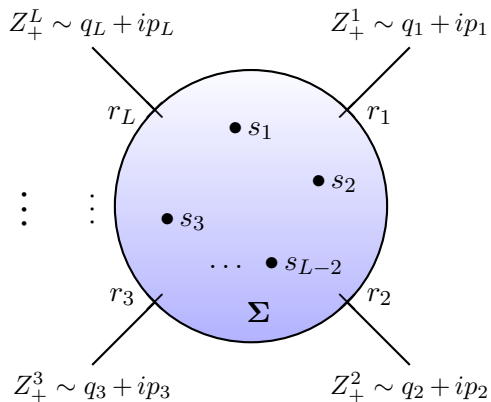
5-brane web picture



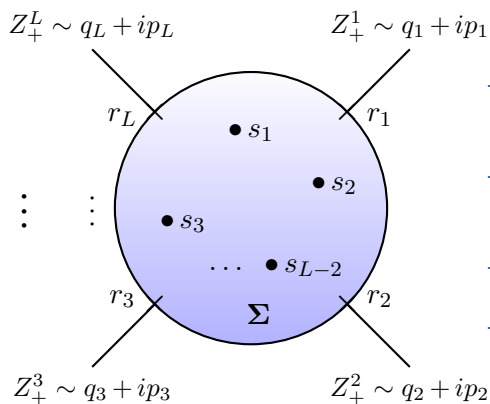
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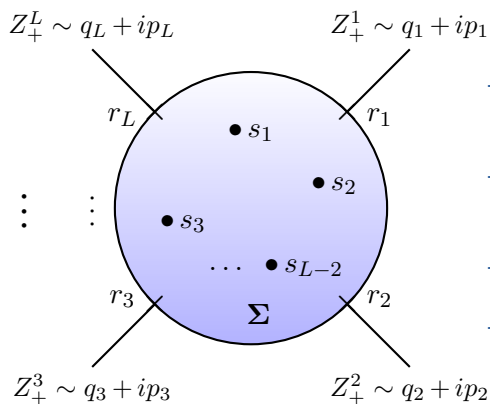


5-brane web picture



- external 5-branes explicitly (p, q) charge conserved
- parametrized by choice of residues mod charge cons.
- $\text{AdS}_6 + 16 \text{ susies} = F(4)$
- need $L \geq 3$, p and q charge

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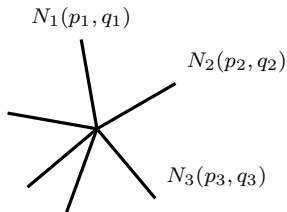
Supergravity solutions for fully localized 5-brane intersections



“Large N ”

Classical supergravity: $Z_+^\ell \sim q_\ell + ip_\ell \in \mathbb{C}$. Limit of string theory with large 5-brane charges \rightarrow effectively continuous.

Large numbers of like-charged coincident 5-branes at each pole



$p_i, q_i \in \mathbb{Z}$, relatively prime

$N_i \gg 1 \forall i$

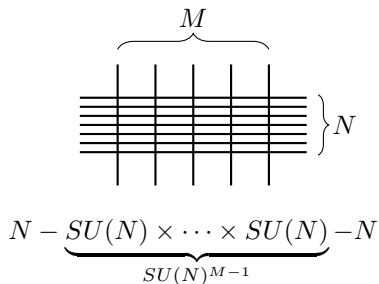
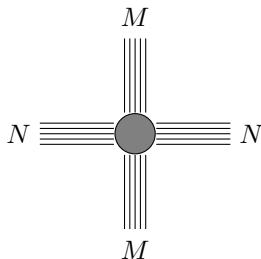
$Z_+^\ell \sim N_\ell q_\ell + ip_\ell$

In particular: generally large D5 and NS5 brane charges.

SCFT/gauge theory connections

NS5/D5 intersection

[Aharony, Hanany, Kol '97]



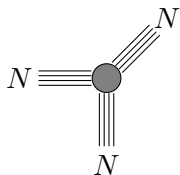
\leftrightarrow 4-pole solution with $Z_+^1 = -Z_+^3 \sim iN$, $Z_+^2 = -Z_+^4 \sim M$

large D5 and large NS5 charge \leftrightarrow large number of nodes in quiver deformation and (at least some) large-rank gauge groups

SCFT/gauge theory connections

5d T_N theories: junction of N D5, N NS5 and N (1, 1) 5-branes

[Benini, Benvenuti, Tachikawa '09]



– Reduce on S^1 to 4d $T[A_{N-1}]$

– IR gauge theory [Bergman, Zafrir '14]

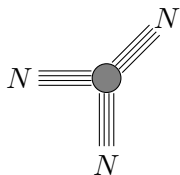
$$N - SU(N - 1) - \cdots - SU(2) - 2$$

\leftrightarrow 3-pole solution with $Z_+^1 \sim iN$, $Z_+^2 \sim N$, $Z_+^3 \sim (1 + i)N$

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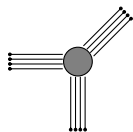
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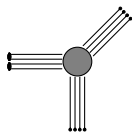
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Potential subtlety: s -rule constraints from 5-branes on 7-branes



vs.



$$SU(4)^3$$

vs.

$$SU(4)^2 \times SU(2) \rightarrow E_7$$

Supergravity sol. \sim unconstrained junctions (details next week).

Global solutions on the disc

Physically regular solutions for $\Sigma = \text{disc}$, $L \geq 3$ poles on $\partial\Sigma$,
 $2L - 2$ real parameters = residues at poles with zero sum.

Poles required by regularity, correspond to (p, q) 5-branes with
 $q + ip \sim Z_+^\ell$. Natural identification with 5-brane junctions.

Recipe for Σ w/ arbitrary numbers of handles and boundaries.

AdS₆ solutions in Type IIB

– More general Riemann surfaces? –

More general Riemann surfaces?

Construction of regular κ^2 works for general Riemann surfaces.
Boundary condition $\mathcal{G}_{\partial\Sigma} = 0$ not automatic.

Two natural options: (i) add handles (ii) add further boundaries

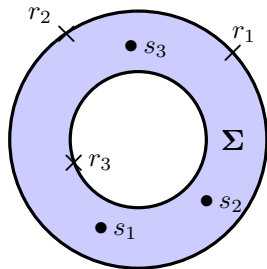
Technically more challenging. Charge distribution on doubled surface + Green's function $\rightarrow \kappa^2$, Green's functions not as simple.

More general Riemann surfaces?

Next-to-simplest option: annulus, no handles, two boundaries

Doubled surface=torus \rightarrow quasi-periodic Jacobi ϑ -functions in Green's function

$L \geq 2$ poles, N "charges"

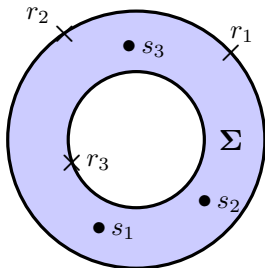


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$$\mathcal{A}_+(w|\tau) = \mathcal{A}_+^0 + \alpha_+ w + \sum_{\ell=1}^L Z_+^\ell \ln \vartheta_1(w - r_\ell|\tau)$$

$$Z_+^\ell = \sigma \frac{\prod_{n=1}^N \vartheta_1(r_\ell - s_n|\tau)}{\prod_{n \neq \ell} \vartheta_1(r_\ell - p_n|\tau)} \exp \left\{ -\frac{2\pi i}{\tau} r_\ell \Lambda_+ - \frac{\pi i}{\tau} \sum_{n=1}^N (r_n^2 + r_n) \right\}$$

More general Riemann surfaces?

Main challenge: $\mathcal{G} = 0$ on two boundaries \Rightarrow non-local condition.

Necessary for $\mathcal{G}|_{\partial\Sigma} = 0$: \mathcal{G} constant along regular boundary segments, no monodromy around poles \sim local conditions.

This is sufficient for a single boundary component:

$$\mathcal{G} = |\mathcal{A}_+|^2 - |\mathcal{A}_-|^2 + \mathcal{B} + \bar{\mathcal{B}} \quad \partial\mathcal{B} = \mathcal{A}_+ \partial\mathcal{A}_- - \mathcal{A}_- \partial\mathcal{A}_+$$

\mathcal{B} defined up to integration constant \rightarrow use it to set $\mathcal{G} = 0$.

Annulus needs extra constraint, connecting both boundaries.

Found no solutions for $L = 2, 3, 4, 5$

X

AdS₆ solutions in type IIB
– solutions with monodromy –

5-brane webs with 7-branes

[DeWolfe, Hanany, Iqbal, Katz '99]

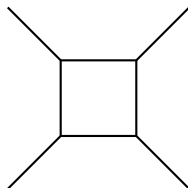
Supergravity fields single-valued in solutions so far. Could allow for $SL(2, \mathbb{R})$ monodromy, induced by simple action on $\mathcal{A}_{\pm} \Rightarrow$ 7-branes

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	0	1	2	3	4	5	6	7	8	9
D5	x	x	x	x	x	x				
NS5	x	x	x	x	x		x			
D7	x	x	x	x	x			x	x	x

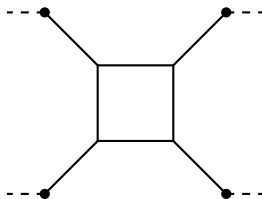


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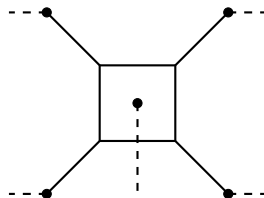
(i) terminate external (p, q) 5-branes on $[p, q]$ 7-branes

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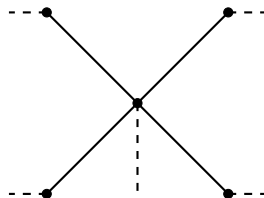
- (i) terminate external (p, q) 5-branes on $[p, q]$ 7-branes
- (ii) add 7-branes into faces of web

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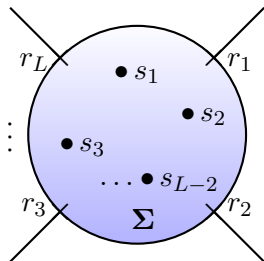


- (i) terminate external (p, q) 5-branes on $[p, q]$ 7-branes
- (ii) add 7-branes into faces of web

In conformal and “near-horizon” limit, (ii) should be represented in corresponding supergravity solution.

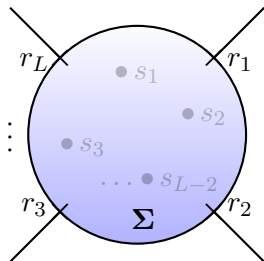
Solutions with monodromy

Construction from disc solutions without monodromy:



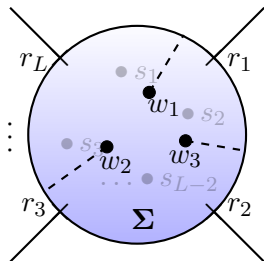
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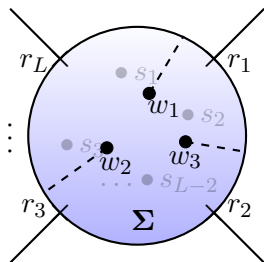
- pick parabolic $SL(2, \mathbb{R})$ matrix

$$M_{[p,q]} = \begin{pmatrix} 1 - pq & p^2 \\ -q^2 & 1 + pq \end{pmatrix}$$

- add punctures w_i , “charge” n_i , orientation of branch cut γ_i

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Encode branch cut structure in $f = \sum_{i=1}^I \frac{n_i^2}{4\pi} \ln \left(\gamma_i \frac{w - w_i}{w - \bar{w}_i} \right)$.

Differentials $\partial \mathcal{A}_{\pm} = \partial \mathcal{A}_{\pm}^s + \eta_{\pm} f (\eta_- \partial \mathcal{A}_{+}^s - \eta_+ \partial \mathcal{A}_{-}^s)$, $\eta_{\pm} = p \mp iq$.

Solutions with monodromy

Differentials $\partial\mathcal{A}_{\pm}$ realize commuting $SL(2, \mathbb{R})$ monodromies $M_{[n_i p, n_i q]}$ around w_i ; sum of residues not necessarily zero.

Locally holomorphic functions \mathcal{A}_{\pm} involve “polylogarithm integrals”
Regular κ^2 by construction, $\mathcal{G}|_{\partial\Sigma} = 0$ still one condition per pole.

Demanding monodromies to lift to $SL(2, \mathbb{R})$ transformations on \mathcal{A}_{\pm} yields further constraints, forcing punctures onto curves in Σ .

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→ regular solutions for disc with punctures and commuting parabolic $SL(2, \mathbb{R})$ monodromies

Solutions with monodromy

At punctures, solution as expected for 7-brane wrapping $\text{AdS}_6 \times S^2$.
Infinitesimal monodromy: recover probe D7 (DBI, κ symmetry) ✓

Additional parameters due to punctures:

charge ✓ branch cut orientation ✓ position in Σ ??

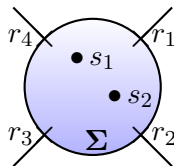
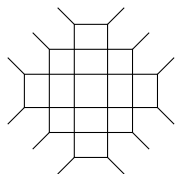
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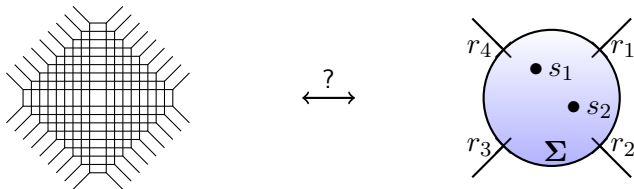
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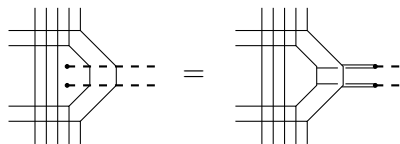
Internal structure of brane web encoded in Σ ?



Dense grid at large N . Position in $\Sigma \leftrightarrow$ choice of face.
Continuous parameter, remains meaningful in conformal limit.

Solutions with monodromy

Hanany/Witten: pull 7-branes out of web \rightarrow 5-brane creation

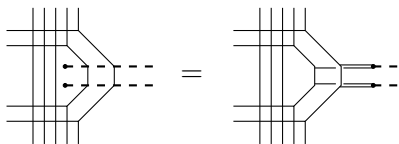


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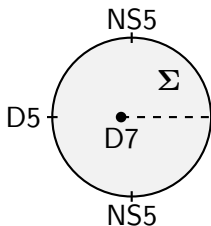
Exception: 7-branes “trapped” by only one NS5 brane



Same as unconstrained 5-brane junction with additional D5s

Solutions with monodromy

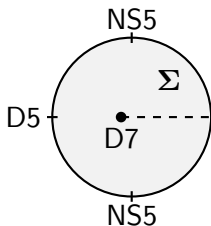
Supergravity picture: family of 3-pole solutions



- two NS5 poles one D5 pole with *fixed* residues
- one D7 puncture, on “equator” charge depends on location

Solutions with monodromy

Supergravity picture: family of 3-pole solutions



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- one D7 puncture, on “equator” charge depends on location

In general, on-shell action (=free energy) different from 4-pole solution w/o monodromy.

As $D7 \rightarrow \partial\Sigma$ w/ fixed 5-brane charges, puncture turns into pole.

Consistent with 5-brane webs just shown. Confirmation: next week.

Summary & Outlook

Summary

Supergravity solutions for fully localized 5-brane junctions in type IIB. Holographic duals for the corresponding 5d SCFTs.

Solutions regular everywhere except for isolated physically meaningful singularities.

Extension to solutions with parabolic $SL(2, \mathbb{R})$ monodromy, for 5-brane junctions with (mutually local) 7-branes.

Outlook

Establish $\text{AdS}_6/\text{CFT}_5$ dualities rigorously: confront holography with QFT results (“stringy” operators, $F(S^5) \rightarrow$ [next week](#)).

AdS/CFT with warped products: spectra, correlation functions, defects, Wilson lines, truncation to 6d $F(4)$ gauged sugra...

Solutions with non-commuting $\text{SL}(2, \mathbb{R})$ monodromies for mutually non-local 7-branes, more general Riemann surfaces, ...

$\text{AdS}_2 \times S^6$: local solution in [\[arXiv:1712.04463\]](#), global solutions?

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Thank you!