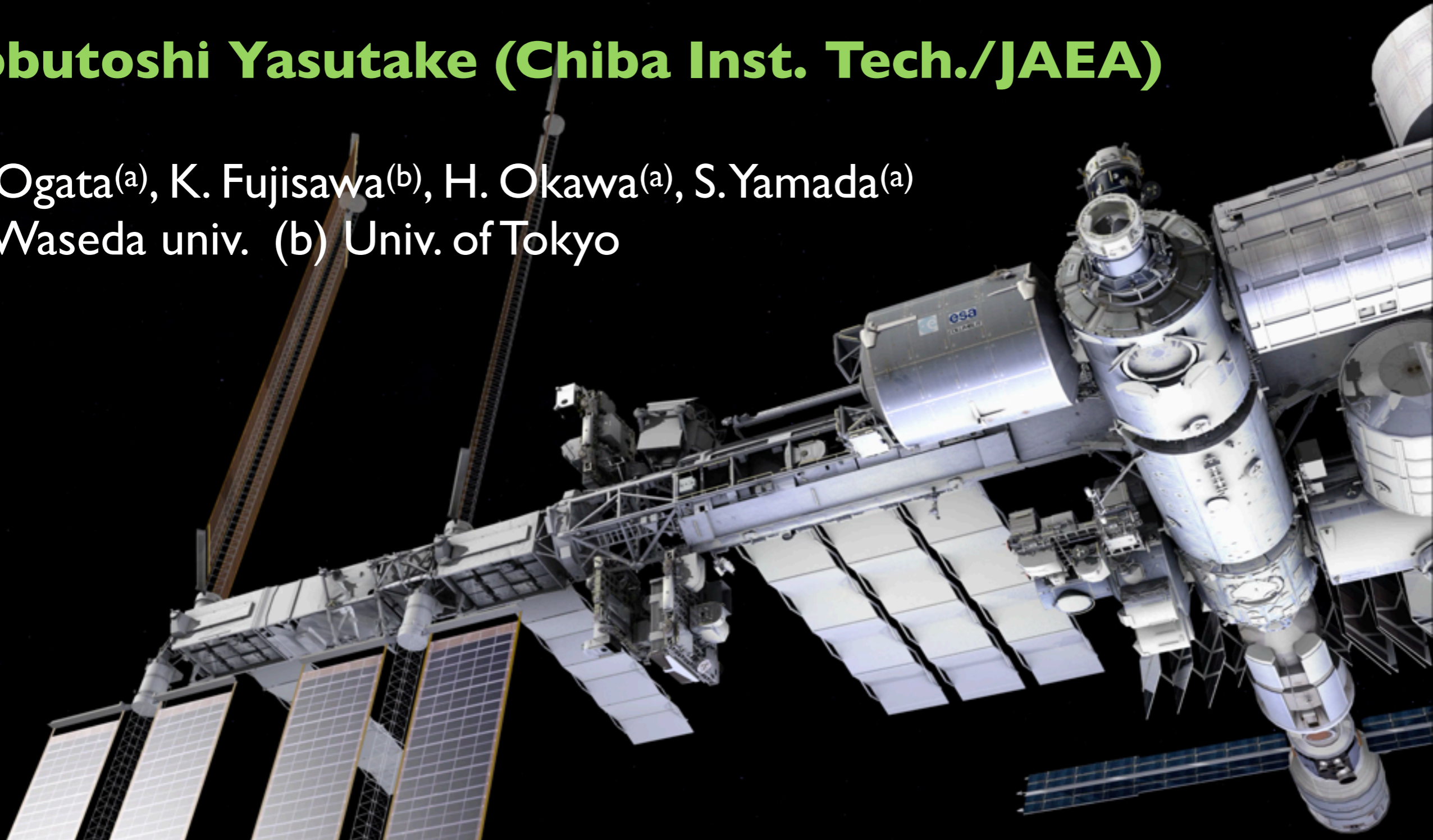


Hydrostatic equilibria of rotating stars in Lagrangian coordinate

Nobutoshi Yasutake (Chiba Inst. Tech./JAEA)

M. Ogata^(a), K. Fujisawa^(b), H. Okawa^(a), S. Yamada^(a)

(a) Waseda univ. (b) Univ. of Tokyo

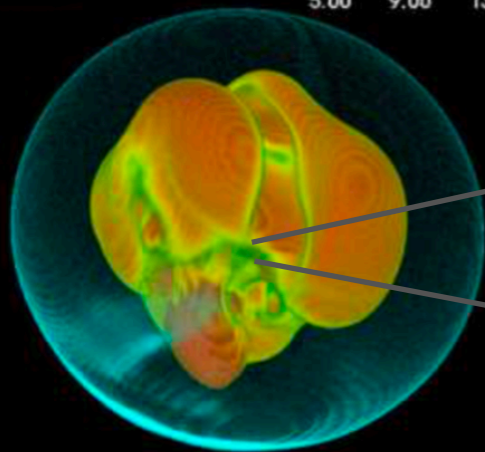


A big problem in astrophysics

There is no scheme of time evolution for "deformed stars" in quasi-equilibrium !!

**Supernovae
dynamical simulation(3D)**

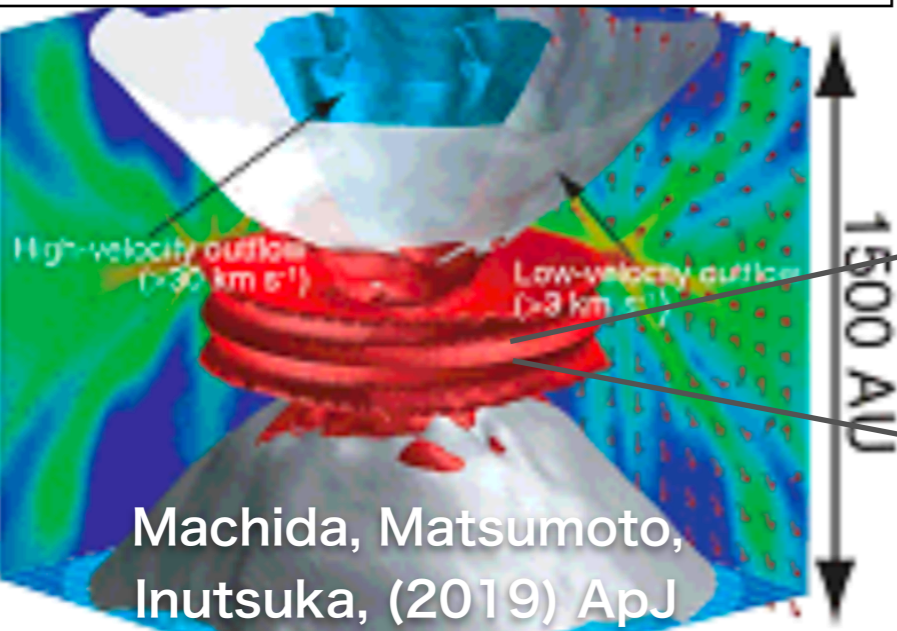
tpb=998 ms 5.00 9.00 13.0 17.0



K. Nakamura, T. Takkiwaki,
K.Kotake (2019) PASJ

10000 km

**Star formations
dynamical simulation(3D)**



Machida, Matsumoto,
Inutsuka, (2019) ApJ

PNSs ($R \sim 50\text{km}$)



$T \sim 30\text{MeV}$

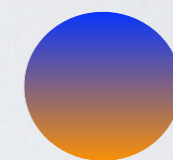
$Y_L \sim 0.3$

C. Jacobi, R. Dedekind, P.L. Dirichlet, B. Riemann,
Chandrasekhar...

evolution (1D)

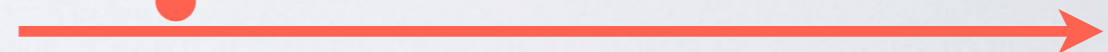
Heyney method (1964) in GR

NSs ($R \sim 10\text{km}$)



$T \sim 0\text{ MeV}$

$Y_v \sim 0$

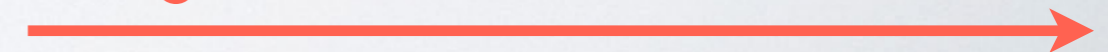


evolution (1D)

Heyney method (1964)

**Massive Stars,
Planets, ... etc**

Proto-stars



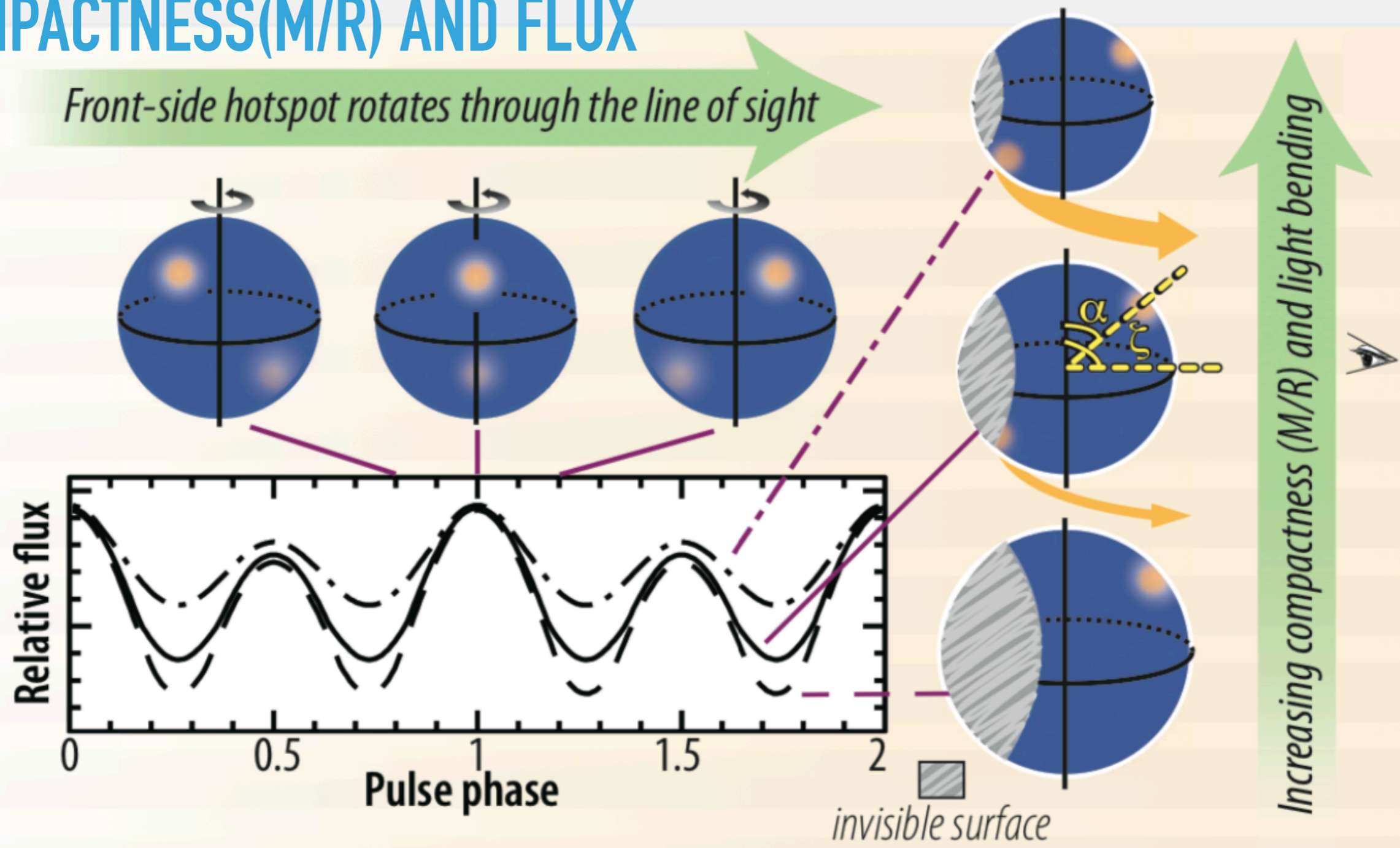
NICER

Neutron star Interior Composition Explorer



COMPACTNESS (M/R) AND FLUX

Thermal Lightcurve Model



(C) KEITH GENDREAW@NASA

Thermal evolutions with magnetic field in 2,3D

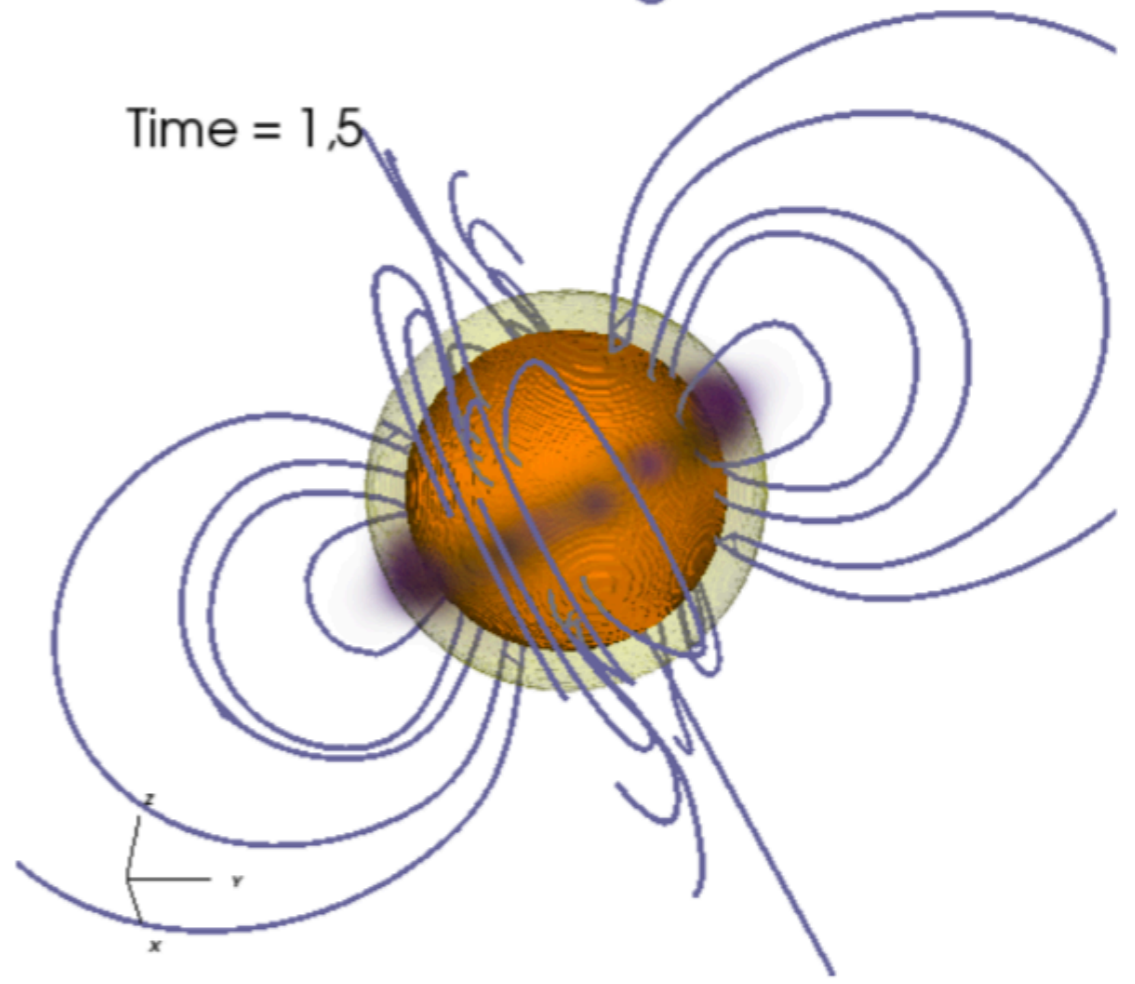
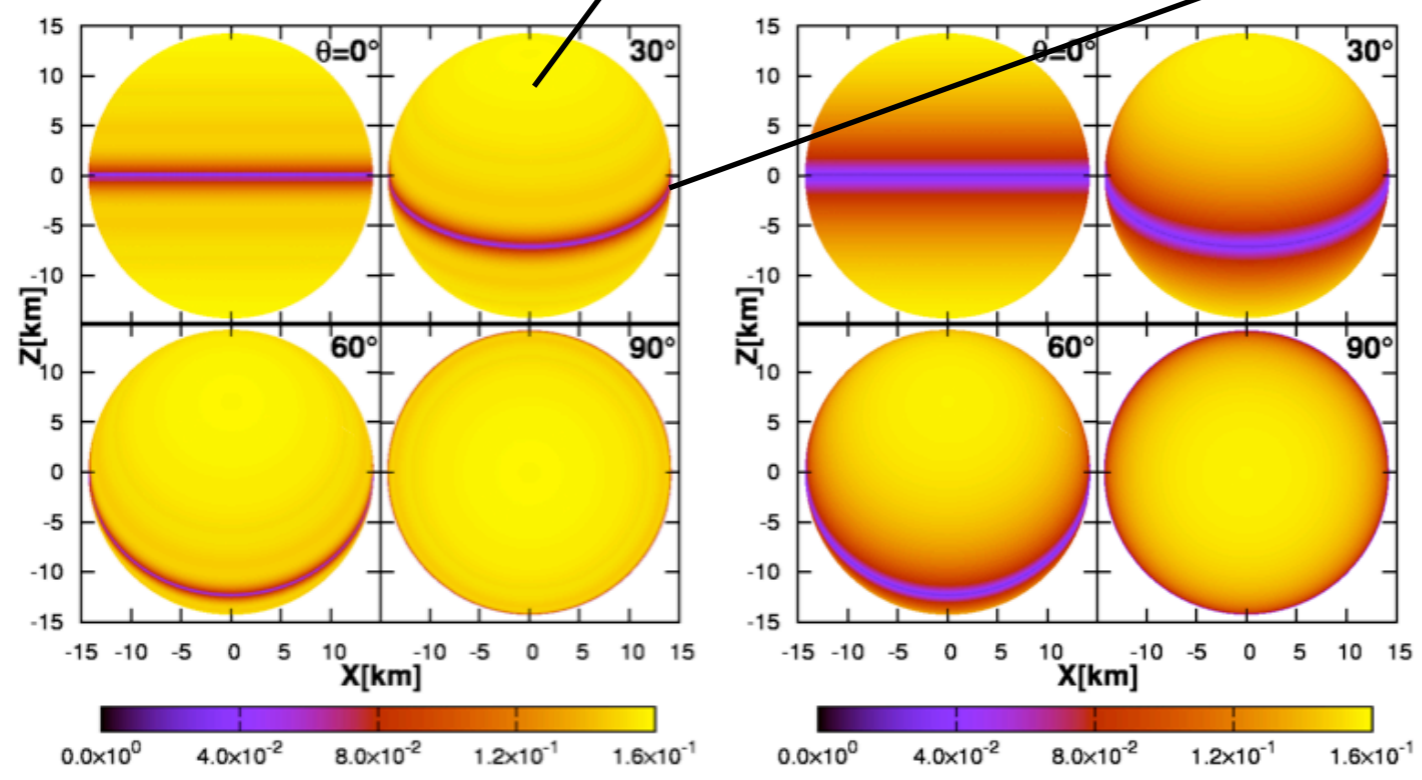
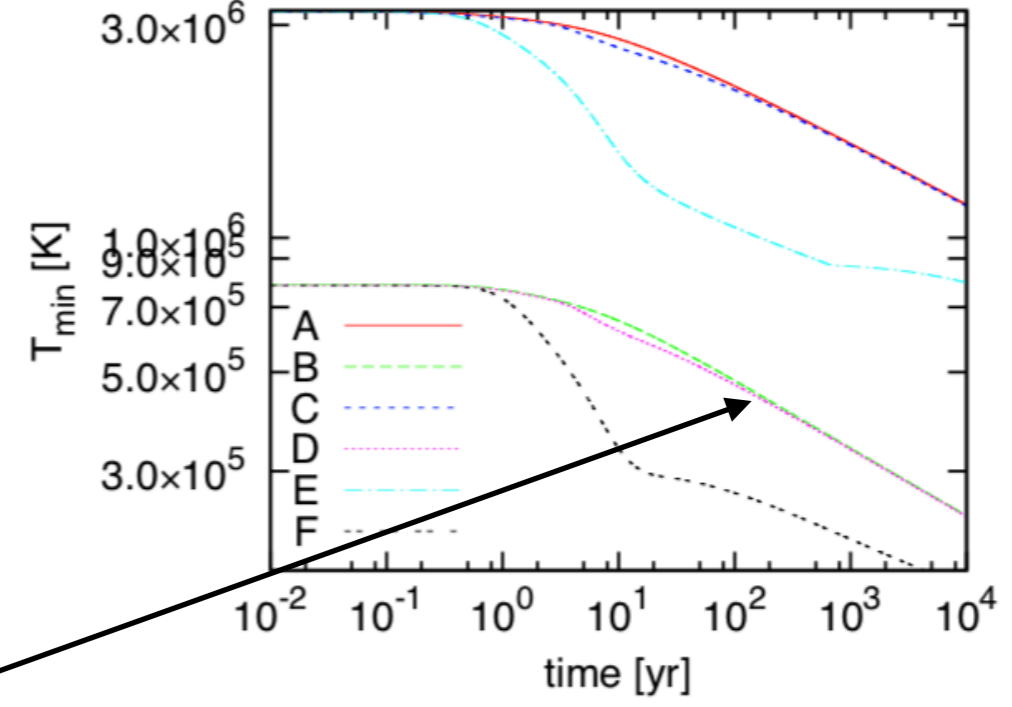
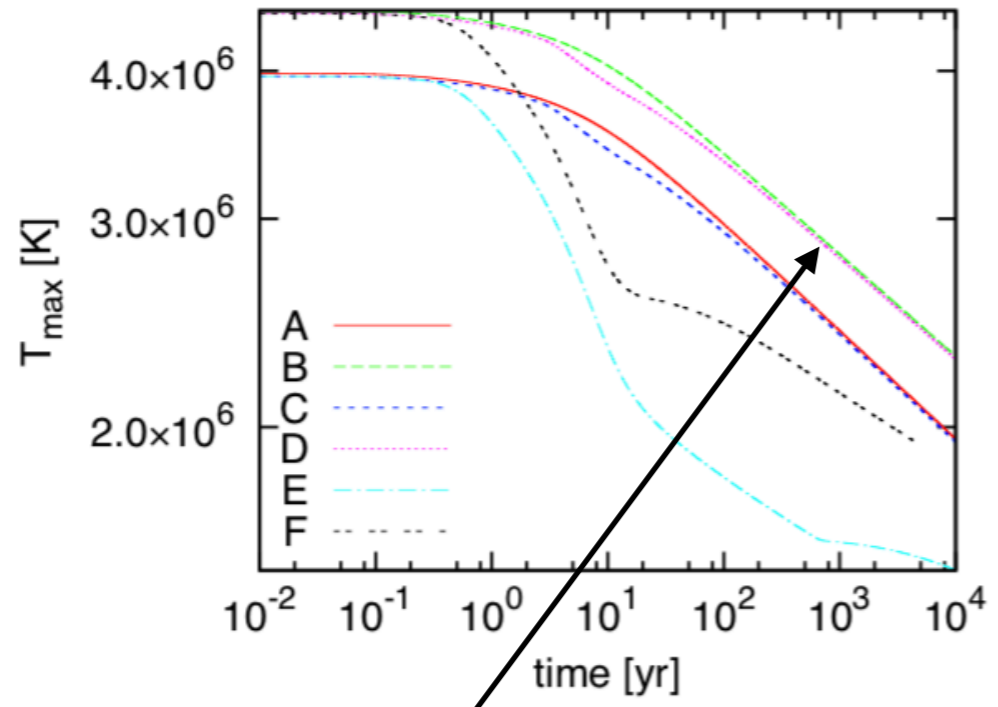
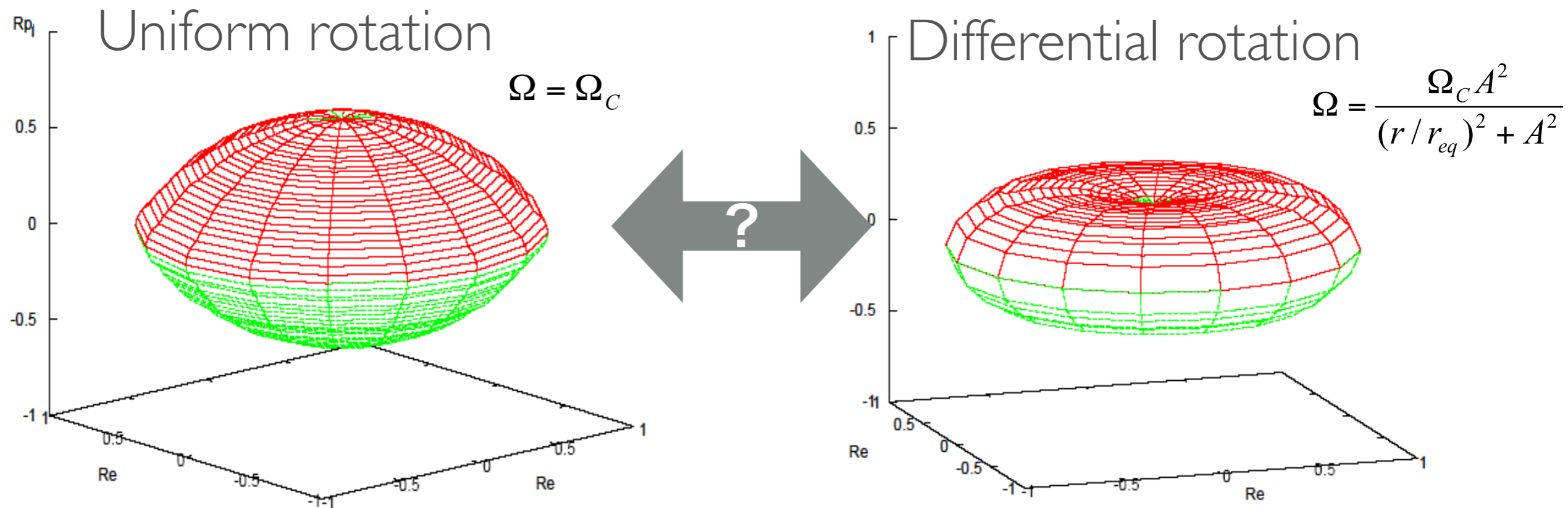


FIG. 7: (Color online) Temperature distribution for model "mSUK" after 10^4 years depended on the inclination angle θ . The unit of color contour is [keV].

Without rotational law, no one can get rotating structures.

Rotating structures of neutron stars in fully GR formulation

NY, Hashimoto, Eriguchi, PTP(2005)



Komatsu, Eriguchi, Hachisu(1989), Cook, Shapiro, Teukolsky(1992)

Shapiro, Teukolsky, Nakamura(1994), Gourgoulhon et al.(1999)

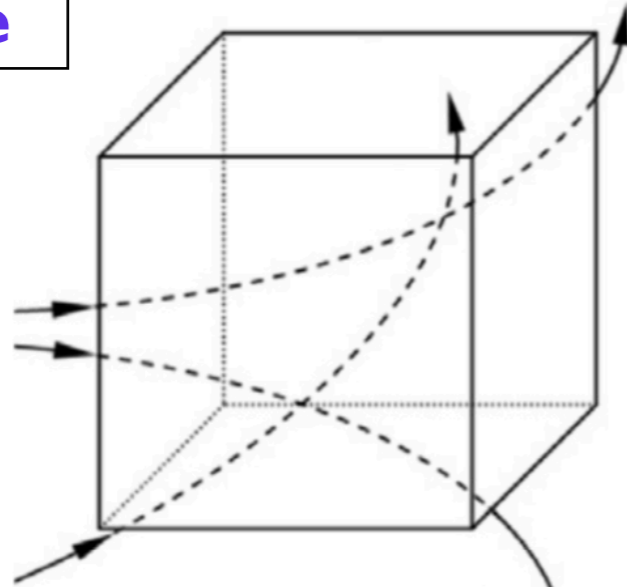
Usui, Uryu, Eriguchi(1999)....Uryu et al.(2017)

also in E. Zhou's talk

SPACE(+TIME) OR PARTICLES

supernovae

Eulerian

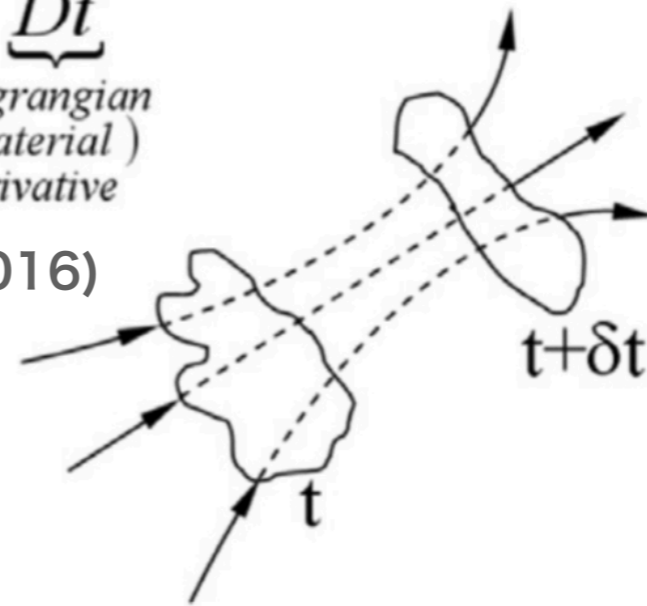


Spatially fixed volume element

$$\underbrace{\frac{\partial}{\partial t} + \mathbf{v} \cdot \nabla}_{\text{Eulerian derivative}} = \underbrace{\frac{D}{Dt}}_{\text{Lagrangian (Material) derivative}}$$

Shadloo et al. (2016)

Lagrangian

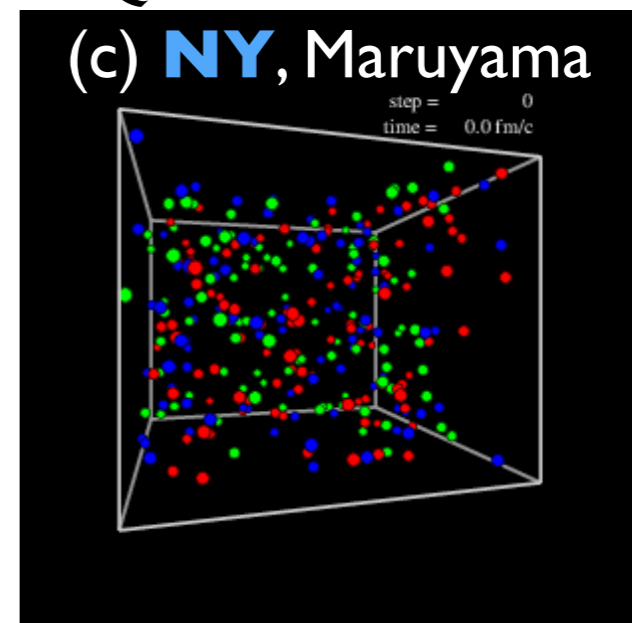


stellar evolutions

Following the motion of the fluid element

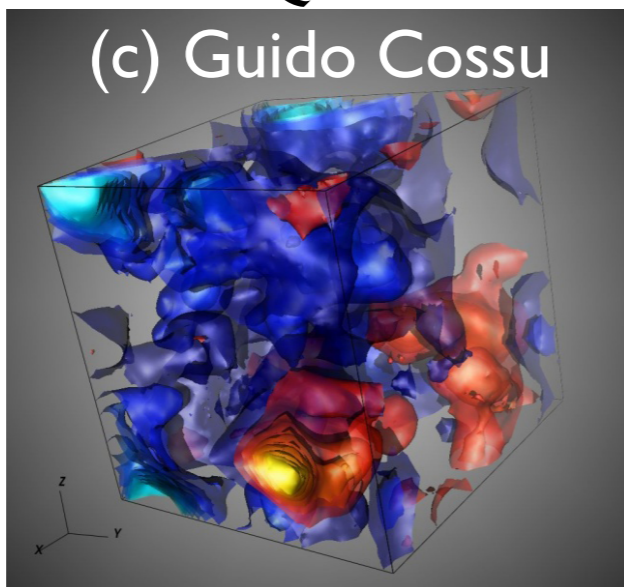
QMD simulation

(c) NY, Maruyama



LQCD

(c) Guido Cossu



c.f.)

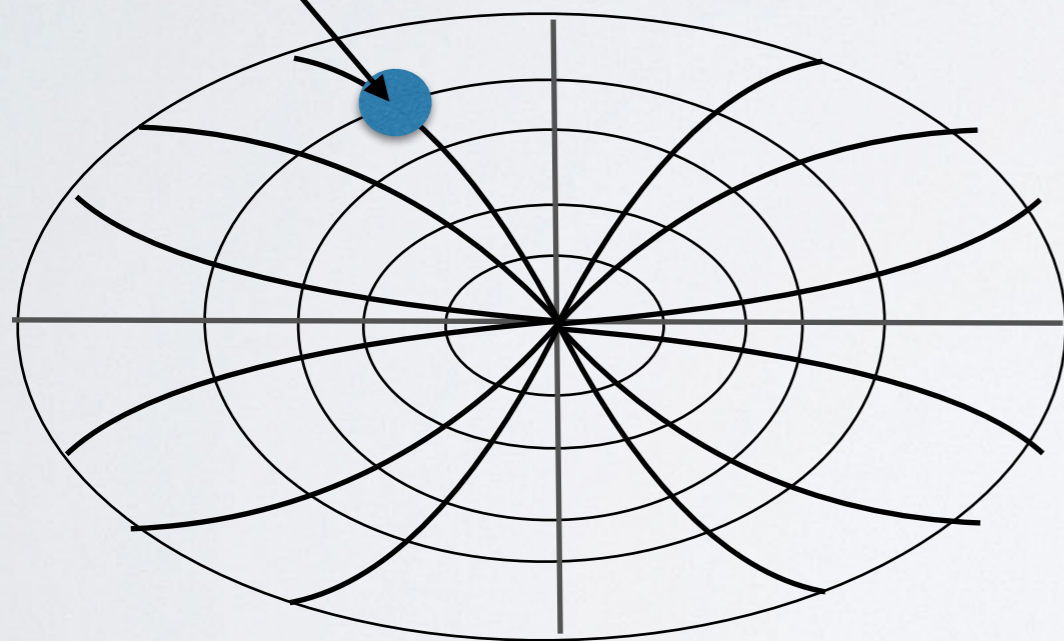
HOW TO SOLVE ?

IN LAGLANGIAN COORDINATE

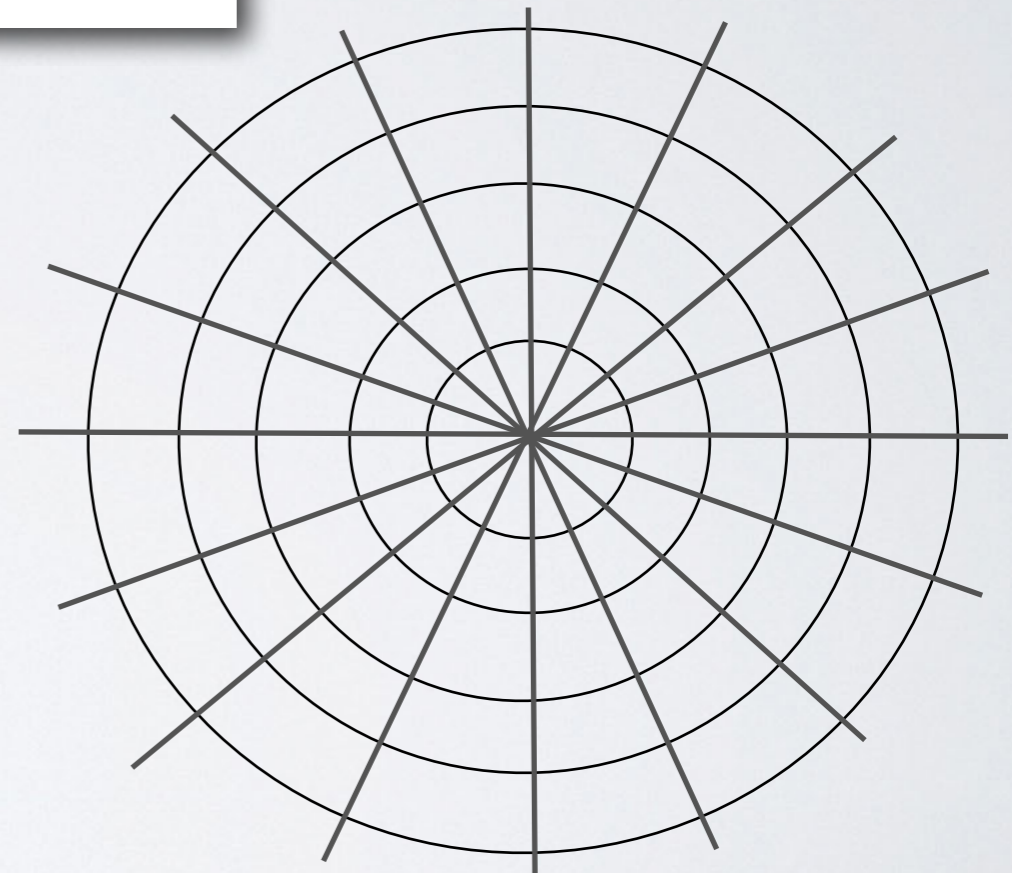
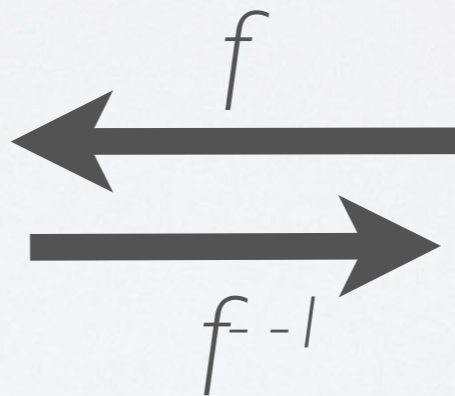
We have to solve this eq. in mass-coordinate for baroclinic cases with arbitrary angular momentum profiles.

$$\nabla P + \rho \nabla \phi - \frac{\rho j^2}{\varpi^3} \mathbf{e}_\varpi = 0, \quad (\text{Newtonian case})$$

(m, j, s)



$$dV = |J| dV'$$



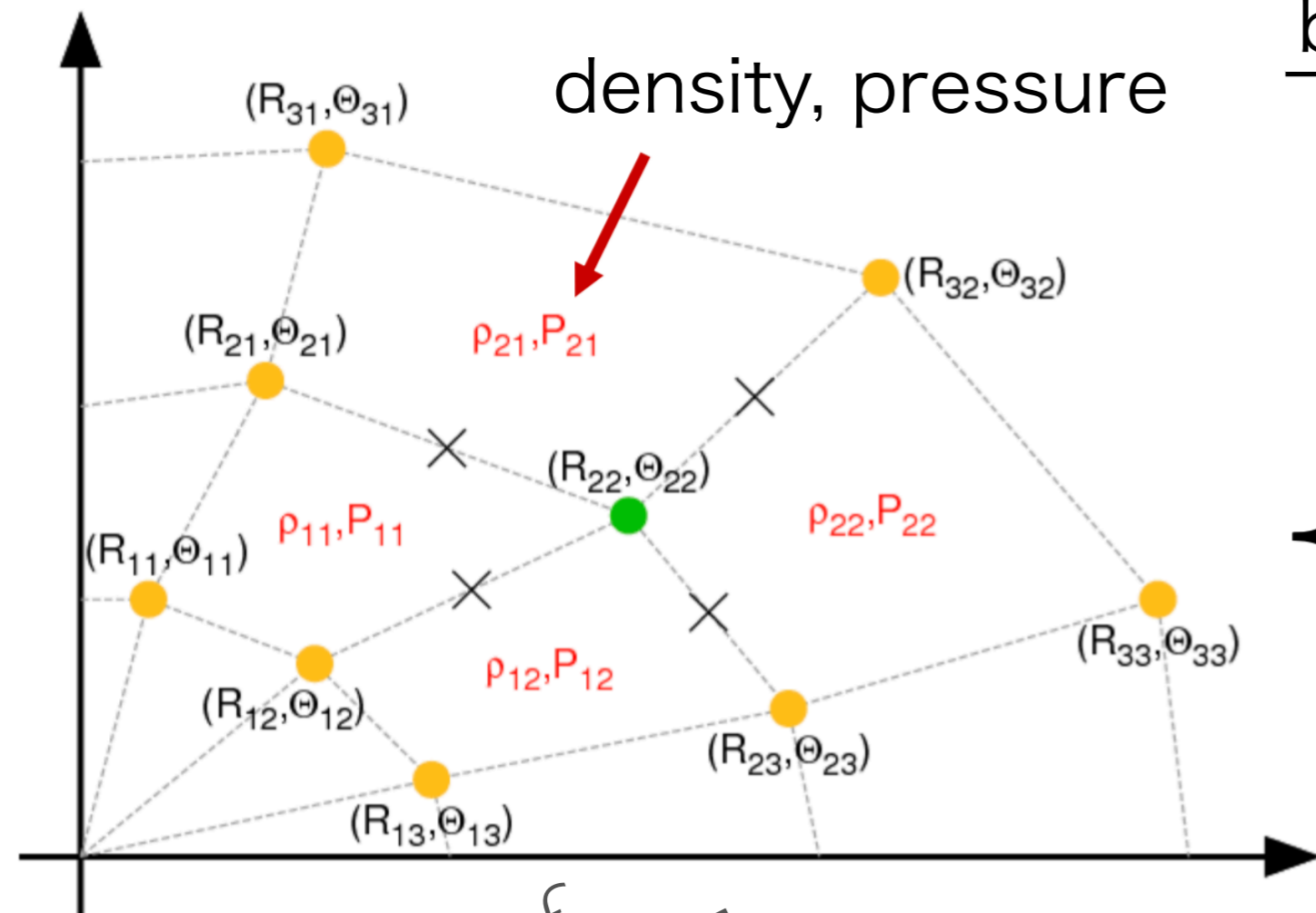
Problems

- ① Uncertainty because of the gauge freedom. (Freedman&Shutz 1978)
- ② Numerical problems cf. Hour-glass problem

① Simple discretization

Ogata, Fujisawa, Okawa, **NY**, Yamada

$$\frac{\partial P}{\partial r} \rightarrow \frac{P_{i+1} - P_i}{r_{i+1} - r_i}$$

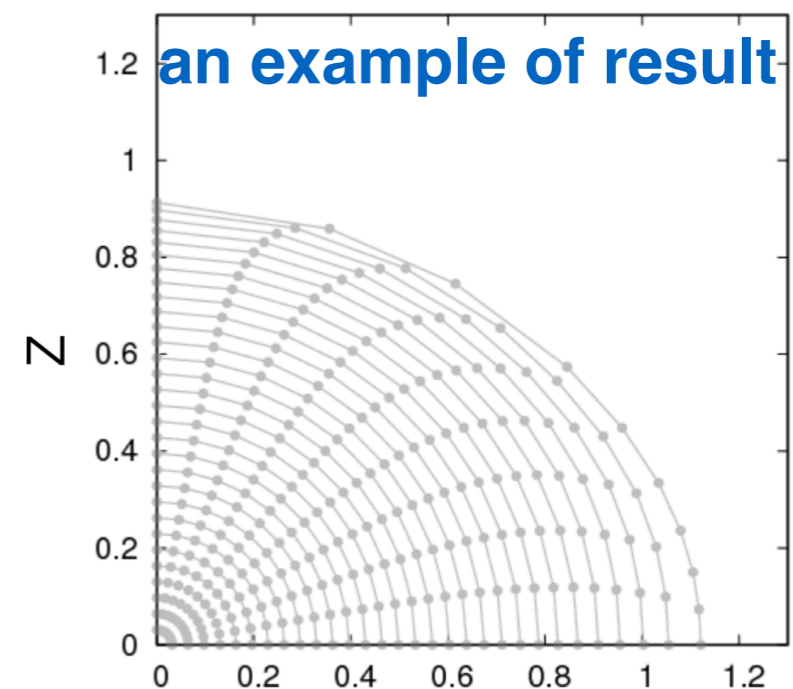
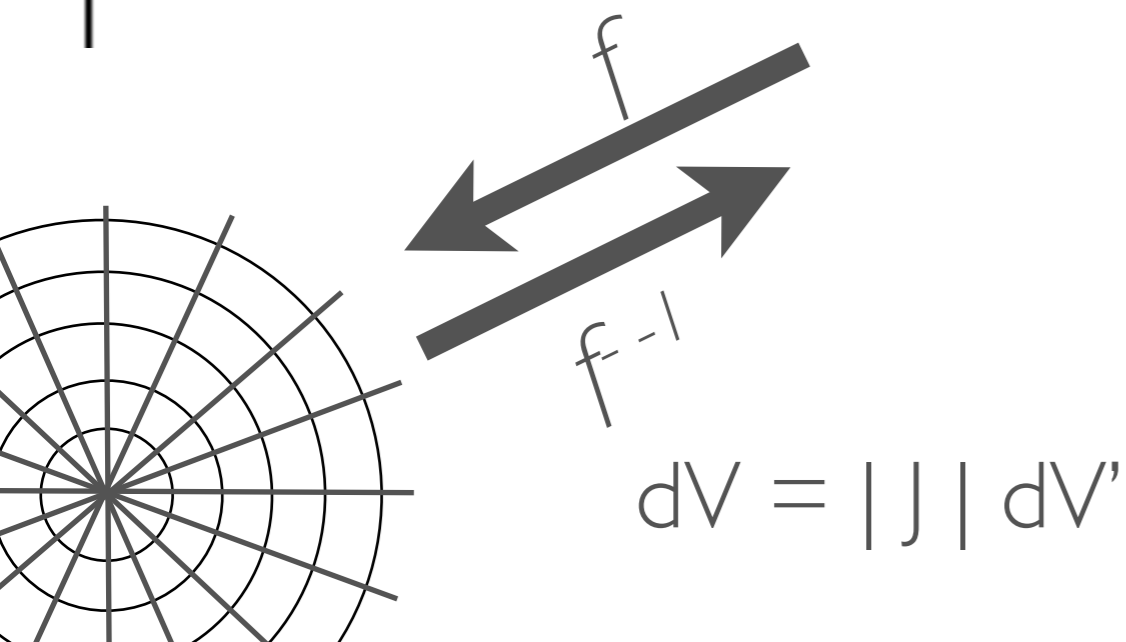


balance equation

$$\frac{1}{\rho} \nabla P - \nabla \phi + \omega^2 (R \sin \theta) \mathbf{e}_\omega = \mathbf{0}$$



$$\begin{cases} F_R = \frac{1}{\rho} \frac{\partial P}{\partial R} - \frac{\partial \phi}{\partial R} + \omega^2 (R \sin \Theta) \sin \Theta = 0 \\ F_\Theta = \frac{1}{\rho R} \frac{\partial P}{\partial \Theta} - \frac{\partial \phi}{R \partial \Theta} - \omega^2 (R \sin \Theta) \cos \Theta = 0 \end{cases}$$



② Spectral method

Fujisawa, Ogata, Okawa, **NY**, Yamada

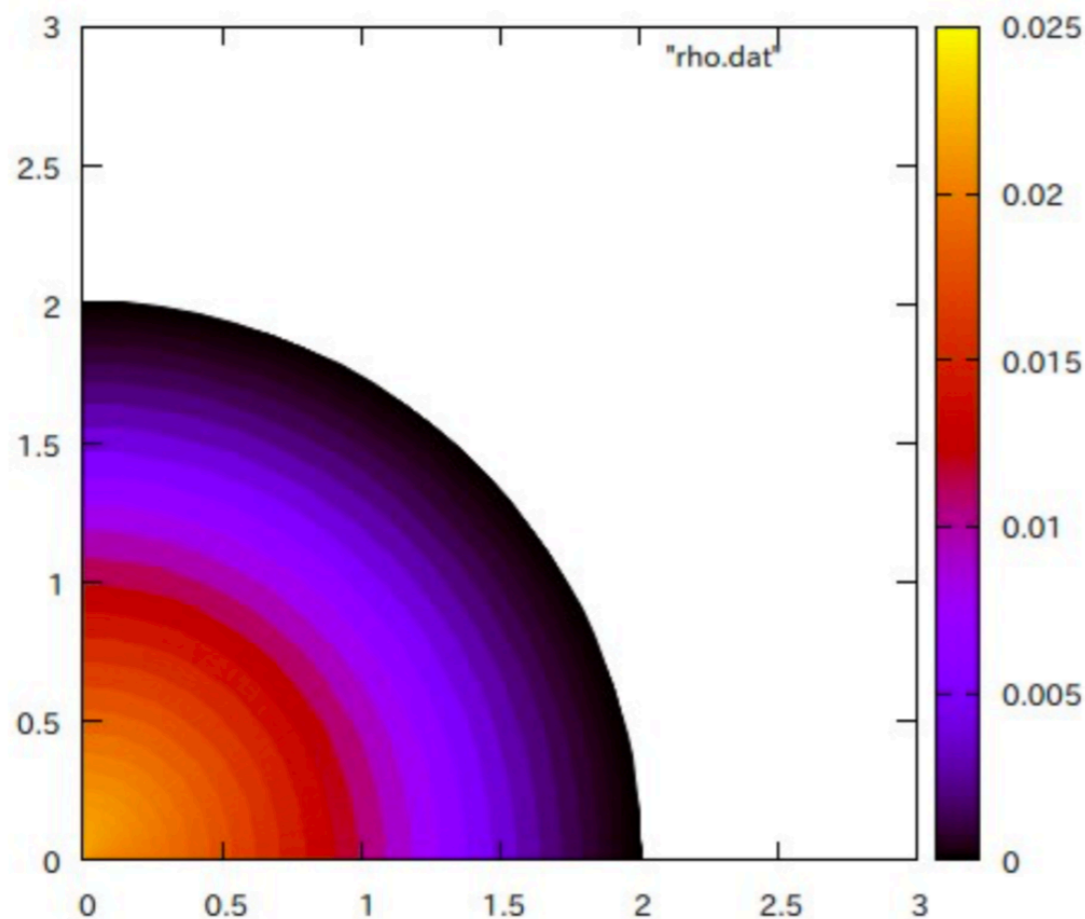
Chebyshev's expansions

$$X(r, \nu) = \sum_{n=0}^{N_{\max}} \sum_{\ell=0}^{L_{\max}} a_{n\ell} T_n(r) T_\ell(\nu)$$

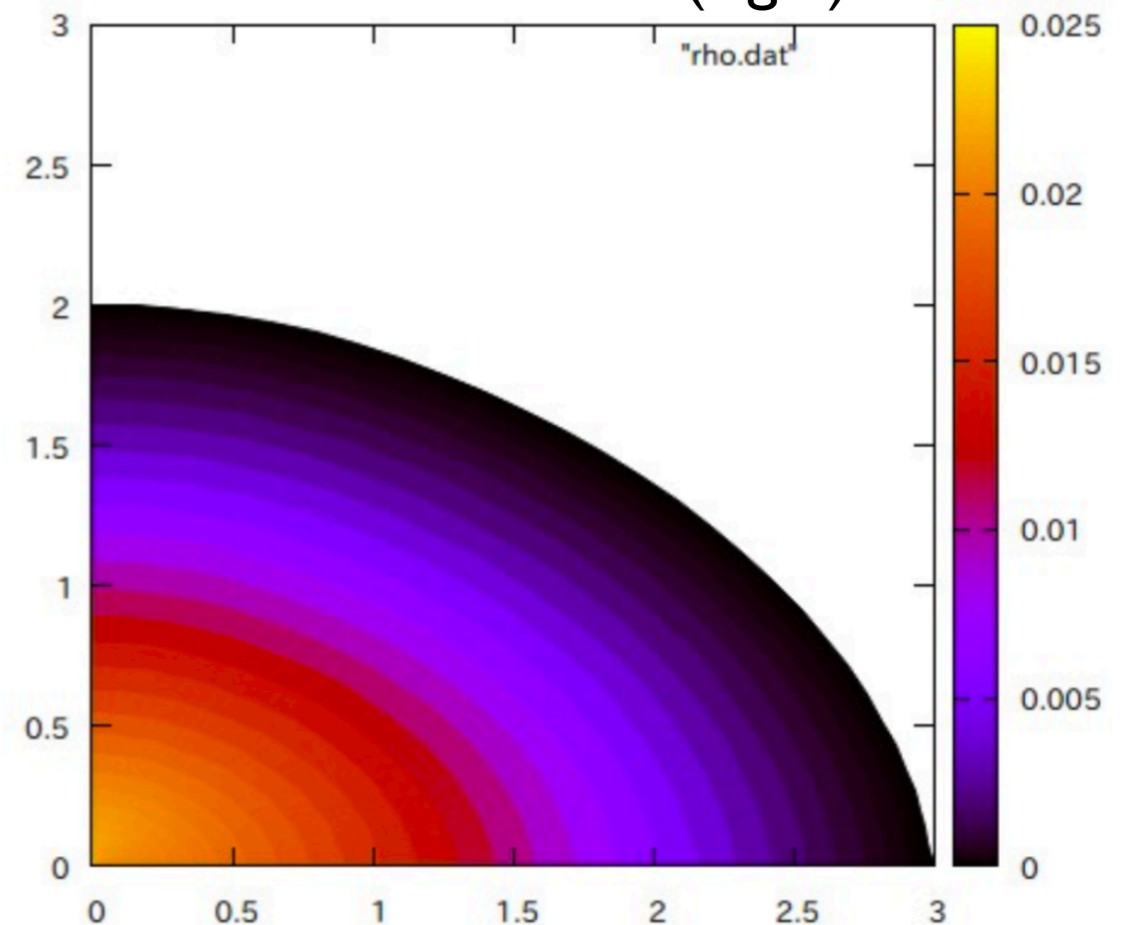
$$X = r, \cos \theta, P, \phi \dots$$

$$\int \int T_n(r) T_\ell(\nu) \left\{ \nabla P + \rho \nabla \phi - \frac{\rho j^2}{\omega^3} \mathbf{e}_\omega \right\} dr d\nu = 0$$

without rotation



with rotation(rigid)



We are ready to conduct evolution calculations !

③ GR case

Okawa, Ogata, Fujisawa, **NY**, Yamada

$\phi(r, \theta) \rightarrow f_0(r, \theta), f_1(r, \theta), f_2(r, \theta), \omega(r, \theta)$
 gravitational potential Metric

Under the axial symmetric metric

$$ds^2 = -f_0^2(r, \theta)dt^2 + f_1^{-2}(r, \theta) (dr^2 + r^2d\theta^2) + f_2^{-2}(r, \theta)r^2 \sin^2 \theta (d\varphi - \omega(r, \theta)dt)^2$$

Highly non-linear equations (Not integral form)

$E_{\mu\nu} (\equiv G_{\mu\nu} - \kappa_G T_{\mu\nu})$ Einstein eqs. ($E_{tt}, E_{\theta\theta}, E_{\varphi\varphi}, E_{t\varphi}$),
 Conservation ($\nabla_\mu T_r^\mu, \nabla_\mu T_\theta^\mu$) and EOS.

Example

$$\begin{aligned} \mathbf{G}_{tt} = & a_1^{(tt)} \left(\frac{\partial^2 f_0}{\partial r^2} + \frac{1}{r^2} \frac{\partial^2 f_0}{\partial \theta^2} \right) + a_2^{(tt)} \left(\frac{\partial^2 f_1}{\partial r^2} + \frac{1}{r^2} \frac{\partial^2 f_1}{\partial \theta^2} \right) + a_3^{(tt)} \left(\frac{\partial^2 f_2}{\partial r^2} + \frac{1}{r^2} \frac{\partial^2 f_2}{\partial \theta^2} \right) \\ & + a_4^{(tt)} \left(\frac{\partial^2 \omega}{\partial r^2} + \frac{1}{r^2} \frac{\partial^2 \omega}{\partial \theta^2} \right) + a_5^{(tt)} \left(\frac{\partial f_0}{\partial r} \frac{\partial \omega}{\partial r} + \frac{1}{r^2} \frac{\partial f_0}{\partial \theta} \frac{\partial \omega}{\partial \theta} \right) \\ & + a_6^{(tt)} \left(\frac{\partial f_2}{\partial r} \frac{\partial \omega}{\partial r} + \frac{1}{r^2} \frac{\partial f_2}{\partial \theta} \frac{\partial \omega}{\partial \theta} \right) + a_7^{(tt)} \frac{\partial f_0}{\partial r} + a_8^{(tt)} \frac{\partial f_1}{\partial r} + a_9^{(tt)} \frac{\partial f_2}{\partial r} + a_{10}^{(tt)} \frac{\partial \omega}{\partial r} \\ & + a_{11}^{(tt)} \frac{\partial f_1}{\partial \theta} + a_{12}^{(tt)} \frac{\partial f_2}{\partial \theta} + a_{13}^{(tt)} \frac{\partial \omega}{\partial \theta} \end{aligned}$$

$$\mathbf{T}_{tt} = \frac{\rho f_0^4 f_2^4 + \rho \omega^2 r^4 \sin^4 \theta (\Omega - \omega)^2 + f_0^2 f_2^2 r^2 \sin^2 \theta (\Omega^2 P + 2\Omega\omega\rho - 2\omega^2\rho)}{f_0^2 f_2^4 + f_2^2 r^2 \sin^2 \theta (\Omega - \omega)^2}$$

W4 methods

A method to solve non-linear equations (arXiv 1809.04495)

H.Okawa, K.Fujisawa, R.Hirai, Y.Yamamoto, NY, H.Nagakura, Yamada

Naive Newton-Raphon methods does not work for multi-dimensional hydrostatic equilibria in Lagrange coordinate [Ref. Friedman & Schutz 1978]. We, then, introduce a new method, named as "W4 method", which is based on the fixed point theorem.

$$\mathbf{F} = \nabla P + \rho \nabla \phi - \frac{\rho j^2}{\omega^3} \mathbf{e}_\varpi$$

$$\ddot{x}_1 + A_1 \dot{x}_1 + B_1 F = 0$$

$$\ddot{x}_2 + A_2 \dot{x}_2 + B_2 F = 0$$

⋮

【tangent vector space on S^1 】

Here, $F = F(x_1, x_2, \dots, y_1, y_2, \dots)$.

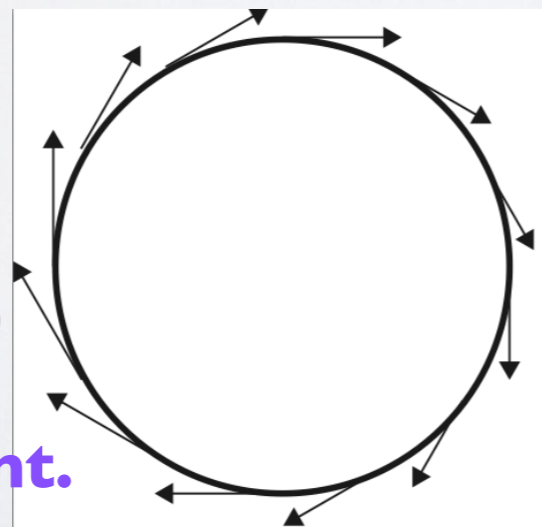
In Lagrange scheme, the variables are "coordinates" of nodes, which have mass, entropy, and angular momenta.

【tangent vector space on S^2 】

Find a fixed point \Leftrightarrow

Find a solution for
1st order differential eq.
(SOR, Newton method etc.)

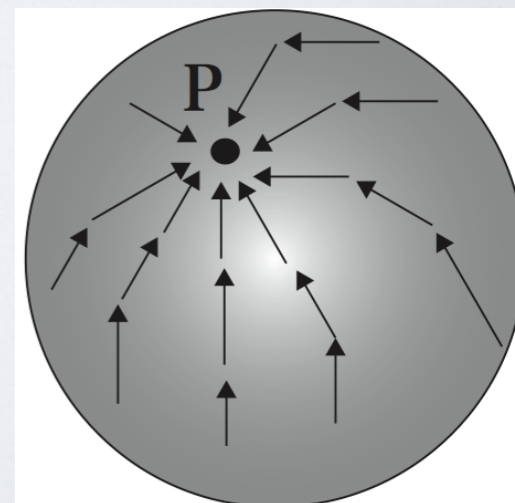
→ **Sometime,
there is no fixed point.**



Find a fixed point \Leftrightarrow

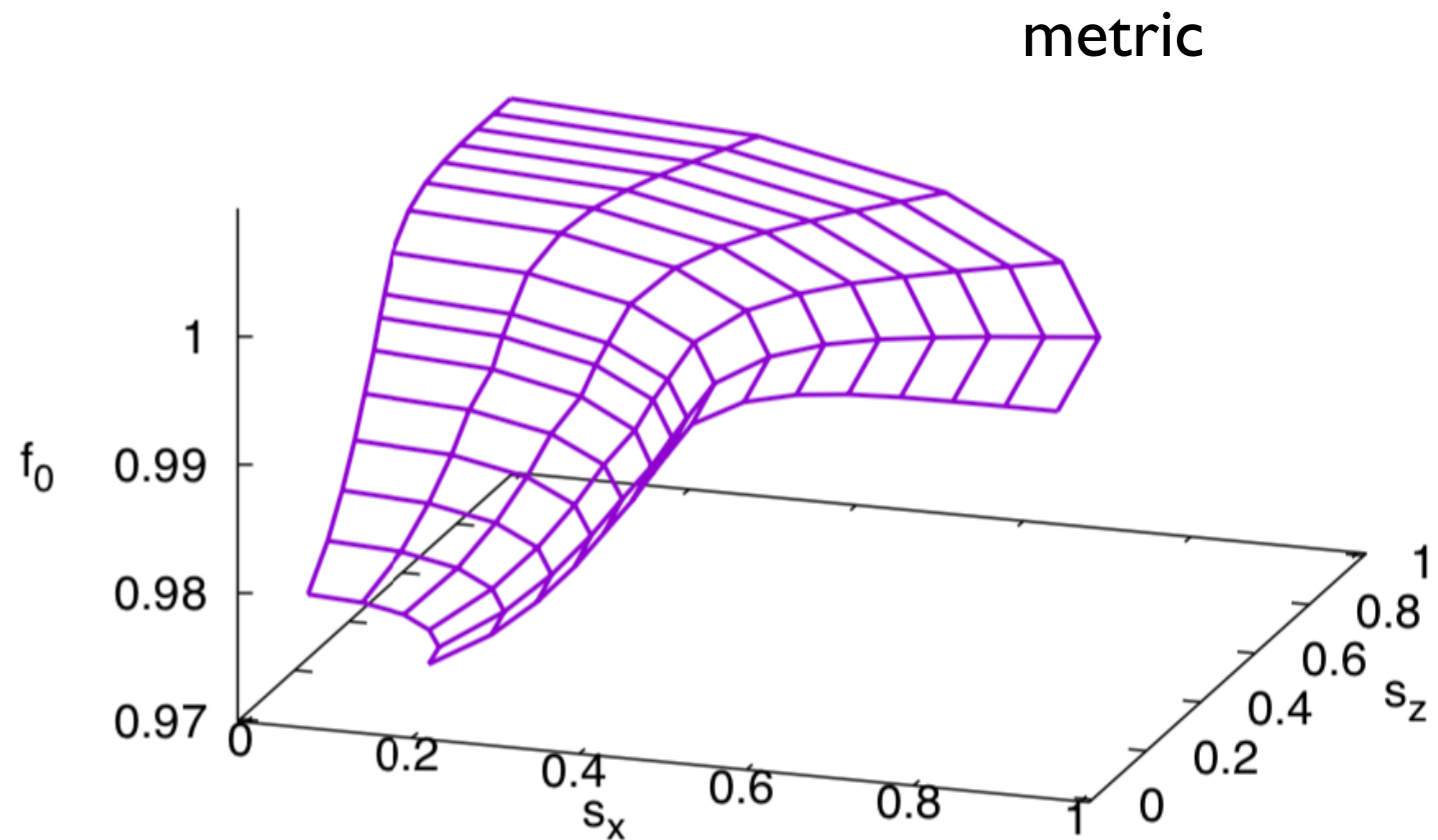
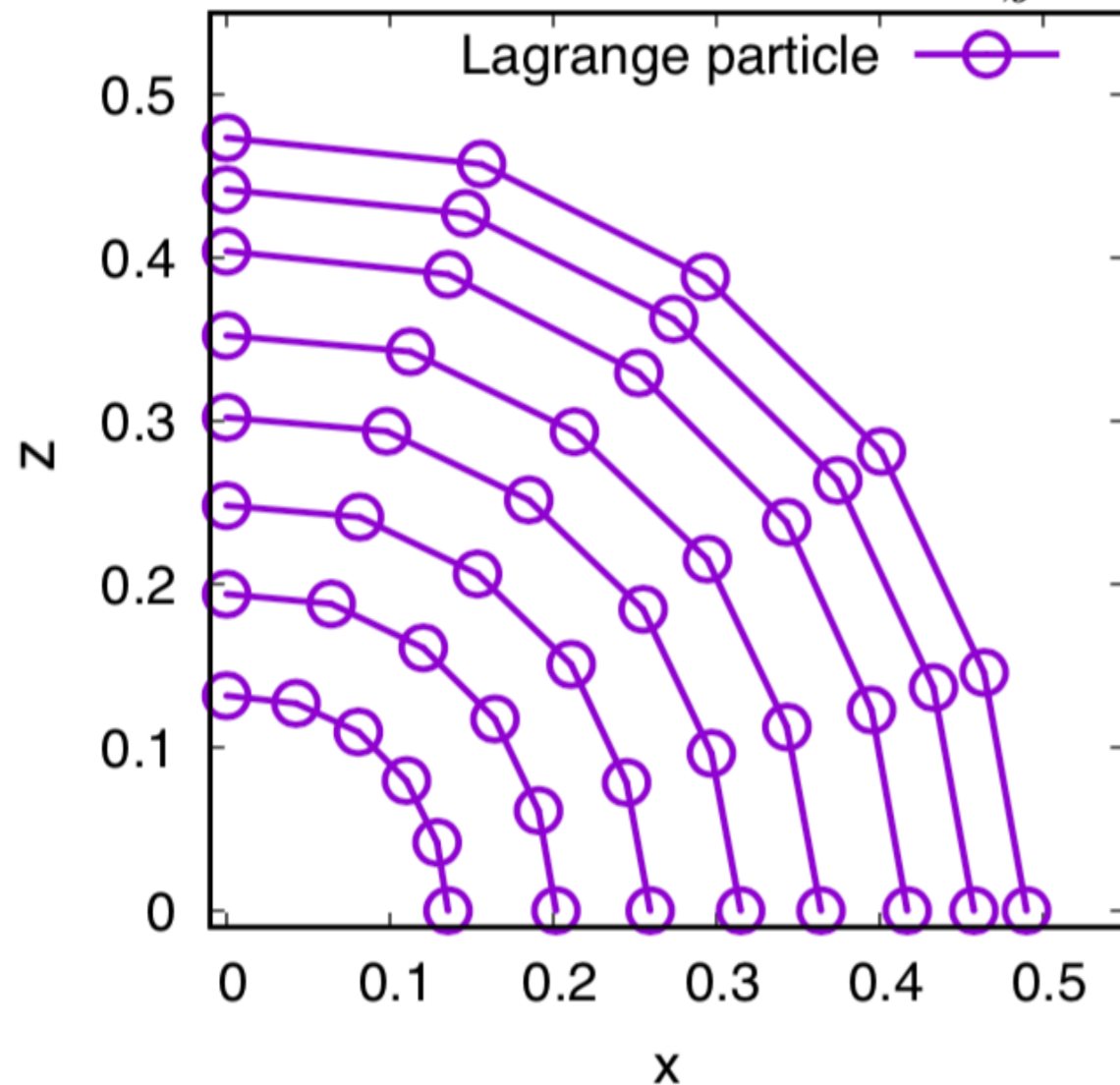
Find a solution for
2nd order differential eq.
(W4, MD calculation etc.)

→ **There must be
fixed points.**



Current status of GR case

with slow rotation(rigid)

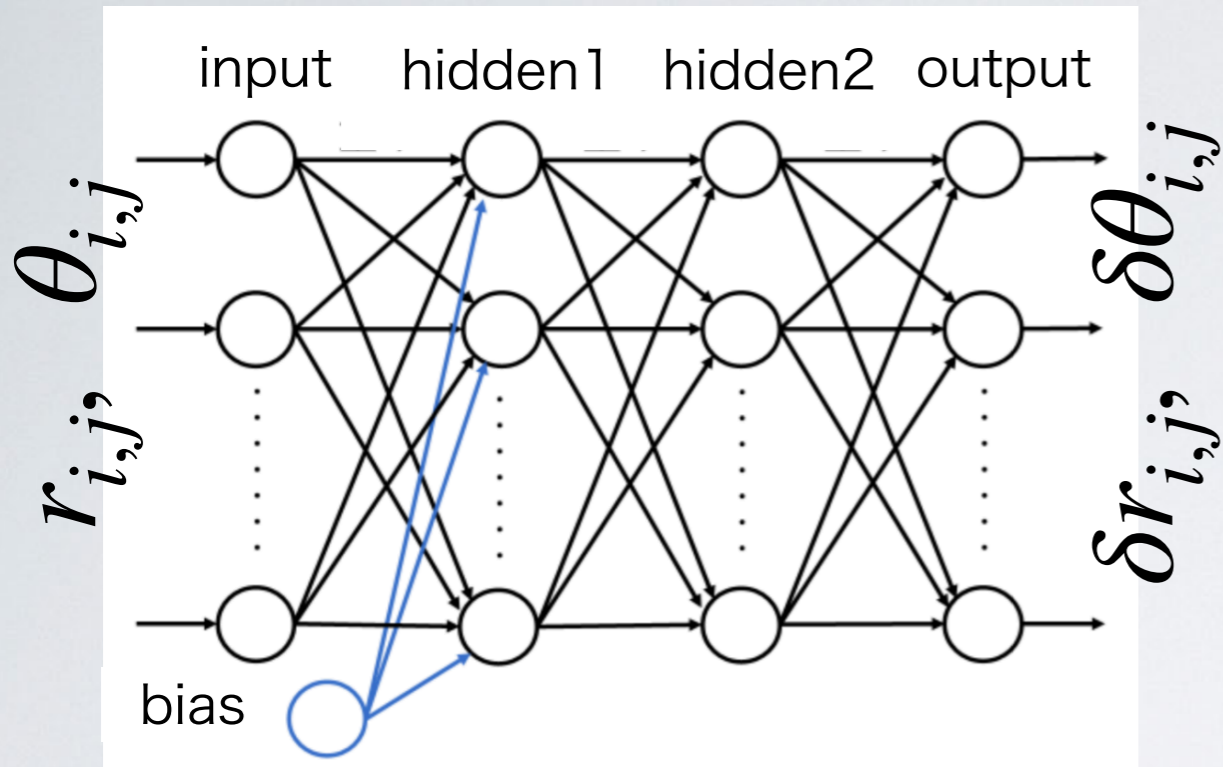


We already have some solutions with rotation.
→ Adaptable to proto-neutron stars: $\tau \sim 20\text{s}$.

$\times 10^2$

cf.) dynamical simulations for supernovae with super computers: $\tau \sim 0.2\text{s}$.

④ NEURAL NETWORK



Equilibrium eq.



Optimization problem

$$E = \sum_{i,j} \left(\nabla P + \rho \nabla \phi - \frac{\rho j^2}{\omega^3} \mathbf{e}_\omega \right)^2$$

Why Neural network?

① **Universal approximation theorem**

→ Useful for “non-linear equations”.

cf. Sign problem

(K. Kashiwa’s talk)

② **Development of Hardware and Software**

cf. cuDNN(Numerical Libraries)

tensorflow, Keras, cafe... etc.

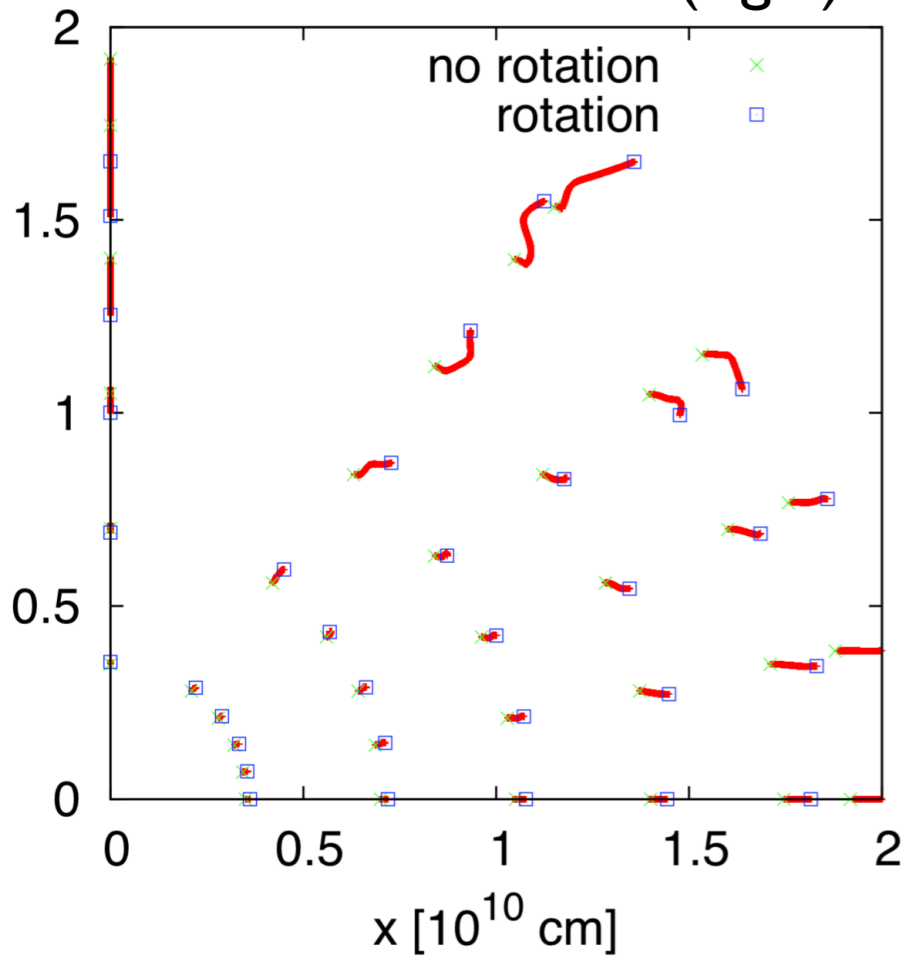
③ **Future high computings method**

such as quantum computers.

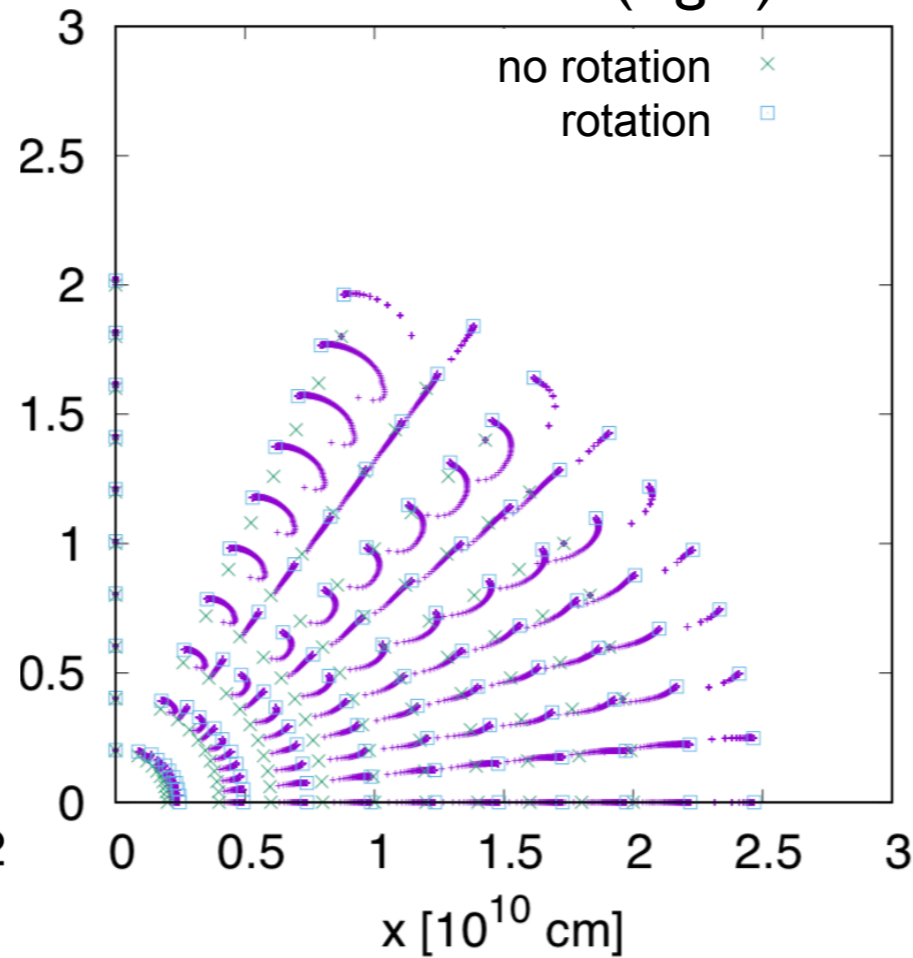
cf. D-wave,..etc.

EXAMPLES OF ROTATING STARS

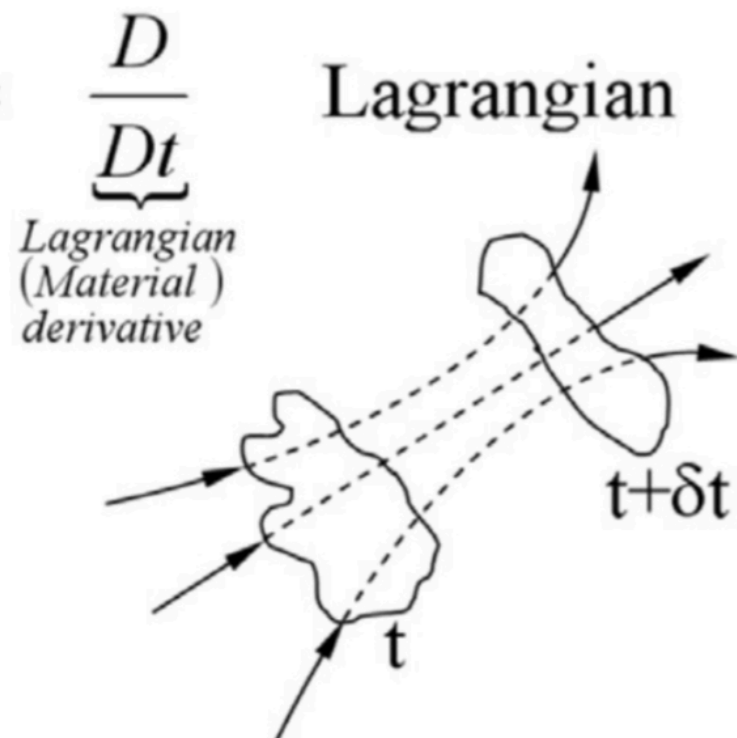
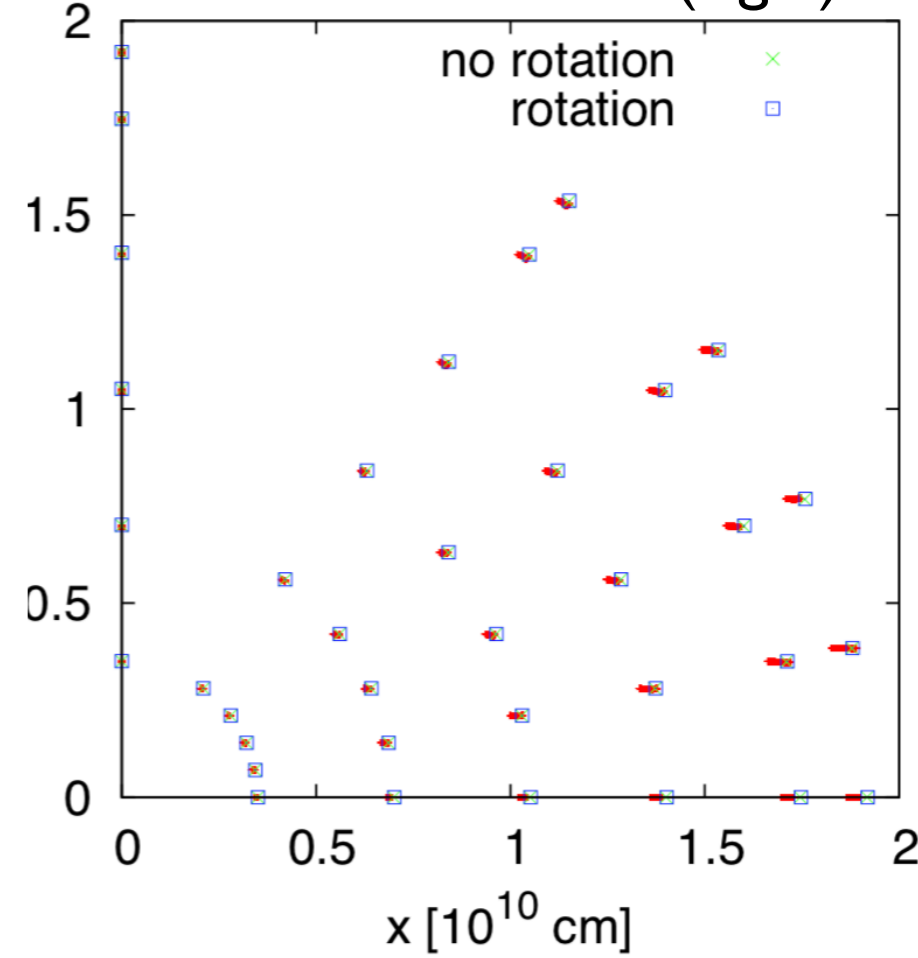
with slow rotation(rigid)



with rotation(rigid)



with slow rotation(rigid)



Blue or red paths show trajectories (vectors in left figure) of Lagrangian mass elements, which are searched by NN.



These trajectories strongly depends on NN, and initial guess. Accuracy is not so good. Optimization doesn't equal to solve equations. High numerical costs are also problems.

SUMMARY & DISCUSSION

It is the most important for stellar evolutions to obtain (baroclinic) hydrostatic equilibria with arbitrary angular-momentum profiles. But they are highly non-linear balance equation.

We have introduced 4 approaches in this talk.

1. Simple discretization

2. Spectral method

3. GR case

4. Neural Network

Optimization \neq Solving non-linear equations

NN

W4 method (arXiv: 1809.04495)

Now, our topic moves to **“stellar evolutions in 2D,3D”**.