



# *RPA approaches and effective interactions*

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*A rendering of the future RAON complex, under construction in Daejeon*



- ❖ Introduction: the general issue
- ❖ Introduction: the specifics
  - Transformed NN interactions via UCOM, SRG  *see also: Hans Feldmeier*
  - “Self-consistent” Second RPA  *see also: Marcella Grasso*
- ❖ Applications of transformed NN interactions in (S)RPA
  - Giant resonances: centroids
  - Giant resonances: fine structure
- ❖ Some comments on...
  - Three-nucleon interactions
  - Ground-state correlations
- ❖ Outlook

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# The (non-relativistic) nuclear many-body problem

## ❖ A interacting nucleons

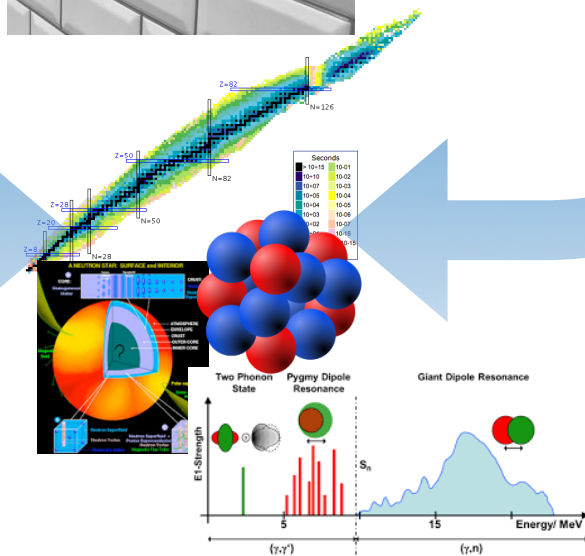
NN(N) interactions



Schroedinger equation

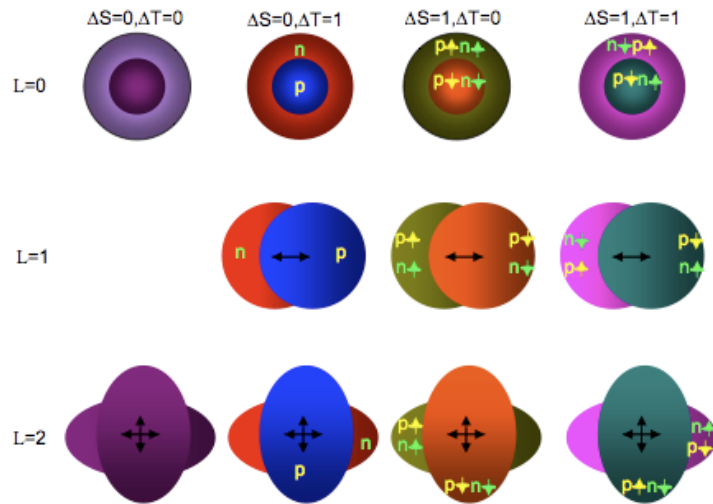
Effective/transformed  
Hamiltonian

Quantum many-body  
method



# Normal modes of vibration

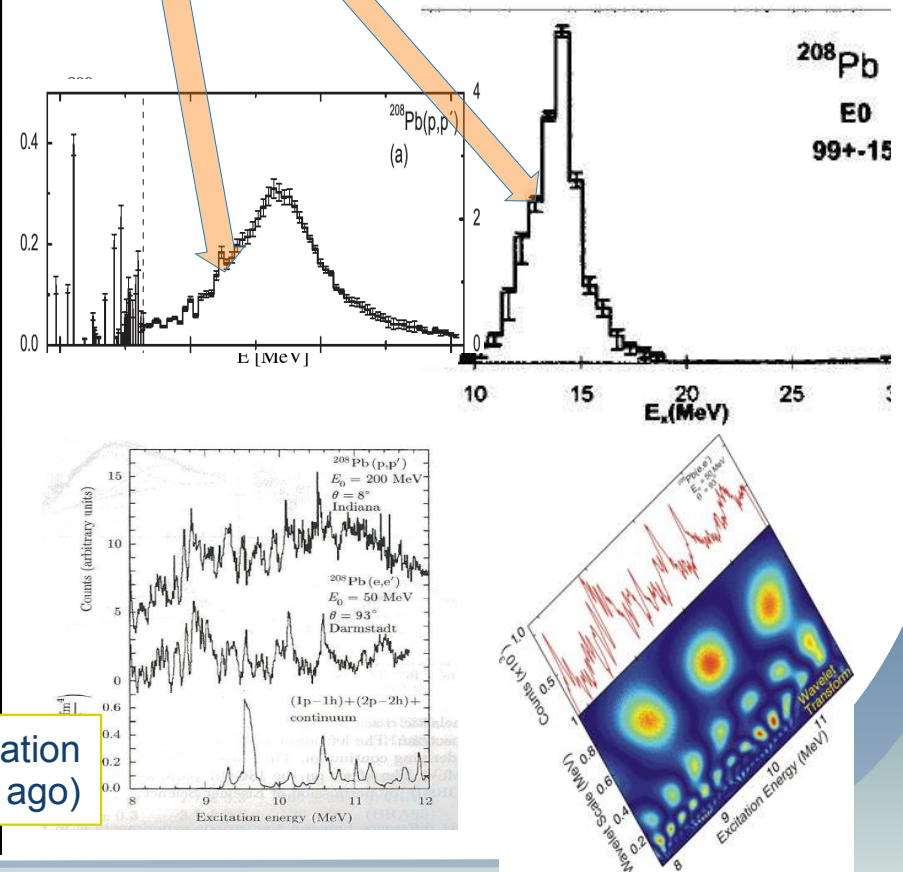
The simplified picture



SRPA explored to explain attenuation (~30yrs ago)

Major part of the EWSR

The reality



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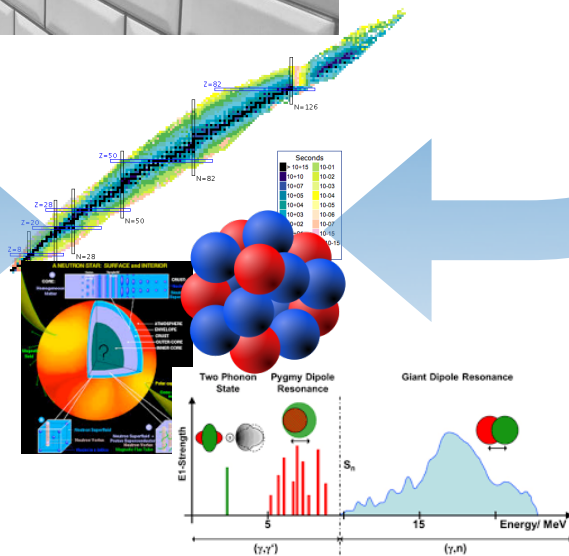
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# The (non-relativistic) nuclear many-body problem

## ❖ $A$ interacting nucleons

### NN(N) interactions

*chiral or directly fitted to  
few-nucleon data*

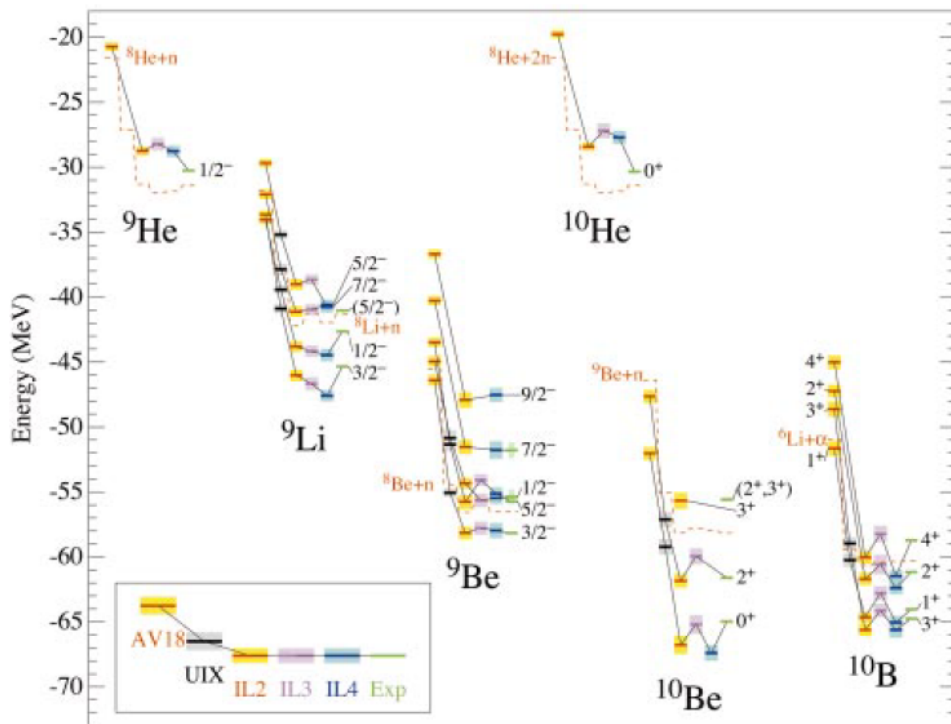


### Schroedinger equation

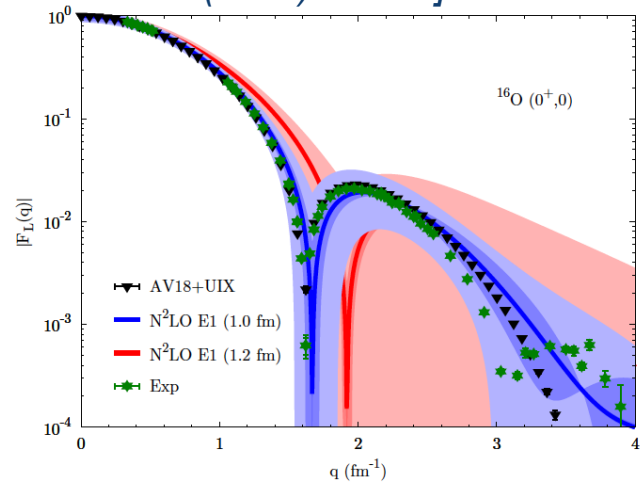
Quantum many-body  
method



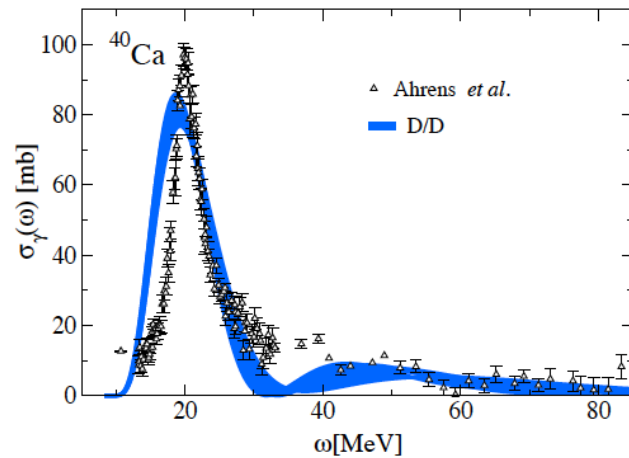
QMC [Pieper et al., PRC66(2002)044310]



QMC [Lonardoni et al., PRC97(2018)044318]



CCM+LIT [Simonis et al, arXiv: 1905.02055]





# The (non-relativistic) nuclear many-body problem

## ❖ $A$ interacting nucleons

NN(N) interactions



Schroedinger equation

Hartree-Fock, RPA  
("Mean field")

standard DFT



*... where a couple of MeV discrepancy in the total mass of heavy nuclei is seen as a failure!*

# The (non-relativistic) nuclear many-body problem

## ❖ A interacting nucleons

NN(N) interactions

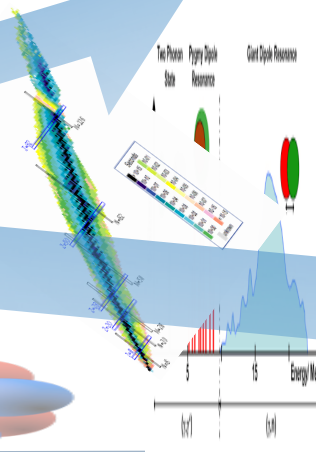
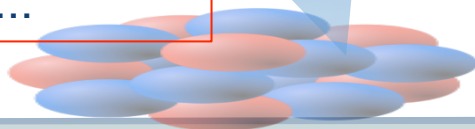


Schroedinger equation

Effective/transformed  
Hamiltonian

Quantum many-body  
method

double counting,  
undercounting,  
divergencies,  
...



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See review by Roth, Neff, Feldmeier, PPNP65(2010)50

## ❖ Unitary correlation operator method (UCOM)

$$\langle \Psi | H | \Psi' \rangle = \langle \Phi | C^\dagger H C | \Phi' \rangle = \langle \Phi | C^{-1} H C | \Phi' \rangle = \langle \Phi | \hat{H} | \Phi' \rangle$$

## ❖ Similarity renormalization group (SRG)

$$\frac{d}{d\alpha} \tilde{H}_\alpha = [\eta_\alpha, \tilde{H}_\alpha] \qquad \frac{d}{d\alpha} U_\alpha = -U_\alpha \eta_\alpha$$

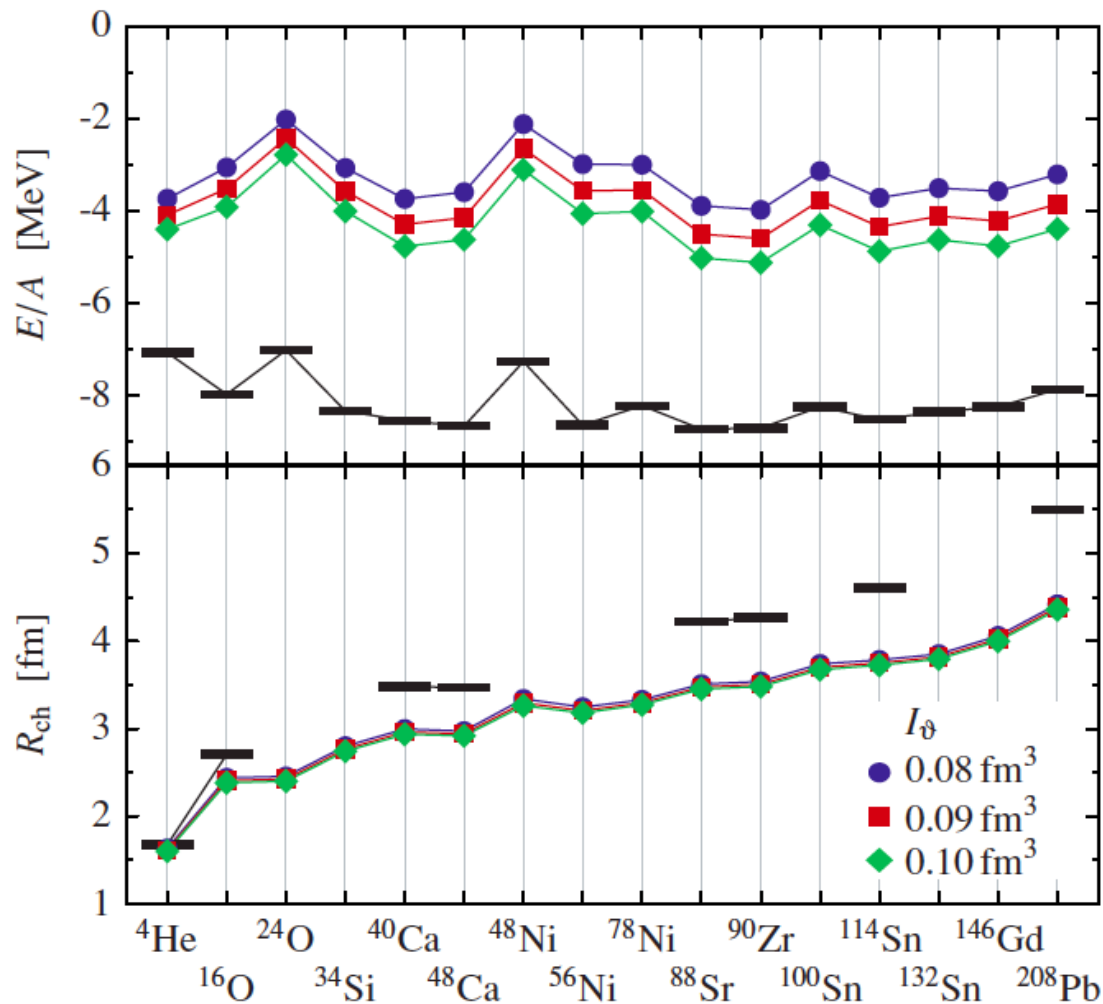
## ❖ Starting Hamiltonian:

- Precisely fit phenomenological (e.g., Argonne V18)
- From chiral EFT

## ❖ Prediagonalized Hamiltonian: to be used within tractable Hilbert spaces – How tractable?

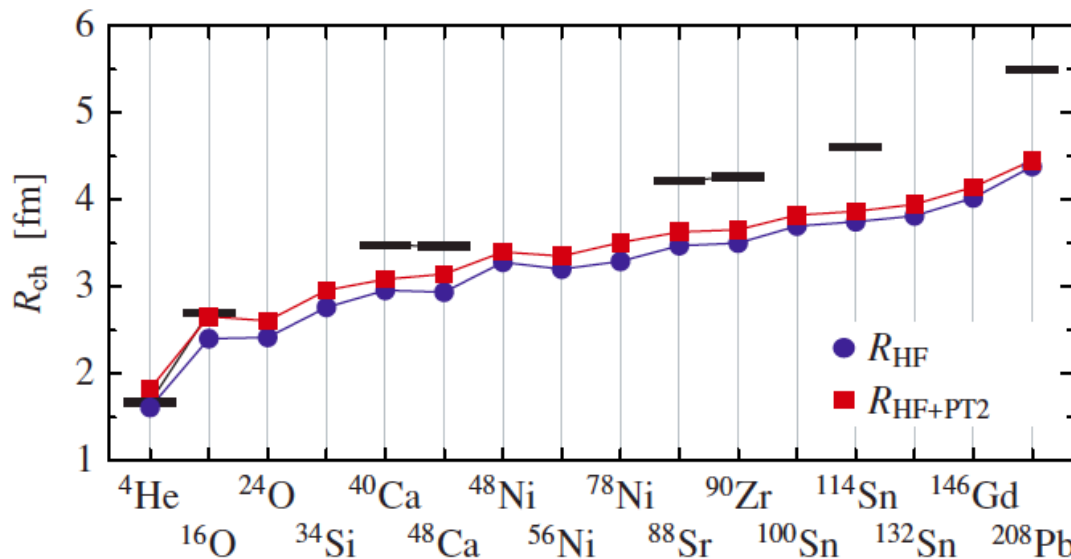
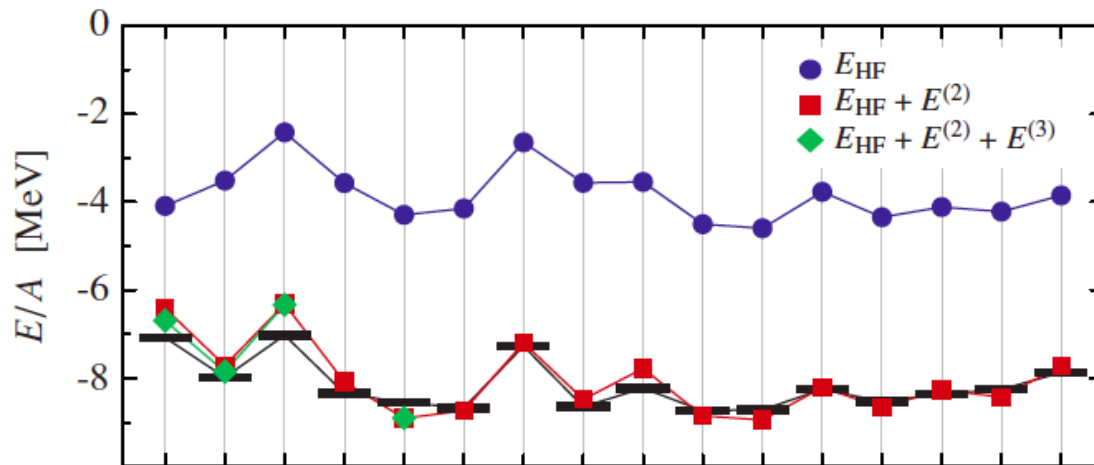
## ❖ AV18+UCOM

[Roth, PP, et al., PRC73(06)044312]

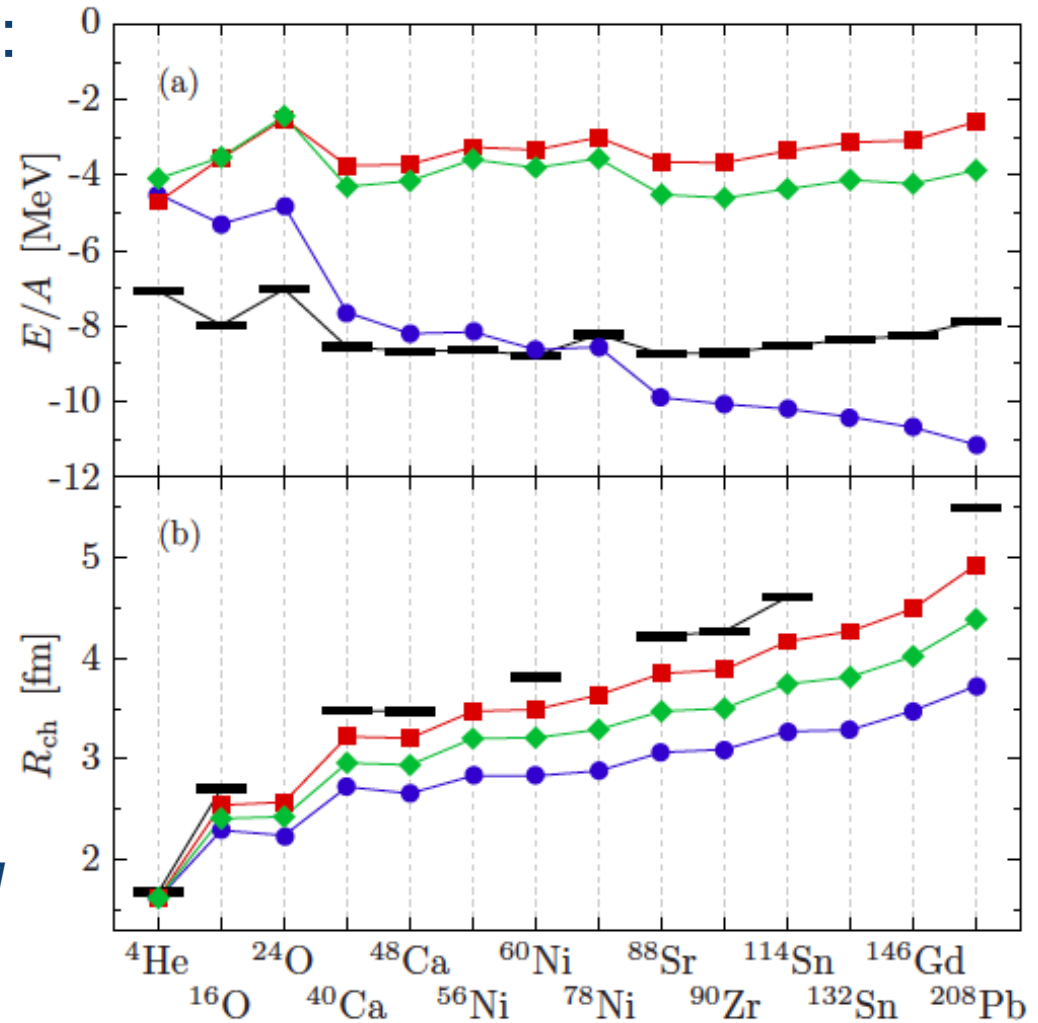


## ❖ AV18+UCOM

[Roth,PP, et al., PRC73(06)044312]



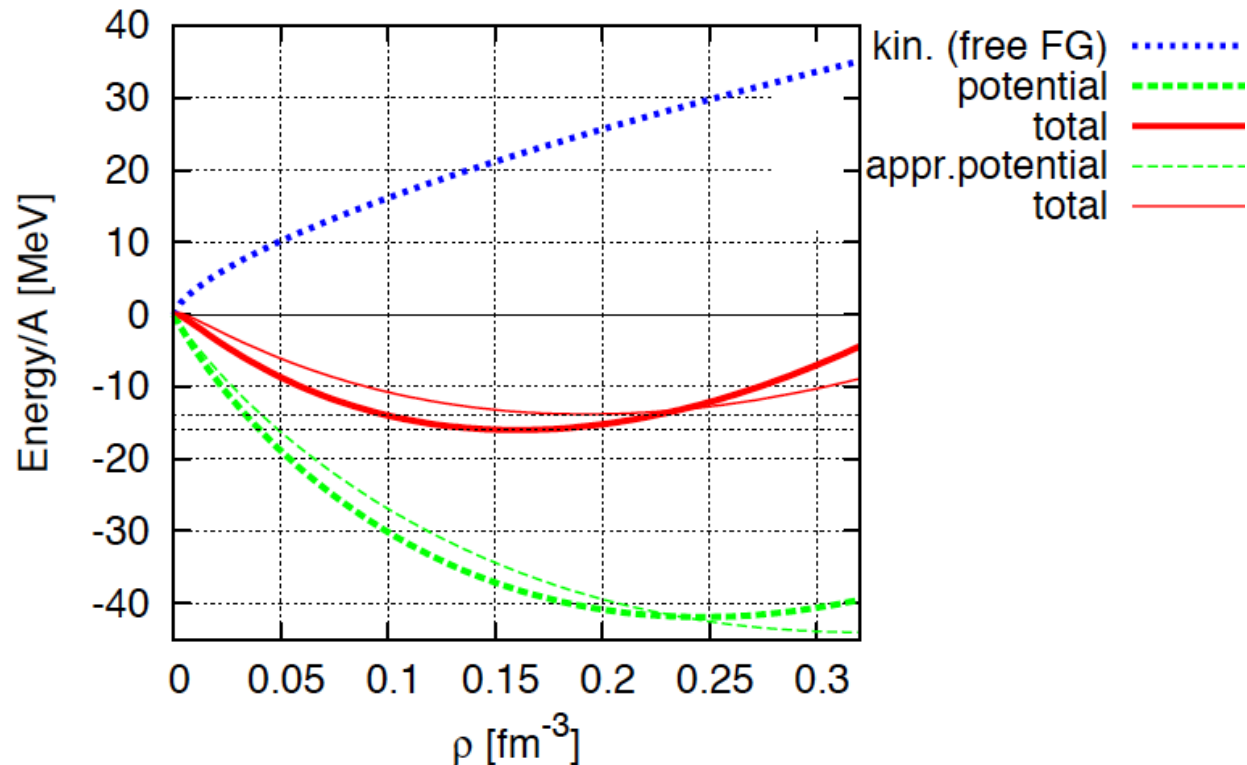
- ❖ AV18 transformed via:
  - original UCOM
  - UCOM-SRG
  - SRG



[Roth et al., PRC77(2008)064003]

- ❖ The total energy is the difference of two comparably large quantities:

SNM



- ❖ Small residual 3N effects can play a critical role



# Standard RPA

- Vibration creation operator:

$$Q_\nu^\dagger = \sum_{ph} X_{ph}^\nu O_{ph}^\dagger - \sum_{ph} Y_{ph}^\nu O_{ph} \quad ; \quad Q_\nu |RPA\rangle = 0 \quad ; \quad Q_\nu^\dagger |RPA\rangle = |\nu\rangle$$

- Standard RPA - the RPA vacuum is approximated by the HF ground state:

$$\langle RPA | \dots | RPA \rangle \rightarrow \langle HF | \dots | HF \rangle \quad ; \quad O_{ph}^\dagger \rightarrow a_p^\dagger a_h$$

- RPA equations in  $ph$ -space:

$$\begin{pmatrix} A & B \\ -B^* & -A^* \end{pmatrix} \begin{pmatrix} X^\nu \\ Y^\nu \end{pmatrix} = \hbar\omega_\nu \begin{pmatrix} X^\nu \\ Y^\nu \end{pmatrix}$$

$$A_{ph,p'h'} = \delta_{pp'}\delta_{hh'}(e_p - e_h) + H_{hp',ph'} \quad ; \quad B_{ph,p'h'} = H_{hh',pp'} \quad ;$$

For some in-medium  
Hamiltonian H

☞ Self-consistent HF+RPA: spurious state and sum rules

- Equivalent to small-amplitude Time-Dependent Hartree-Fock
- Linear Response Theory

# Second RPA

- Vibration creation operator: Includes  $2p2h$  configurations

$$Q_{\nu}^{\dagger} = \sum_{ph} X_{ph}^{\nu} O_{ph}^{\dagger} - \sum_{ph} Y_{ph}^{\nu} O_{ph} + \sum_{p_1 h_1 p_2 h_2} \mathcal{X}_{p_1 h_1 p_2 h_2}^{\nu} O_{p_1 h_1 p_2 h_2}^{\dagger} - \sum_{p_1 h_1 p_2 h_2} \mathcal{Y}_{p_1 h_1 p_2 h_2}^{\nu} O_{p_1 h_1 p_2 h_2}$$

- The SRPA vacuum is approximated by the HF ground state:

$$\langle \text{SRPA} | \dots | \text{SRPA} \rangle \rightarrow \langle \text{HF} | \dots | \text{HF} \rangle$$

- SRPA equations in  $ph \oplus 2p2h$ -space:

$$\left( \begin{array}{cc|cc} A & \mathcal{A}_{12} & B & 0 \\ \mathcal{A}_{21} & \mathcal{A}_{22} & 0 & 0 \\ \hline -B^* & 0 & -A^* & -\mathcal{A}_{12}^* \\ 0 & 0 & -\mathcal{A}_{21}^* & -\mathcal{A}_{22}^* \end{array} \right) \begin{pmatrix} X^{\nu} \\ \mathcal{X}^{\nu} \\ Y^{\nu} \\ \mathcal{Y}^{\nu} \end{pmatrix} = \hbar\omega_{\nu} \begin{pmatrix} X^{\nu} \\ \mathcal{X}^{\nu} \\ Y^{\nu} \\ \mathcal{Y}^{\nu} \end{pmatrix}$$

$$A_{ph,p'h'} = \delta_{pp'} \delta_{hh'} (e_p - e_h) + H_{hp',ph'} ; \quad B_{ph,p'h'} = H_{hh',pp'}$$

$\mathcal{A}_{12}$ : interactions between  $ph$  and  $2p2h$  states

$\mathcal{A}_{22}$ :  $\delta_{p_1 p'_1} \delta_{h_1 h'_1} \delta_{p_2 p'_2} \delta_{h_2 h'_2} (e_{p_1} + e_{p_2} - e_{h_1} - e_{h_2})$  + interactions among  $2p2h$  states

For some in-medium Hamiltonian H

# Second RPA

- Vibration creation operator: Includes  $2p2h$  configurations

$$Q_\nu^\dagger = \sum_{ph} X_{ph}^\nu O_{ph}^\dagger - \sum_{ph} Y_{ph}^\nu O_{ph} + \sum_{p_1 h_1 p_2 h_2} \mathcal{X}_{p_1 h_1 p_2 h_2}^\nu O_{p_1 h_1 p_2 h_2}^\dagger - \sum_{p_1 h_1 p_2 h_2} \mathcal{Y}_{p_1 h_1 p_2 h_2}^\nu O_{p_1 h_1 p_2 h_2}$$

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$$A_{ph,p'h'} = \delta_{pp'} \delta_{hh'} (e_p - e_h) + H$$

$\mathcal{A}_{12}$ : interactions between  $ph$

$$\mathcal{A}_{22}: \delta_{p_1 p_1'} \delta_{h_1 h_1'} \delta_{p_2 p_2'} \delta_{h_2 h_2'} (e_{p_1} +$$

**linear response with collision term**

❖ Solutions appear in adjoint pairs

❖ If the stability condition is satisfied,

$$\langle 0|[F^\dagger, [H, F]]|0\rangle = \sum_{ab} F_a^* M_{ab} F_b \geq 0$$

solutions are real

❖ The EWSR is preserved  $\left( \sum_n (E_n - E_0) |\langle 0|V|n\rangle|^2 = \int d^3r \rho_0(\vec{r}) \frac{(\vec{\nabla}V)^2}{2m} \right)$

$$\sum_{n: N_n=1} |\langle 0|O|n\rangle|^2 E_{n0} = \frac{1}{2} \langle 0|[O, H, O]|0\rangle$$

❖ ... and therefore spurious transitions have zero energy (restoration of symmetries)

$|HF\rangle = \text{vacuum of } ph \text{ states} \rightarrow \text{stability of RPA for physical states}$

- ❖ Solutions still appear in adjoint pairs
- ❖ However, the stability condition is violated and spurious states appear at non-zero energy...
- ❖ Note that the rhs is still zero:

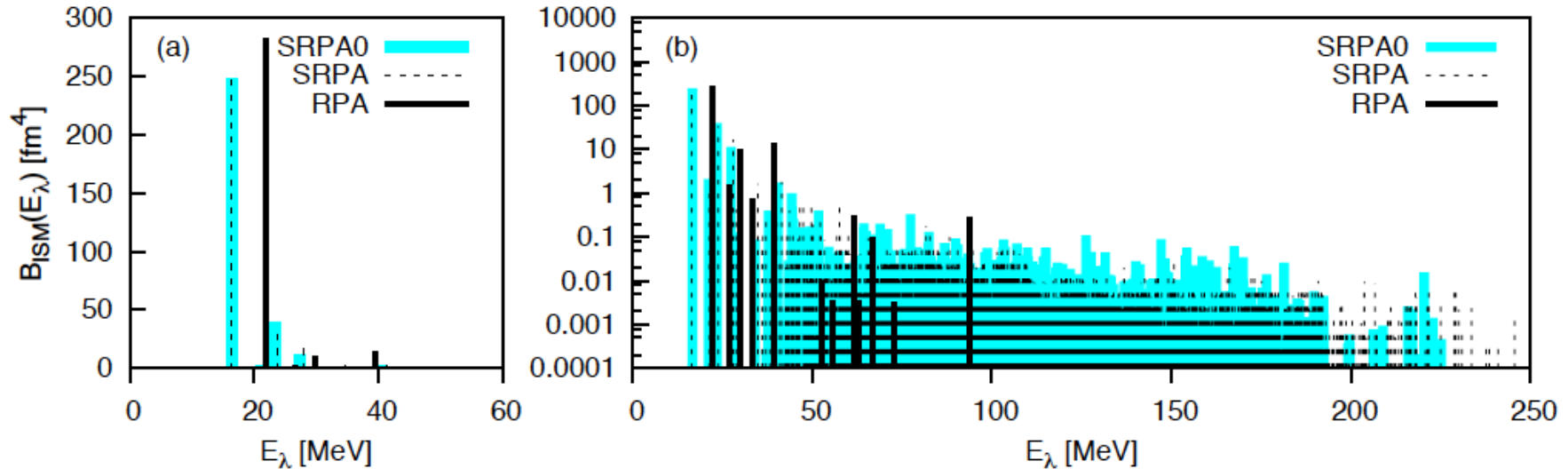
$$\sum_{n: N_n=1} |\langle 0|O|n\rangle|^2 E_{n0} = \frac{1}{2} \langle 0|[O, H, O]|0\rangle$$

- ❖ Does this mean that the EWSR is not conserved?
- ❖ **The EWSR is conserved:** the difference is that there are positive-energy states with negative norm / negative-energy states with positive norm
- ❖ Has been verified numerically

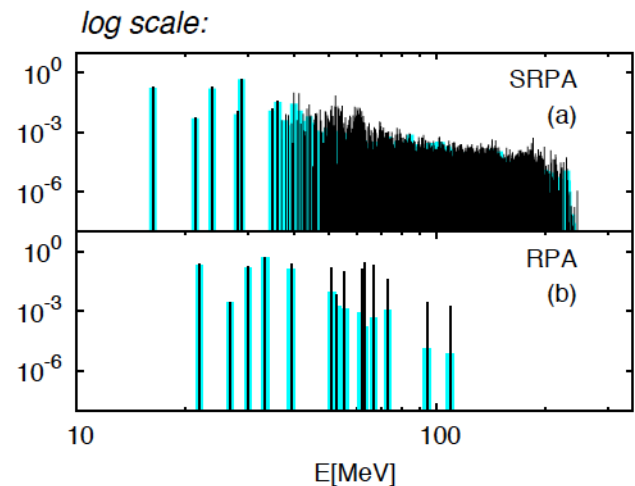
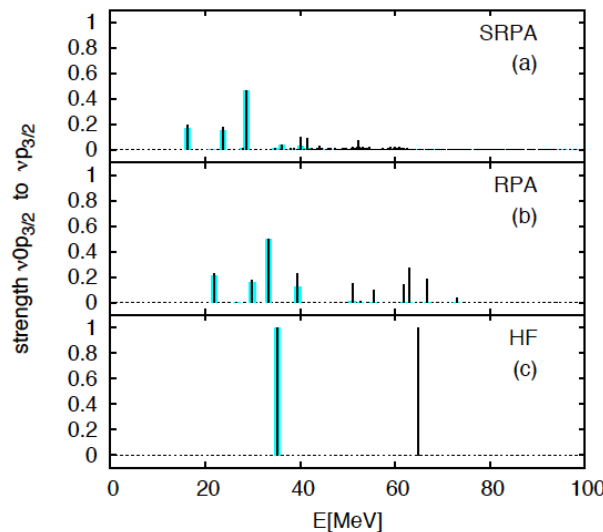
# RPA vs SRPA: a numerical demonstration

first "self-consistent" SRPA calculations:  
[PP & Roth, PRC81(2010)024317]

## ❖ Strength distribution



## ❖ ph states:

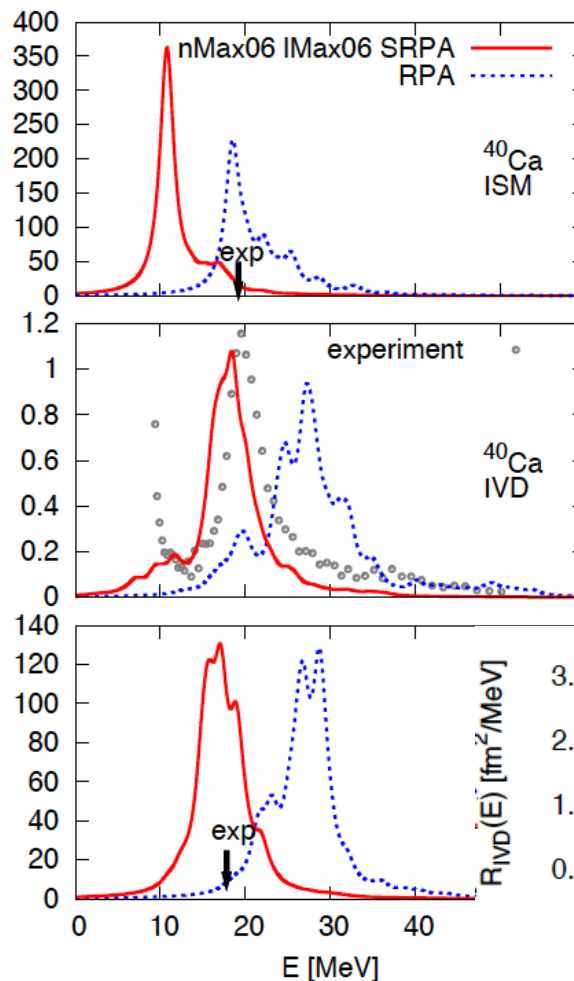
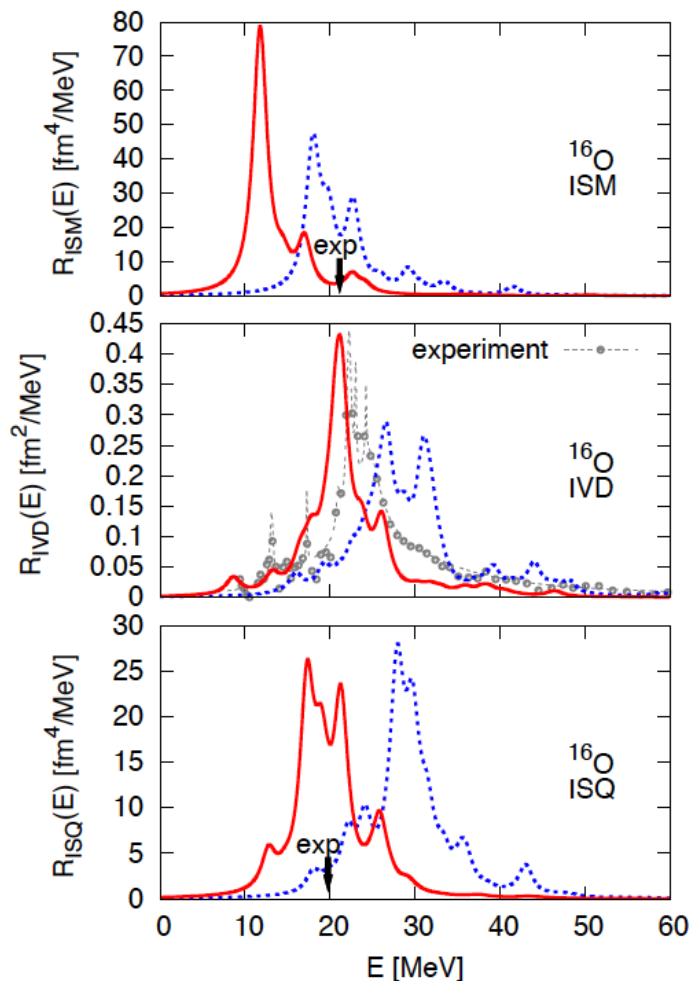


<sup>16</sup>O, 7 shells

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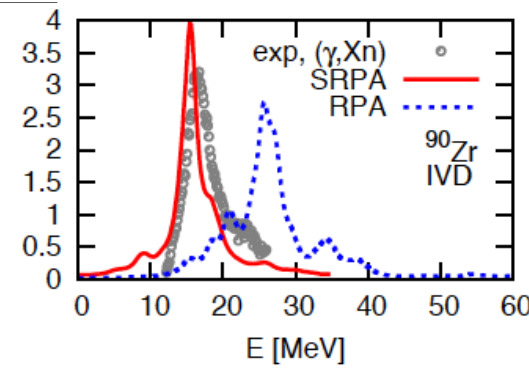
## ❖ RPA vs SRPA with UCOM

[PP & Roth, PLB671(2009)356]



Residual 3N effects?

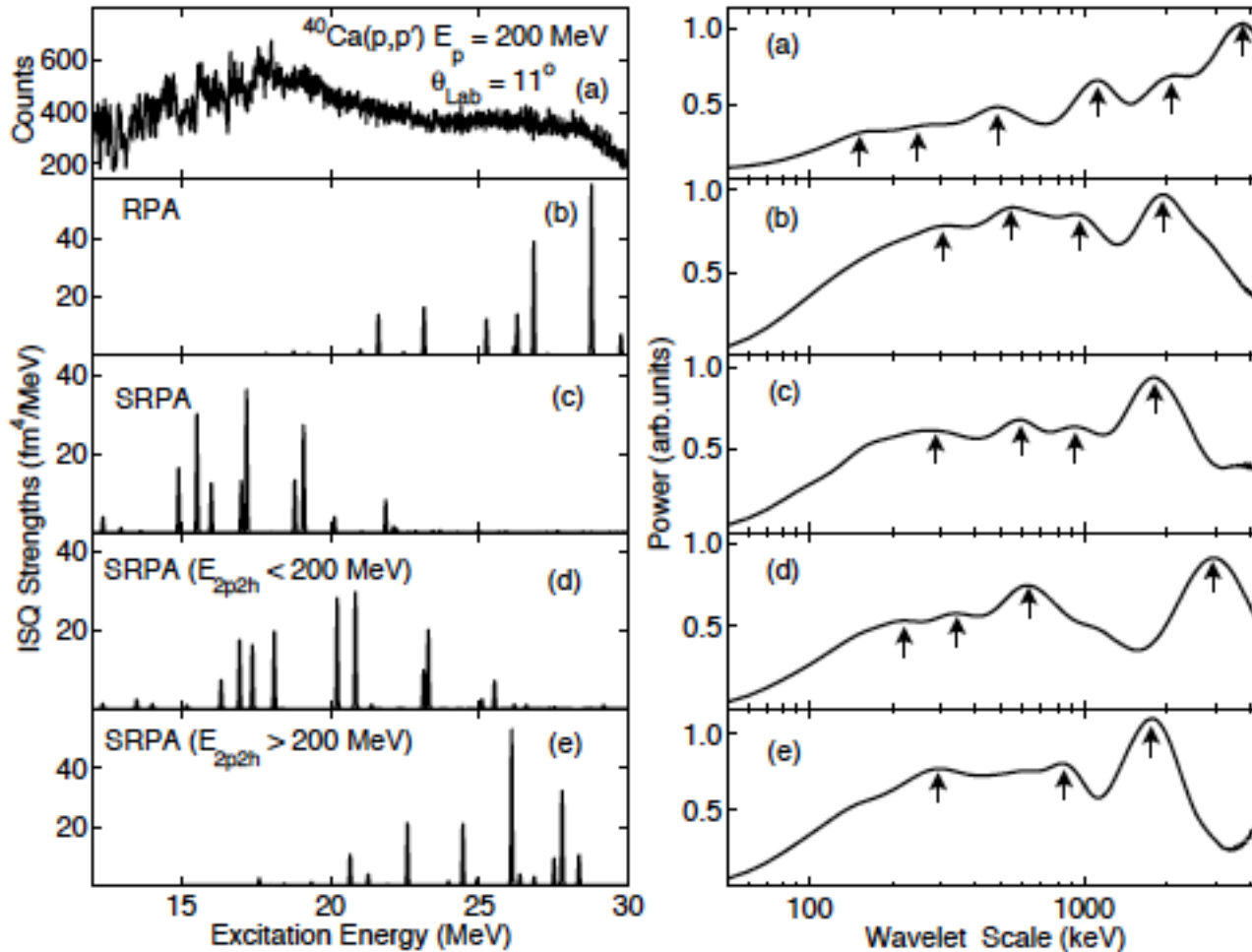
Dressing of orbitals



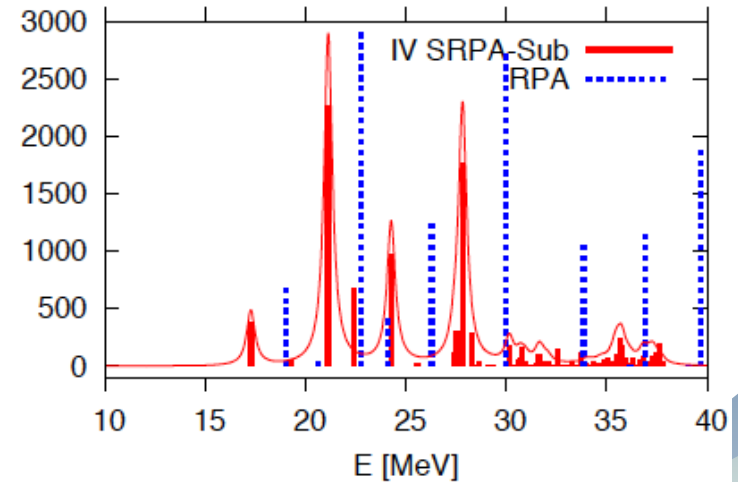
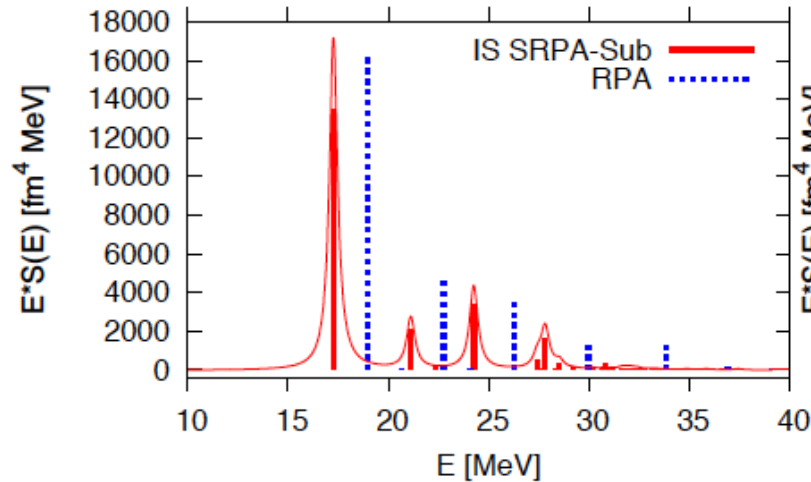
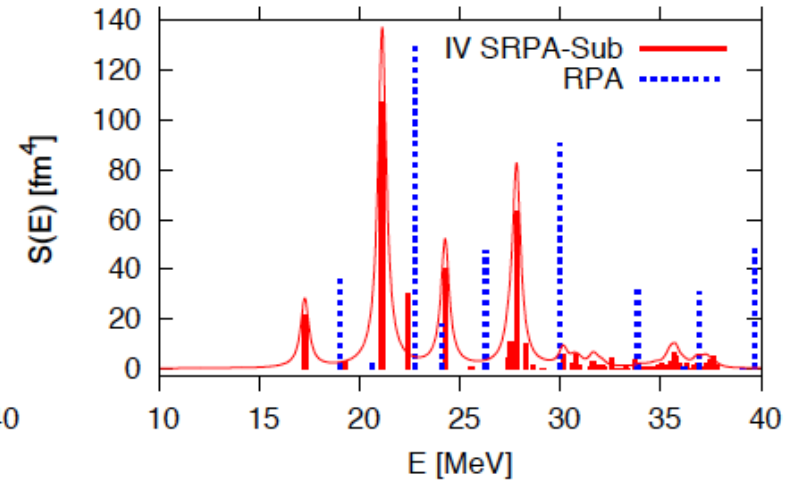
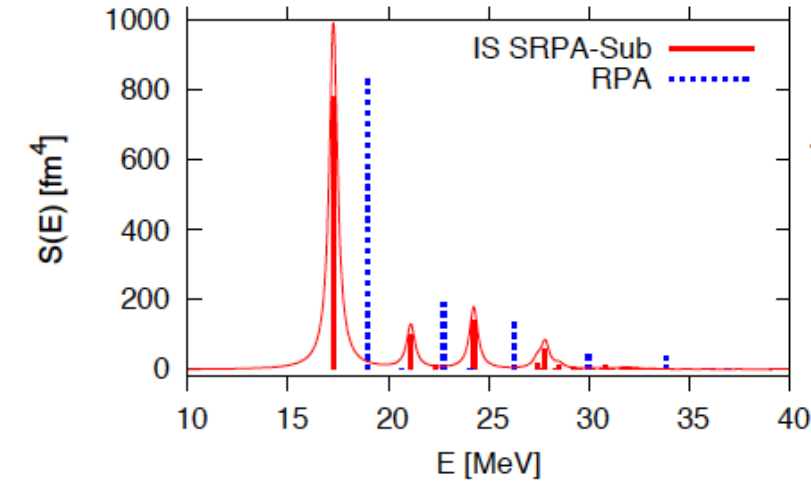


# $^{40}\text{Ca} (p,p')$ and UCOM-SRPA

Fine structure of the isoscalar giant quadrupole resonance in  $^{40}\text{Ca}$  due to Landau damping?  
[Usman,...PP,...et al,PLB698(2011)]



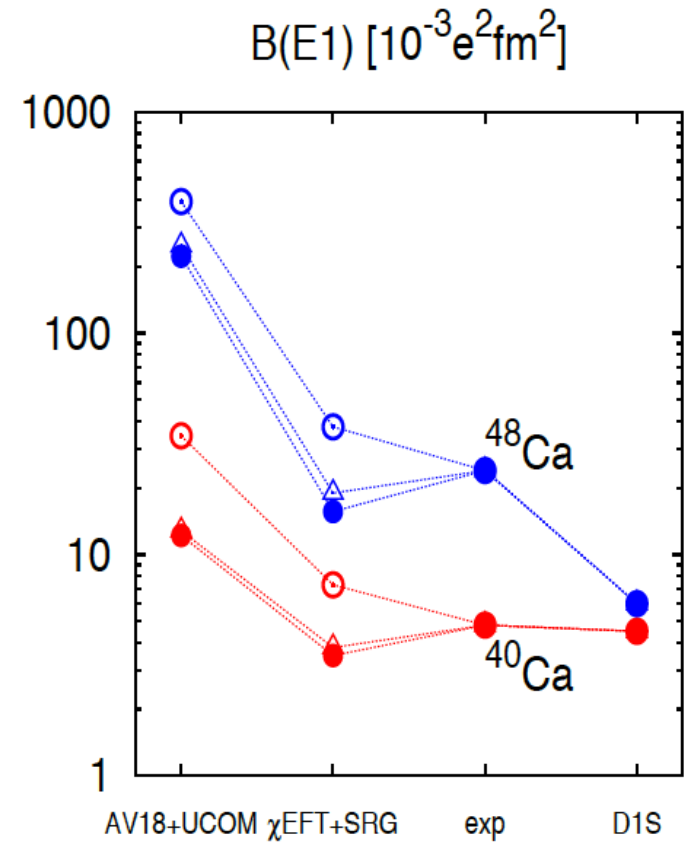
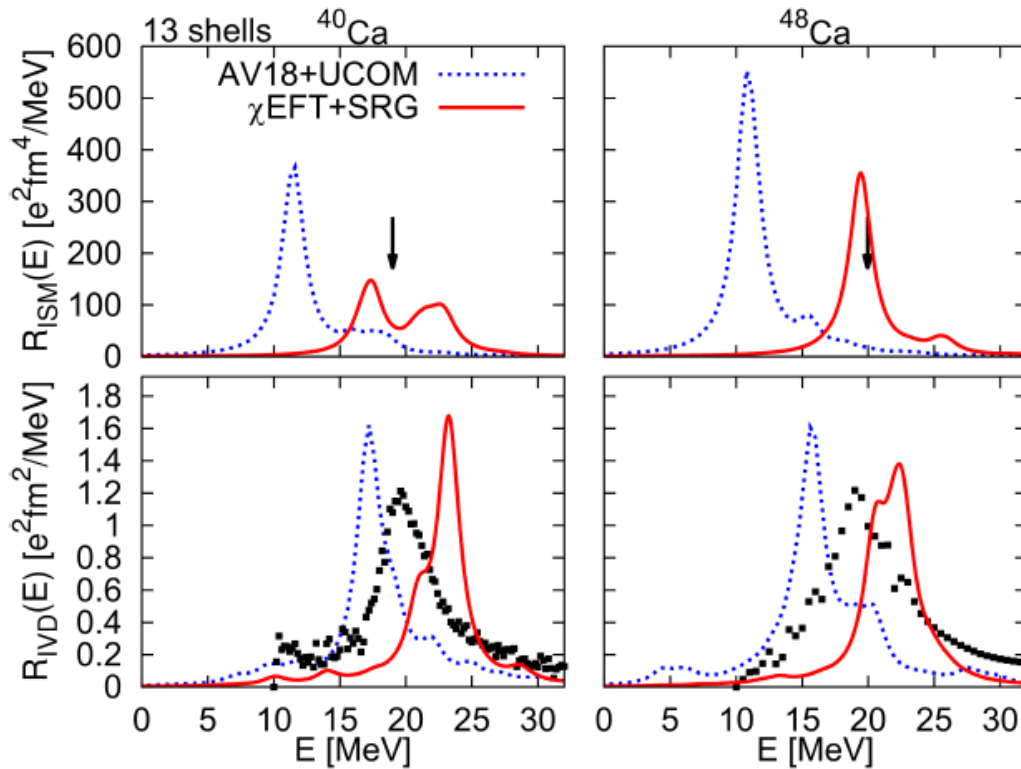
## ❖ Monopole response of $^{40}\text{Ca}$



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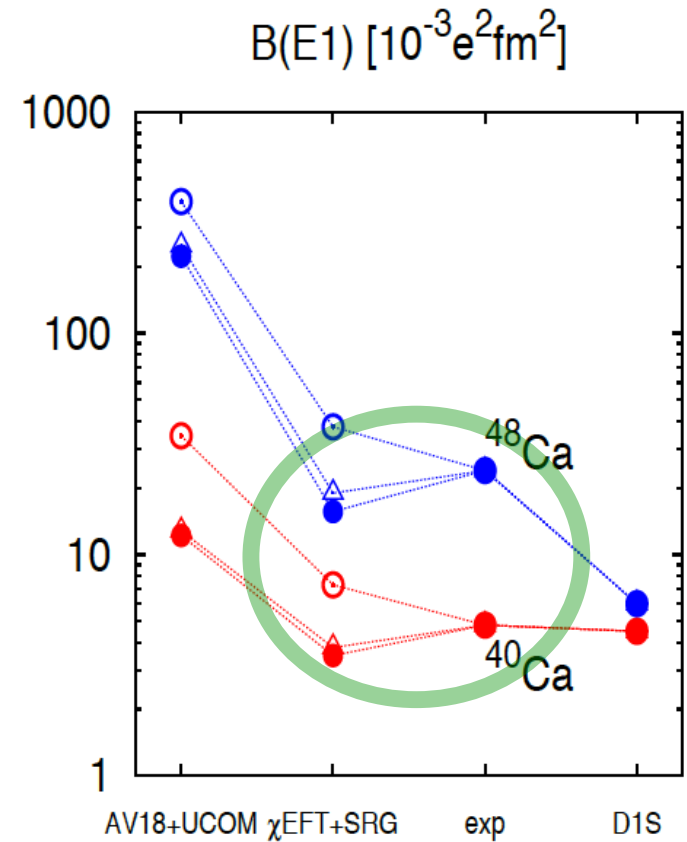
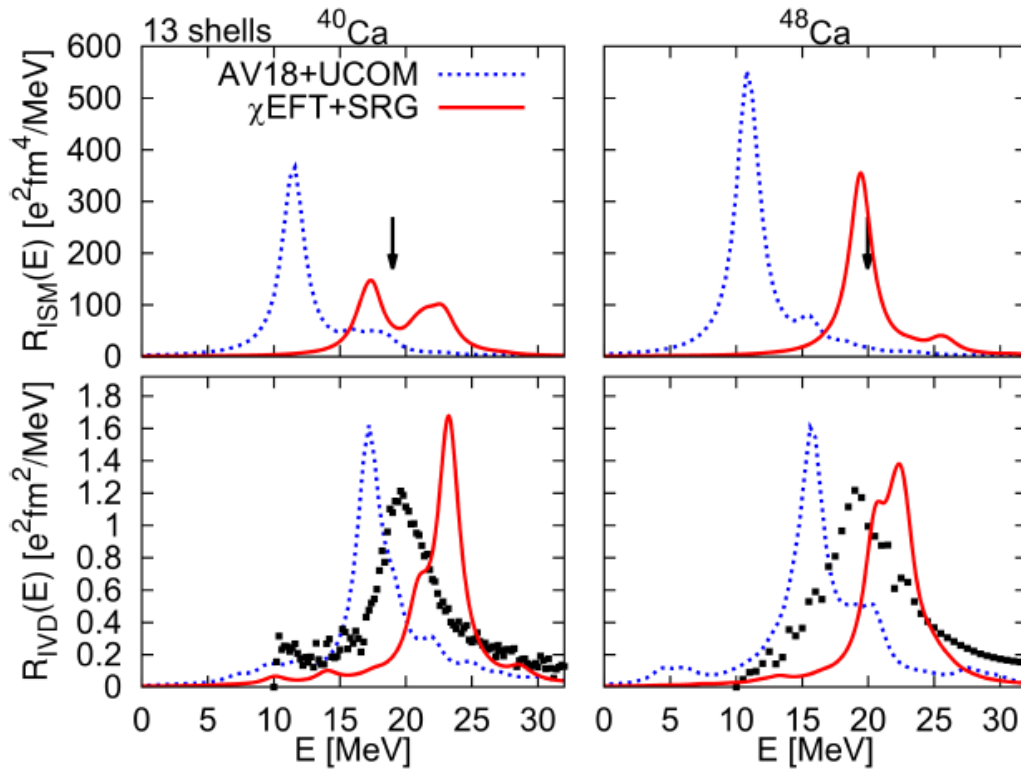
# Chiral NN+NNN and extended RPA

PP, Trippel, Roth, APPB48(Zakopane2016)



# Chiral NN+NNN and extended RPA

PP, Trippel, Roth, APPB48 (Zakopane 2016)



- ❖ Phenomenological functionals: a density-dependent term is required for saturation
- ❖ Without it, the ground state collapses to large density
- ❖ In the present approaches, most nuclei saturate already at the 2N level; the 3N only improves the saturation point.

❖ ***3N force an order of magnitude weaker than a typical density-dependent interaction, e.g.:***

$$t_3 \text{ (UCOM-SRG)} * \rho_0 \div t_3 * \rho_0^\alpha \text{ (D1S)} \sim 0.08$$

- A handle on the “acceptable” strength of a (in-medium) DDI vs. momentum dependence etc.?

P. PAPAKONSTANTINOU AND R. ROTH

PHYSICAL REVIEW C 81, 024317 (2010)

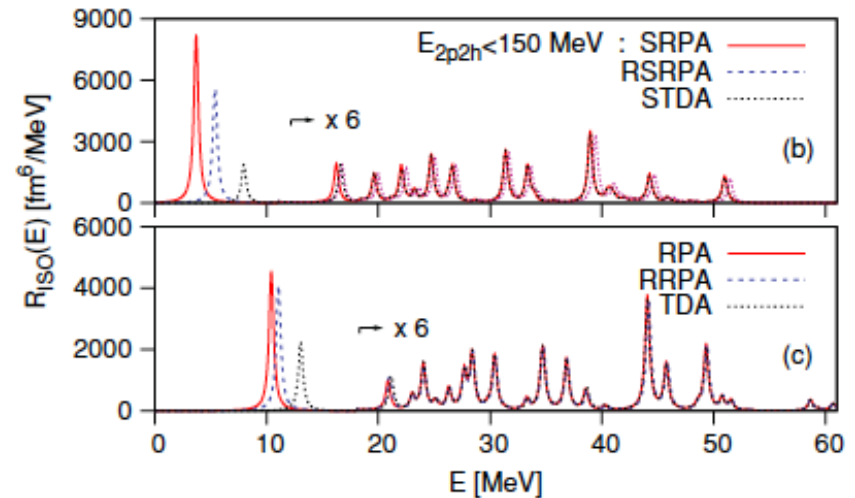
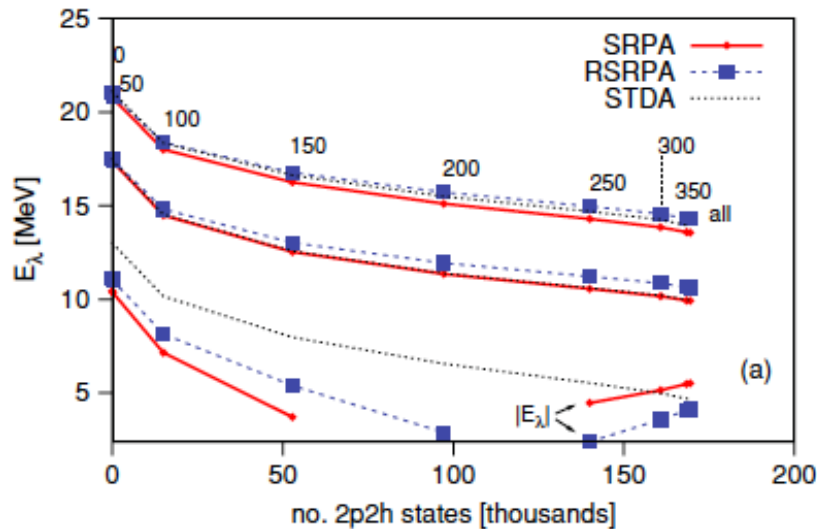


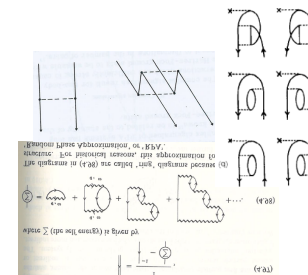
FIG. 9. (Color online) Exploring the influence of ground-state correlations: Isoscalar  $3^-$  response of  $^{16}\text{O}$  ( $\epsilon_{\text{max}} = 14$ ,  $\ell_{\text{max}} = 10$ ) within RPA or SRPA and “renormalized” RPA or SRPA, (RRPA, RSRPA; see text), as well as TDA and STDA (SRPA with  $B = 0$ ). (a) Evolution of the three lowest (closest to zero) eigenvalues as we increase the number of  $2p2h$  configurations considered, starting from zero and up to all  $2p2h$  states available in the single-particle space,  $N_2$ . The numbers assigned to the various points indicate the corresponding  $2p2h$ -energy cutoff imposed,  $E_{2p2h,\text{max}}$ , in MeV. The energy of the lowest eigenstate becomes imaginary for larger  $N_2$  in (R)SRPA; then its amplitude is indicated. Note that the STDA and SRPA results for the second state almost coincide. (b) Strength function for  $E_{2p2h,\text{max}} = 150$  MeV in SRPA, RSRPA, and STDA ( $\Gamma = 2$  MeV). (c) Strength function in RPA, RRPA, and TDA.

See also: PP et al., PRC75,014310

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  - Phenomenological
  - Realistic (chiral)
- ❖ Outlook



- ❖ Model space and Hamiltonian meeting each other
- ❖ Collective phenomena (GRs) and their decays: SRPA represents a smart (tailored) choice of model space.
- ❖ Unitarily transformed interactions are perturbative and do not produce double counting
- ❖ Missing:
  - 3N effects: However, only a small correction is required
  - GSC: could affect substantially less collective states – M1, GT, etc.
  - How dressed should the particles be?
- ❖ A consistent framework is required

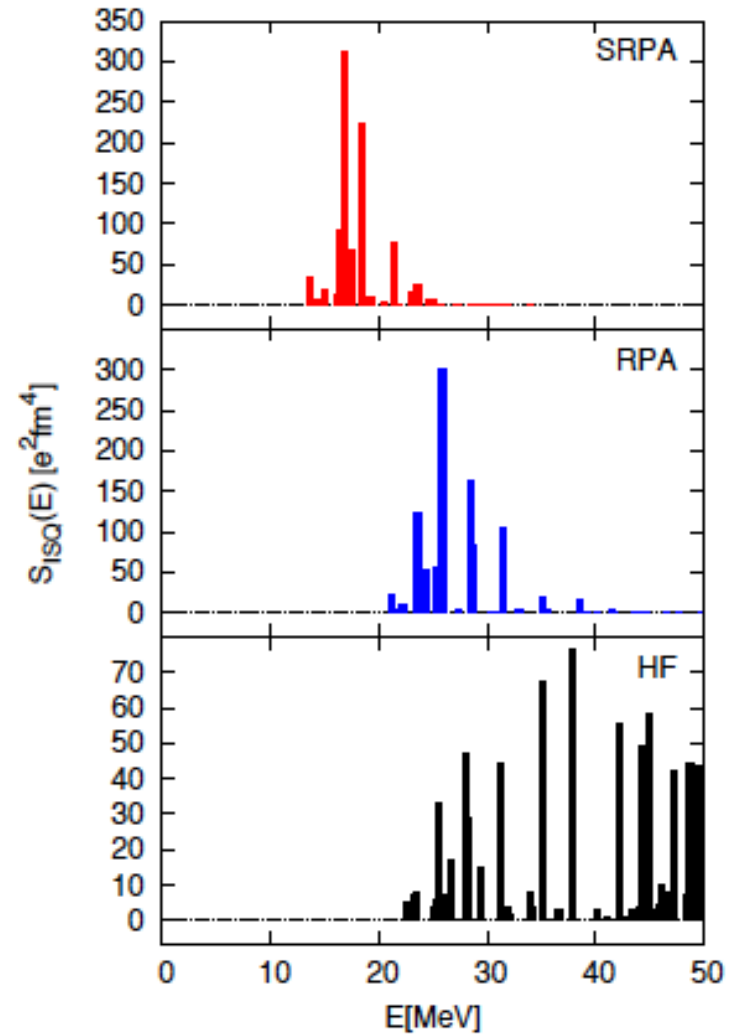
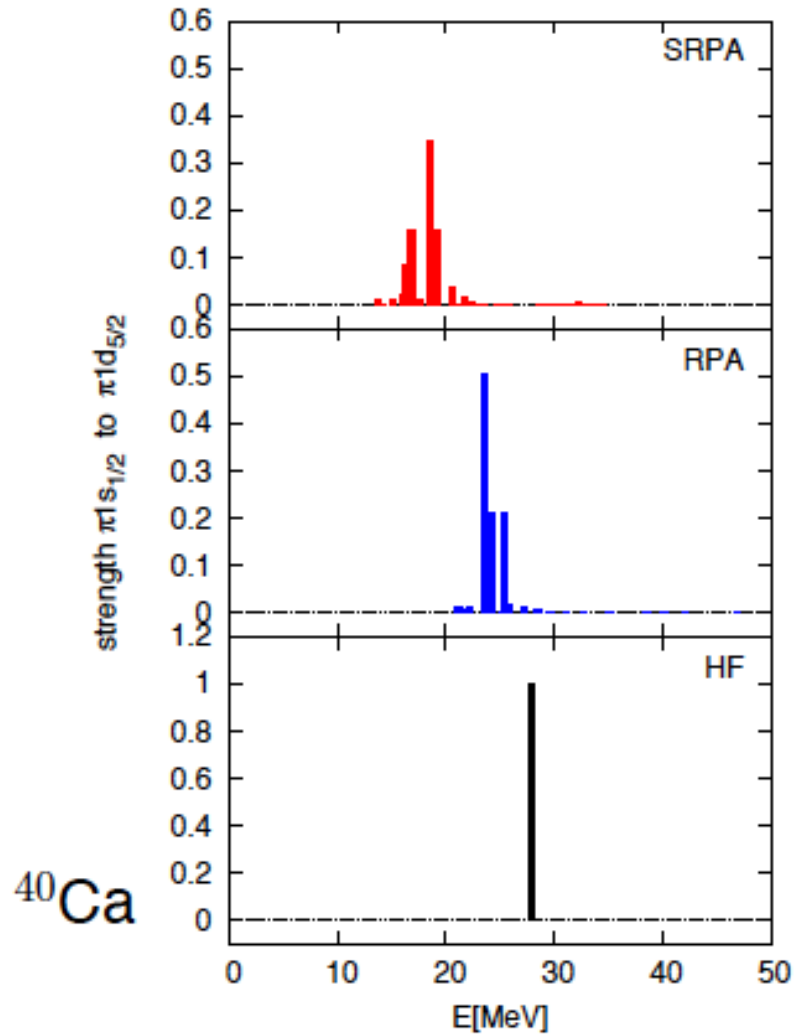


Thank you!

감사합니다~

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and the National Research Foundation (NRF) of Korea  
(2013M7A1A1075764).

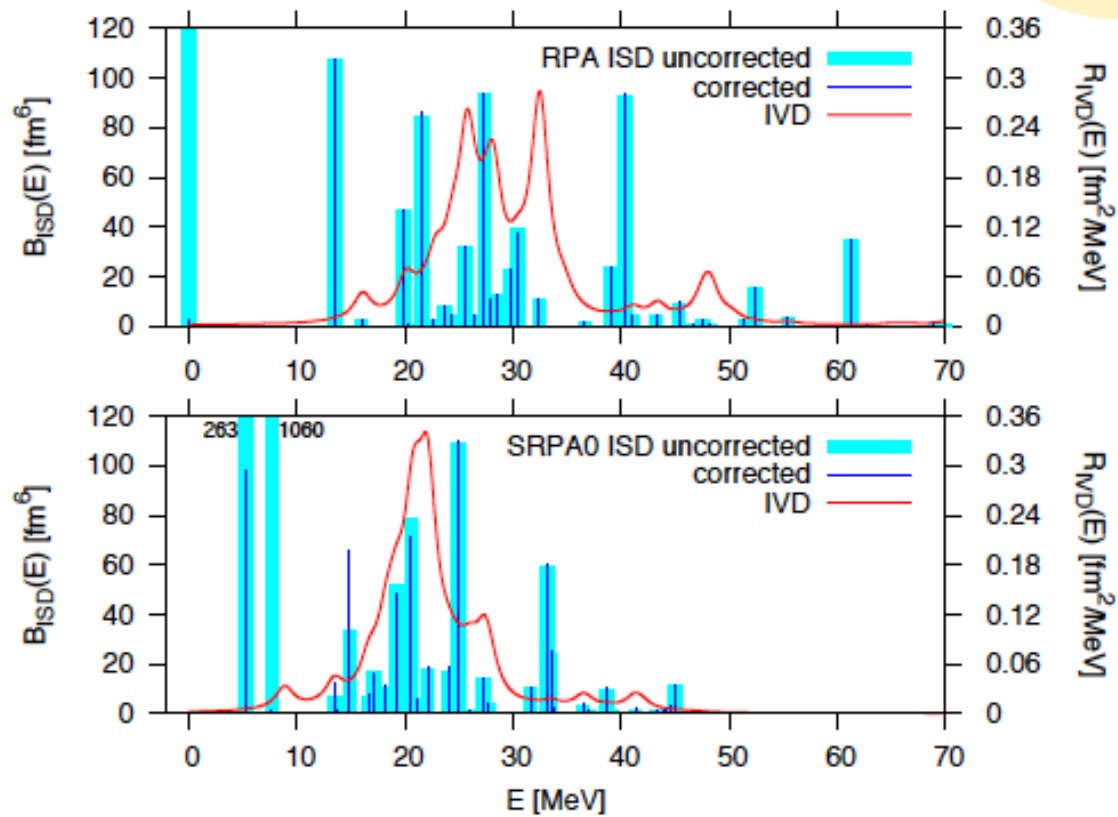
# Fragmentation of $ph$ strength



# Spurious states!

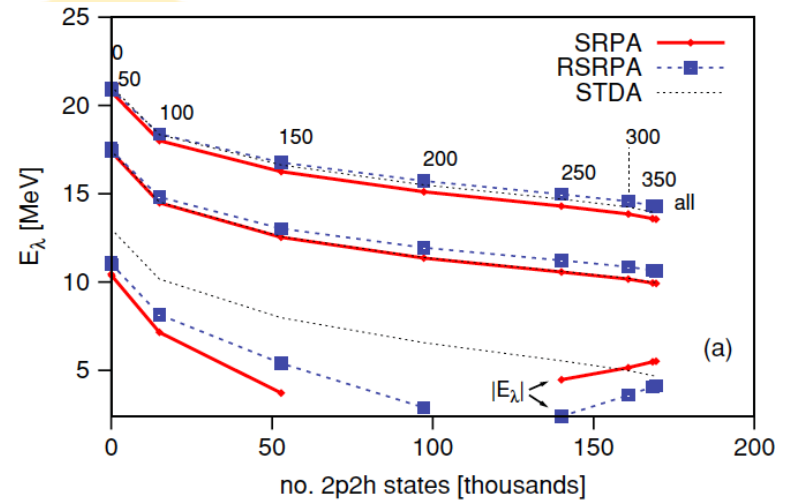
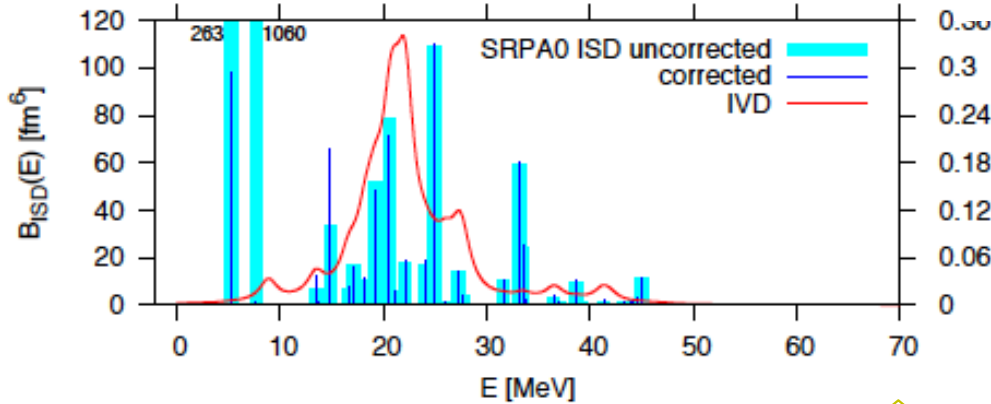
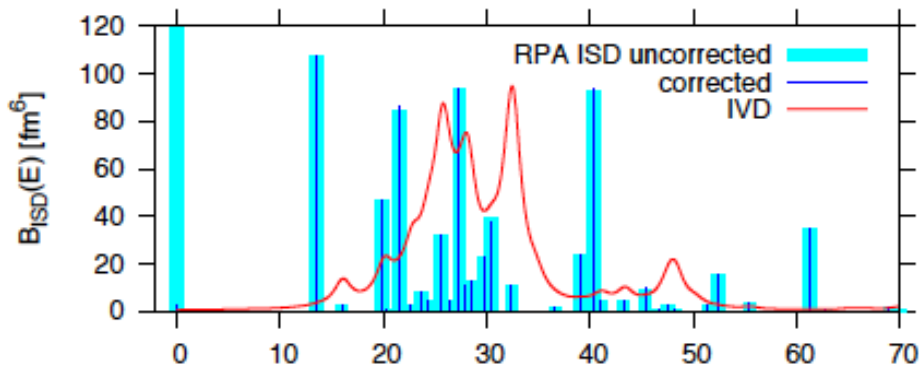
ISD corrected radial operator  $r^3 - \frac{5}{3}\langle r^2 \rangle r$  vs  $r^3$

$^{16}\text{O}$   
 $N_{\text{max}} = 12$



# Spurious states and instabilities in SRPA

ISD corrected radial operator  $r^3 - \frac{5}{3}\langle r^2 \rangle r$  vs  $r^3$



Instability of low-energy states

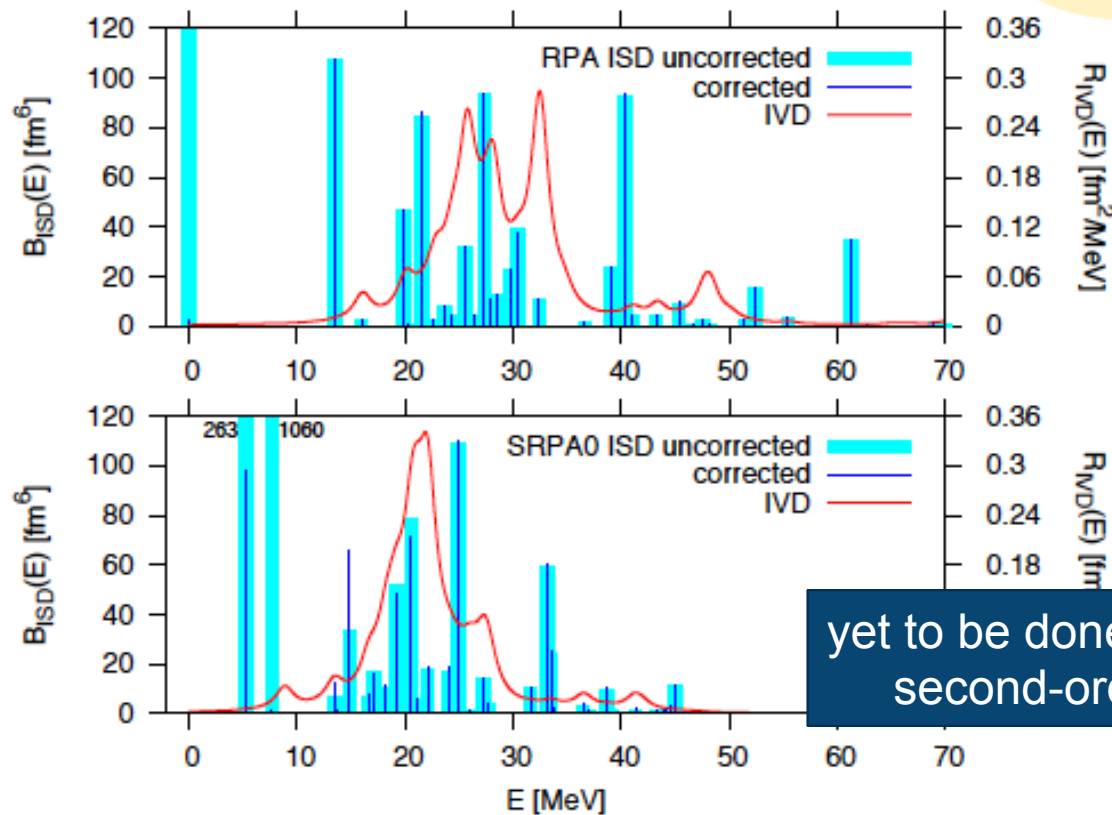


spurious admixtures

# Spurious states!

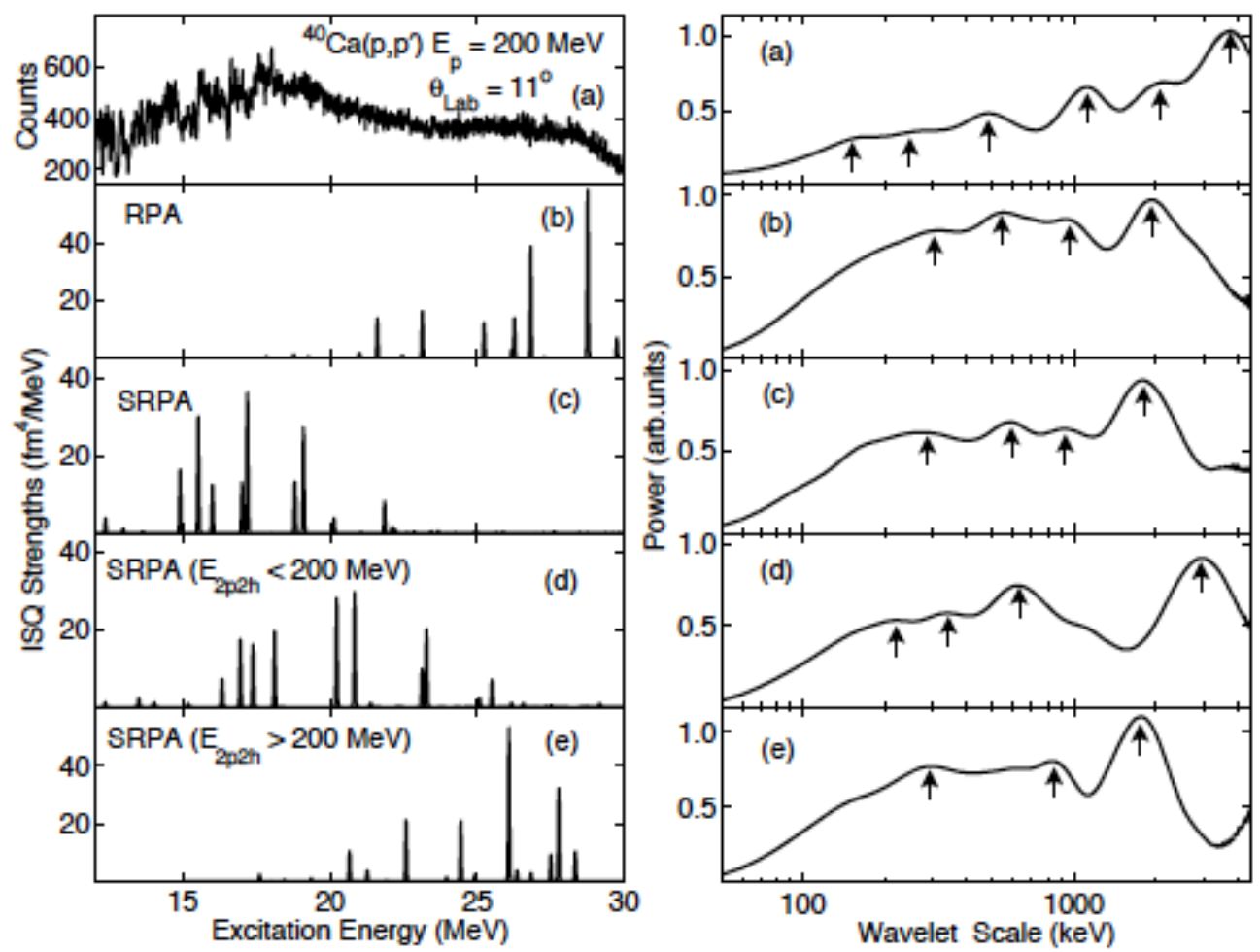
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$^{16}\text{O}$   
 $N_{\text{max}} = 12$



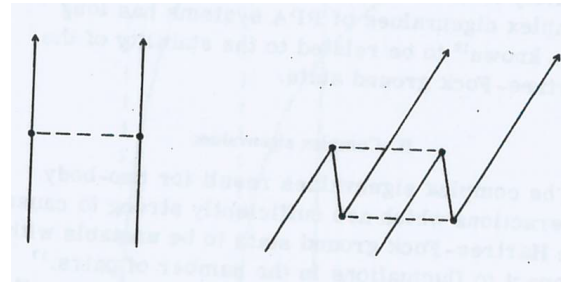
yet to be done: self-consistent second-order formalism

# $^{40}\text{Ca} (p,p')$ and UCOM-SRPA



## ❖ pp-RPA

- From A to A±2 system
- Pairing interaction
- Ladder diagrams



## ❖ ph-RPA

- Same number of particles
- Phonon excitations
- Ring diagrams

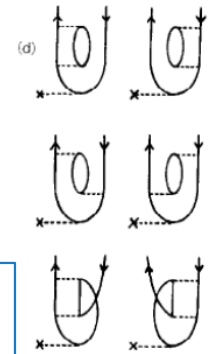
## ❖ ph-2p2h-RPA (Second RPA)

- Extension of ph-RPA
- Reference state

- HF reference state
- Spherical systems (J good quantum number)
- Everything antisymmetrized

where  $\Sigma$  (the self energy) is given by

The diagrams in (4.98) are called 'ring' diagrams because of their ring-like structure. For historical reasons, this approximation for  $G$  is called the 'Random Phase Approximation' or 'RPA'.





## ❖ Ground-state correlations

- Correction to the Hartree-Fock energy via the backward amplitudes
- Correction real in the absence of phase transitions (e.g. superfluidity in pp-RPA)

## ❖ Excited states

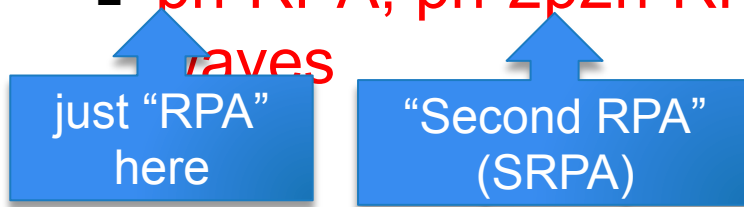
- pp-RPA: 2-particle transfer, spectroscopy
- ph-RPA, ph-2p2h-RPA: Vibrational states, sound waves

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## ❖ Excited states

- pp-RPA: 2-particle transfer, spectroscopy
- **ph-RPA, ph-2p2h-RPA: Vibrational states, sound waves**



# First SRPA results

- Softened realistic Hamiltonian (AV18, UCOM)
- No arbitrary truncations
- Explored: Energetic shift; EWSR; ph fragmentation; etc.

