

# **KIDS energy density functional for nuclei and nuclear matter**

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- **KIDS**

**K**orea: **I**BS-**D**aegu-**S**ungkyunkwan

- **Collaboration**

- **First work**

P. Papakonstantinou, Y. Lim (IBS), T.-S. Park (Sungkyunkwan Univ.)

- **Nuclear structures**

H. Gil (Kyungpook) ,P. Papakonstantinou, (IBS), Y. Oh (Kyungpook)

- **Astrophysics and Heavy-ion collision**

M. Kim (Pusan), Y. -M. Kim (UNIST), Y. Kim (IBS), C.-H. Lee

(Pusan),

S. Jeon (McGill)

# Outline

EFT for two nucleons

Building a model

Application

Summary

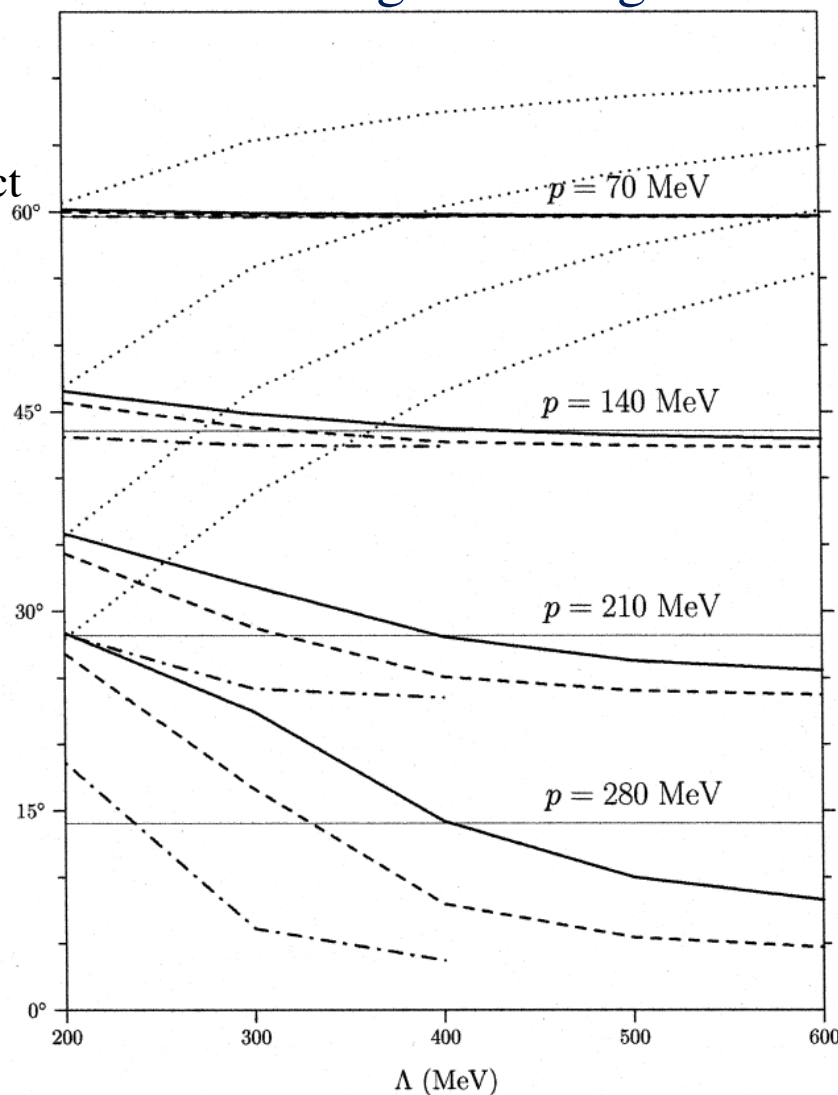
EFT for two nucleons

- Pionful EFT for the np scattering (CHH, T.-S. Park, D.-P. Min, PLB 473 (2000))

NN potential up to NLO following Weinberg counting rule

LO: OPE, 4N contact

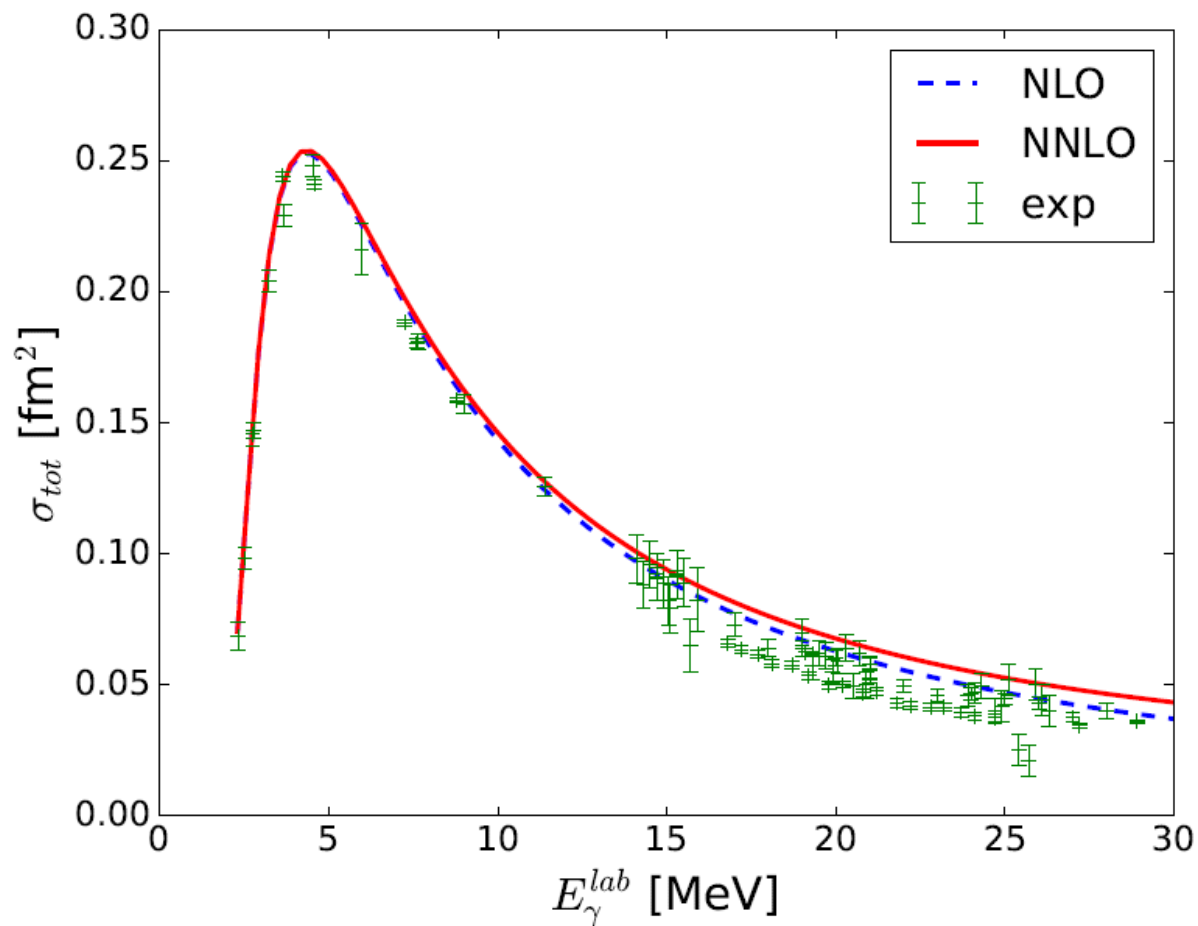
NLO: TPE, 4N contact



- np phase shifts
- LO: dotted
  - reproduce experiment only at small momentum
  - strongly cut-off dependent
- NLO: solid
  - less dependent on cut-off
  - good agreement at large momenta

- Pionless EFT for the  $np \rightarrow dg$  (Y.-H. Song, S.-i. Ando, CHH, PRC 96 (2017))

Pionless EFT: Valid for  $p \ll mp$  ( $p \sim mp$ ,  $E_\gamma \sim 10$  MeV)



- Can we have similar thing for the nuclear matter and nuclei?
  - Systematic and converging expansion
  - Uncertainty quantification
- Could be ‘Yes’

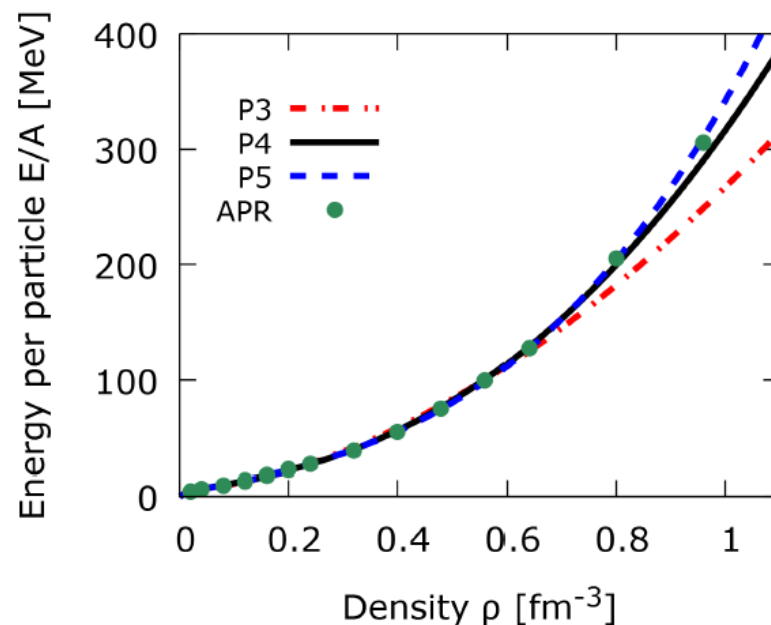


FIG. 5: Energy per particle of pure neutron matter with models P3, P4 and P5 presented in Table IV. Here, the symmetric EoS parameters  $\alpha_i$  are fixed as model S3b in Table I.

(H. Gil, Y.-M. Kim, CHH, P. Papakonstantinou, Y. Oh, PRC 100 (2019))

Building a model



- Scale in nuclear matter

- Momentum scale:  $k_F$  (270 MeV for saturated symmetric matter)
- $k_F/m_\rho$  : less than 1 even at  $r = 8r_0$  ( $k_F \propto r^{1/3}$ )
- Expand the energy density in powers of  $k_F/m_\rho$  ( $m_\rho$ : rho-meson mass)

- Rules

- Rule1: Expand EDF for nuclear matter in powers of  $k_F/m_\rho$
- Rule2: Fit the parameters to the well-known nuclear matter properties
- Rule3: Keeping the parameters unchanged, apply the model to nuclei

- KIDS Ansatz

$$\mathcal{E}(\rho, \delta) = \mathcal{T}(\rho, \delta) + \sum_{i=0}^3 c_i(\delta) \rho^{1+i/3}, \quad \rho = \rho_n + \rho_p$$

- Fitting

$$c_i(\delta) = \alpha_i + \delta^2 \beta_i, \quad \delta \equiv (\rho_n - \rho_p) / \rho$$

- $c_0(0), c_1(0), c_2(0)$ :  $r_0, E/A, K_0$  (assume  $c_3(0) = 0$ )
- $c_i(1)$  : APR PNM EoS (14 data in  $r = 0.02 - 0.96 \text{ fm}^{-3}$ )

- Fitting result

$$\chi^2(\delta) = \sum_j \exp\{-\beta\rho_j/\varrho_0\} \left( \frac{\mathcal{E}(\rho_j) - D_j}{\mathcal{T}(\rho_j)} \right)^2 ; \quad \beta \geq 0$$

- **Fitting to APR PNM EoS**

	<b>c0</b>	<b>c0 , c1</b>	<b>c0 - c2 (P3)</b>	<b>c0 - c3 (P4)</b>	<b>c0 - c5 (P5)</b>
<b>c2</b>	0.071632	0.001566	0.000529	0.000138	0.000115

- Fitting improves with more terms
- Improvement saturates at P5 (5 terms for PNM)
- **Double check: Fitting to QMC PNM EoS** (J. Carlson et al., Rev, Mod. Phys. 87

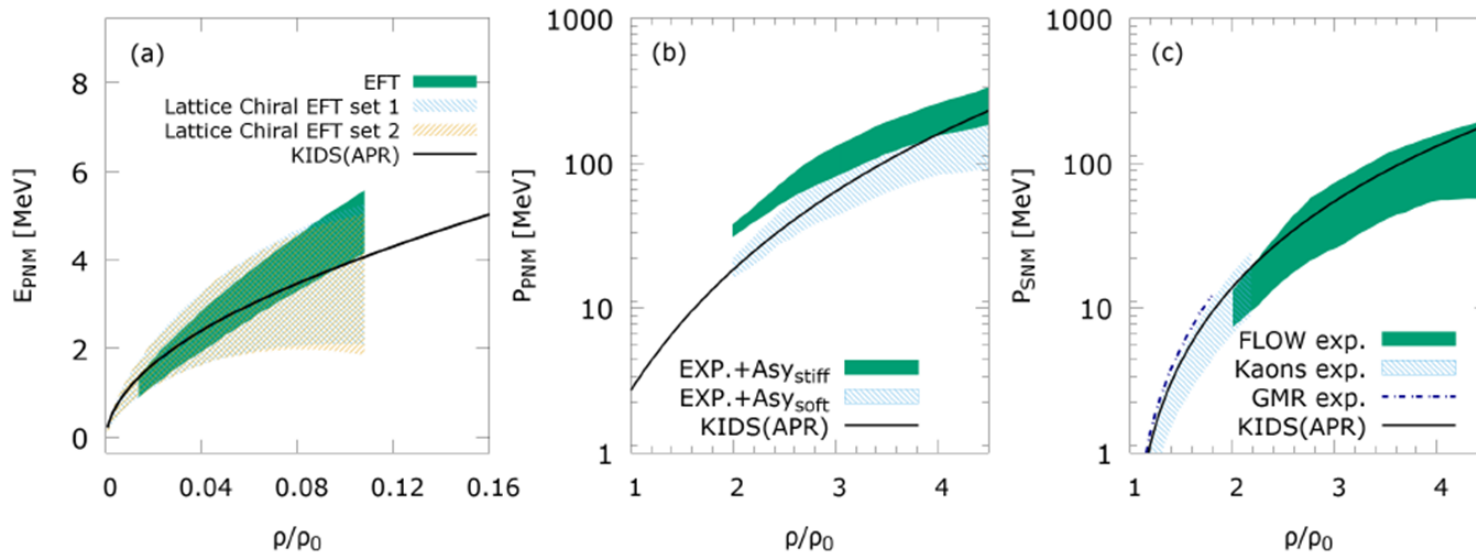
(2015))

<b>P4</b>	<b>P5</b>	<b>P6</b>

# • Nuclear matter properties: P4 parameters

Dutra et al., PRC 85 (2012)

	K0 (MeV)	-Q0 (MeV)	J (MeV)	L (MeV)	Kt (MeV)	S(r0/2)/J	3PPNM/(L r0)
KIDS	240.00*	372.65	32.75	49.10	-377.06	0.667	1.03
Exp./Em p.	200-260	200-1200	30-35	40-76	-760,-37 2	0.57-0.86	0.90-1.10



Lattice chiral EFT: E .Epelbaum et al., EPJA40 (2009)

# Application

Nuclei

Neutron star

- Skyrme type force

$$\begin{aligned}
 v_{ij} = & (t_0 + y_0 P_\sigma) \delta(\mathbf{r}_i - \mathbf{r}_j) + \frac{1}{6} \sum_{n=1}^{N-1} (t_{3n} + y_{3n} P_\sigma) \rho^{n/3} \delta(\mathbf{r}_i - \mathbf{r}_j) \\
 & + \frac{1}{2} (t_1 + y_1 P_\sigma) [\delta(\mathbf{r}_i - \mathbf{r}_j) \mathbf{k}^2 + \mathbf{k}'^2 \delta(\mathbf{r}_i - \mathbf{r}_j)] + (t_2 + y_2 P_\sigma) \mathbf{k}' \cdot \delta(\mathbf{r}_i - \mathbf{r}_j) \mathbf{k} \\
 & + iW_0 \mathbf{k}' \times \delta(\mathbf{r}_i - \mathbf{r}_j) \mathbf{k} \cdot (\sigma_i + \sigma_j),
 \end{aligned}$$

- Energy density

$$\begin{aligned}
 \mathcal{E} = & \frac{\hbar^2}{2m} \tau + \frac{3}{8} t_0 \rho - \frac{1}{8} (t_0 + 2y_0) \rho \delta^2 + \frac{1}{16} \sum_{n=1}^{N-1} t_{3n} \rho^{1+n/3} - \frac{1}{48} \sum_{n=1}^{N-1} (t_{3n} + 2y_{3n}) \rho^{1+n/3} \delta^2 \\
 & + \frac{1}{64} (9t_1 - 5t_2 - 4y_2) \frac{(\nabla \rho)^2}{\rho} - \frac{1}{64} (3t_1 + 6y_1 + t_2 + 2y_2) \frac{(\nabla \rho \delta)^2}{\rho} \\
 & + \frac{1}{8} (2t_1 + y_1 + 2t_2 + y_2) \tau - \frac{1}{8} (t_1 + 2y_1 - t_2 - 2y_2) \sum_q \frac{\rho_q \tau_q}{\rho} \\
 & + \frac{1}{2} W_0 \left( \frac{\mathbf{J} \cdot \rho}{\rho} + \sum_q \frac{\mathbf{J}_q \cdot \nabla \rho_q}{\rho} \right), \quad \tau = \frac{3}{5} \left( \frac{6\pi^2}{\nu} \right)^{2/3} \rho^{5/3}
 \end{aligned}$$

$$\text{KIDS EDF: } \mathcal{E}(\rho, \delta) = \mathcal{T}(\rho, \delta) + \sum_{i=0}^{N-1} c_i(\delta) \rho^{1+i/3},$$

- Transformation of coefficients

$$t_0 = \frac{8}{3}c_0(0), \quad y_0 = \frac{8}{3}c_0(0) - 4c_0(1),$$

$$t_{3n} = 16c_n(0), \quad y_{3n} = 16c_n(0) - 24c_n(1), \quad (n \neq 2)$$

$$t_{32} = 16c_2(0) - \frac{3}{5} \left( \frac{3}{2}\pi^2 \right)^{2/3} \theta_s \equiv 16c_2(0)(1 - k)$$

$$y_{32} = 16c_2(0) - 24c_2(1) + \frac{3}{5}(3\pi^2)^{2/3} \left( 3\theta_\mu - \frac{\theta_s}{2^{2/3}} \right) \equiv [16c_2(0) - 24c_2(1)](1 - k')$$

with

$$\theta_s \equiv 3t_1 + 5t_2 + 4y_2 = \frac{5}{3}(3\pi^2)^{-2/3}16c_2(0)k$$

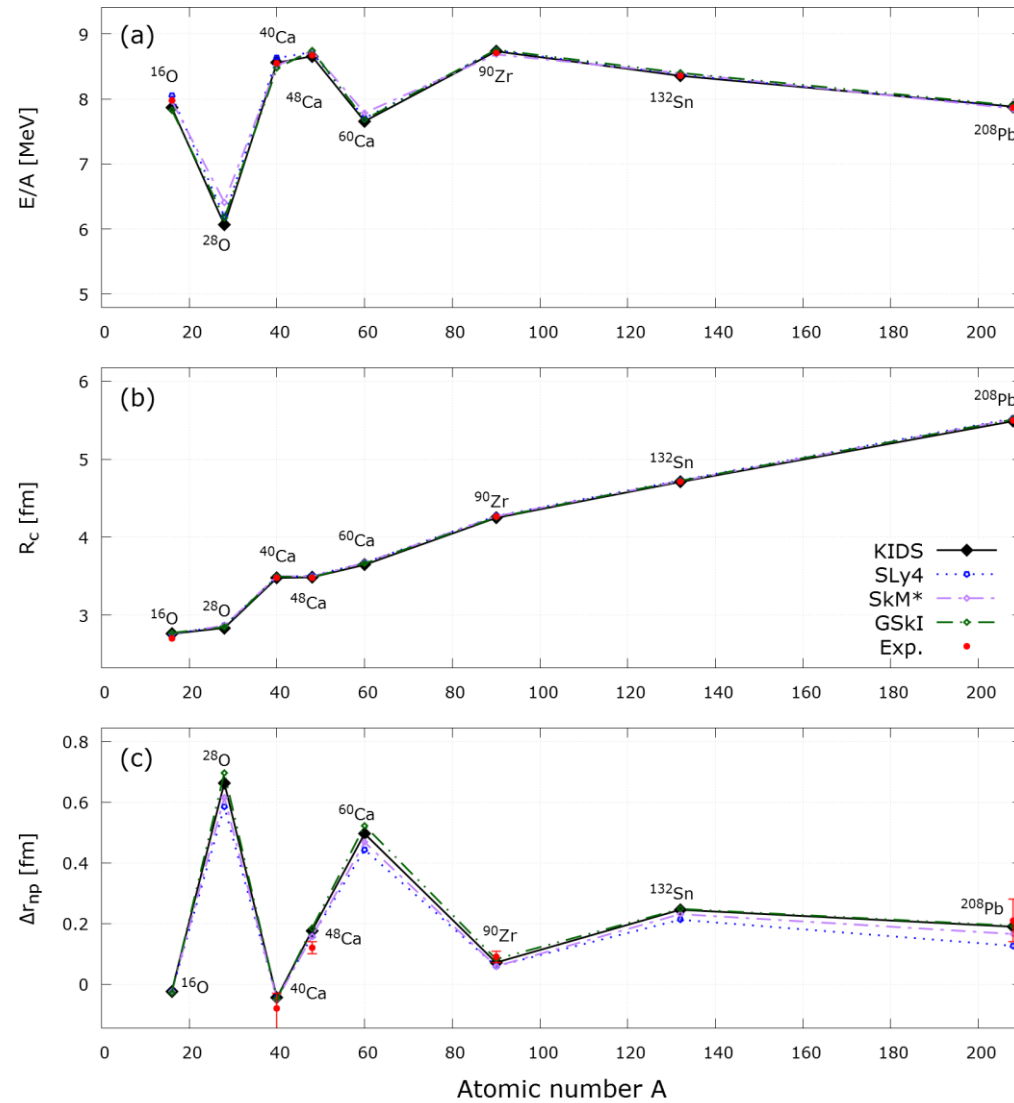
$$\theta_\mu \equiv t_1 + 3t_2 - y_1 + 3y_2 = -\frac{5}{9}(3\pi^2)^{-2/3}[16c_2(0) - 24c_2(1)]k' + \frac{\theta_s}{3 \cdot 2^{2/3}}.$$

$k, k'$ : fraction of gradient terms in the  $c_2$  term (r2/3)

- Two parameter

fitting

- Assume  $k=k'$
- Assume  $y_1=y_2=0$
- Remaining parameters  $k, W_0$
- Fit to  $E/A$  and  $R_c$  of  $^{40}\text{Ca}$ ,  $^{48}\text{Ca}$ ,  $^{208}\text{Pb}$



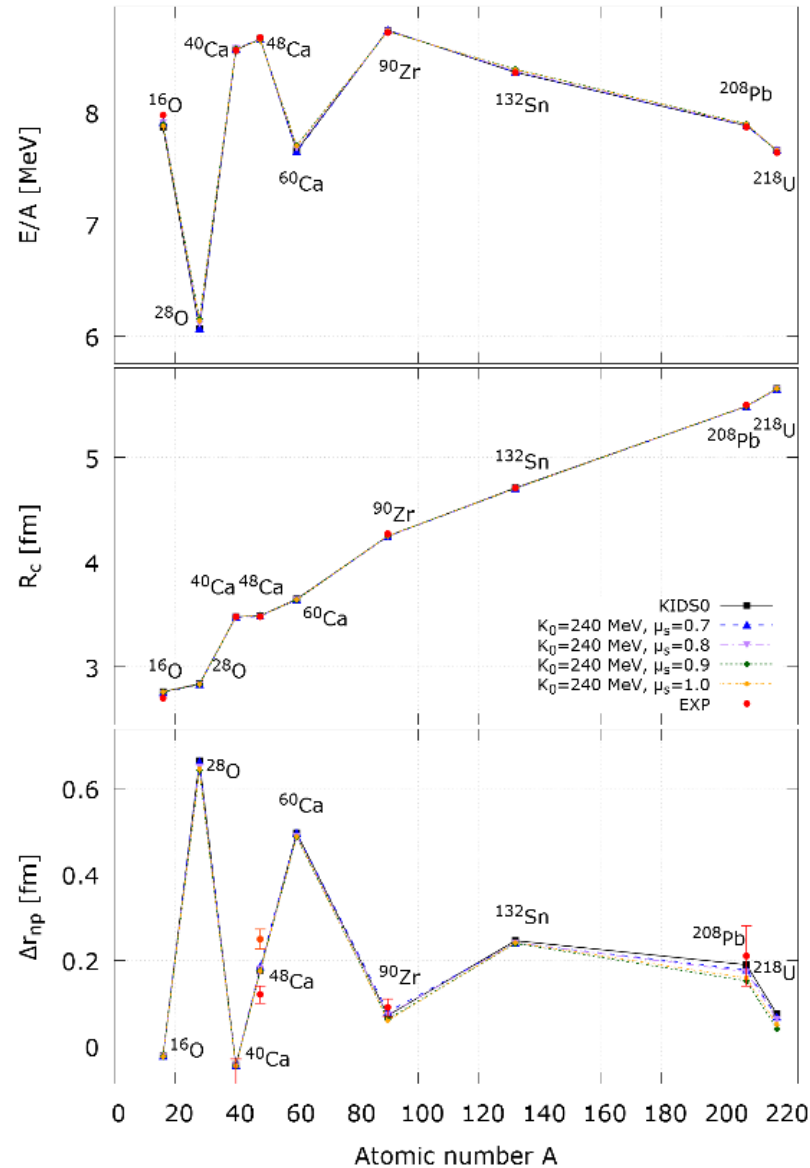


- Dependence on the effective mass

- Assume non-zero  $y_1, y_2$
- Fit to specific isoscalar
- effective mass
- Isoscalar effective mass

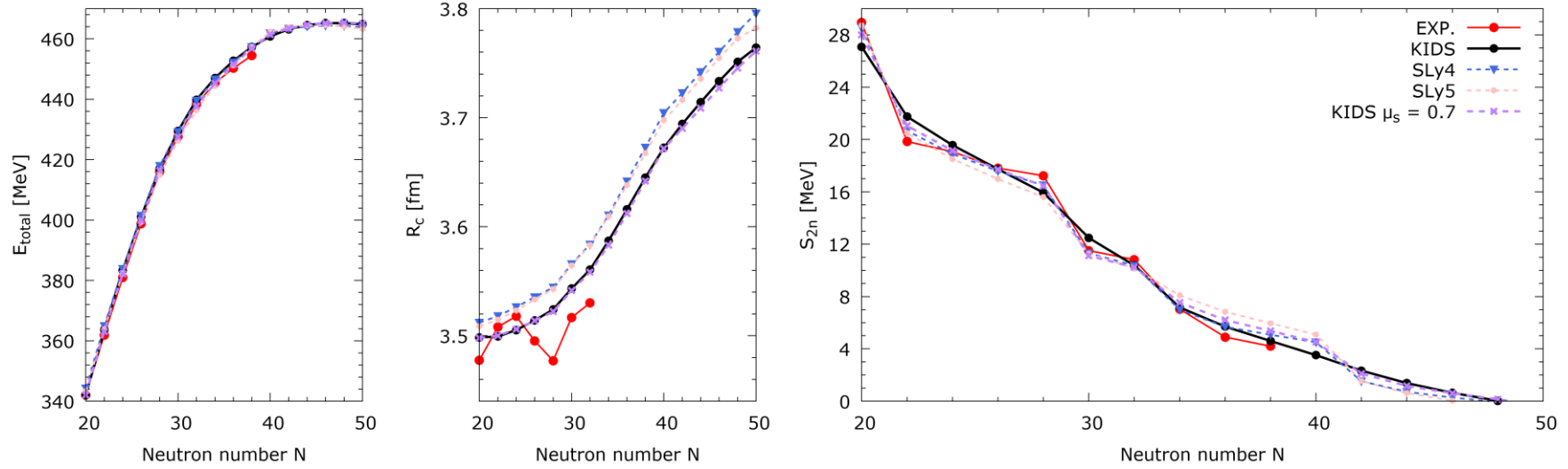
$$\mu_s^{-1} \equiv (m_{IS}^*/m)^{-1} = 1 + \frac{m}{8\hbar^2} \rho \theta_s$$

$$\theta_s \equiv 3t_1 + 5t_2 + 4y_2$$

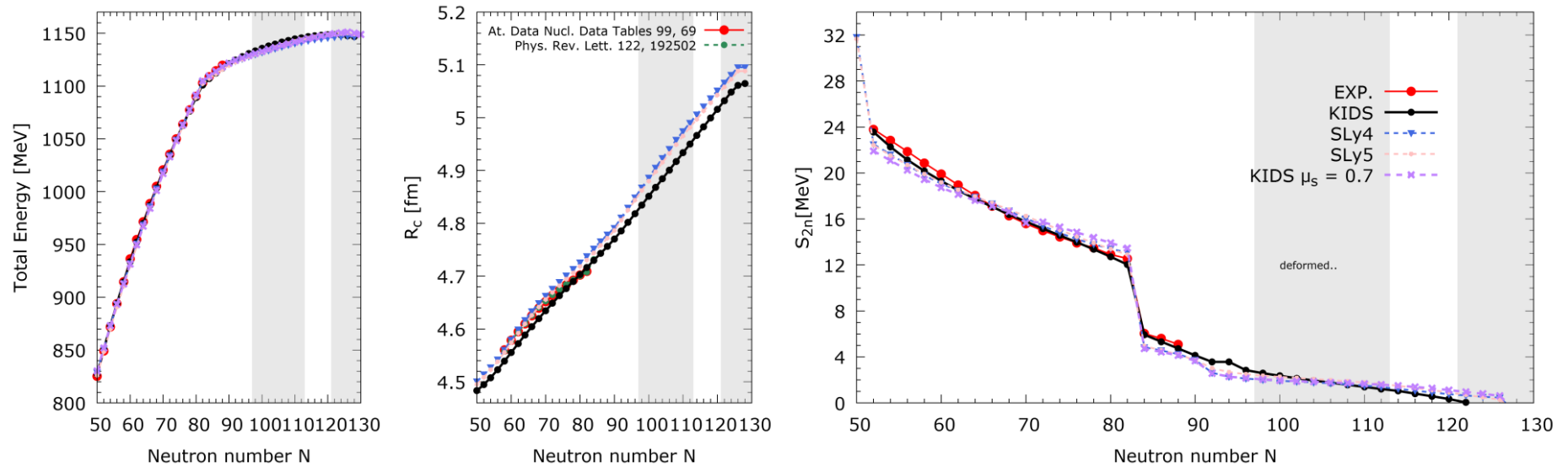


- Isotopes with pairing force

Ca isotope ( $Z = 20$ )

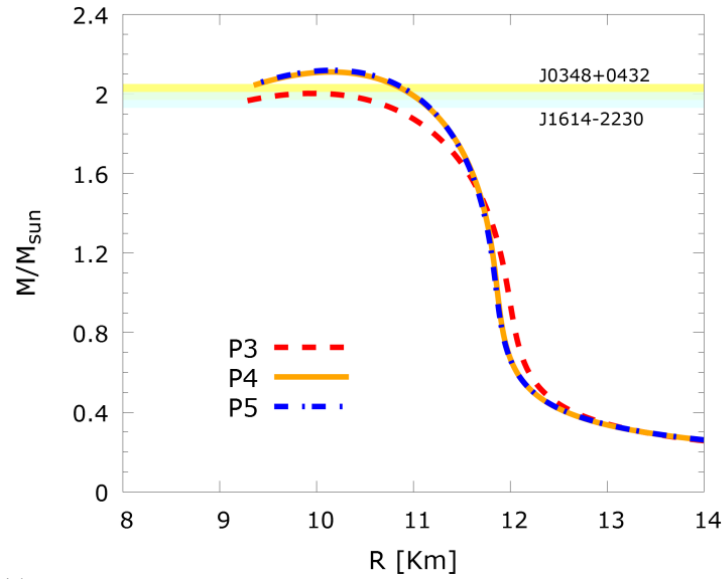


Sn isotope ( $Z = 50$ )



- Neutron star

- Mass-radius



- Tidal deformability

	GSkI	SLy4	SkI4	SGI	KIDS
$R_{1.4M_{\odot}}$ [km]	11.94	11.82	12.46	12.77	11.79
$k_2(1.4M_{\odot})$	0.079	0.077	0.092	0.097	0.076
$\lambda(1.4M_{\odot})$ [ $10^{36} \text{g cm}^2 \text{s}^2$ ]	1.906	1.770	2.772	3.292	1.737
$\Lambda(1.4M_{\odot})$	337.2	312.9	490.9	583.0	307.5

Recent analysis (B. P. Abbott et al. arXiv:180511581v1 [ $\xi \Lambda(1.4M_{\odot}) = 190_{-120}^{+390}$ ])

# Summary

- Link low energy EFT to nuclear matter and structures

- Novel EDF constructed with simple rules



- Nuclear matter properties agree well with exp./emp. data

- Most updated data of neutron stars well reproduced

- Nuclear properties reproduced over wide range of mass number

- Effective mass controlled without altering bulk properties



- Odd-odd, odd-even nuclei: tensor force, deformation

- Mass table whole nuclear chart

- Response to external probe