

# Hyperon-resonance productions with photon and meson beams

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Contents based on

S.i.N. [PRD96, 076021 \(2017\)](#)

J.K.Ahn, S.i.N. [PRD98, 114012 \(2018\)](#)

S.i.N., A.Hosaka [arXiv:1902.09106 \[hep-ph\]](#)

J.K.Ahn, S.B.Yang, K.S.Choi, S.i.N. *In preparation*



Internationalization  
Cooperation  
Synergy

# What are we interested in?

Hyperons of terra incognita:  $\Lambda(1405)$  and  $\Xi(1690)$

New types of hadrons

Structures seem different from others

Rarely known physical properties

How can we probe the puzzles via experiments?

Theoretical guides and estimations for Future

What are we interested in?

$\Lambda(1405)$  at LEPS/SPring-8 and CLAS/JLab

$$\gamma^{(*)} p \rightarrow K^+ \pi^+ \Sigma^- \quad \text{S.i.N. PRD96, 076021 (2017)}$$

$$\gamma^{(*)} p \rightarrow K^{*+} \pi^0 \Sigma^0 \quad \text{S.i.N., A.Hosaka arXiv:1902.09106 [hep-ph]}$$

$\Xi(1690)$  at Hadron hall/J-PARC

$$K^- p \rightarrow K^+ K^- \Lambda \quad \text{J.K.Ahn, S.i.N. PRD98, 114012 (2018)}$$

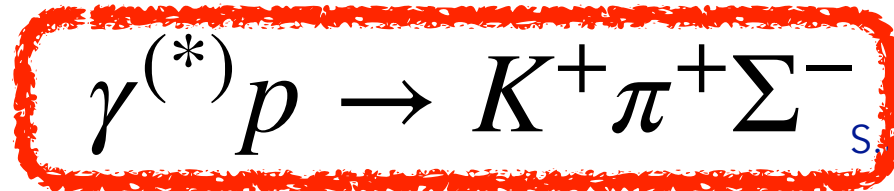
S=0 and S=-1 resonances at Belle/KEK (Dr. S.B.Yang)

$$\Lambda_c^+ \rightarrow \pi \bar{K} p$$

J.K.Ahn, S.B.Yang, K.S.Choi, S.i.N. In preparation

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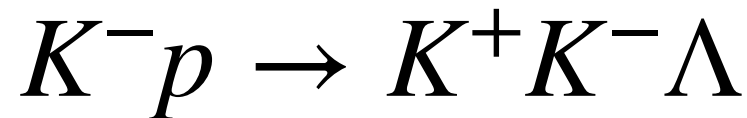


S.i.N. PRD96, 076021 (2017)



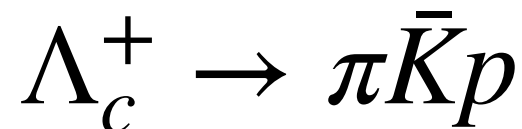
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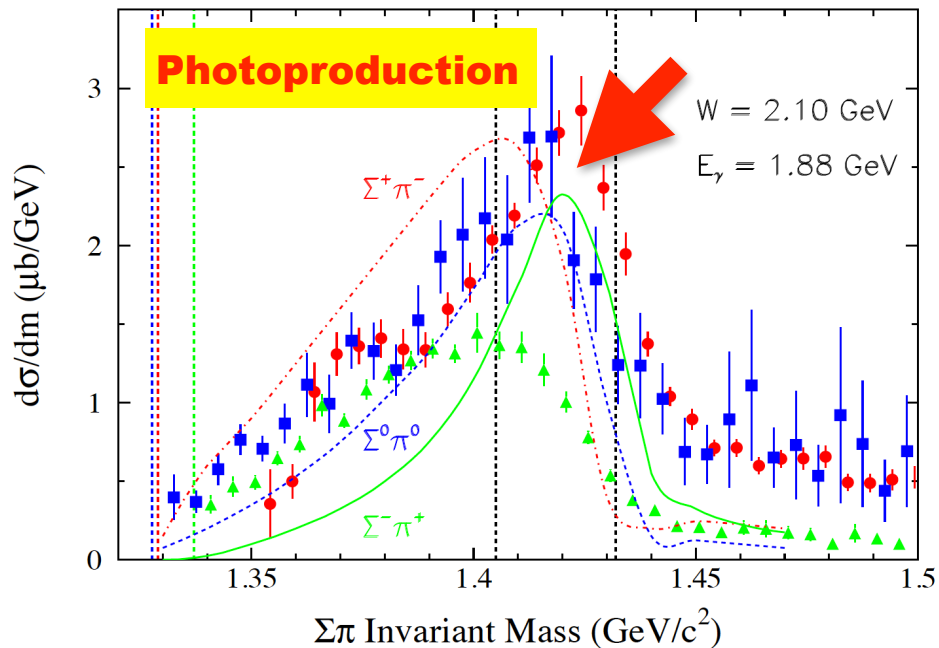
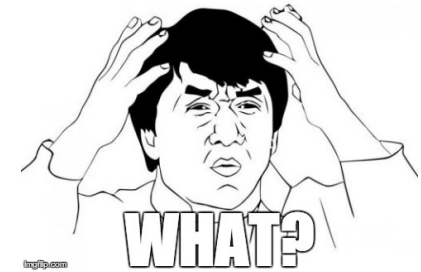


J.K.Ahn, S.B.Yang, K.S.Choi, S.i.N. In preparation

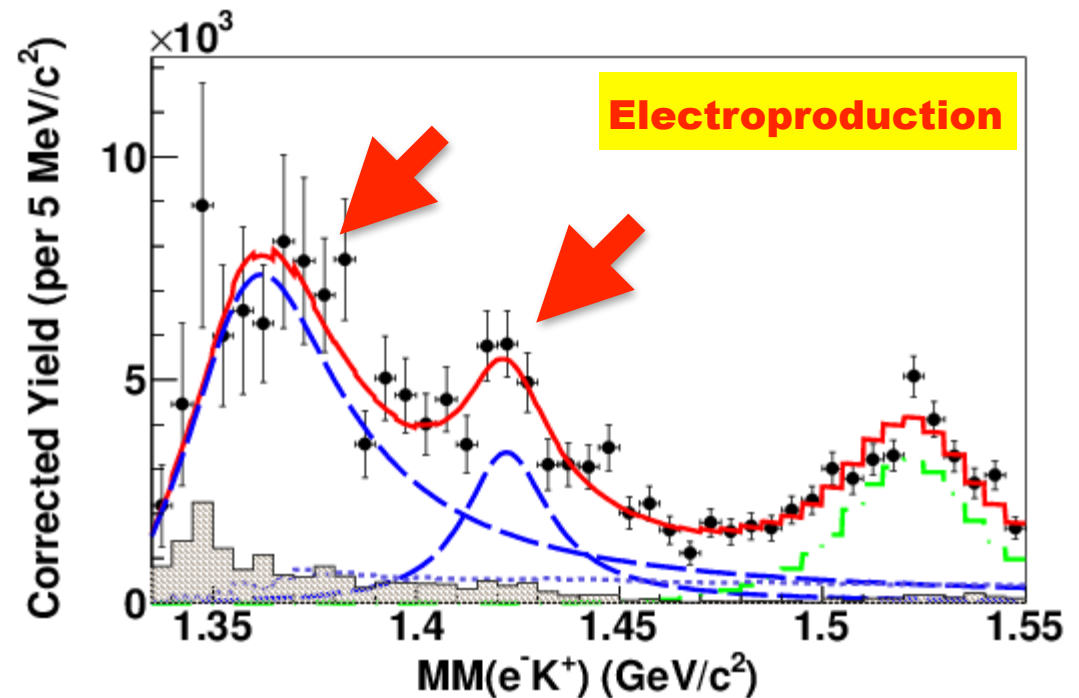
# Motivation

Obvious difference between  $\Lambda(1405)$  invariant masses

from photo- and electro-production in experiments: **Why??**



K. Moriya et al. [CLAS], PRC87, no. 3, 035206 (2013)



H. Y. Lu et al. [CLAS], PRC88, 045202 (2013)

Single peak vs. Double peak



## Possible explanations (real vs. virtual photon)

Virtual photon has scalar component  $\rightarrow$  Increasing t-ch contribution?

In terms of strong interaction,  
No differences for two poles  
(even more, numerically small...)

Virtual photon probes  $\Lambda(1405)$  EM form factor  $\rightarrow$  New EM excitation?

Hadron EM form factors play the roles  
then, nontrivial interference occurs?

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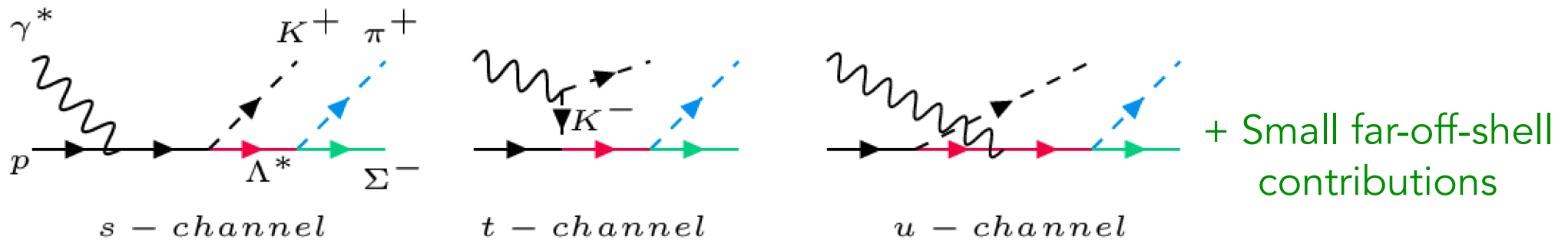
Virtual photon probes  $\Lambda(1405)$  EM form factor → New EM excitation?

Hadron EM form factors play the roles  
then, nontrivial interference occurs?

**We will take this scenario!!!**

## Theoretical framework

Effective Lagrangian approach in Born approximation at tree level



Relevant strong couplings from ChUM

$B$	Mass	Width	$g_{KNB}$	$g_{\pi\Sigma B}$	$\kappa_B$
$p$	938.272 MeV	$\sim 0$ MeV	—	—	2.79
$L$	1368 MeV	100 MeV	$1.2 + 1.7i$	$-2.5 - 1.5i$	0.30
$H$	1423 MeV	50 MeV	$-2.5 + 0.94i$	$0.42 - 1.4i$	0.41



## Theoretical framework

Scattering amplitudes conserving WT identity for arb.  $Q^2$  values

$$i\mathcal{M}_{\text{total}} = \sum_{\Lambda^*=H,L} \sum_x i\mathcal{M}_x^{\Lambda^*}, \quad i\mathcal{M}_x^{\Lambda^*} = ieg_{KN\Lambda^*} g_{\pi\Sigma\Lambda^*} \frac{\bar{u}(k_5)[\not{q}_{4+5} + M_{\Lambda^*}^2]\Gamma_x^{\Lambda^*} u(k_2)}{M_{\pi^+\Sigma^-}^2 - M_{\Lambda^*}^2 - i\Gamma_{\Lambda^*} M_{\Lambda^*}},$$

$$\Gamma_s^{\Lambda^*} = \frac{\not{k}_1 + \not{k}_2 + M_N}{s - M_N^2} \left[ \not{\epsilon} + (F_1^p - 1) \left[ \not{\epsilon} + \frac{(\epsilon \cdot k_1)\not{k}_1}{Q^2 + \delta} \right] - \frac{\kappa_p F_2^p}{2M_N} (\not{\epsilon}\not{k}_1 - \epsilon \cdot k_1) \right] F_c,$$

$$\Gamma_t^{\Lambda^*} = \frac{1}{t - M_K^2} \left[ \epsilon \cdot (2k_3 - k_1) + (F_K - 1) \left[ \epsilon \cdot (2k_3 - k_1) + \frac{(\epsilon \cdot k_1)[k_1 \cdot (2k_3 - k_1)]}{Q^2 + \delta} \right] \right] F_c,$$

$$\Gamma_u^{\Lambda^*} = F_u \left[ F_1^{\Lambda^*} \left[ \not{\epsilon} + \frac{(\epsilon \cdot k_1)\not{k}_1}{Q^2 + \delta} \right] - \frac{\kappa_{\Lambda^*} F_2^{\Lambda^*}}{2M_N} (\not{\epsilon}\not{k}_1 - \epsilon \cdot k_1) \right] \frac{\not{k}_2 - \not{k}_3 + M_{\Lambda^*}}{u - M_{\Lambda^*}^2},$$

Strong form factors satisfying WT identity as well

$$i\mathcal{M}_{\text{total}}^{\text{dressed}} = (i\mathcal{M}_{Es} + i\mathcal{M}_{Et} + i\mathcal{M}_{Eu})F_c + i\mathcal{M}_{Ms}F_s + i\mathcal{M}_{Mt}F_t + i\mathcal{M}_{Mu}F_u.$$

$$F_c = 1 - (1 - F_s)(1 - F_t),$$

S.i.N. et al., JKPS59, 2676 (2011).

# Theoretical framework

Photon polarization in cm frame

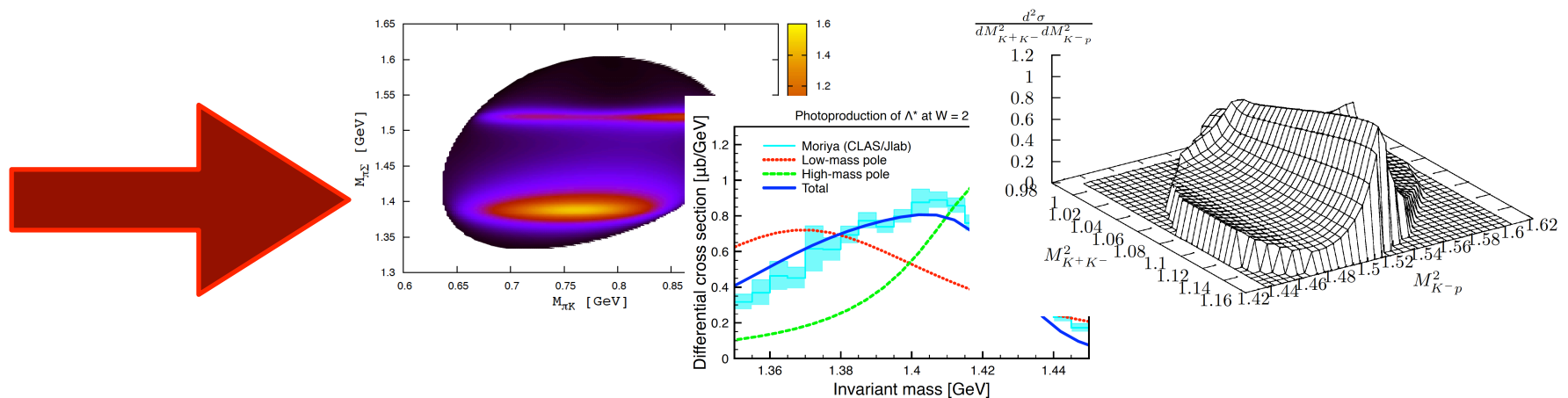
$$\varepsilon_x = (0, \sqrt{1 + \varepsilon}, 0, 0), \quad \varepsilon_y = (0, 0, \sqrt{1 - \varepsilon}, 0), \quad \varepsilon_z = \frac{\sqrt{2\varepsilon}}{\sqrt{|Q^2|}} (k, 0, 0, E_1).$$

Transverse-polarization parameter  $\varepsilon = (0.3 \sim 0.7)$  at CLAS, so that

we choose it  $\varepsilon = 0.5$

We compute double differential cross section  $W=2.4$  GeV

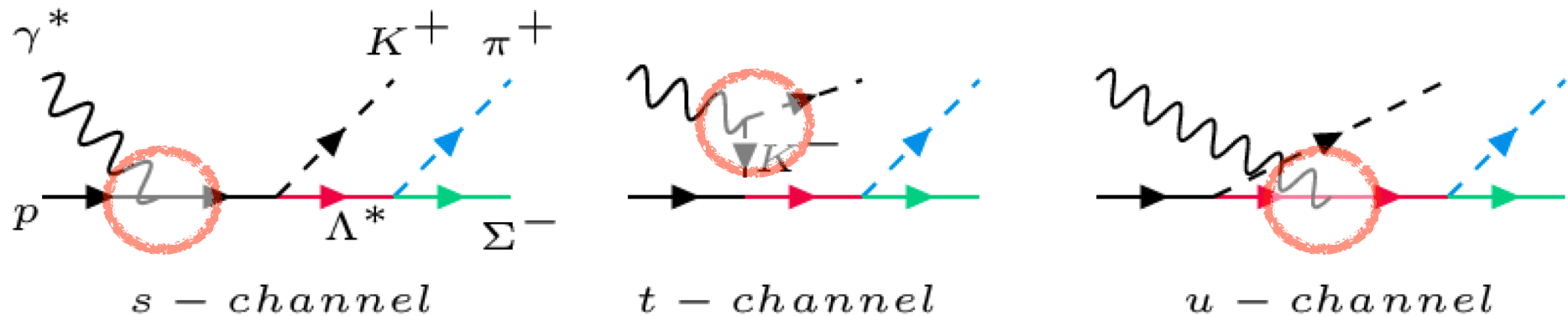
$$\frac{d^2\sigma_{\gamma^{(*)}p \rightarrow K^+\pi^+\Sigma^-}}{dM_{K^+\pi^+}dM_{\pi^+\Sigma^-}} = \frac{1}{4|\vec{k}_{\gamma^{(*)}}||E_{\gamma^{(*)}} - E_N|} \frac{1}{128\pi^4 s} \int M_{K^+\pi^+} M_{\pi^+\Sigma^-} d\cos\theta_{K^+} d\phi_{\Sigma^-} \frac{1}{2n} \sum |\mathcal{M}_{\gamma^{(*)}p \rightarrow K^+\pi^+\Sigma^-}|^2,$$



## Theoretical framework

Hadron EM form factors for electroproduction

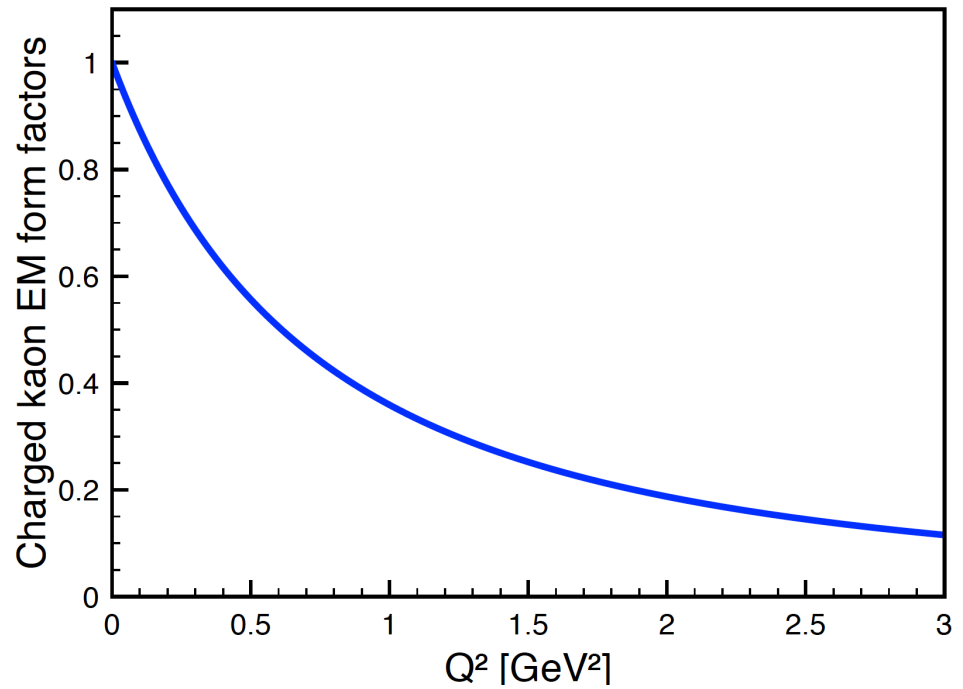
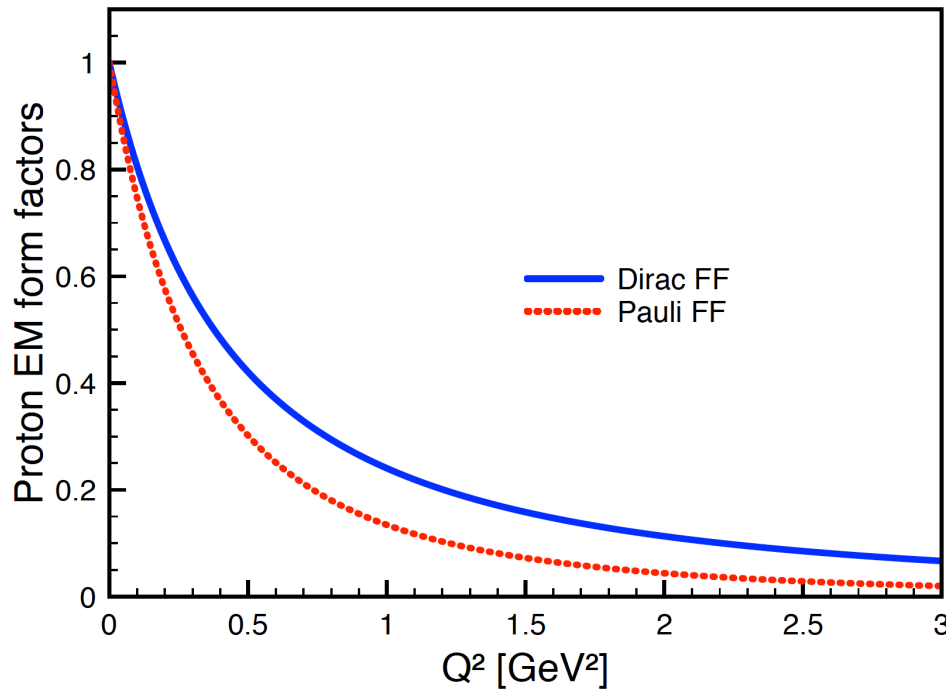
Where hadron EM form factors appear?



Proton, Kaon, and  $\Lambda(1405)$ , three EM form factors necessary

# Theoretical framework

EM form factors (FF) for proton and kaon: **Well known**



$$G_E^p(Q^2) \simeq G_D(Q^2), \quad G_M^p(Q^2) \simeq \mu_p G_D(Q^2), \quad G_D(Q^2) = \left[ \frac{1}{1 + Q^2 \langle r^2 \rangle_E^p / 12} \right]^2,$$

J. J. Kelly, Phys. Rev. C 70, 068202 (2004).

$$F_K(Q^2) = \frac{1}{[1 + Q^2 / (0.845 \text{ GeV})^2 + Q^4 / (1.270 \text{ GeV})^4]}, \quad \text{W. Jaus, Phys. Rev. D 44, 2851 (1991).}$$

## Theoretical framework

EMFF for  $\Lambda(1405)$ : Less known so far. So How to model it?

How can we construct  $\Lambda(1405)$  EMFF???

1) It's neutral so possibly similar structure to neutron EMFF

2) EM charge rms radii relates to EMFF

$$\langle r^2 \rangle_E^{H,L} = -6 \frac{dG_E^{H,L}(Q^2)}{dQ^2} \Big|_{Q^2=0}, \quad \langle r^2 \rangle_M^{H,L} = -\frac{6}{\mu_{H,L}} \frac{dG_M^{H,L}(Q^2)}{dQ^2} \Big|_{Q^2=0},$$

**cf) Galster parameterization**

$$G_E^n(Q^2) \simeq -\frac{a\mu_n\tau}{1+b\tau} G_D(Q^2).$$

$$G_E^{H,L}(Q^2) = -\frac{\langle r^2 \rangle_E^{H,L}}{6} Q^2 F_K(Q^2) \left[ \frac{1}{1 + Q^2 \langle r^2 \rangle_M^{H,L} / 12} \right]^2, \quad G_M^{H,L}(Q^2) \approx \mu_{H,L} \left[ \frac{1}{1 + Q^2 \langle r^2 \rangle_M^{H,L} / 12} \right]^2$$

2) EM information of  $\Lambda(1405)$  from ChUM: EM charge rms radii

M.M.Kaskulov, P.Grabmayr, EPJA 19, 157 (2004).

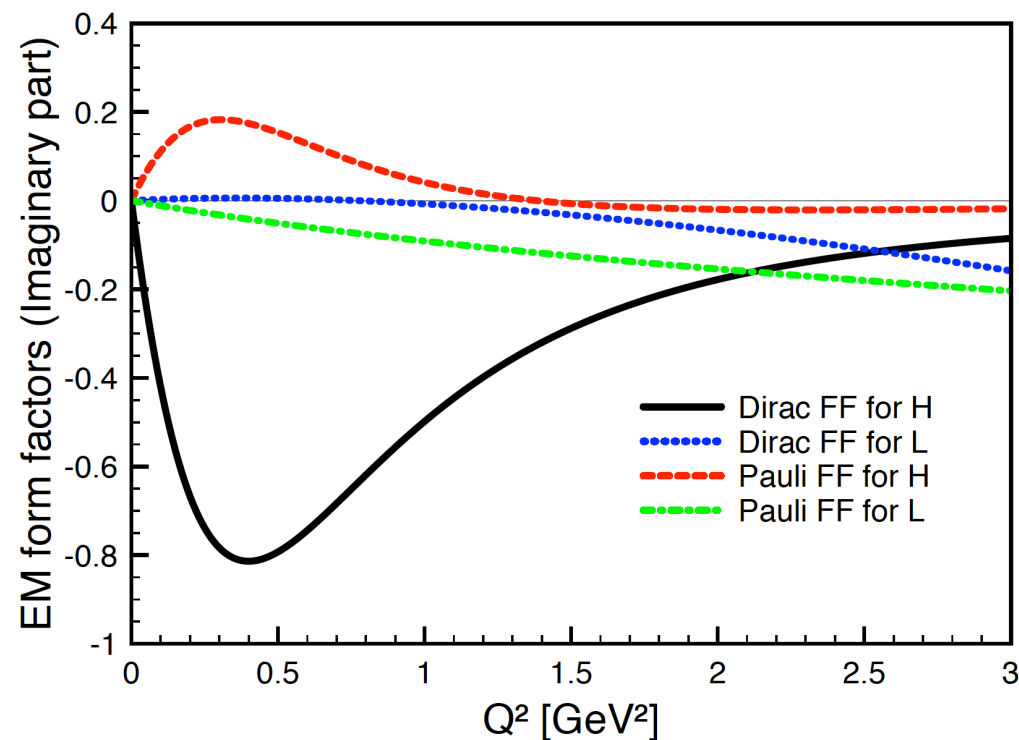
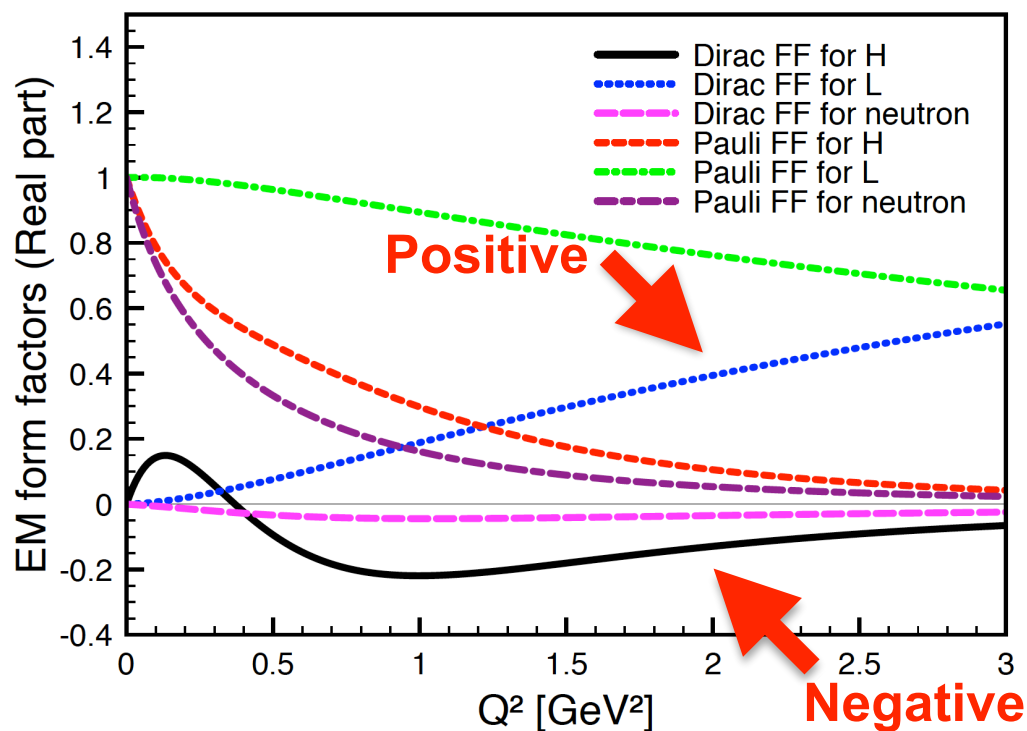
**Neutron EMFF + Charge rms radii from ChUM  $\approx$   $\Lambda(1405)$  EMFF**

# Theoretical framework

EMFF for  $\Lambda(1405)$ : **Less known so far. So How to model it?**

T. Sekihara, T. Hyodo and D. Jido, Phys. Lett. B 669, 133 (2008).

$\langle r^2 \rangle_E^H$	$\langle r^2 \rangle_M^H$	$\langle r^2 \rangle_E^L$	$\langle r^2 \rangle_M^L$	$\langle r^2 \rangle_E^n$
$-3.365 + 7.783i$	$6.859 - 10.455i$	$0.462 - 0.051i$	$-0.334 + 0.539i$	$-2.877 \pm 0.077$

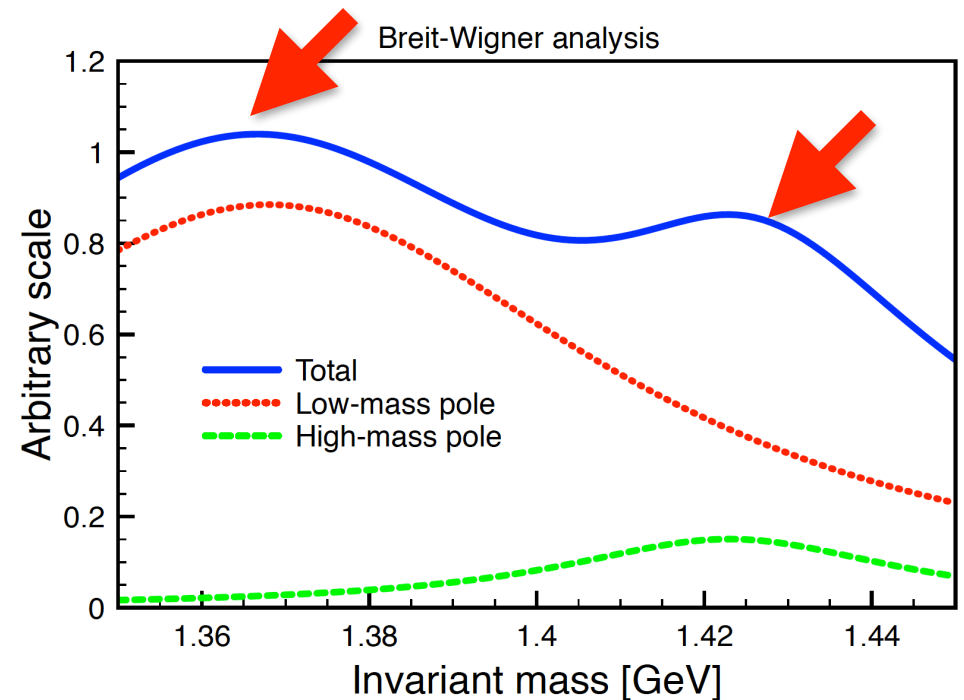
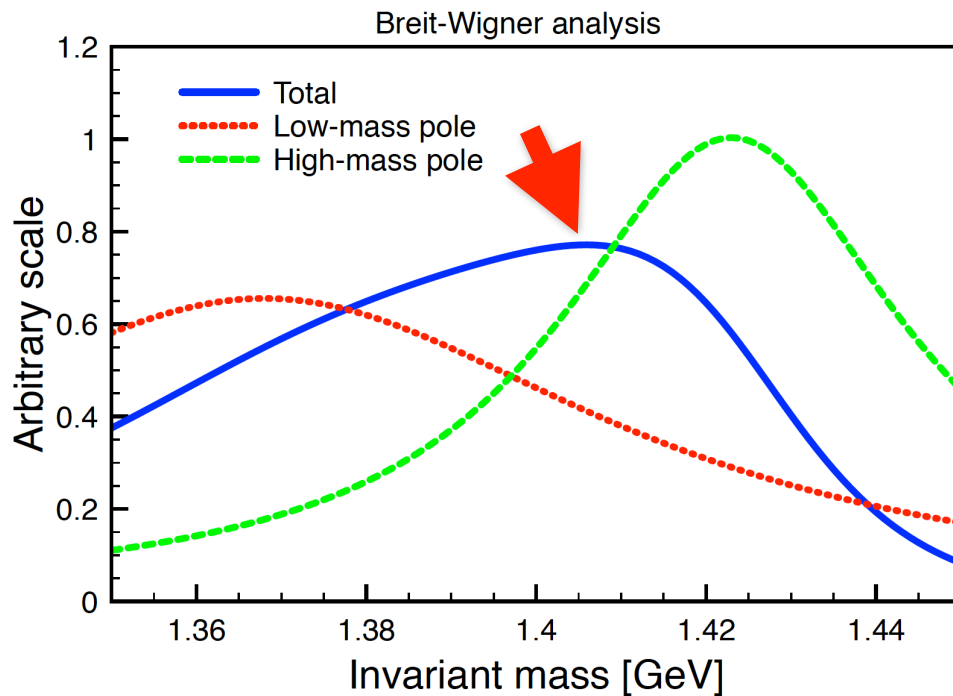


# Theoretical framework

## Invariant mass plots with Breit-Wigner type distributions (test)

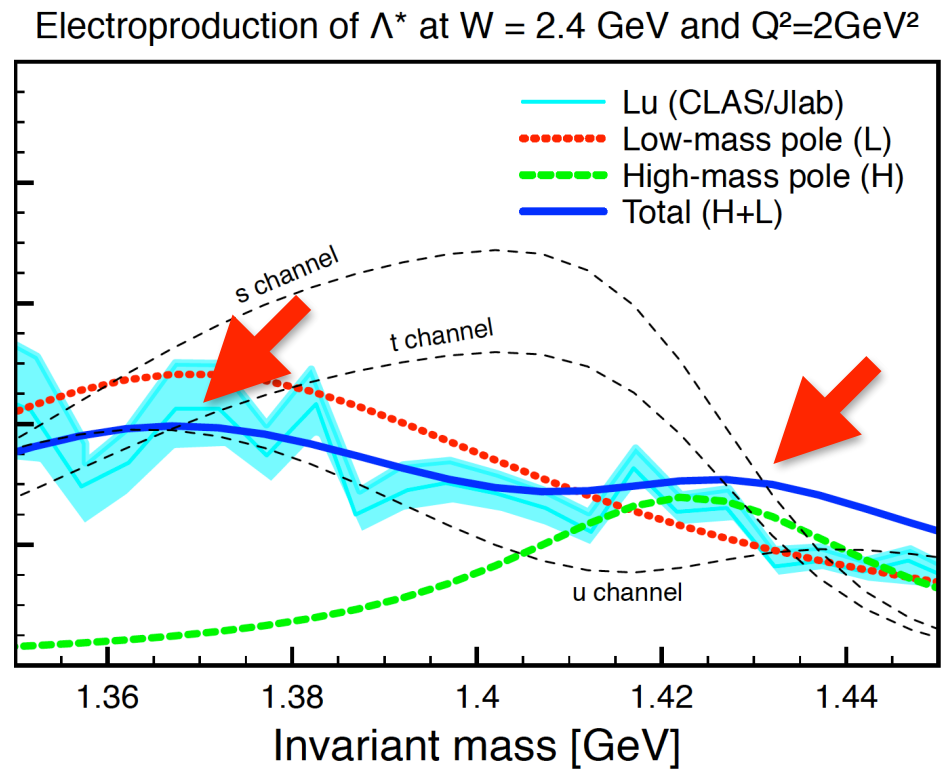
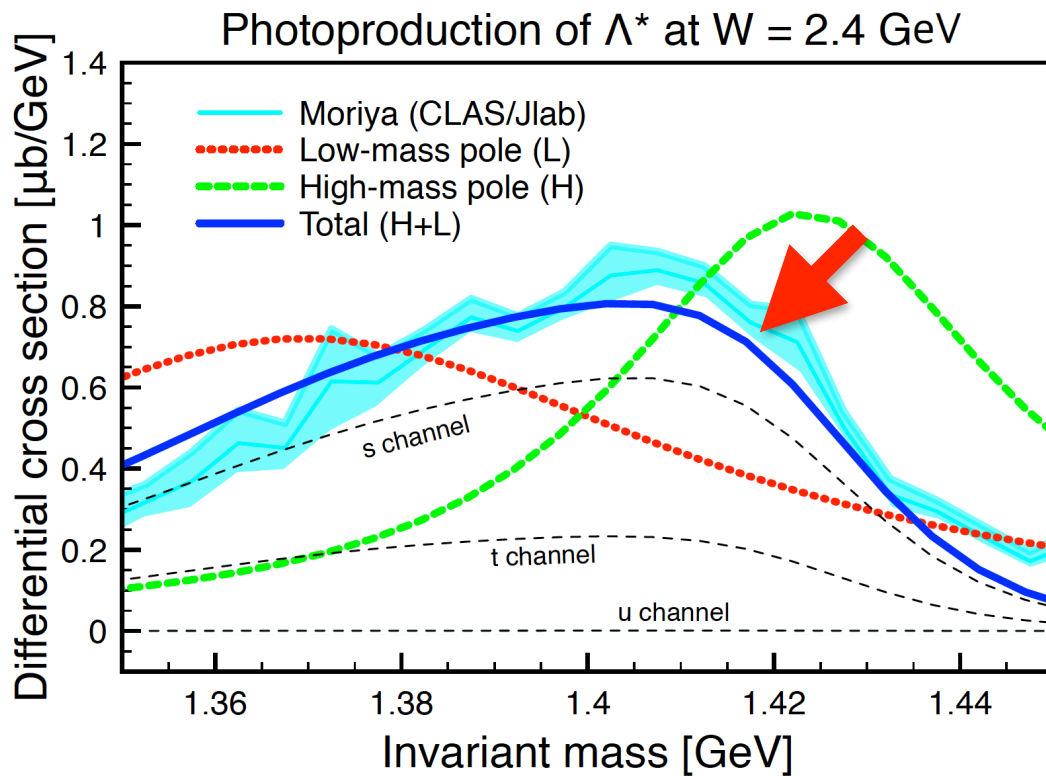
$$i\mathcal{M}_H = \frac{g_{KNH} g_{\pi\Sigma H} A_H}{M_{\pi^+\Sigma^-}^2 - M_H^2 - i\Gamma_H M_H}, \quad i\mathcal{M}_L = \frac{g_{KNL} g_{\pi\Sigma L} A_L}{M_{\pi^+\Sigma^-}^2 - M_L^2 - i\Gamma_L M_L}.$$

$$g_{KNL} g_{\pi\Sigma L} F_1^L = \underline{0.660 - 0.549i}, \quad g_{KNH} g_{\pi\Sigma H} F_1^H = \underline{-0.583 - 2.363i}.$$



# Theoretical framework

Invariant mass plots: Full calculations



**Destructive interference for photoproduction  $\Rightarrow$  Single pole**

vs.

**Constructive interference for electroproduction  $\Rightarrow$  Double pole**



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S.i.N. PRD96, 076021 (2017)

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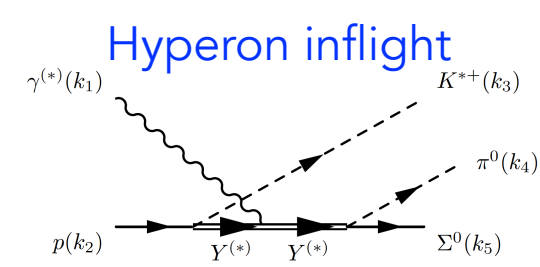
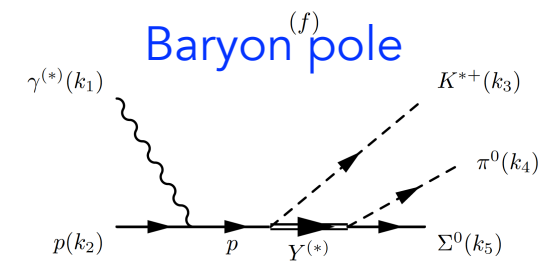
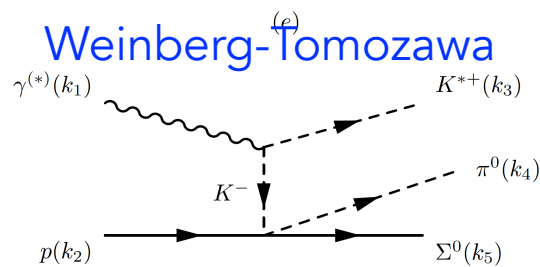
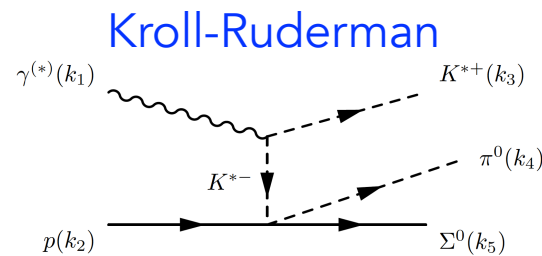
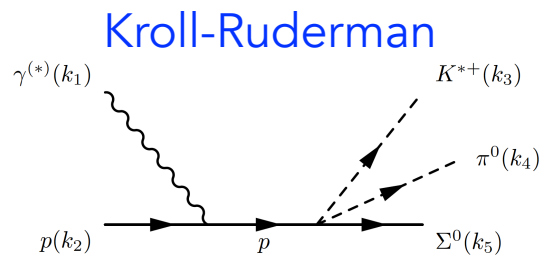
New questions and challenges

- 1) Is this observation universal for  $\Lambda(1405)$ ?
- 2) Then, what about other meson productions?
- 3) Theoretically,  $K^*N\Lambda(1405)$  coupling rarely known
- 4) It is notorious for  $\Lambda(1405)$  reaction studies
- 5) Now is the time for  $\gamma^{(*)} p \rightarrow K^{*+} \pi^0 \Sigma^0$

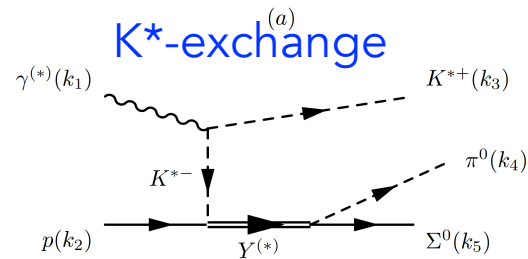
# Theoretical framework

## Effective Lagrangian approach at tree-level

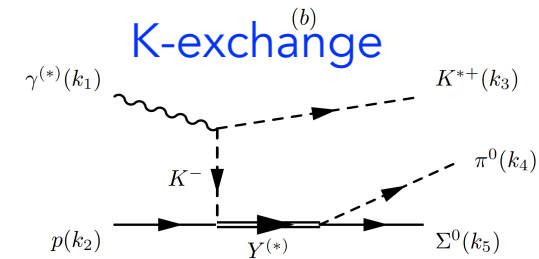
### Born approximation



(g)



(c)

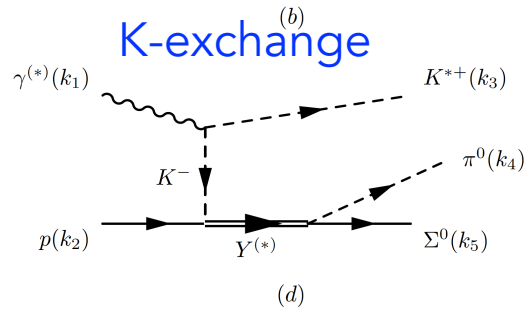
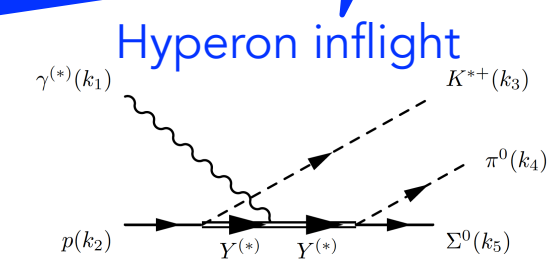
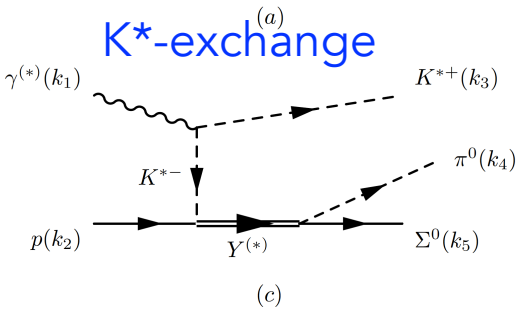
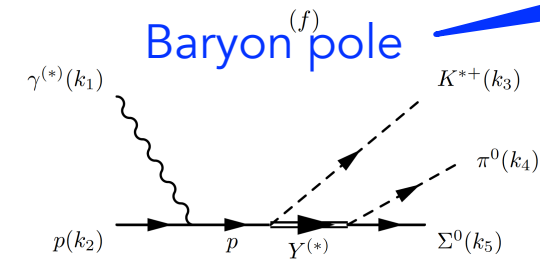
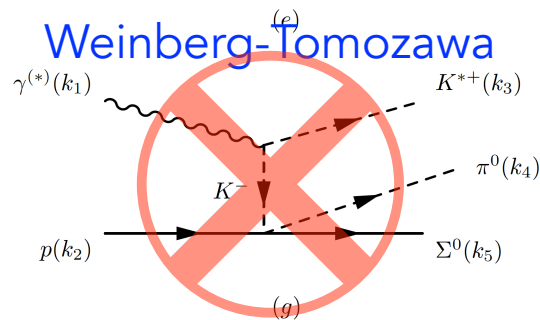
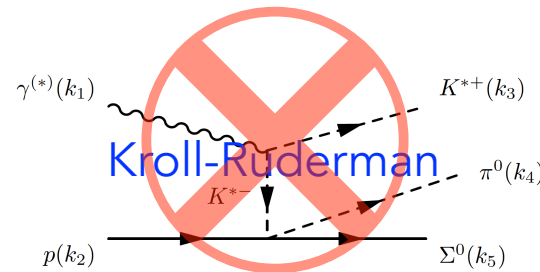
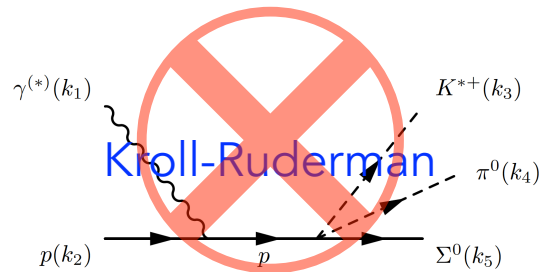


(d)

# Theoretical framework

Effective Lagrangian approach at tree-level

Born approximation

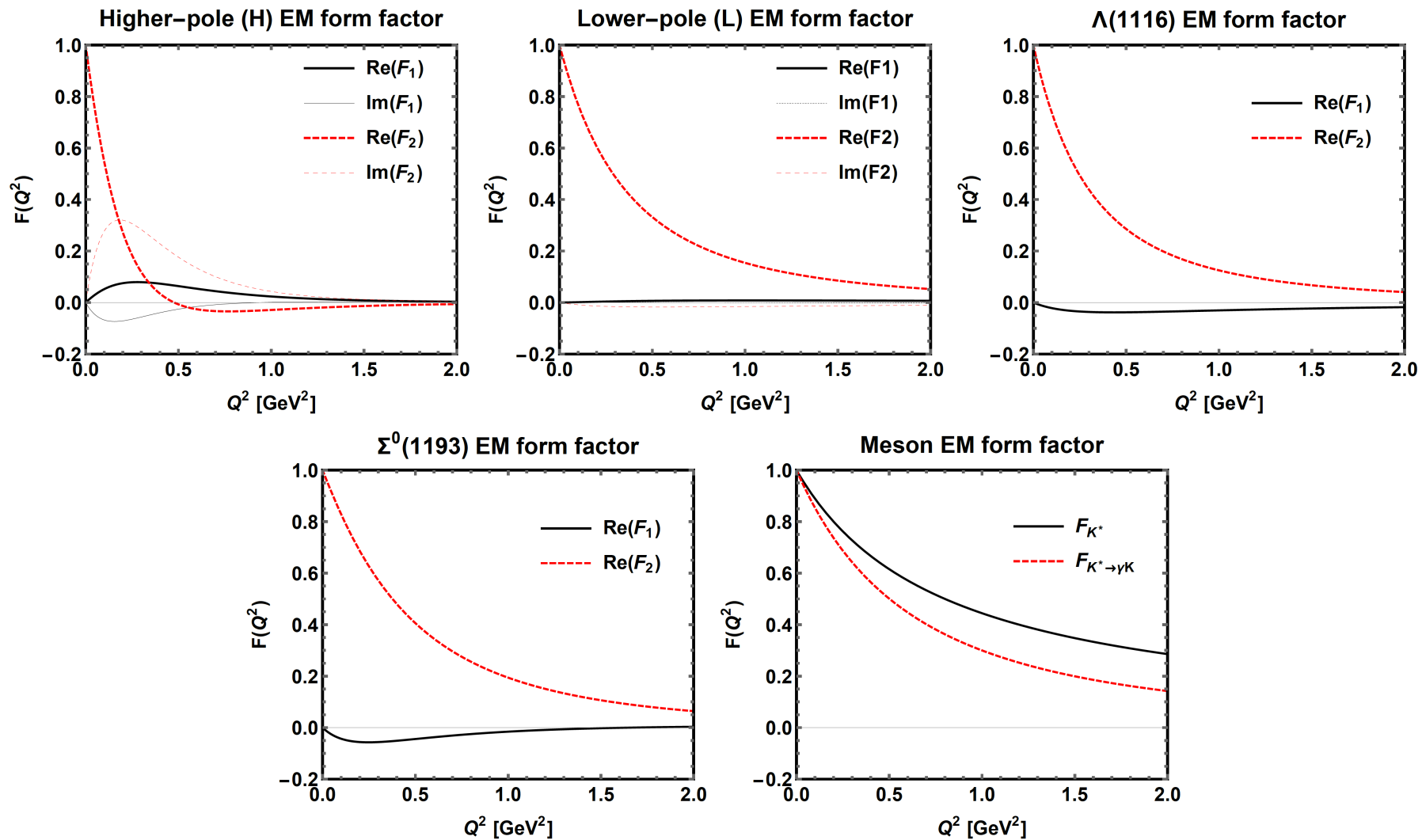


Background

# Theoretical framework

## Hyperon EM form factors from exp. and theo.

$$\frac{\langle r_E^2 \rangle_H \quad \langle r_M^2 \rangle_H \quad \langle r_E^2 \rangle_L \quad \langle r_M^2 \rangle_L}{-0.131 + 0.303i \quad 0.267 - 0.407i \quad 0.018 + 0.002i \quad -0.013 + 0.021i}$$

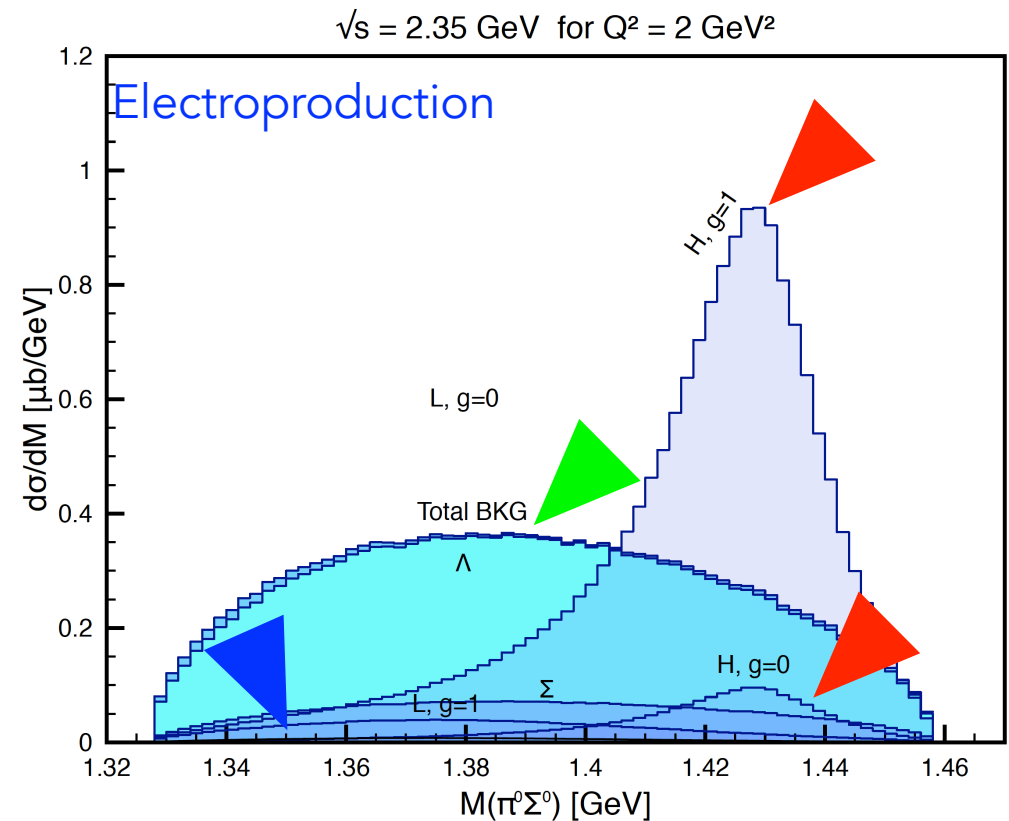
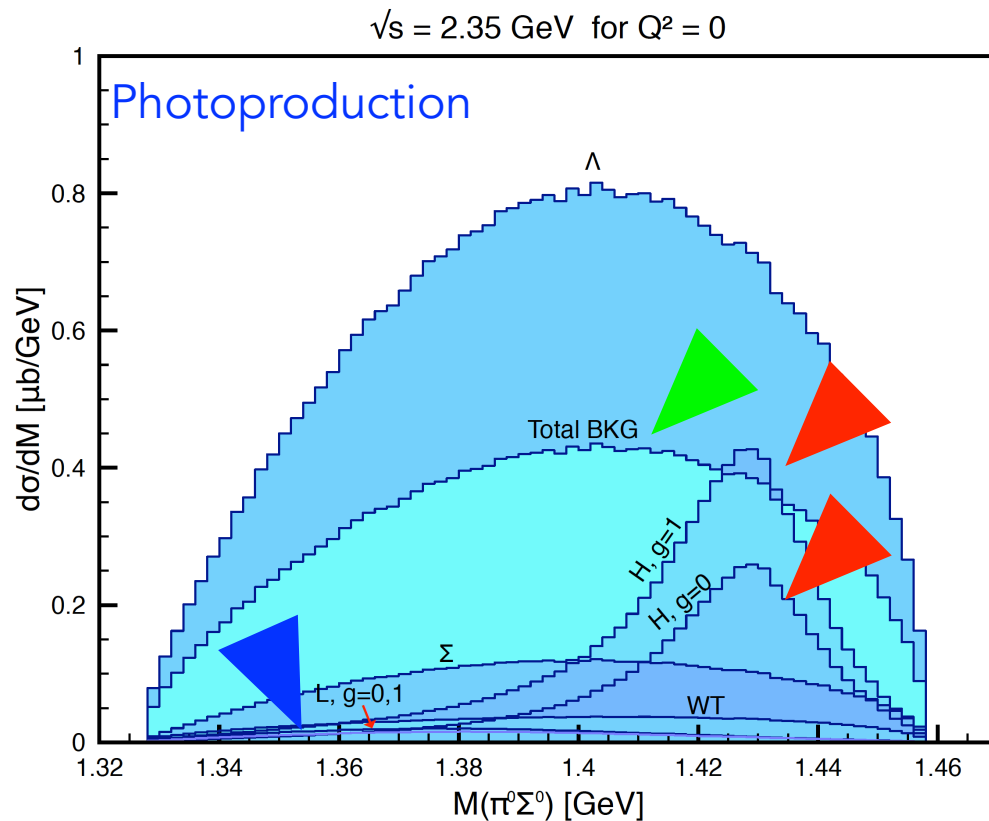


# Numerical results

Each contribution for  $Q^2=0$  and  $Q^2>0$  GeV<sup>2</sup> at  $E_{cm}=2.35$  GeV

	$g_{KN\Lambda_H}$	$g_{K^*N\Lambda_H}$	$g_{\pi\Sigma\Lambda_H}$	$g_{KN\Lambda_L}$	$g_{K^*N\Lambda_L}$	$g_{\pi\Sigma\Lambda_L}$
$g=0$	$2.4 + 1.1i$	$0.0 + i0.0$	$-0.2 - 1.4i$	$1.4 - 1.6i$	$0.0 + 0.0i$	$-2.3 + 1.4i$
$g=1$	$2.4 + 1.1i$	$0.1 - 0.9i$	$-0.2 - 1.3i$	$1.4 - 1.6i$	$-0.4 + 0.1i$	$-2.3 + 1.5i$

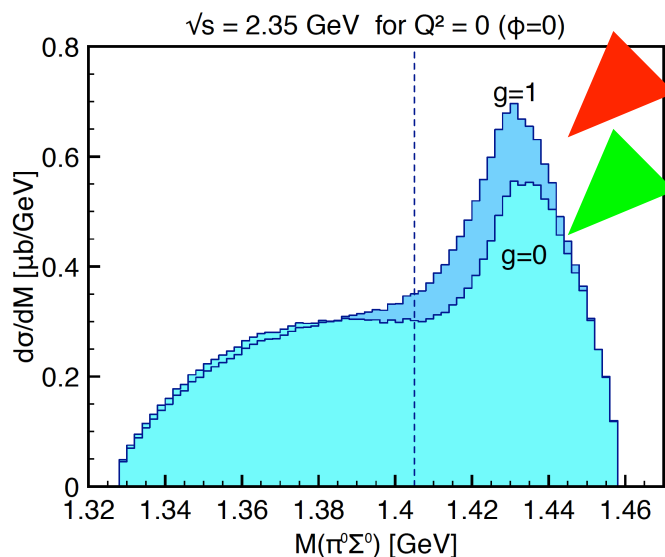
$M_H$	$\Gamma_H$	$M_L$	$\Gamma_L$
1430 MeV	30 MeV	1376 MeV	126 MeV



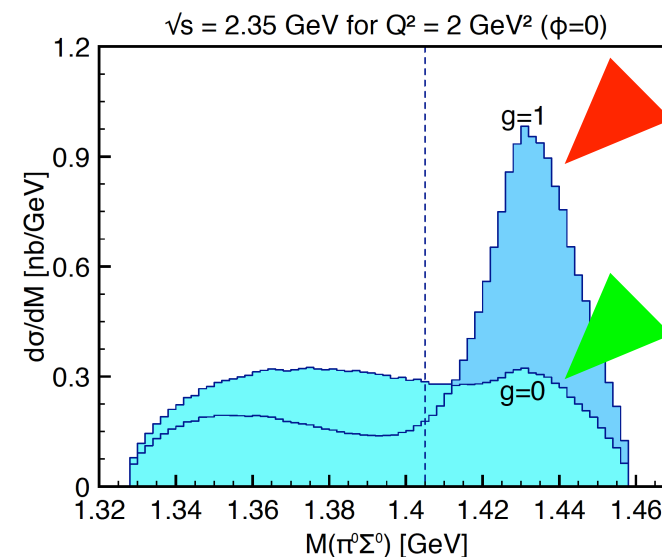
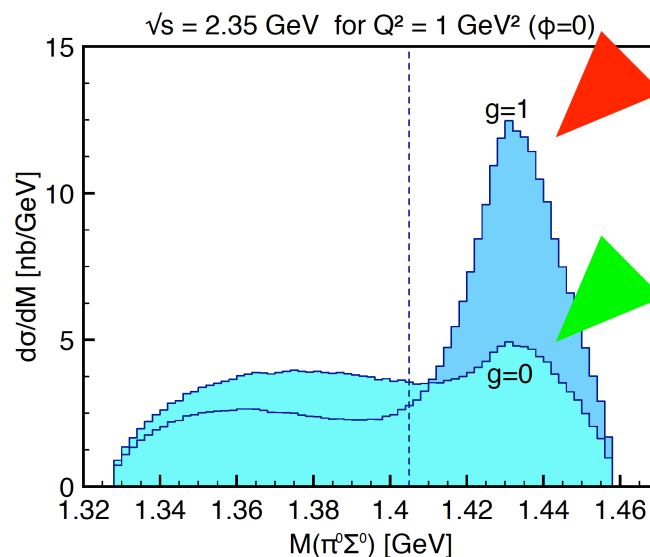
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Photoproduction



Electroproduction



With  $K^*$  exchange

vs.

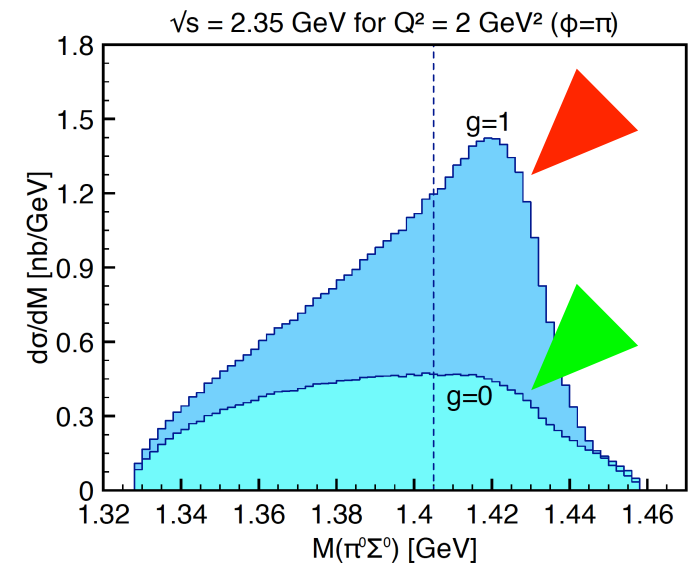
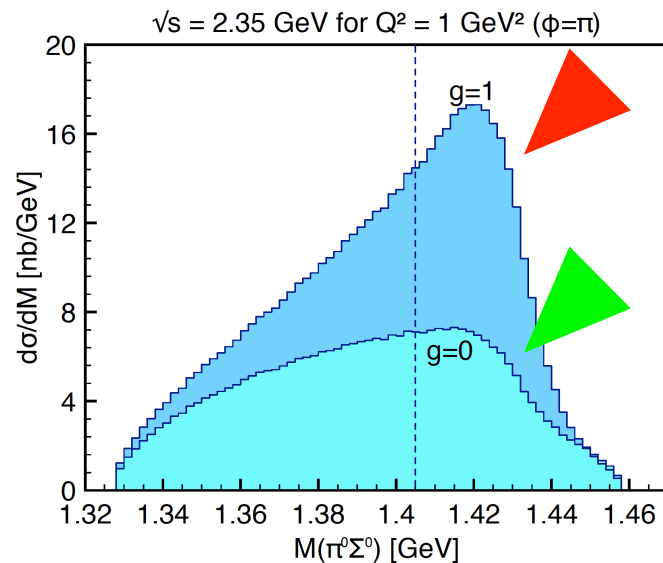
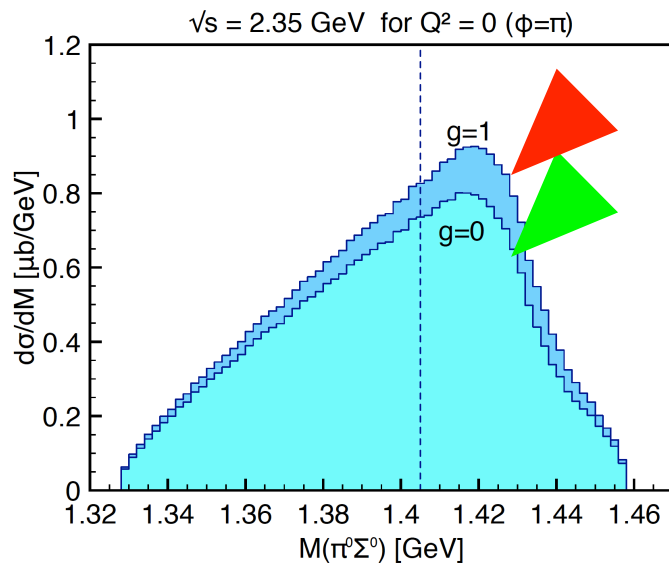
Without  $K^*$  exchange

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Photoproduction

Electroproduction



With  $K^*$  exchange

vs.

Without  $K^*$  exchange



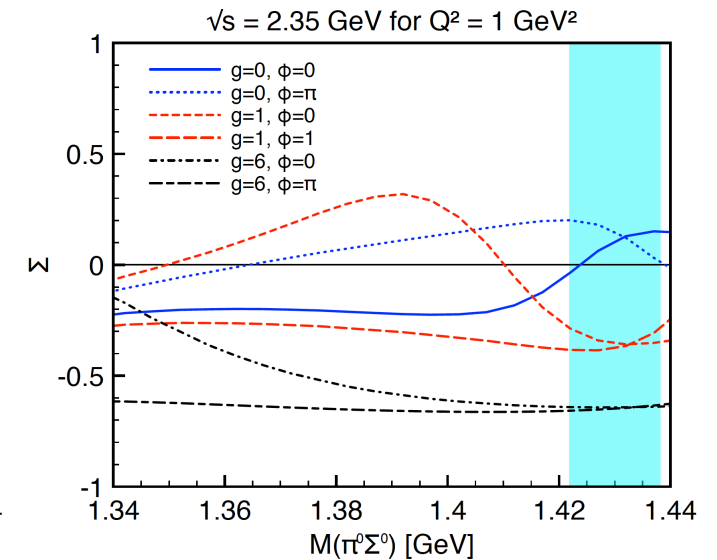
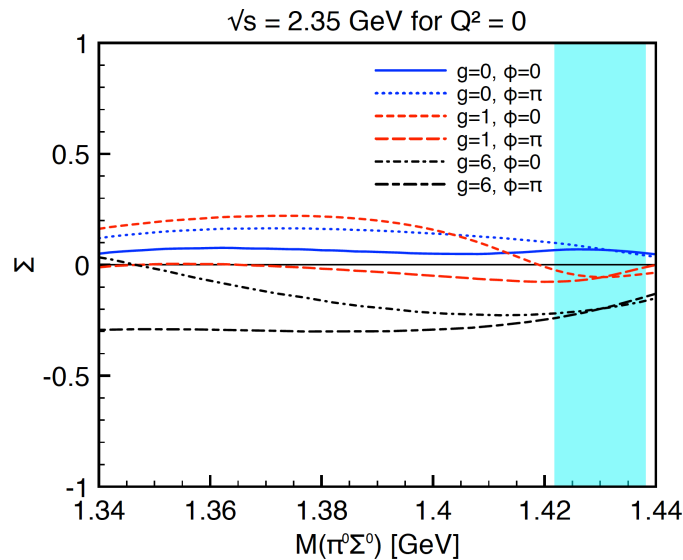
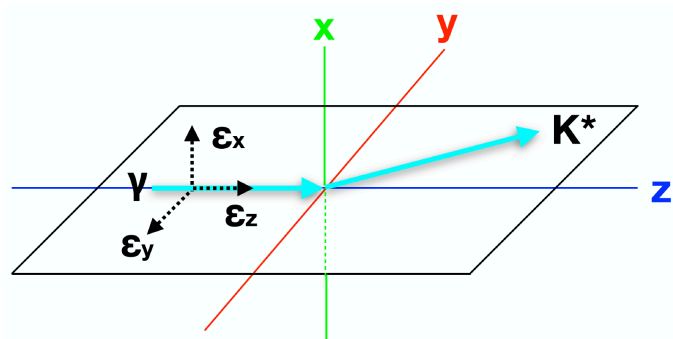
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$$\Sigma(M_I) = \frac{d\sigma_x/dM_I - d\sigma_y/dM_I - d\sigma_z/dM_I}{d\sigma_x/dM_I + d\sigma_y/dM_I + d\sigma_z/dM_I},$$

Photoproduction

Electroproduction



## Summary, conclusion, and perspectives

$K^*$  photo and electroproductions are studied

Higher-pole and BKG dominate:

No single or double bump observed in  $K$  productions

If  $K^*N \Lambda^*$  coupling is negligible, no peak observed in electroproduction, whereas photoproduction shows it at  $M_{INV} \sim 1430$  MeV

If  $K^*N \Lambda^*$  coupling is sizable, similar invariant mass shapes observed in both at  $M_{INV} \sim 1430$  MeV

## Summary, conclusion, and perspectives

These observations based on ChUM, indicating that

Two-pole structure of  $\Lambda^*$

Furthermore, we can extract information of  $K^*N$   $\Lambda^*$  coupling

How to pin down two-pole structure of  $\Lambda^*$  from  $K^*$  photoproduction?: LEPS2 experiment

Polarization will help!!!! (works in progress)

What are we interested in?

$\Lambda(1405)$  at LEPS/SPring-8 and CLAS/JLab

$$\gamma^{(*)} p \rightarrow K^+ \pi^+ \Sigma^-$$

S.i.N. PRD96, 076021 (2017)

$$\gamma^{(*)} p \rightarrow K^{*+} \pi^0 \Sigma^0$$

S.i.N., A.Hosaka arXiv:1902.09106 [hep-ph]

$\Xi(1690)$  at Hadron hall/J-PARC

$$K^- p \rightarrow K^+ K^- \Lambda$$

J.K.Ahn, S.i.N. PRD98, 114012 (2018)

S=0 and S=-1 resonances at Belle/KEK (Dr. S.B.Yang)

$$\Lambda_c^+ \rightarrow \pi \bar{K} p$$

J.K.Ahn, S.B.Yang, K.S.Choi, S.i.N. In preparation

# Motivation

$S=-2$   $\Xi$  hyperon spectra have not been established

Theoretical contradictions such as  $\Xi(1620)$ ,  $\Xi(1690)$

BaBar experiment assigned  $\Xi(1690, 1/2^-)$  and  $\Xi(1820, 3/2^-)$  almost confirmed:  $\Xi(1620, ??)$

$N(1535, 1/2^-) \sim N(1530, 3/2^-)$ : degenerated spin partner

$\Lambda(1405, 1/2^-) < \Lambda(1520, 3/2^-)$ :  $\sim 115$  MeV difference

$\Xi(1690, 1/2^-) < \Xi(1820, 3/2^-)$ :  $\sim 130$  MeV difference

$\Xi(1690, 1/2^-)$  can be molecular?

# Motivation

$\Xi(1690, 1/2^-)$  from  $\Lambda_c \rightarrow \Lambda K_s^0 K^+$  in Belle/Belle-II, LHCb

Interference  $\Xi(1690, 1/2^-)$  and  $a_0(980)$  unavoidable

$K^- p \rightarrow K^+ K^- \Lambda$  can control interference by  $E_{cm}$

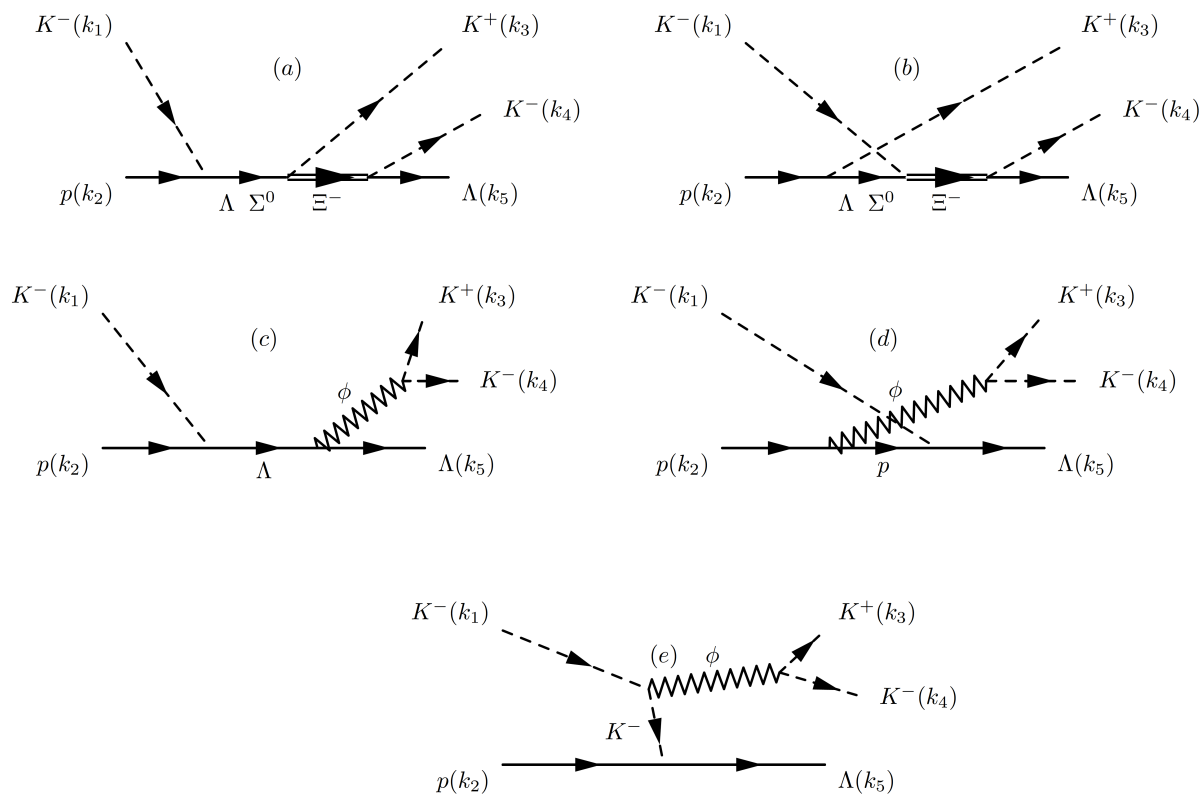
Narrow  $\varphi(1020)$  can be isolated easily

Old  $K^- p \rightarrow K^+ K^- \Lambda$  BC data focusing on  $\Xi(1820)$

Experimental and theoretical guides for  $\Xi(1690, 1/2^-)$   
in  $K^- p \rightarrow K^+ K^- \Lambda$  at J-PARC

# Theoretical framework

Effective Lagrangian approach at tree-level Born approx.



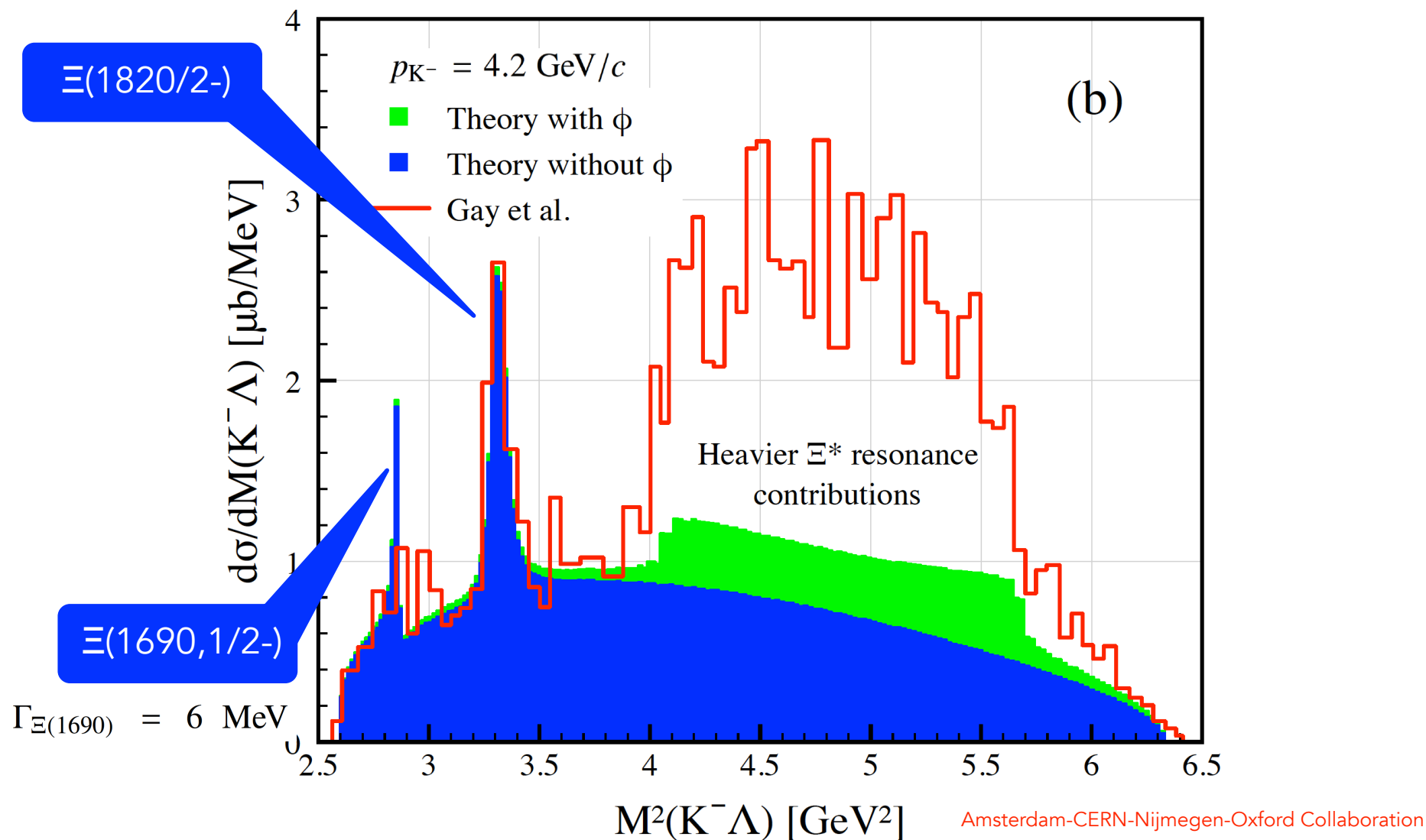
$$F_{\Xi, \phi}(x) = \frac{\Lambda_{\Xi, \phi}^4}{\Lambda_{\Xi, \phi}^4 + (x - M_x^2)^2}, \quad \text{for } x = (s, u, t),$$

$$i\mathcal{M}_{\text{total}} = ie^{i\pi/2} \mathcal{M}_{\Xi_{3/2}} + ie^{i3\pi/2} \mathcal{M}_{\Xi_{1/2}} + i\mathcal{M}_{\phi}$$

$(M, J^P)$	$\Lambda(1116, 1/2^+)$	$\Lambda(1405, 1/2^-)$	$\Lambda(1520, 3/2^-)$	$\Lambda(1670, 1/2^-)$	$\Sigma(1193, 1/2^+)$	$\Sigma(1385, 3/2^+)$
$N(938, 1/2^+)$	-13.24 [22]	0.91 [21]	-10.90 [1]	0.30 [20]	3.58 [22]	-3.22 [21]
$\Xi(1322, 1/2^+)$	3.52 [22]	0.91 [21]	3.27 [21]	-0.18 [20]	-13.26 [22]	-3.22 [21]
$\Xi(1532, 3/2^+)$	4.08	-	-	-	3.22	-
$\Xi(1690, 1/2^-)$	-0.3 [8]	-	-	-	1.8 [8]	-
$\Xi(1820, 3/2^-)$	6.10 [5]	-	-	-	8.00 [5]	-

# Numerical results

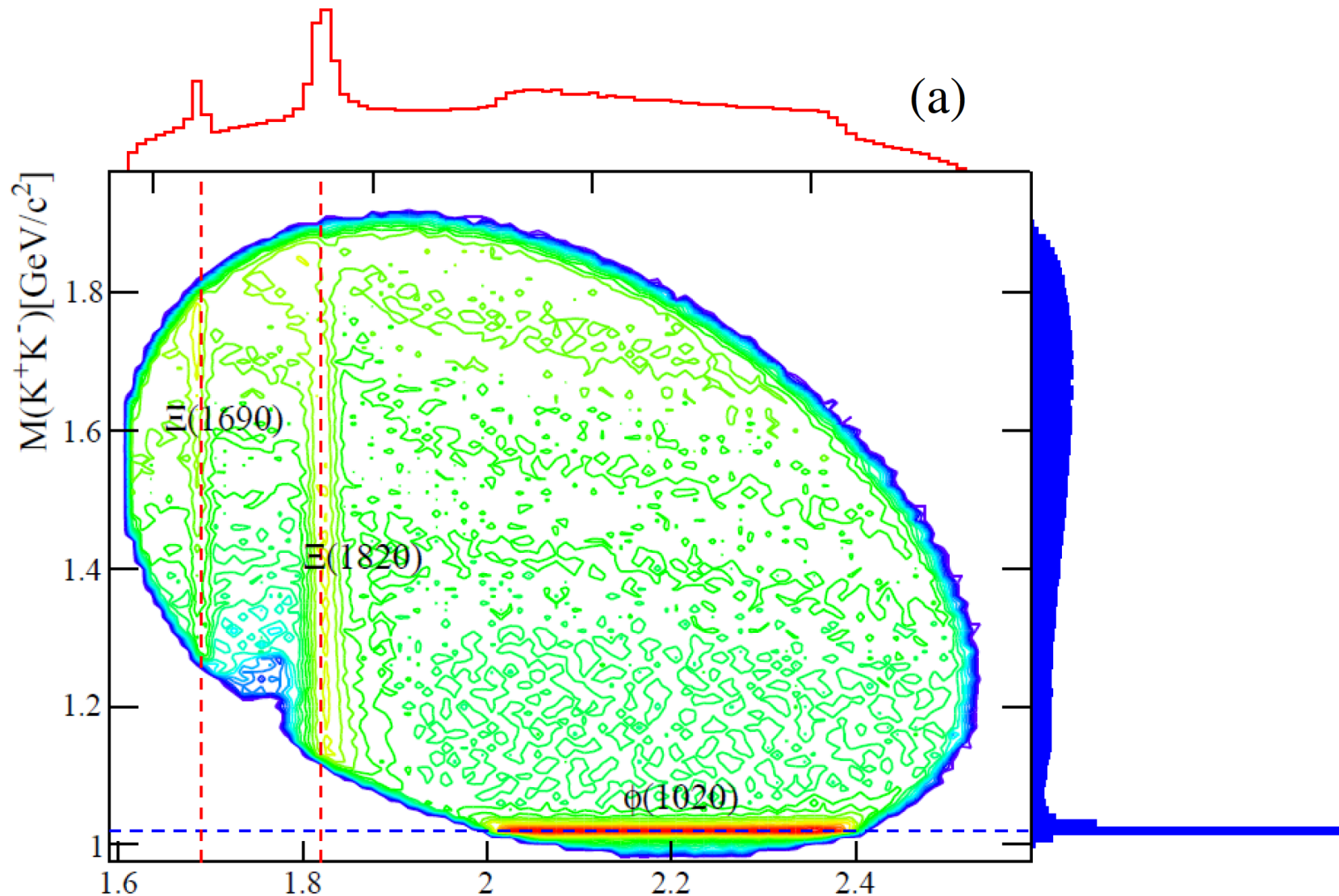
Invariant mass plot: Significant  $\phi$  contribution





# Numerical results

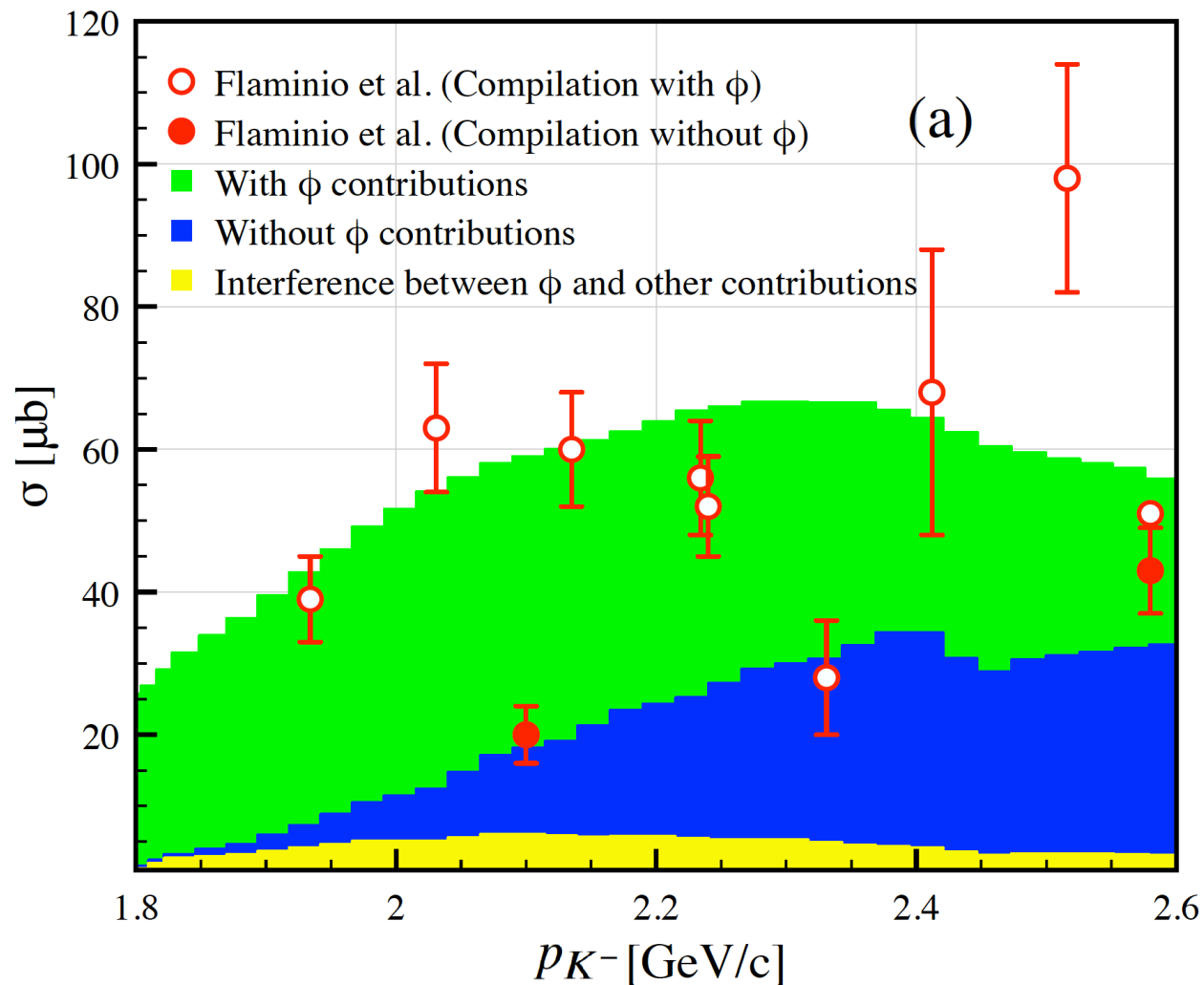
Dalitz plot: Negligible interference (merits of production)



# Numerical results

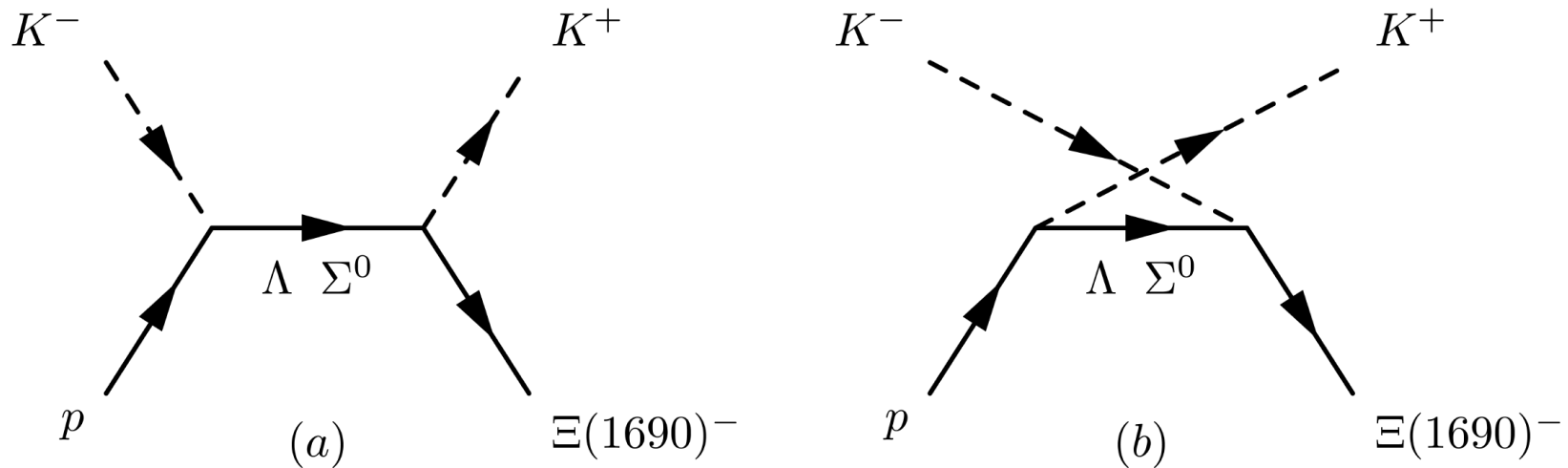
All the model parameters determined then...

Qualitative agreement with other data!!!



# Numerical results

The model parameters applied to  $K^- p \rightarrow K^+ \Xi(1690)^-$

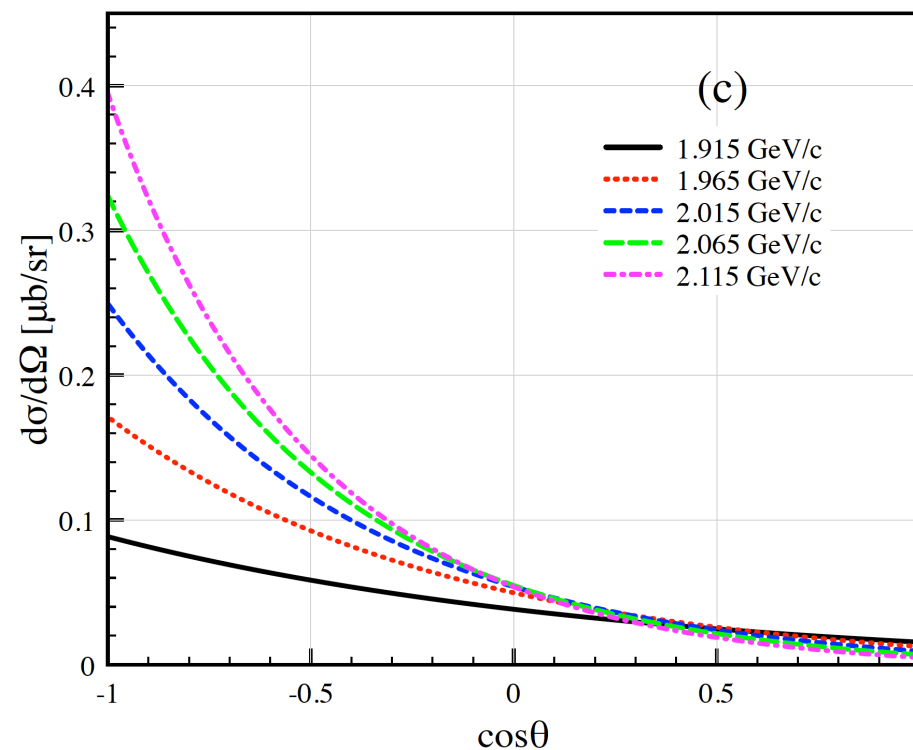
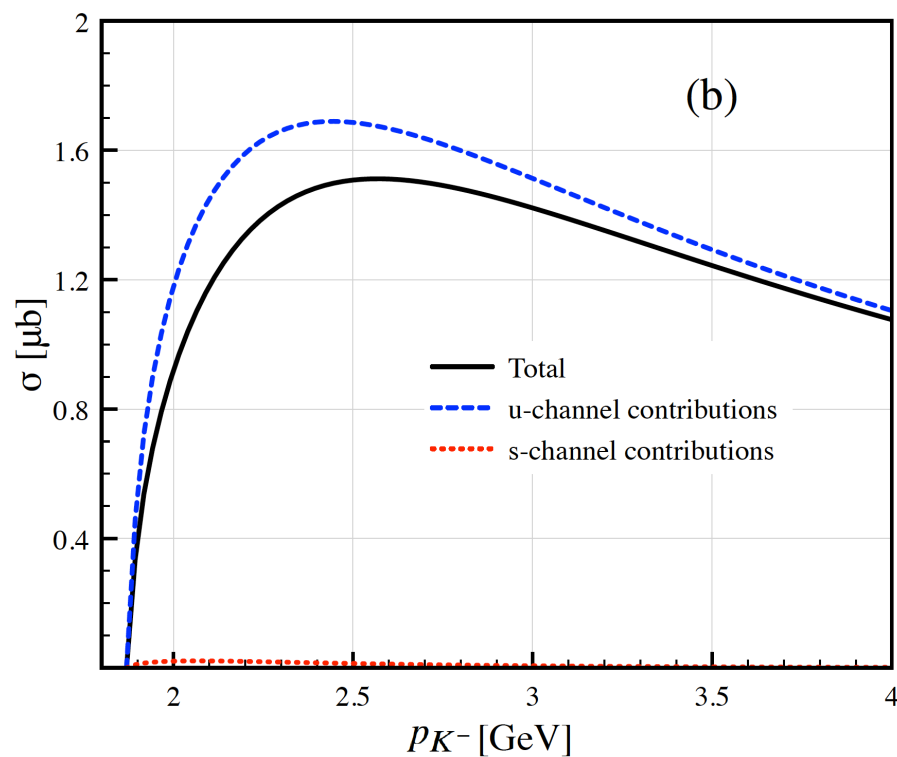


No  $S=-2$  meson makes problem much easier

No hyperon Regge: Is it necessary?!?!

# Numerical results

The model parameters applied to  $K^- p \rightarrow K^+ \Xi^-(1690)$



u-channel dominates  $\sim 1 \mu\text{b}$

Measurable at J-PARC:  $p_{K^-} = 1.878 \text{ GeV}$  at threshold

## Numerical results

Strong backward scattering of  $K^+$  in cm frame due to u-channel dominance

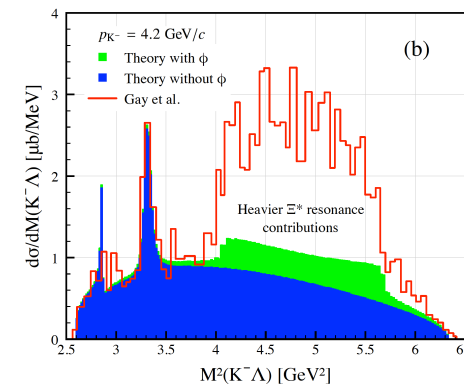
Experimentally  $\Xi^-(1690)$  decays into  $\Sigma^- K^0$ ,  $\Sigma^0 K^-$ ,  $\Xi^- \pi^0$ ,  $\Xi^0 \pi^-$  : Only one neutral decay particle

Easy to reconstruct mass of  $\Xi^-(1690)$

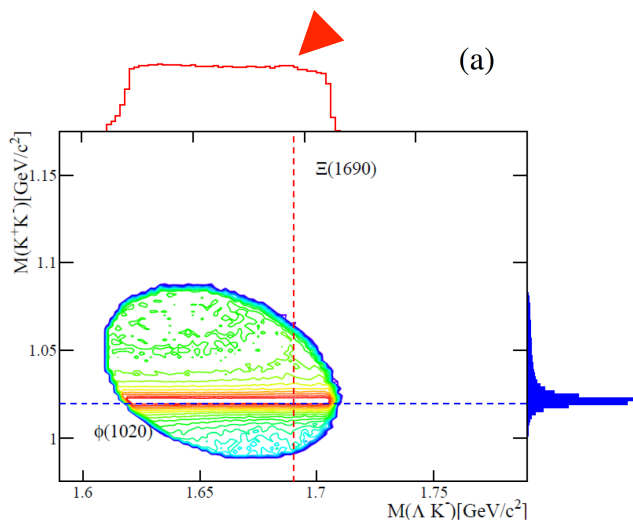
For future experiment at high-intensity  $K^-$  beam at J-PARC, theory calculations for Dalitz plots are employed to simulate experiment via event generators

# Numerical results

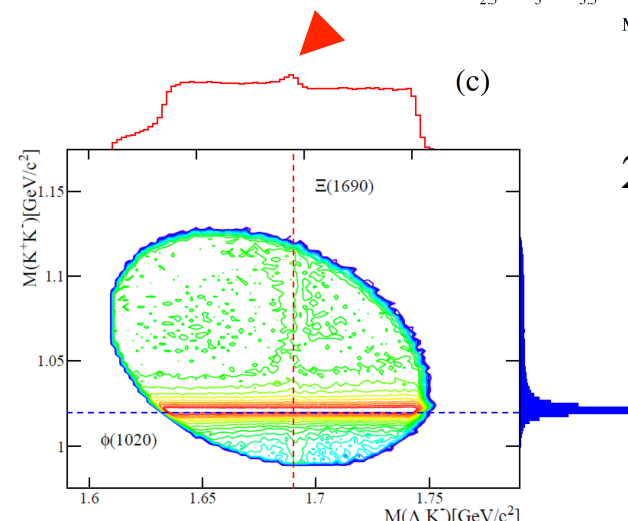
## Simulated Dalitz plots based on theory



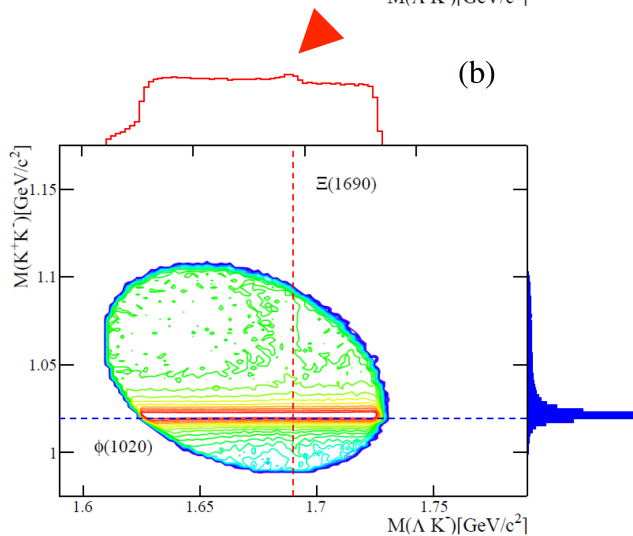
1.915 GeV



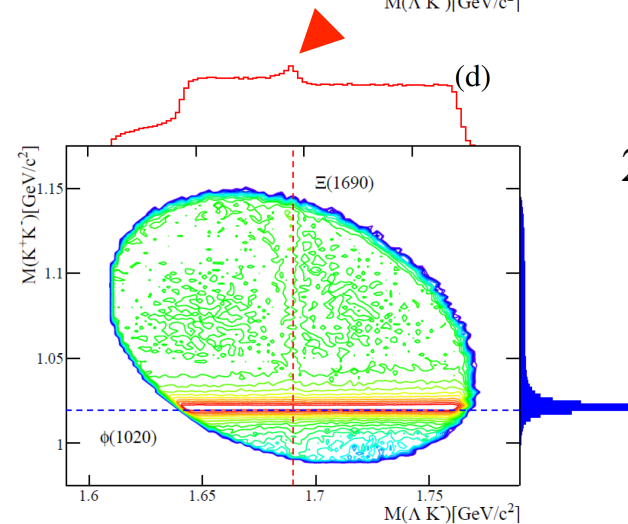
2.015 GeV



1.965 GeV

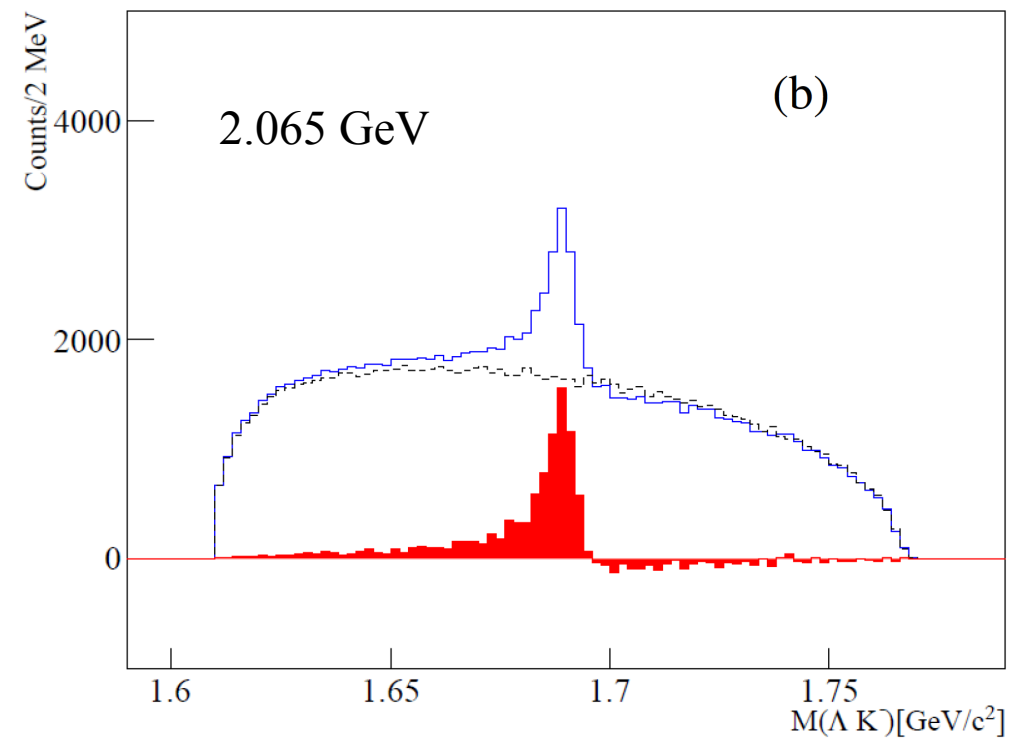
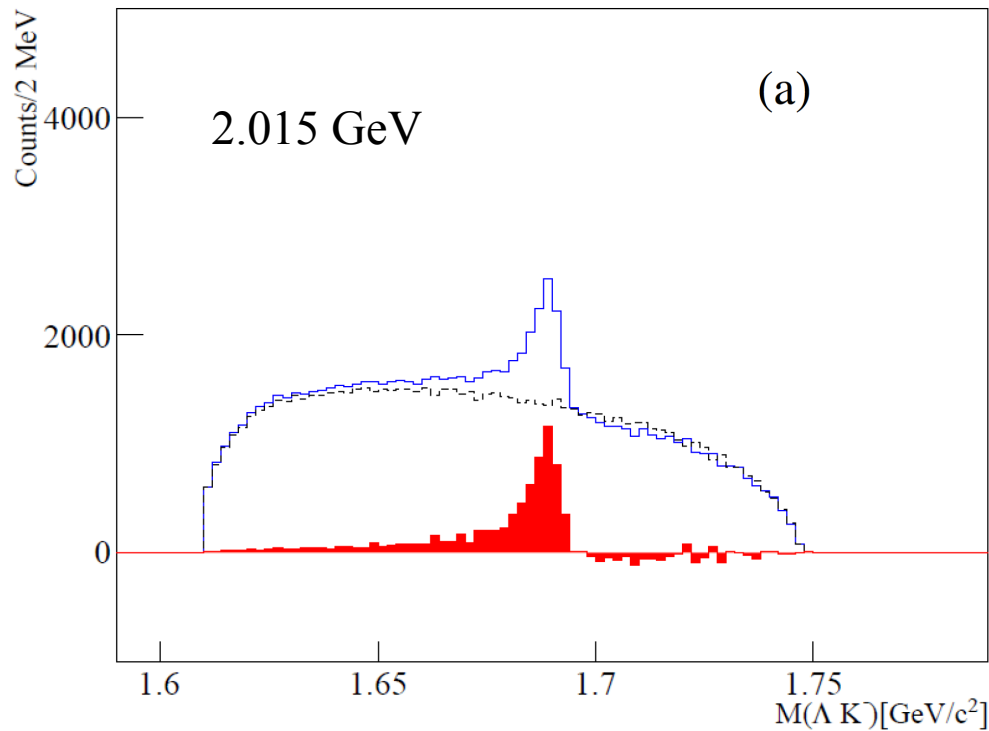
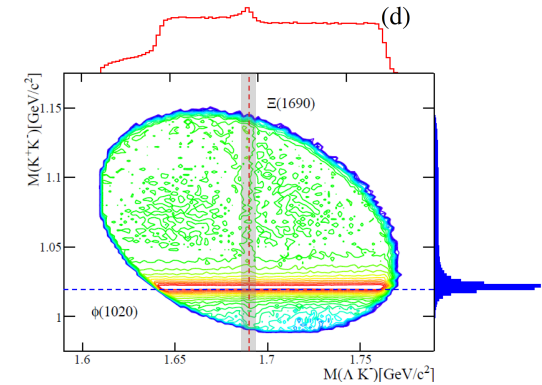


2.065 GeV



# Numerical results

## Simulated Dalitz plots excluding $\phi(1020)$

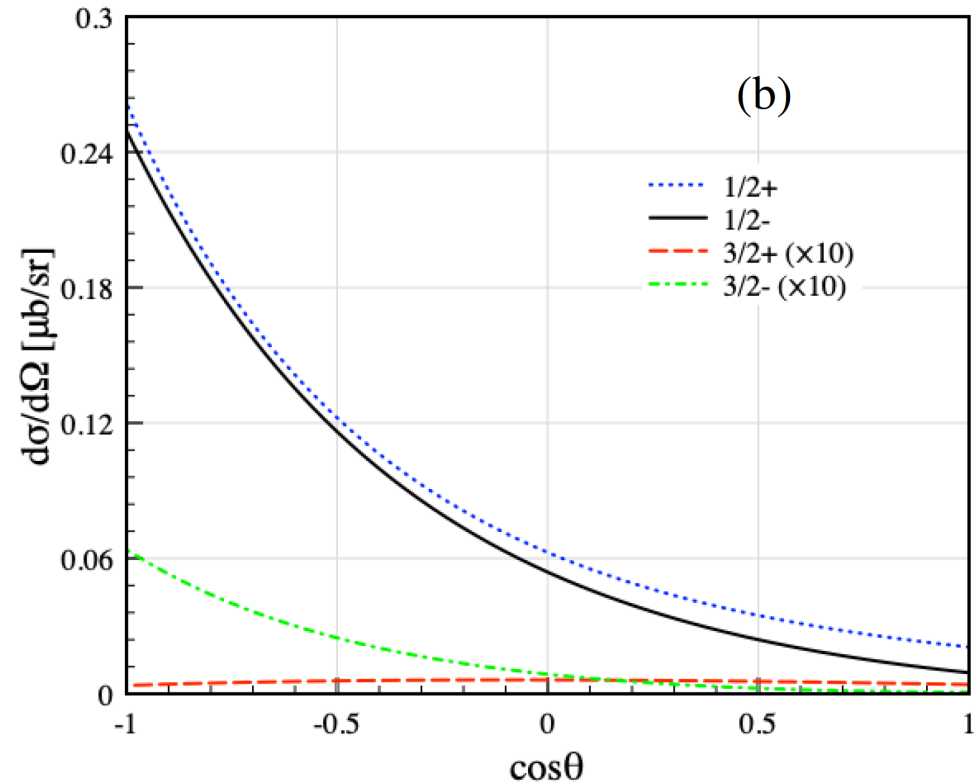
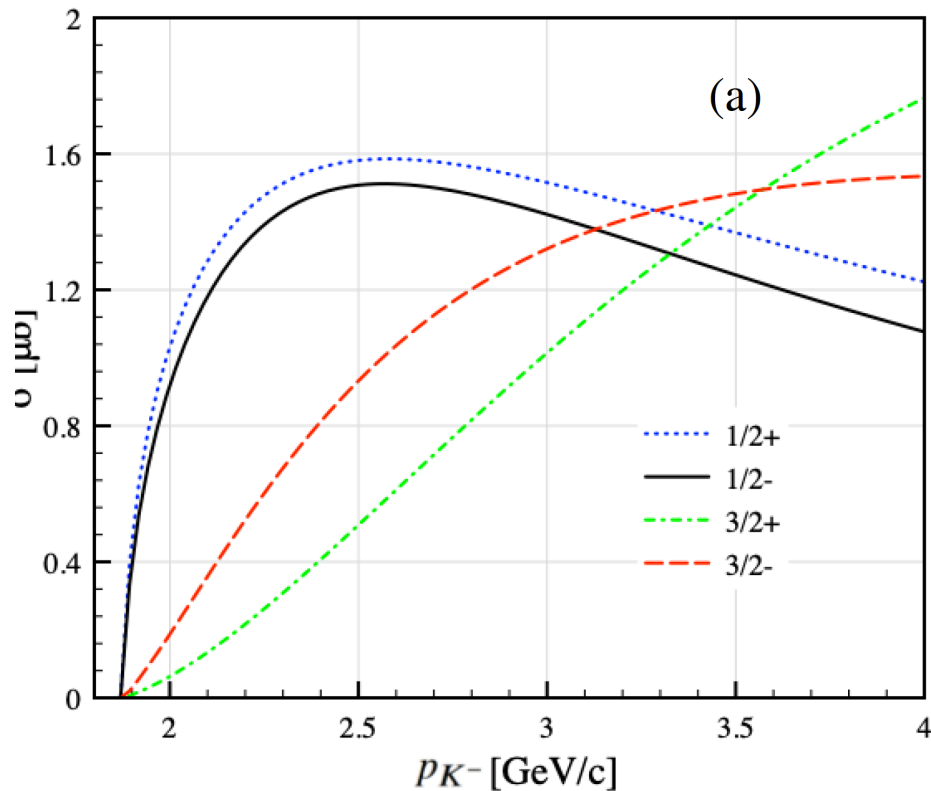


Clear peak with  $\Xi(1690)$  and interference

# Numerical results

## Different spin-parity states of $\Xi(1690)$

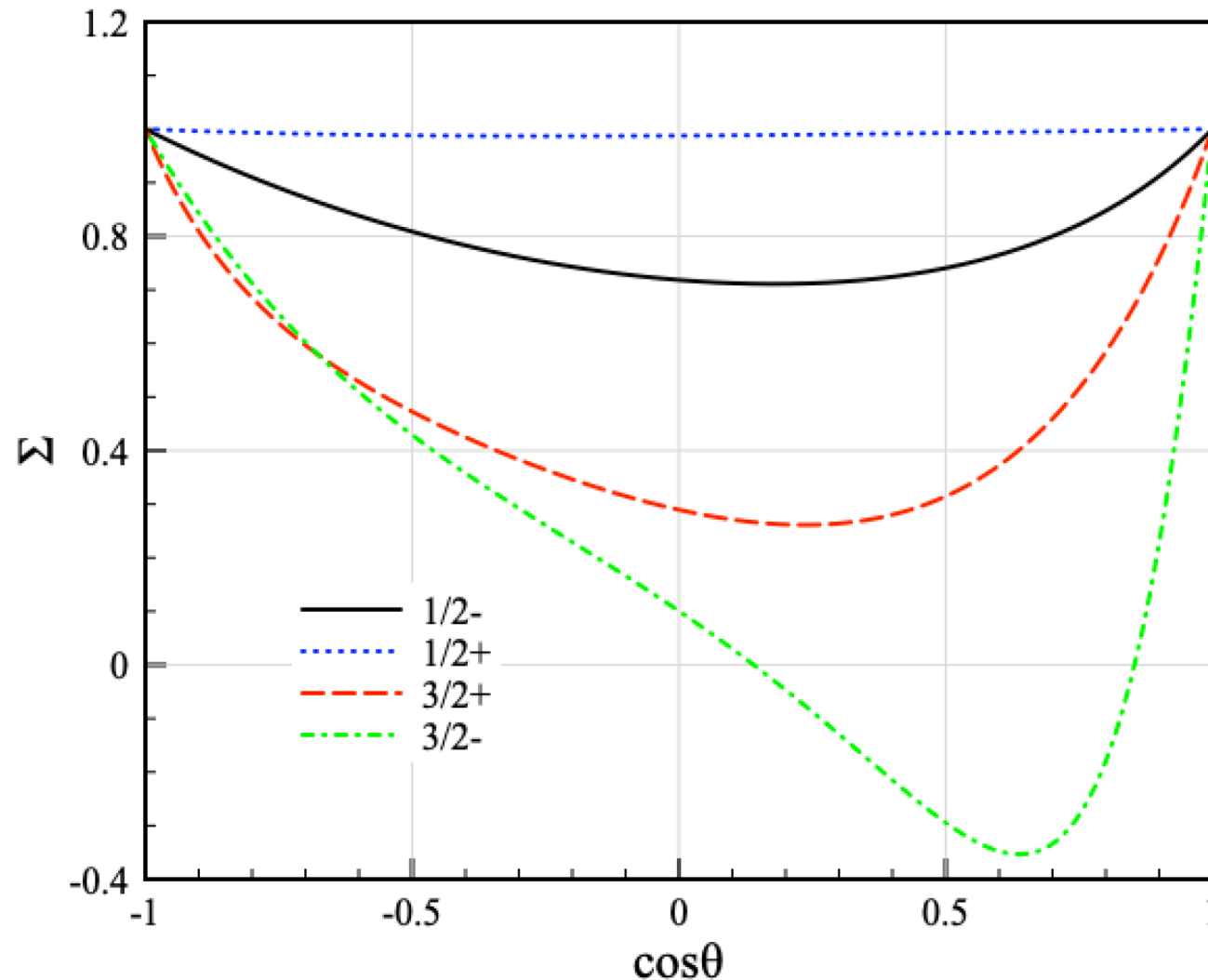
$$\frac{\Gamma_{\Xi(1690) \rightarrow \bar{K}\Lambda}}{\Gamma_{\Xi(1690)}} = 0.271, \quad \frac{\Gamma_{\Xi(1690) \rightarrow \bar{K}\Sigma}}{\Gamma_{\Xi(1690)}} = 0.533.$$





# Numerical results

Double polarization asymmetry  $\Sigma(s_R) = \frac{d\sigma_{\uparrow}/d\Omega - d\sigma_{\downarrow}/d\Omega}{d\sigma_{\uparrow}/d\Omega + d\sigma_{\downarrow}/d\Omega}$ .



## Summary, conclusion, and perspectives

$\Xi^-(1690)$  production with K- beam studied  
Effective Lagrangian method at tree-level Born approx.  
All the parameters determined by Dalitz process

$K^- p \rightarrow K^+ \Xi^-(1690)$  shows  $\sigma \sim 1 \mu\text{b}$  and backward scattering enhancement due to u-channel contribution

Simulated Dalitz and invariant mass plot via event generators based on theory: Clean peak for  $\Xi^-(1690)$  after  $\varphi$  exclusion

Reliable guides for future J-PARC experiments

# Thank you for your attention!

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