

# New era of dark matter research and direct detection experiments

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# What is Dark Matter?

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What **particle** is dark matter?

- Mass?
- (Non-gravitational) Interactions?

DM - SM

DM - DM

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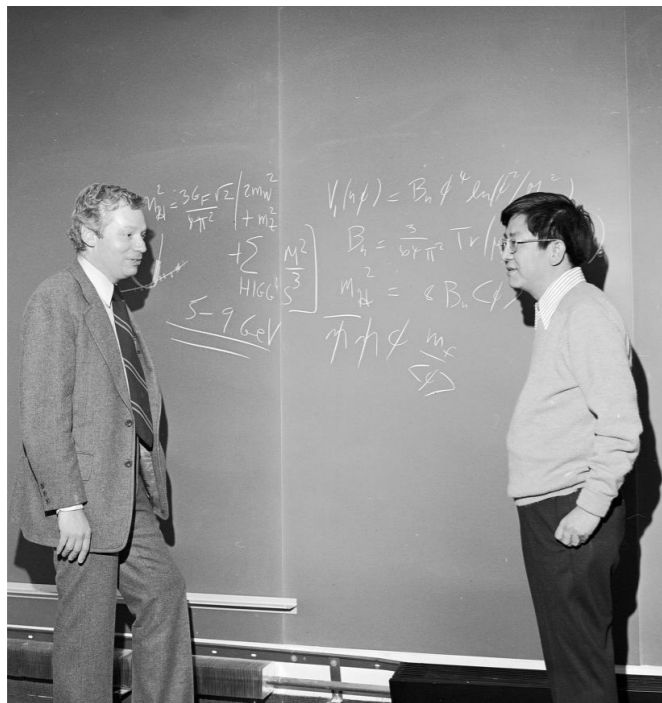
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Preferred candidate so far was

## Weakly Interacting Massive Particle (WIMP)



- Weak scale mass:  $O(1 \sim 100) \times$  proton mass
- Weak interaction with the SM particles:  
about  $< 10^{-12}$  (in cross section) smaller than EM

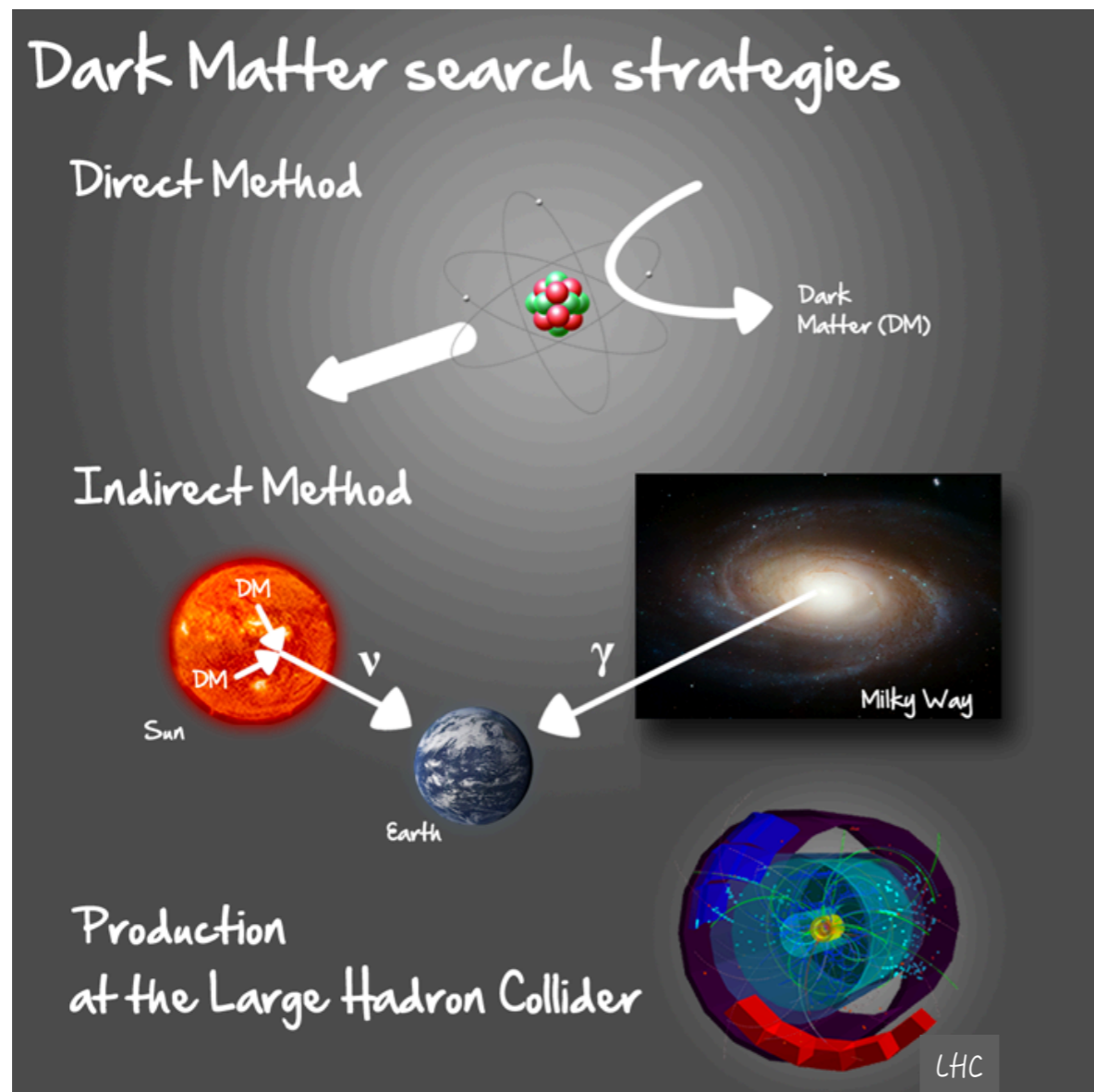
Byproduct of many BSM theories  
for resolving the hierarchy problem

# What is Dark Matter?

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- (Non-gravitational) Interactions?

DM - SM  $\rightarrow$  i) Observation  
ii) Amount of DM  
DM - DM

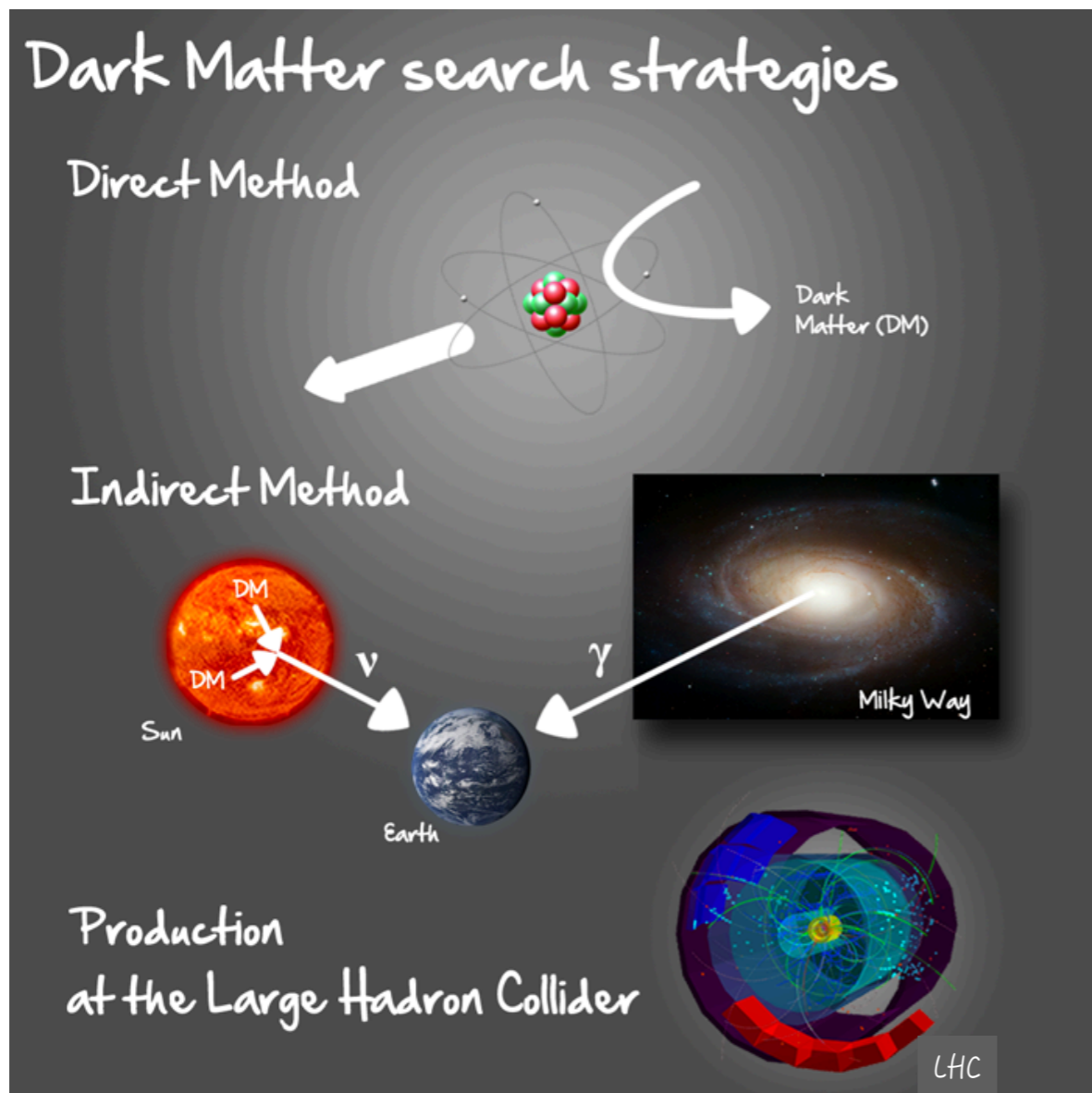


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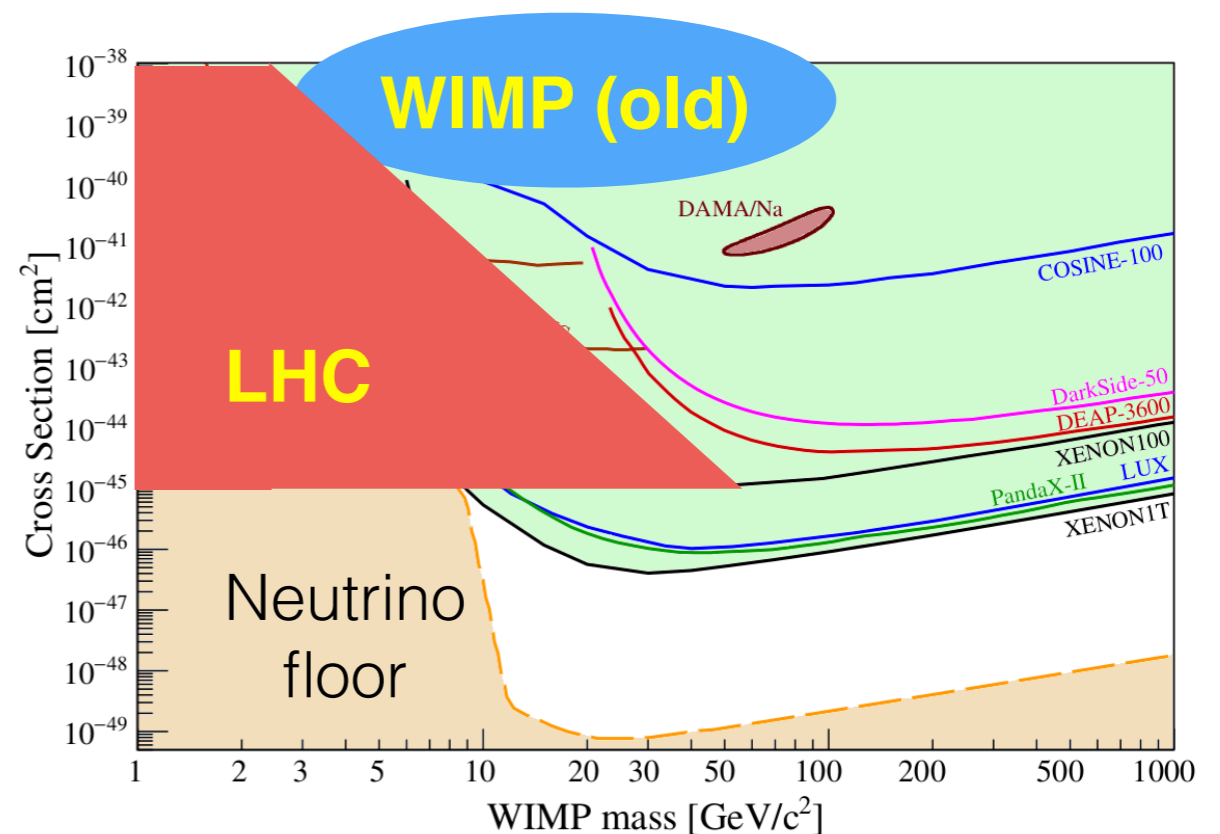
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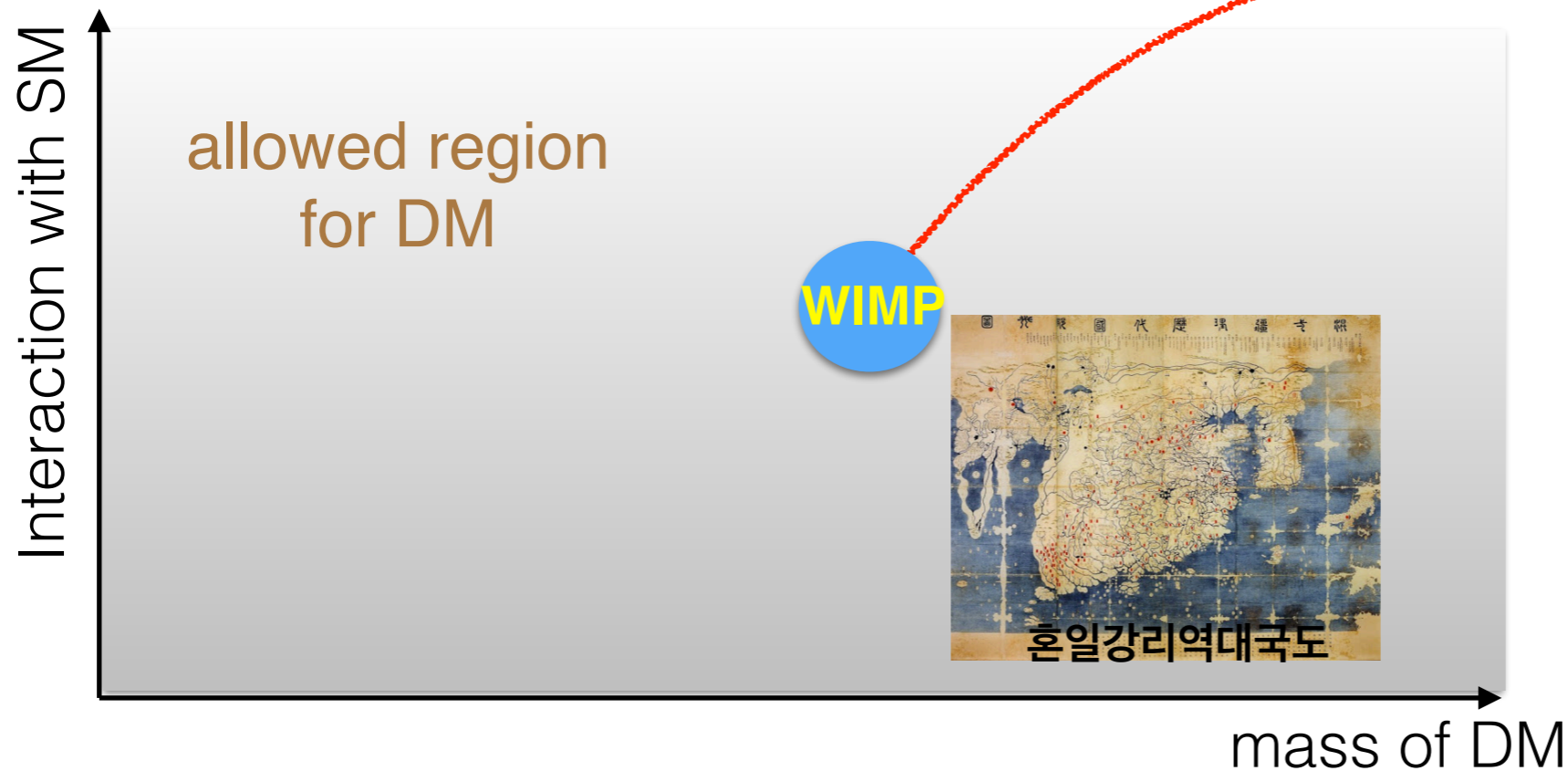


WIMP strongly constrained!

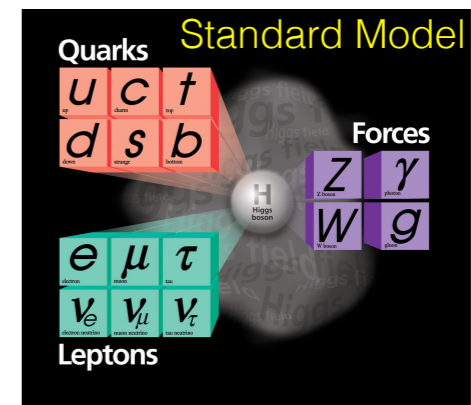


# Dark World beyond WIMP

WIMP may be a theoretical bias.

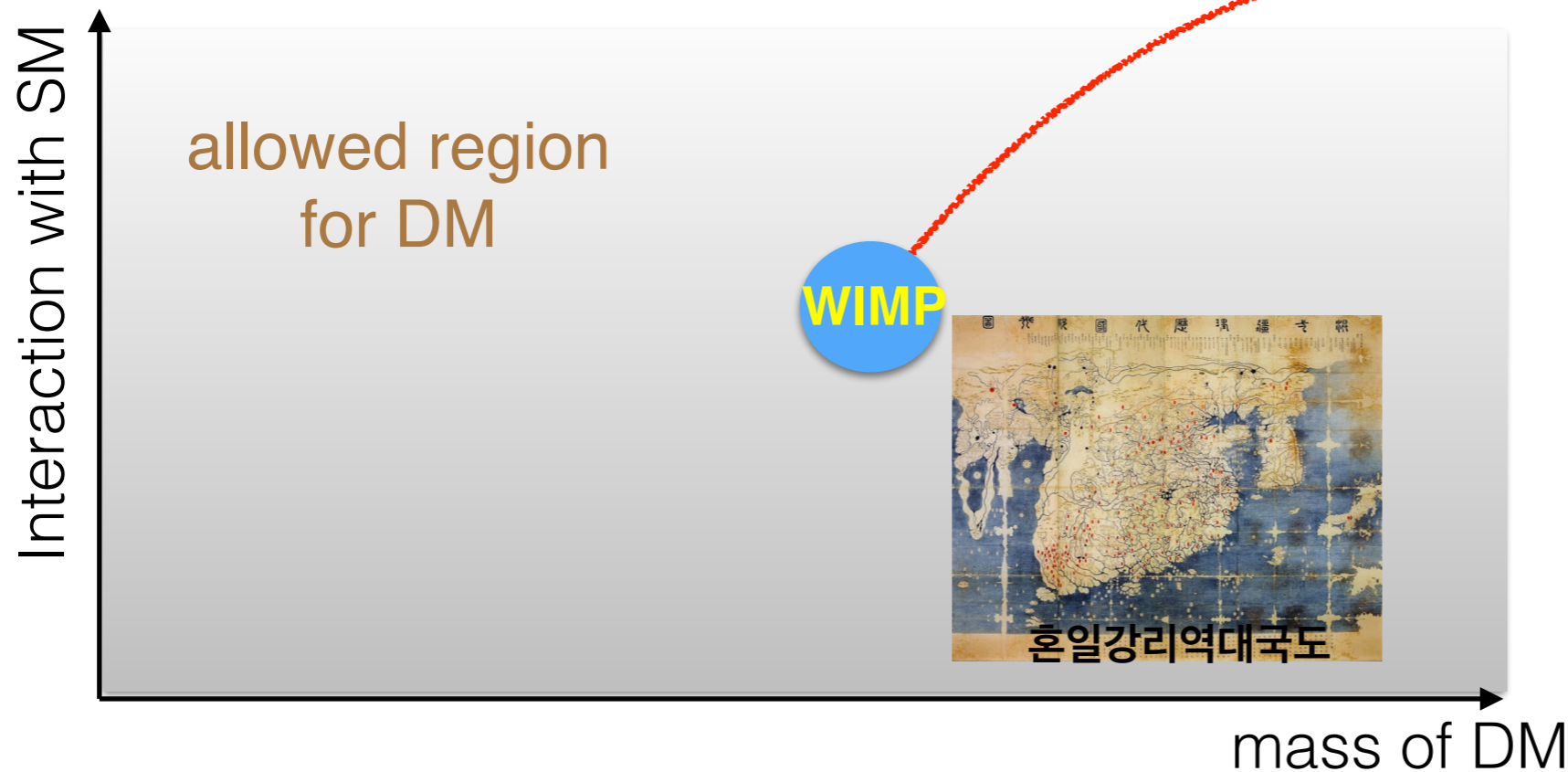


- Small region
- Oversimplification compared to

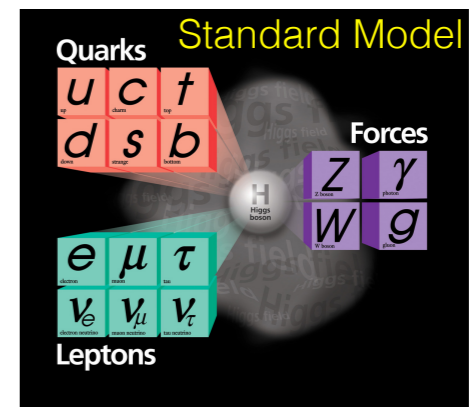


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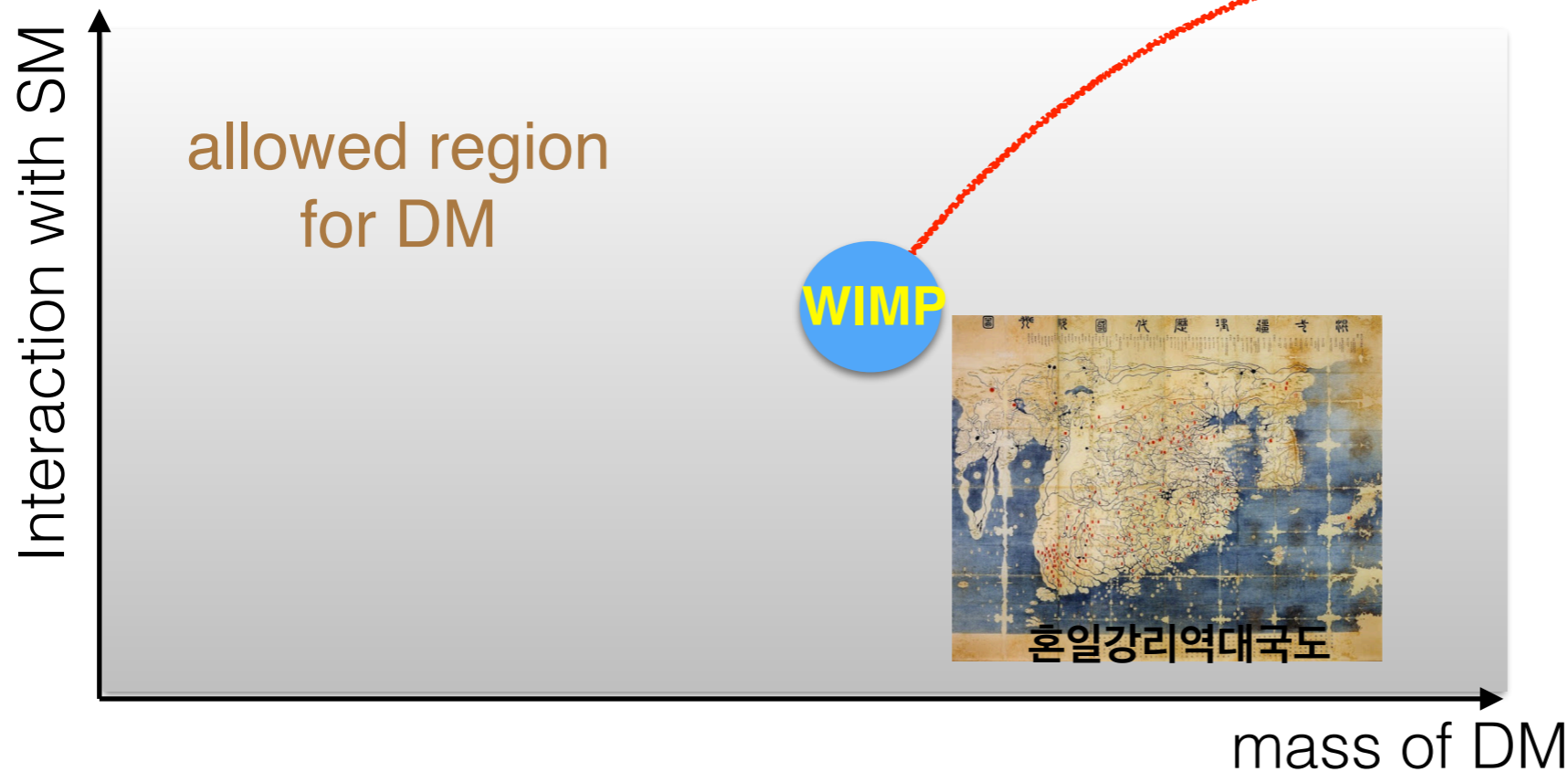


- Models with fast moving light DM are being focused recently.  
(including the light DM reflected by some energetic objects or emitted)
- Boosted Dark Matter (BDM) where the light DM is produced at the present universe by the unique dark sector structure is in the frontline.

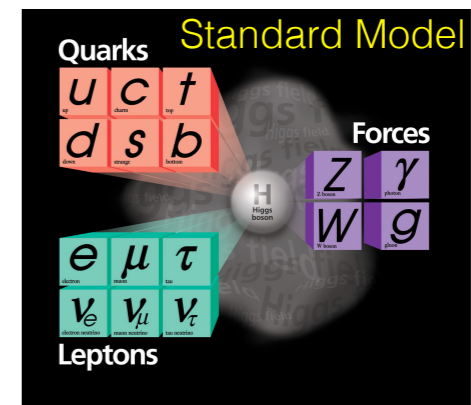


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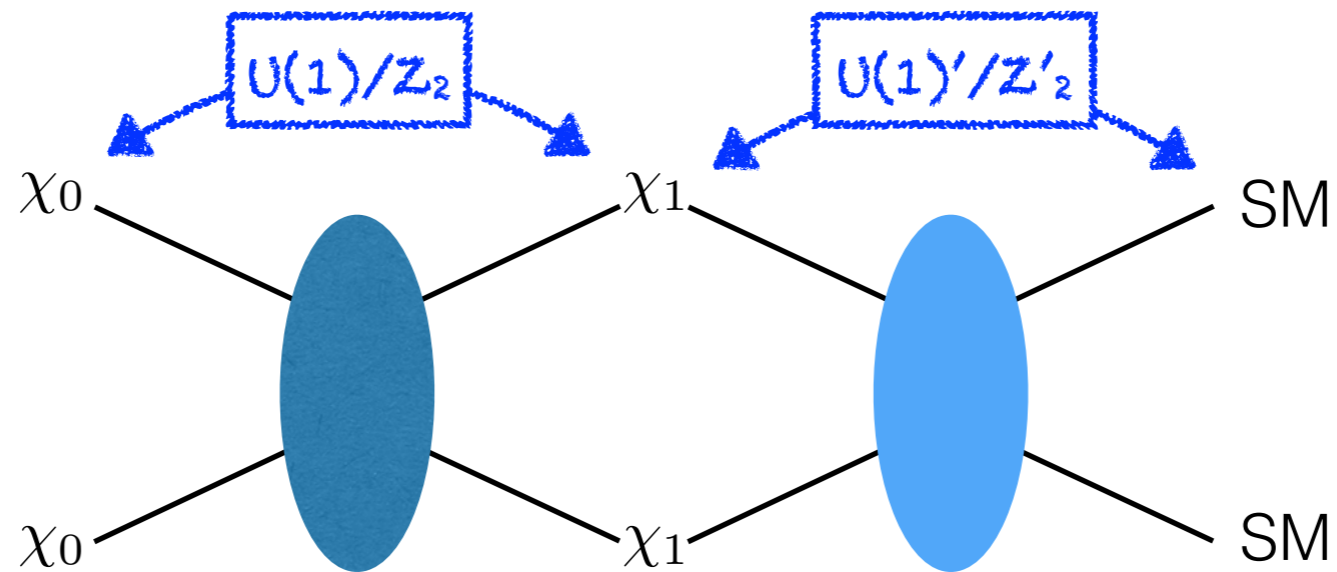


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# Multi-component Boosted DM (BDM)

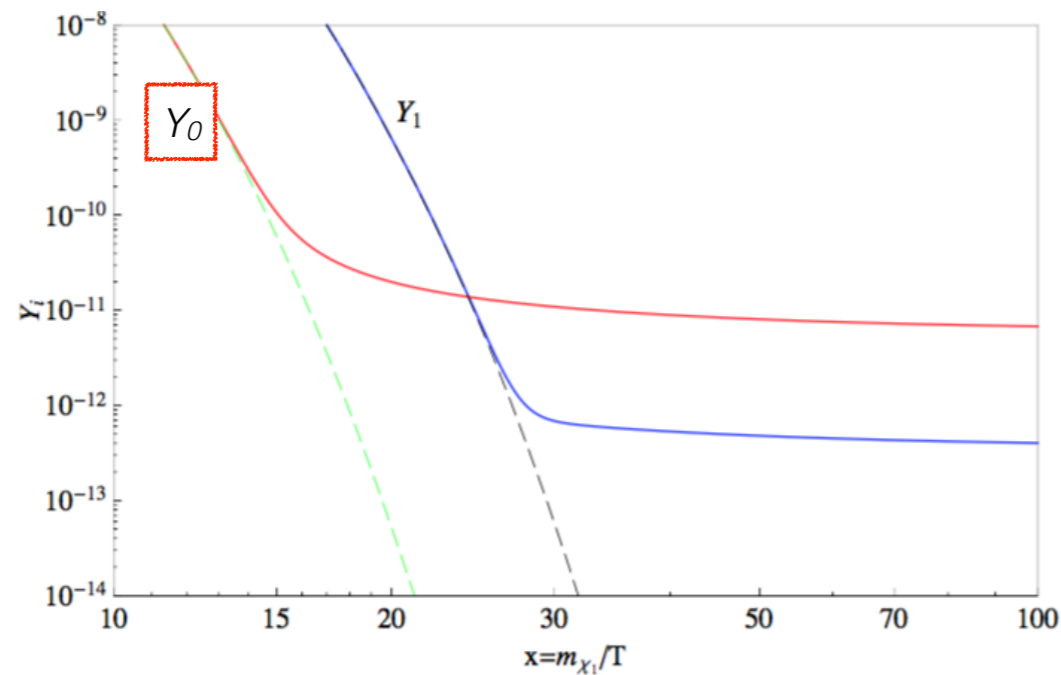
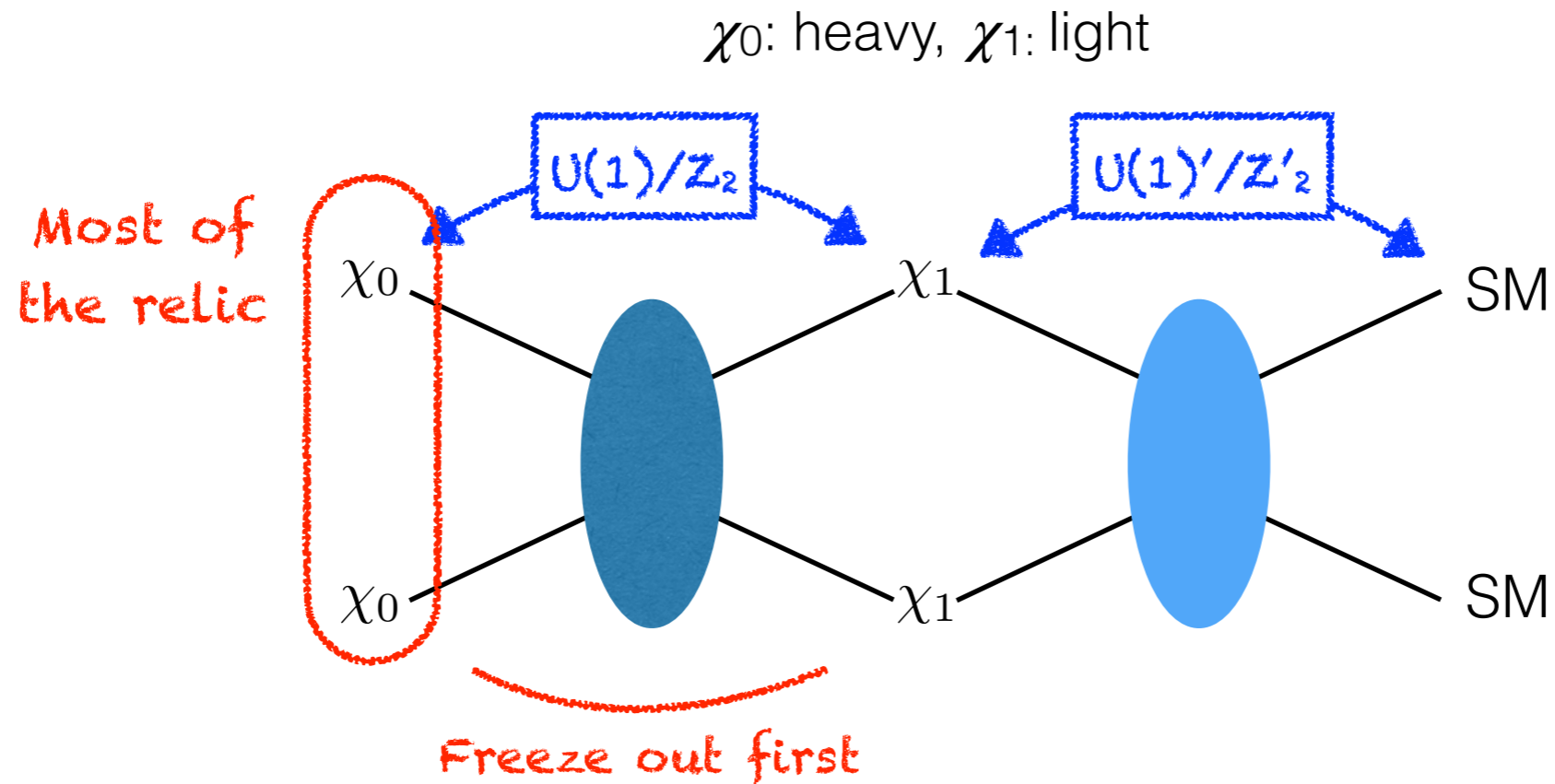
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$\chi_0$ : heavy,  $\chi_1$ : light



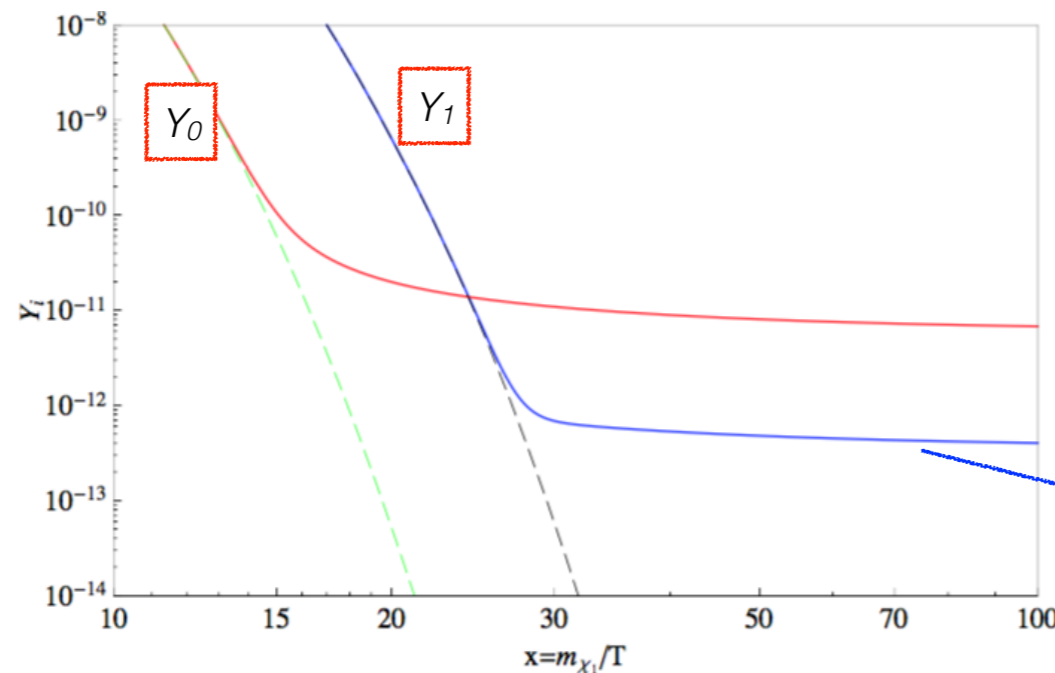
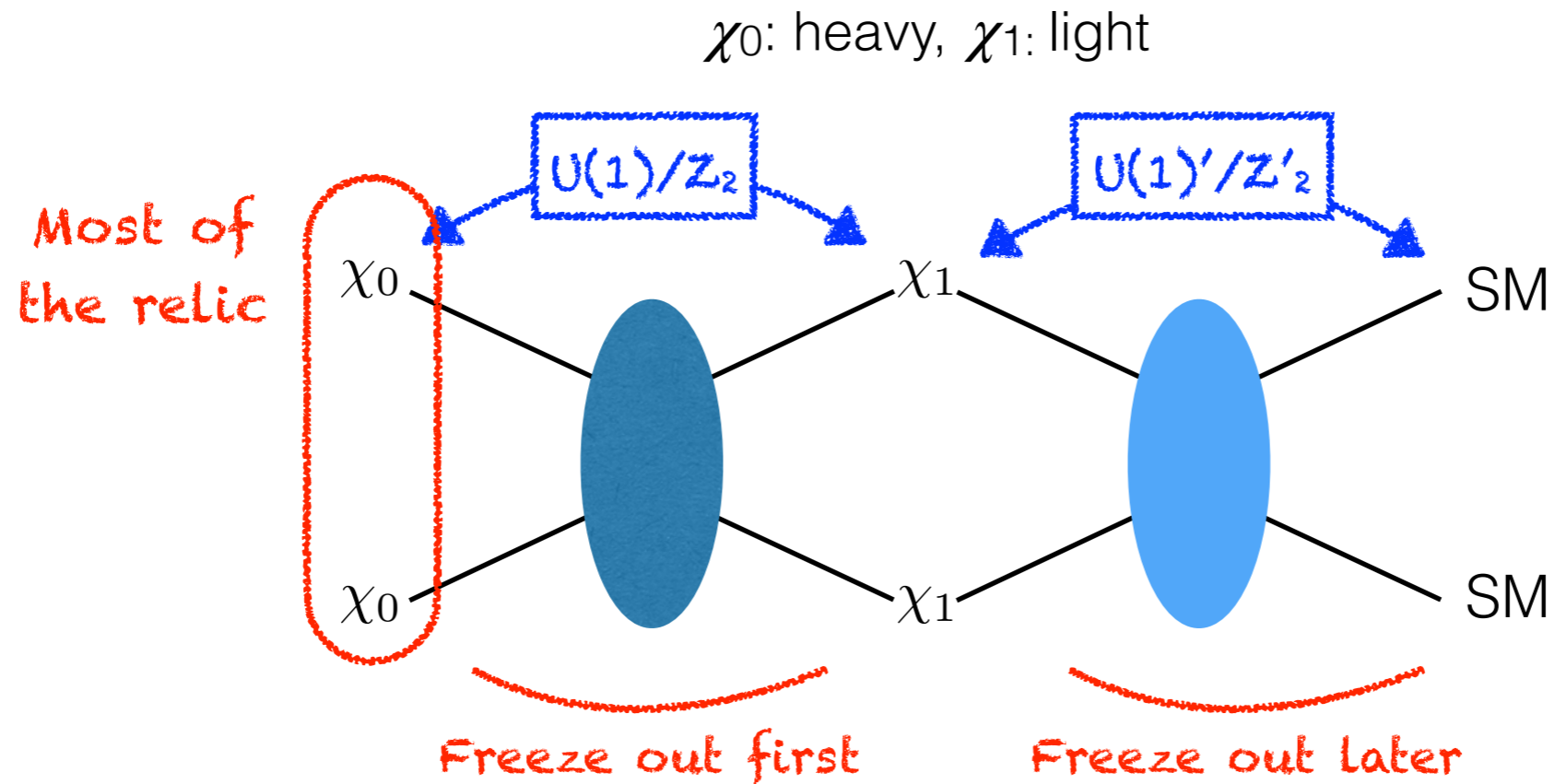
Agashe, Cui, Necib, Thaler, 1405.7370

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Agashe, Cui, Necib, Thaler, 1405.7370

# Multi-component Boosted DM (BDM)



Agashe, Cui, Necib, Thaler, 1405.7370

Belanger, Park, 1112.4491

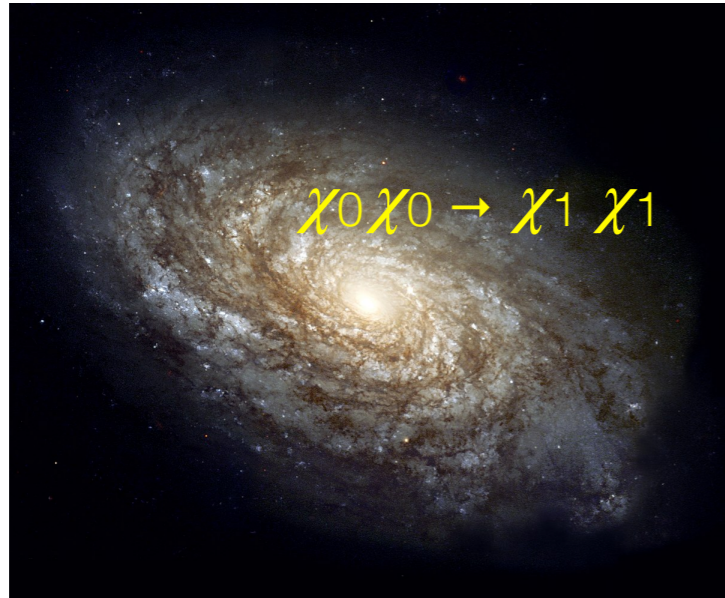
Assisted freeze-out mechanism

non-relativistic relic  $\chi_1$  (negligible)

$$Y_0 \gg Y_1$$

# Multi-component BDM

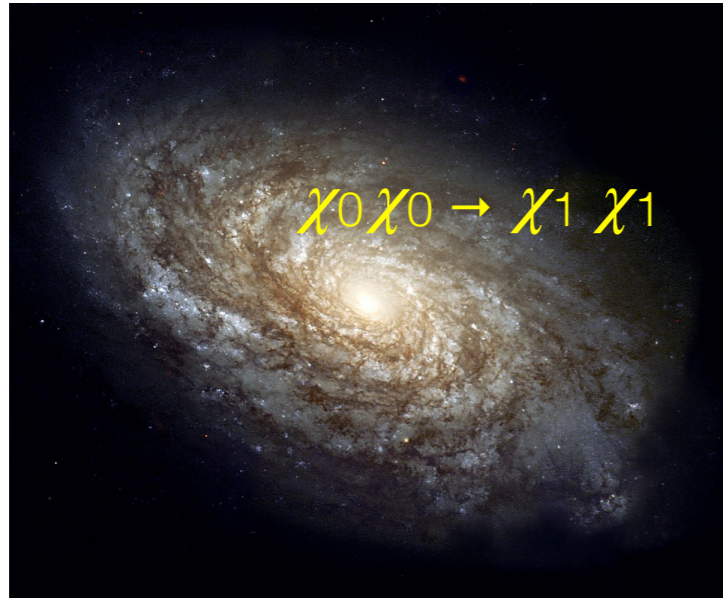
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- $\chi_0$ : accumulated  
(GC, Sun, dSphs)
- $\chi_0\chi_0 \rightarrow \chi_1\chi_1$  (current universe) **relativistic**
  - ※ relic  $\chi_1$  is non-relativistic

# Multi-component BDM

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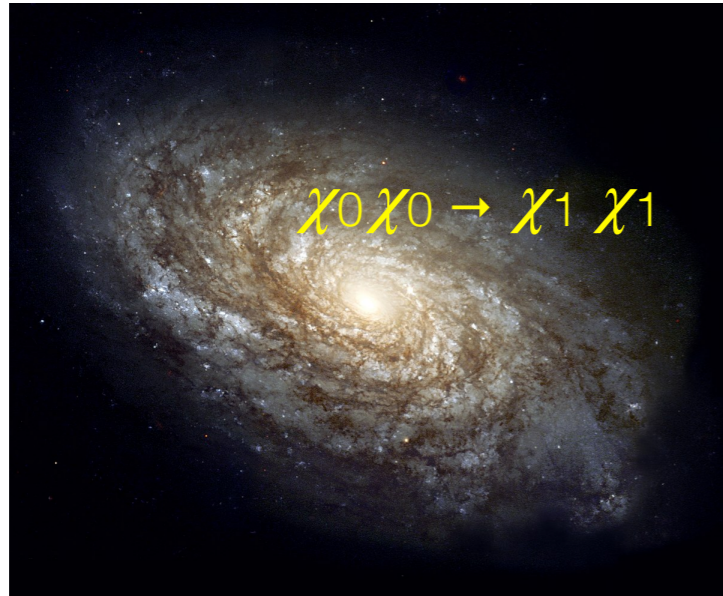


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Observe  $\chi_1$  scattering off target with  $E_1 > E_{th}$   
(indirect detection of  $\chi_0$ )

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$$\text{Flux of } \chi_1 \simeq 1.6 \times 10^{-8} \text{ cm}^{-2} \text{ s}^{-1} \times \left( \frac{\langle \sigma v \rangle_{0 \rightarrow 1}}{5 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}} \right) \times \left( \frac{100 \text{ GeV}}{m_0} \right)^2$$

Assume: NFW



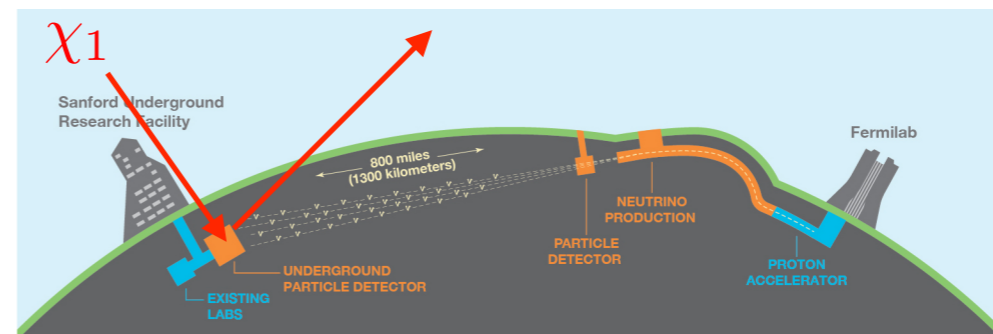
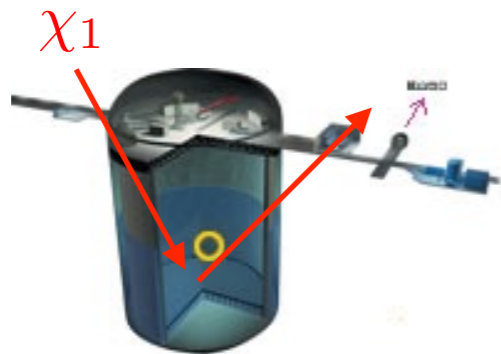
Fixed  $\sim 1$  if s-channel annihilation dominates

10,000 times smaller than the flux of atmospheric  $\nu$  if  $m_0 \sim 100 \text{ GeV}$

# Huge detector if $m_{\chi_0} \approx O(10 \text{ GeV})$

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Flux: small & Energy of  $\chi_1$ : large  $\longrightarrow$  Large volume  $\nu$  experiments

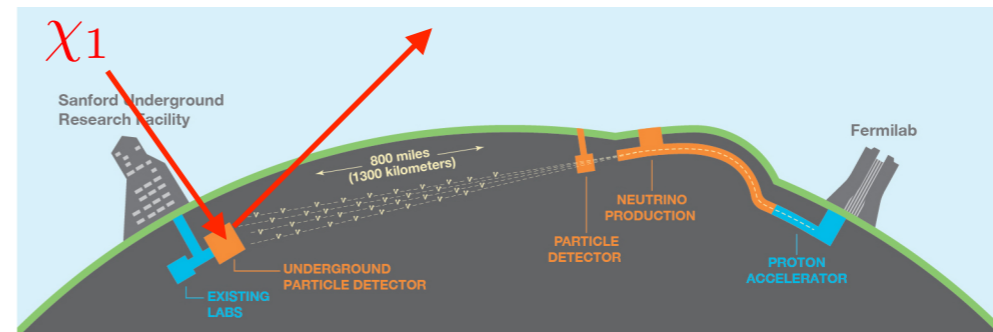
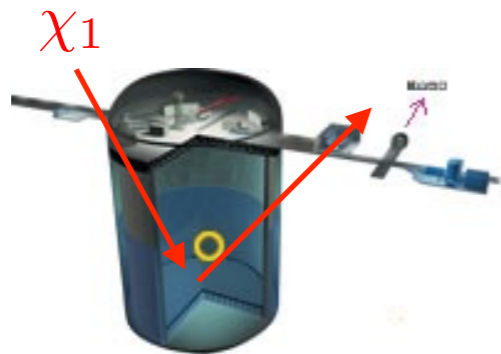




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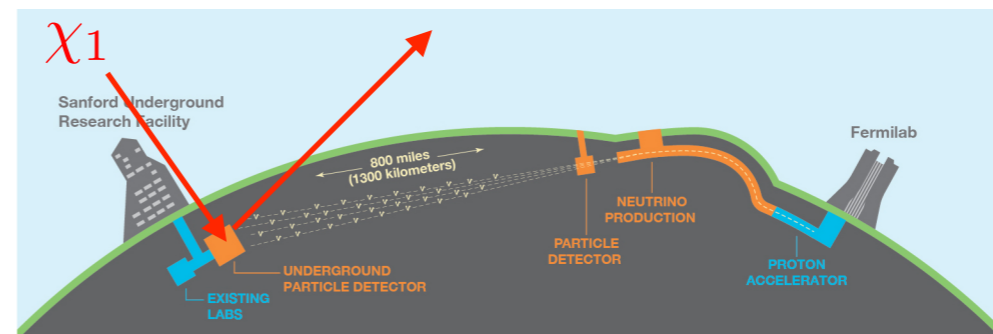
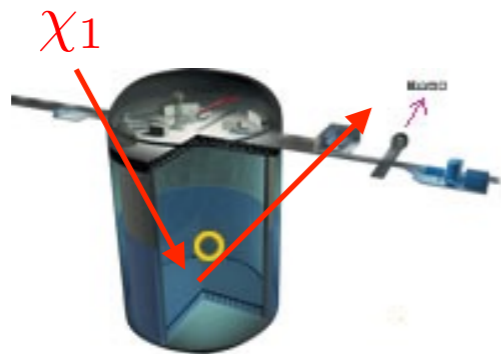


Subtraction of  
major background ( $\nu$ )

Important for all cosmogenic  
BSM signals

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Subtraction of  
major background ( $\nu$ )

Important for all cosmogenic  
BSM signals

- Directional information:  
e.g., GC, Sun, dSphs
- Signal with unique feature

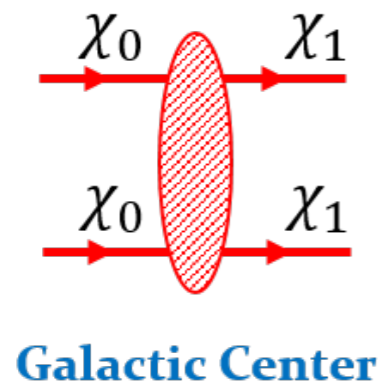
Open up novel possibilities of BDM search in many experiments

# Inelastic BDM (iBDM)

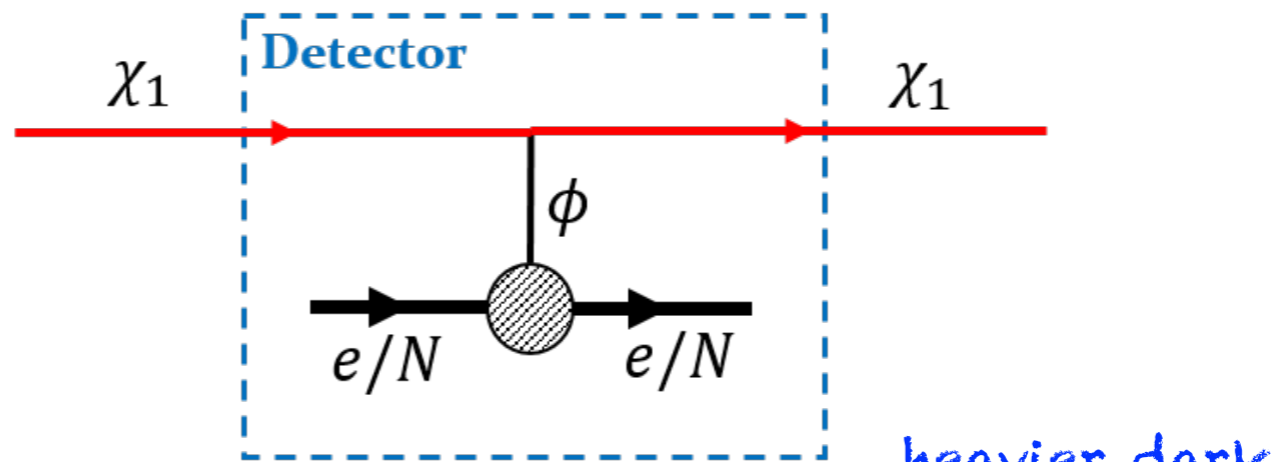
$\chi_0$ : heavy DM

$\chi_1$ : light BDM

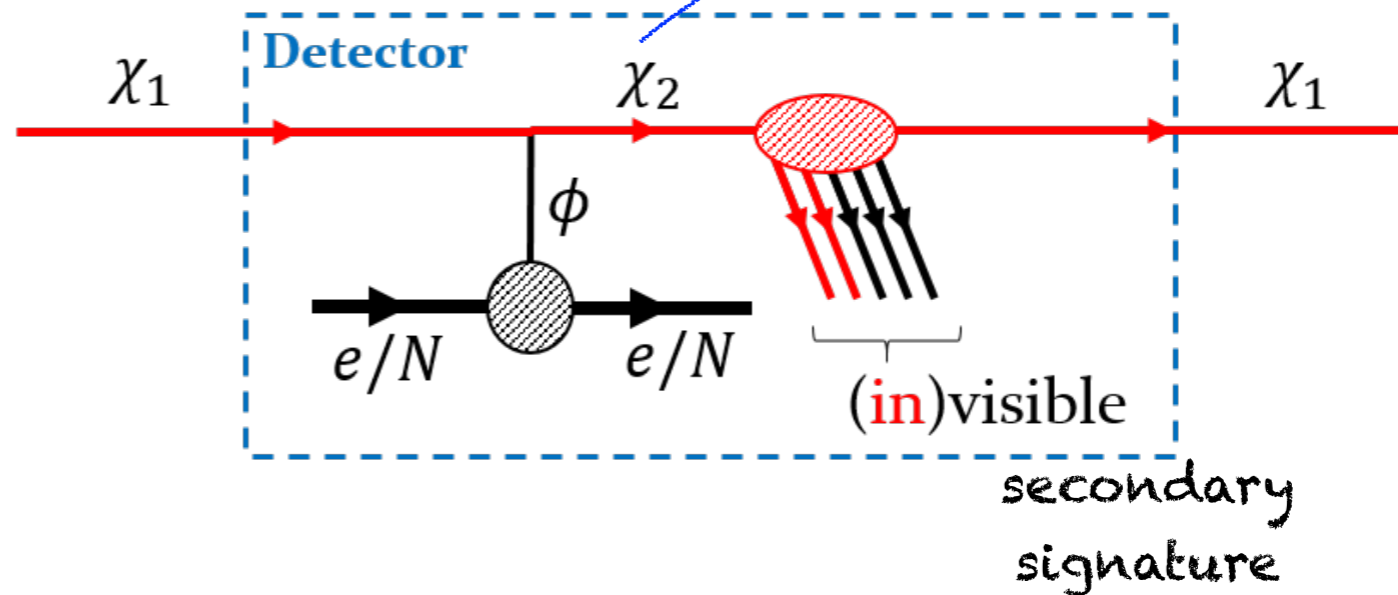
$\chi_2$ : excited state



(a) Elastic scattering (eBDM)



(b) Inelastic scattering (iBDM)



Kim, Park, **SS**, PRL 119, 161801 (2017)

Giudice, Kim, Park, **SS**, PLB 780, 543 (2018)

# Signals inside a fiducial volume

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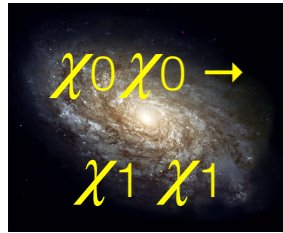
iBDM

Fiducial volume

$p_e$

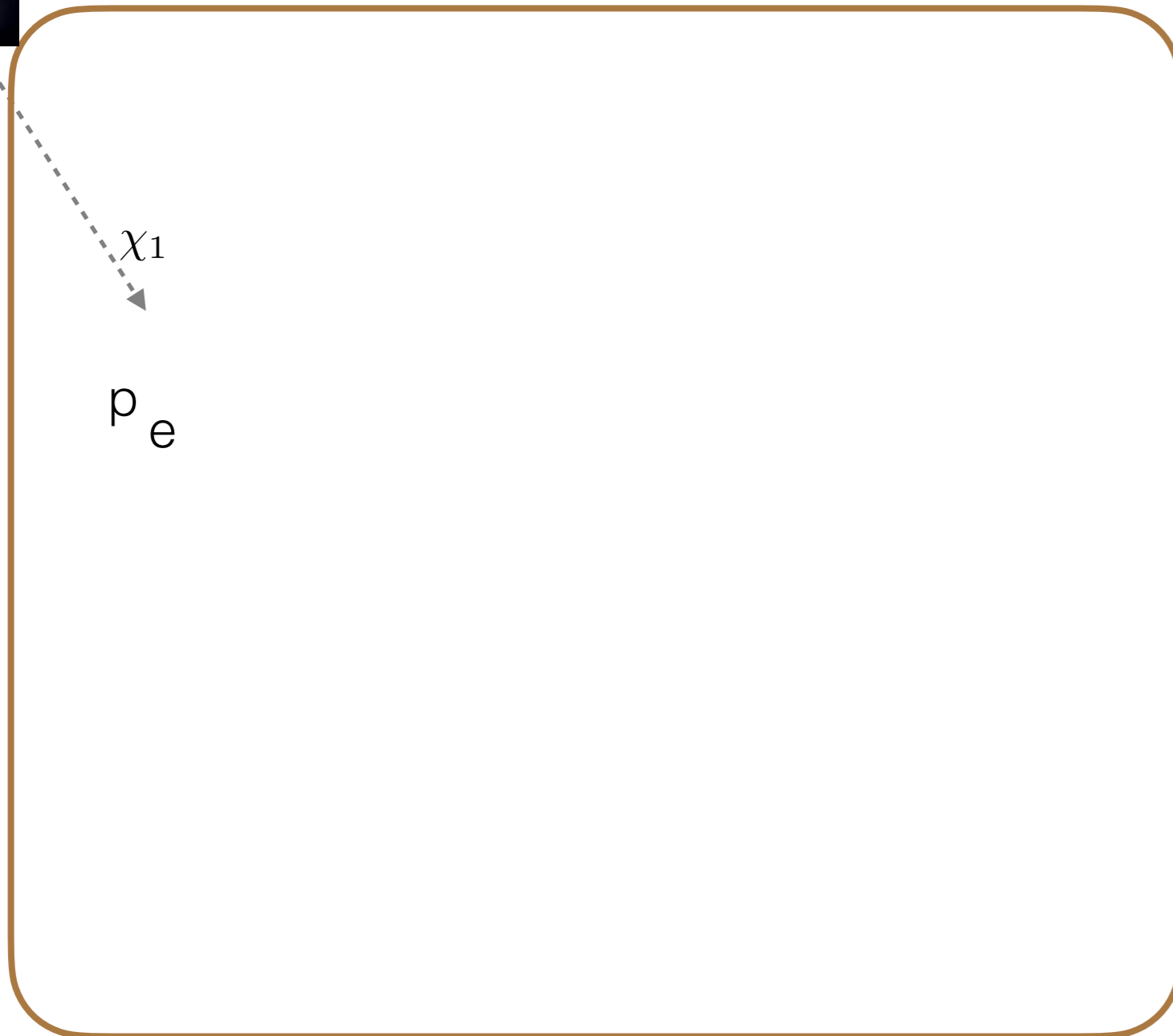
$\chi_1$ : light BDM,  $\chi_2$ : excited state

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iBDM

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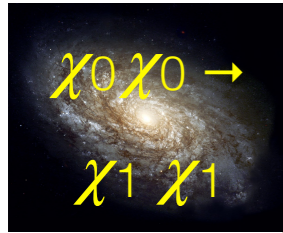


$\chi_1$

$p_e$

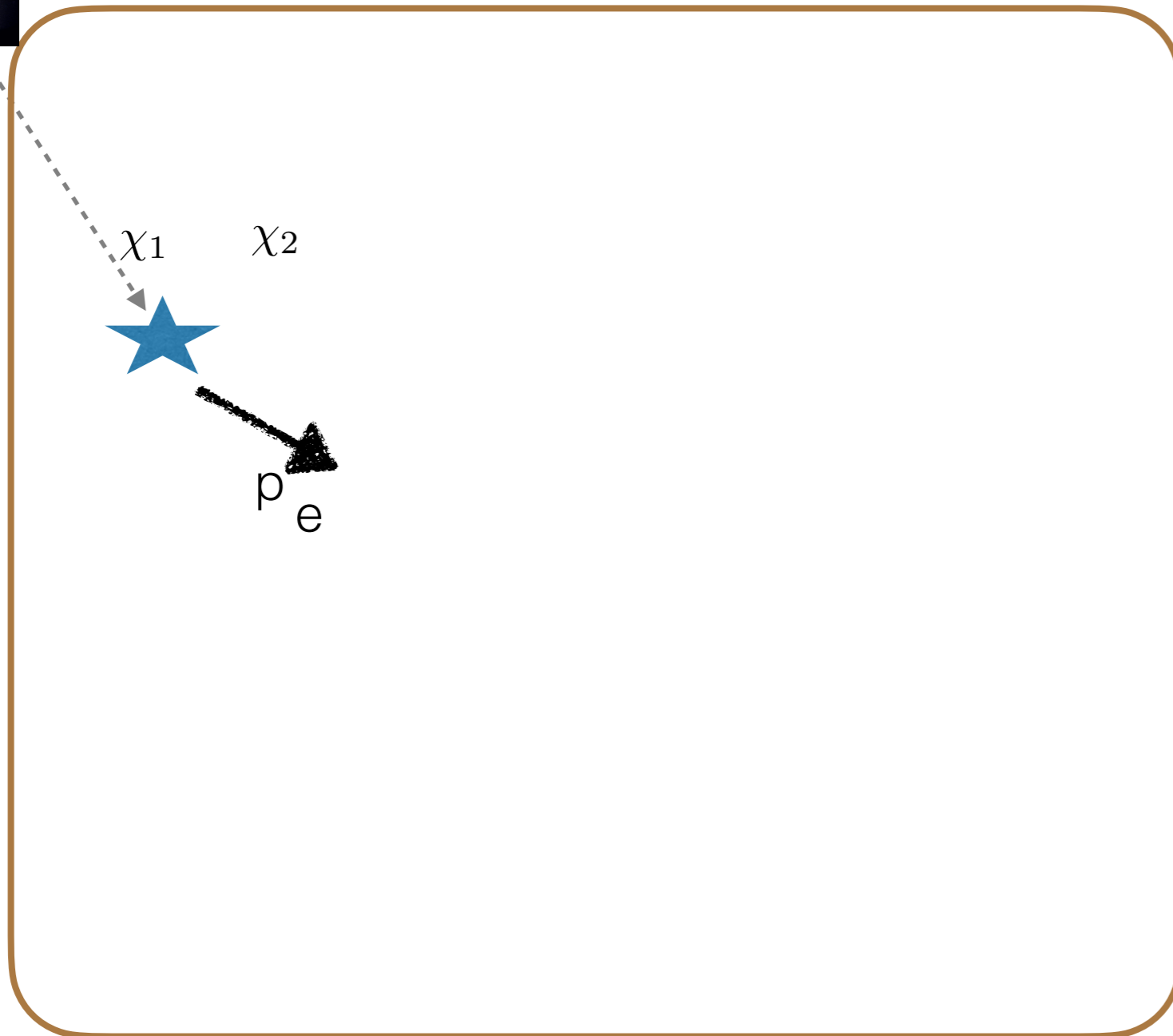
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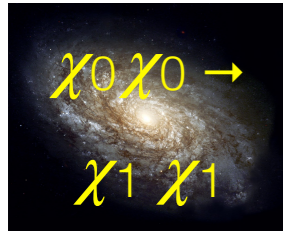
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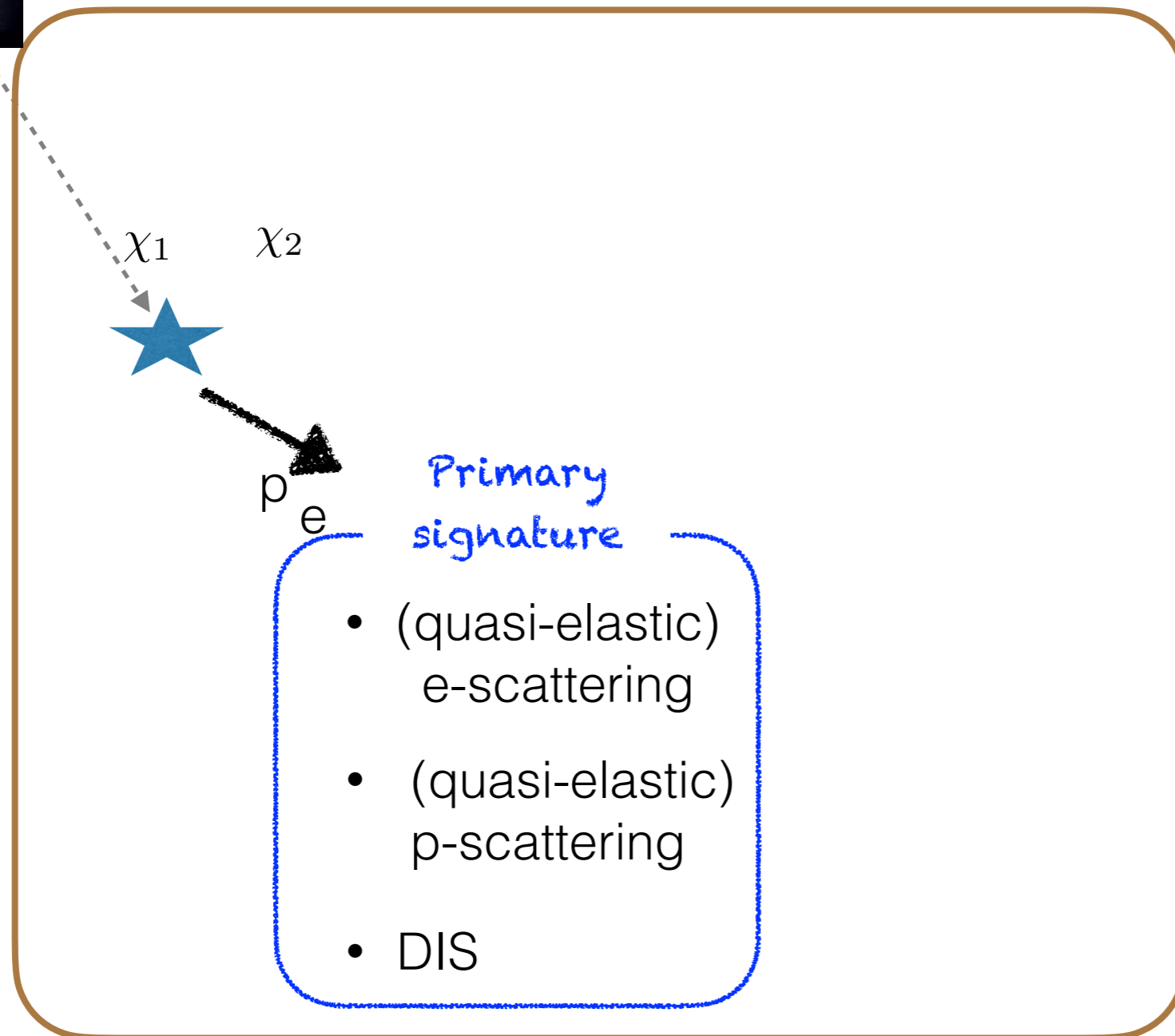
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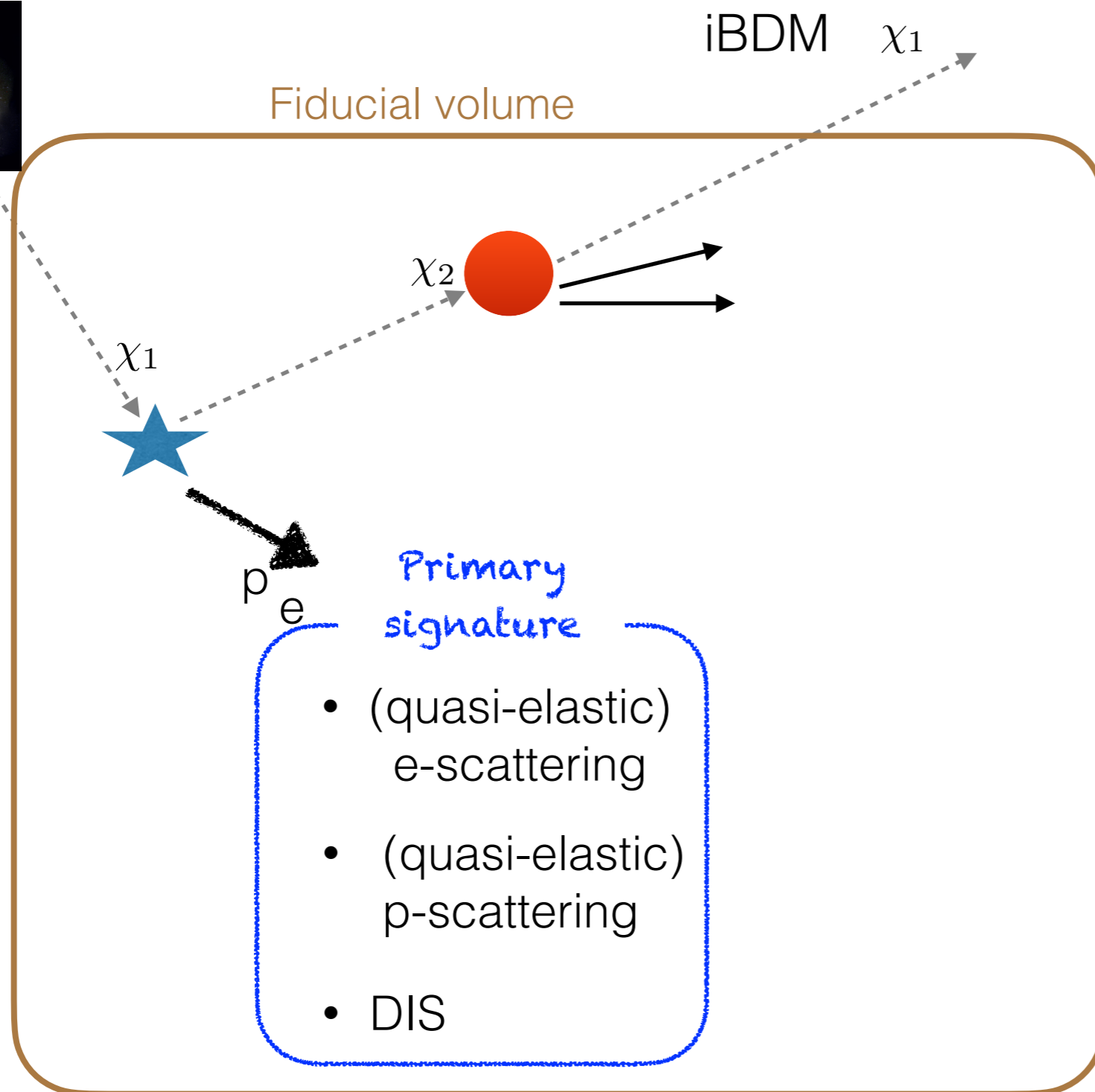
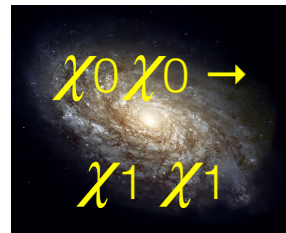
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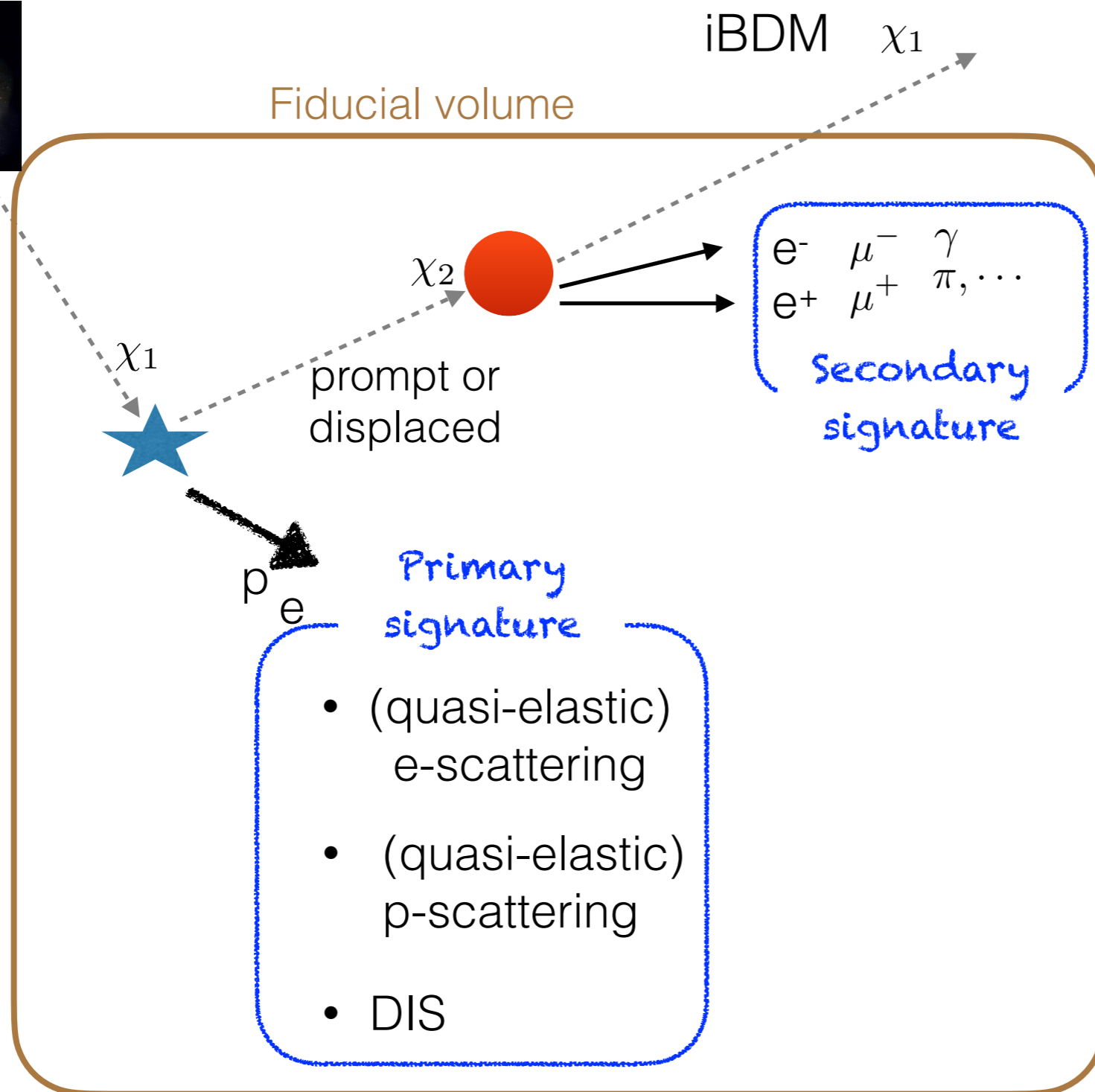
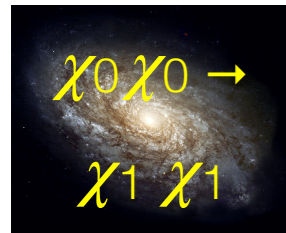
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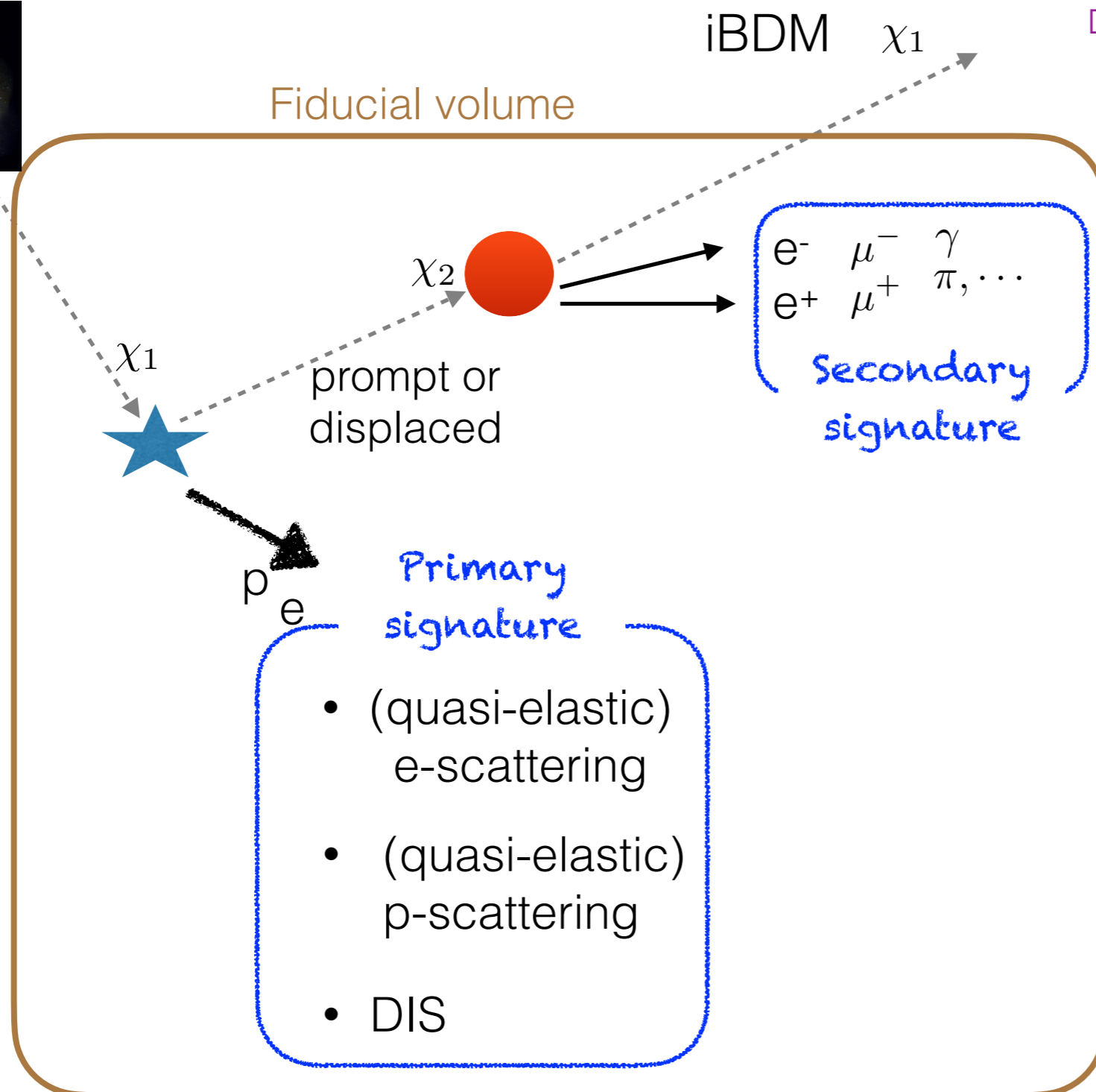
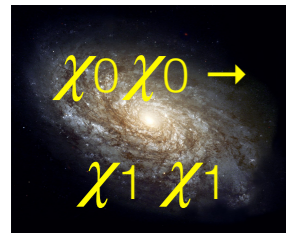


# Signals inside a fiducial volume



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# Signals inside a fiducial volume



De Roeck, Kim, Moghaddam, Park, **SS**,  
Whitehead, 2005.08979

- Tracks pop-up inside fiducial volume (reject the tracks extending outside the fiducial volume)
- Two signatures can be separated beyond the position and angular resolutions. (1cm for all and 1° for e, 5° for p)
- dE/dx for the merged e-recoil signal events

$\chi_1$ : light BDM,  $\chi_2$ : excited state

*Zero - O(10) bkg. (conservative)*

# New method in eBDM search: darkstrahlung

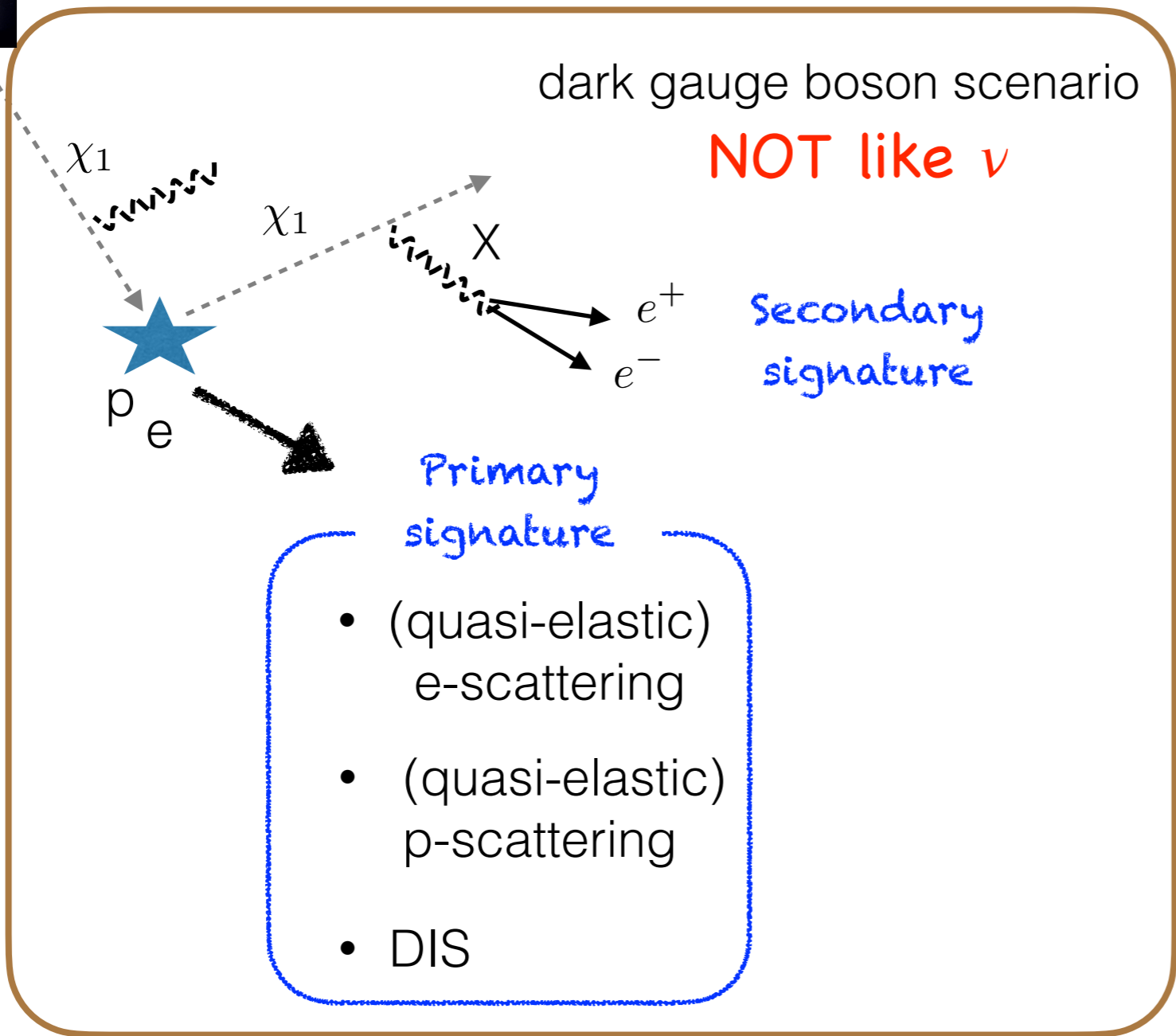
Kim, Park, **SS**, PRD2019, arXiv:1903.05087



Fiducial volume

dark gauge boson scenario

**NOT like  $\nu$**



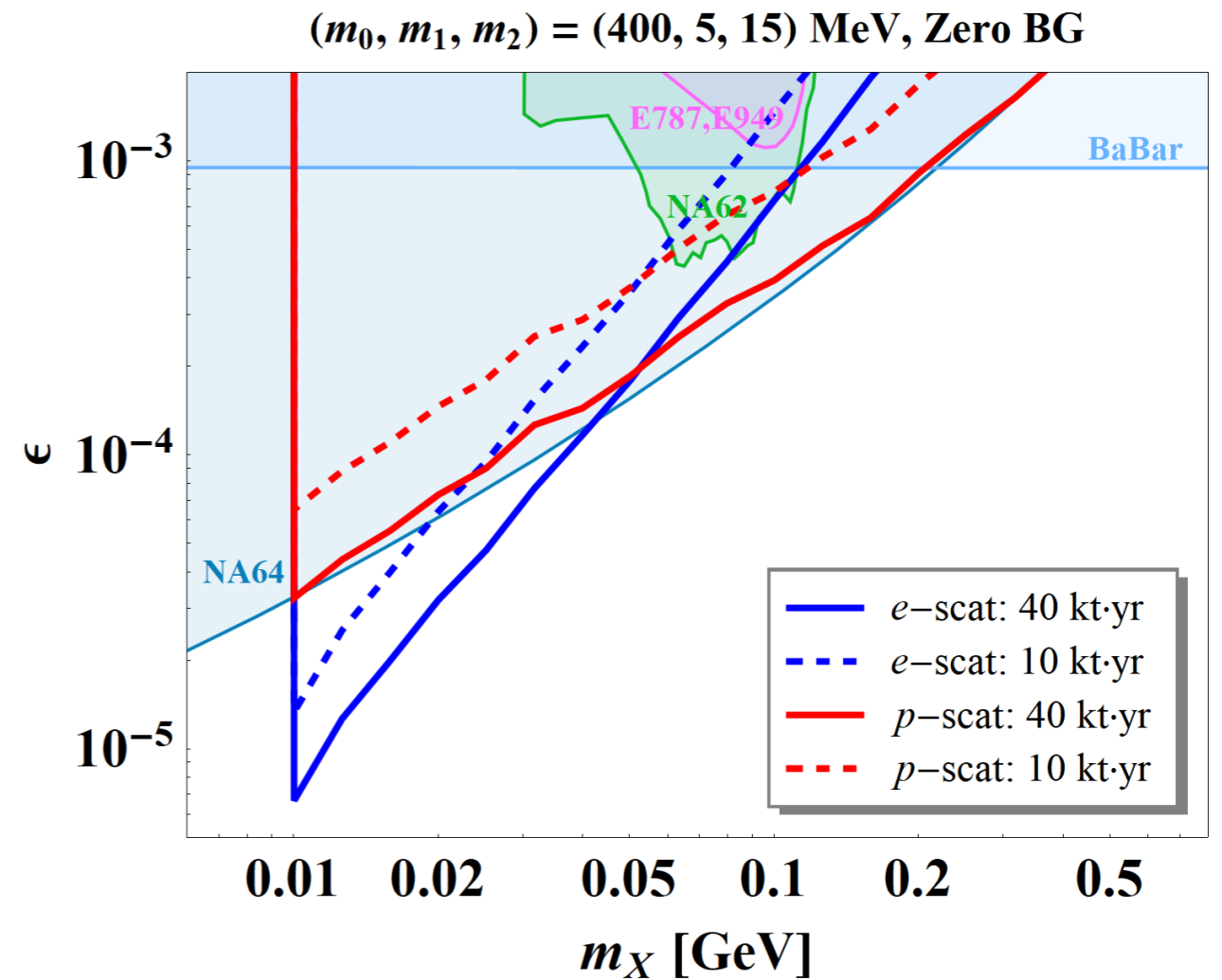
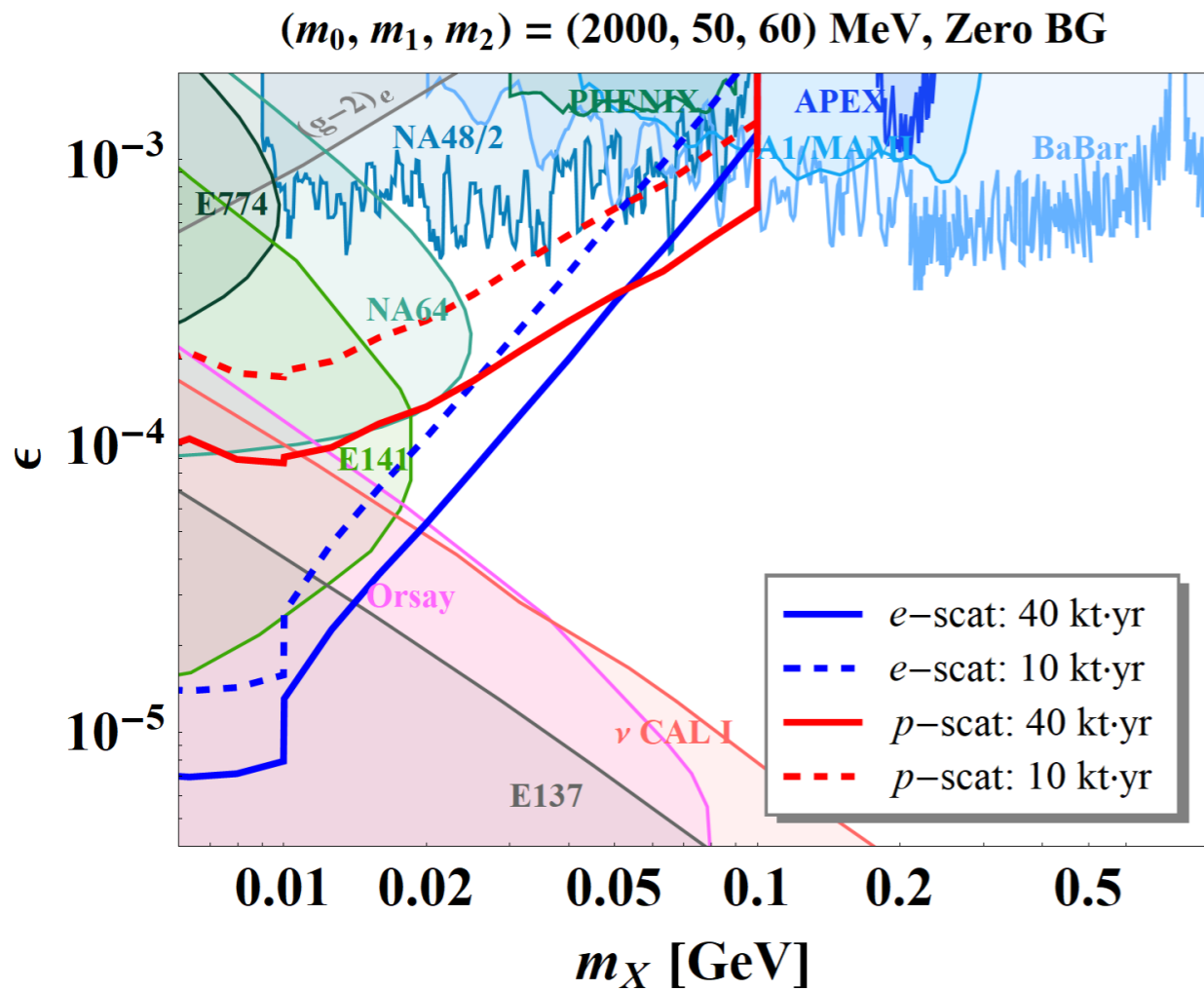
eBDM: elastically scattering BDM

- Different from DM  $\rightarrow \nu \nu$
- NLO but O(10-20%) of LO possible (impossible for beam produced DM)
- Efficient for large  $N_{BG}$  (cosmogenic BSM signal)

$\chi_1$ : light BDM

# Experimental Sensitivities

## DUNE



# Target experiments

$$\text{Flux of } \chi_1 \simeq 1.6 \times 10^{-8} \text{ cm}^{-2} \text{ s}^{-1} \times \left( \frac{\langle \sigma v \rangle_{0 \rightarrow 1}}{5 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}} \right) \times \left( \frac{100 \text{ GeV}}{m_0} \right)^2$$

Assume: NFW

- $m_{\chi_0} \gtrsim O(10 \text{ GeV})$ : Large volume  $\nu$ -experiments:

DUNE (TDR 2002.03005), SK (1711.05278, PRL 2018), HK/KNO, IceCube, ..

- $m_{\chi_0} \lesssim O(1 \text{ GeV})$ : 1. Moderate volume  $\nu$ -experiments:

ProtoDUNE Chatterjee, De Roeck, Kim, Moghaddam, Park, **SS**,  
Whitehead, Yu, 1803.03264, PRD 98, 075027 (2018)

SBNP Kim, Kong, Park, **SS**, 1804.07302, JHEP 1808, 155 (2018)

- 2. Ton-scale DM direct detection experiments:

Giudice, Kim, Park, **SS**, PLB 780, 543 (2018)

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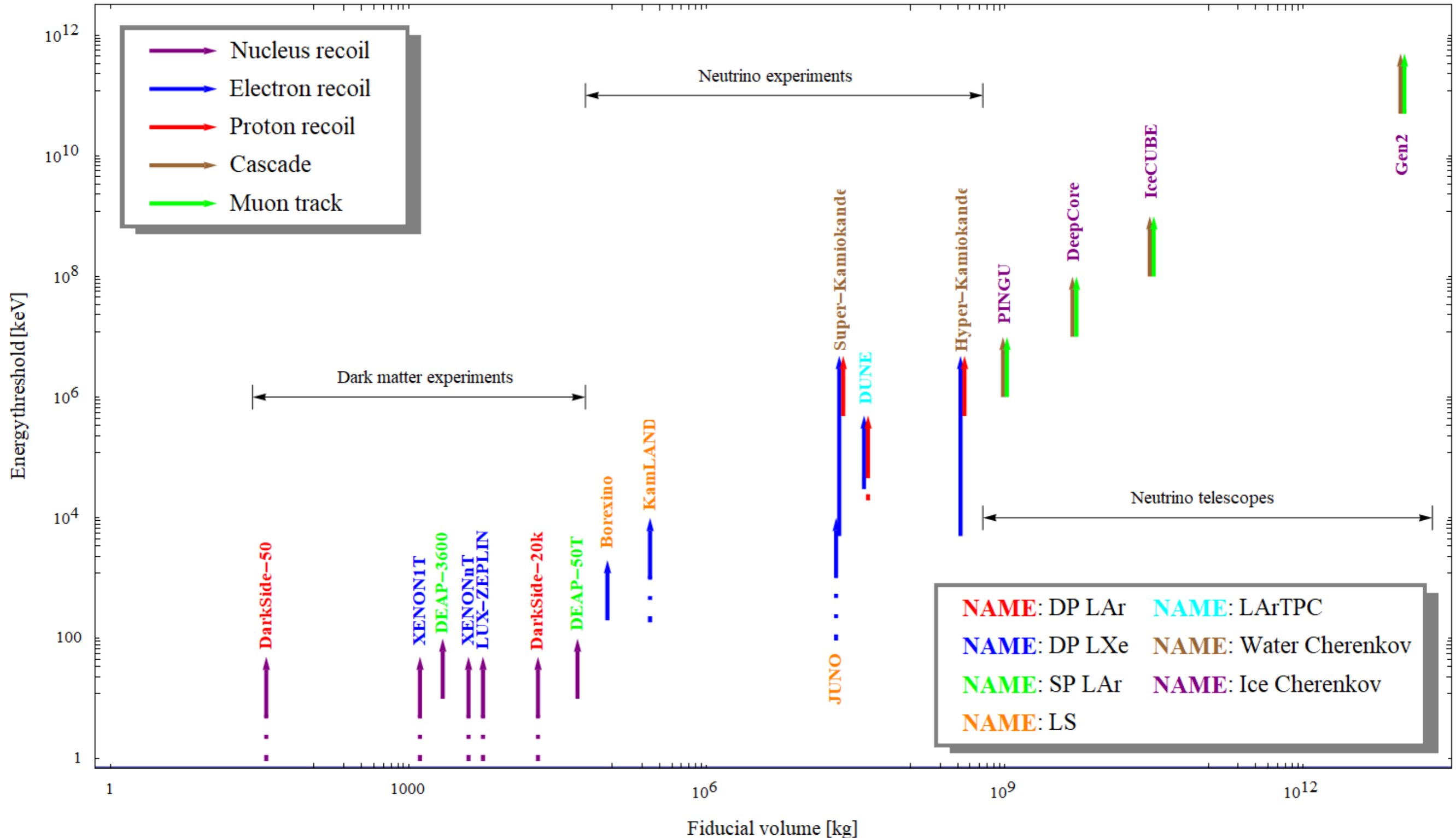
First proposal of high e-recoil  
in DM direct detection exp.  
such as XENON1T



2. Ton-scale DM direct detection experiments:

Giudice, Kim, Park, **SS**, PLB 780, 543 (2018)

# Search for BDM

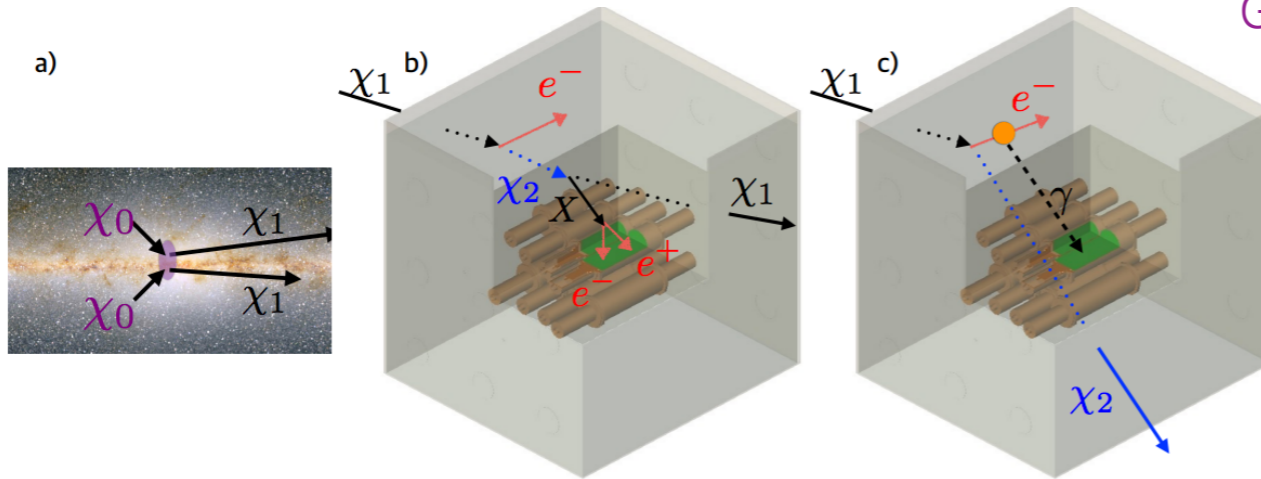


# COSINE-100 result

COSINE-100, PRL 2019

Based on theoretical study

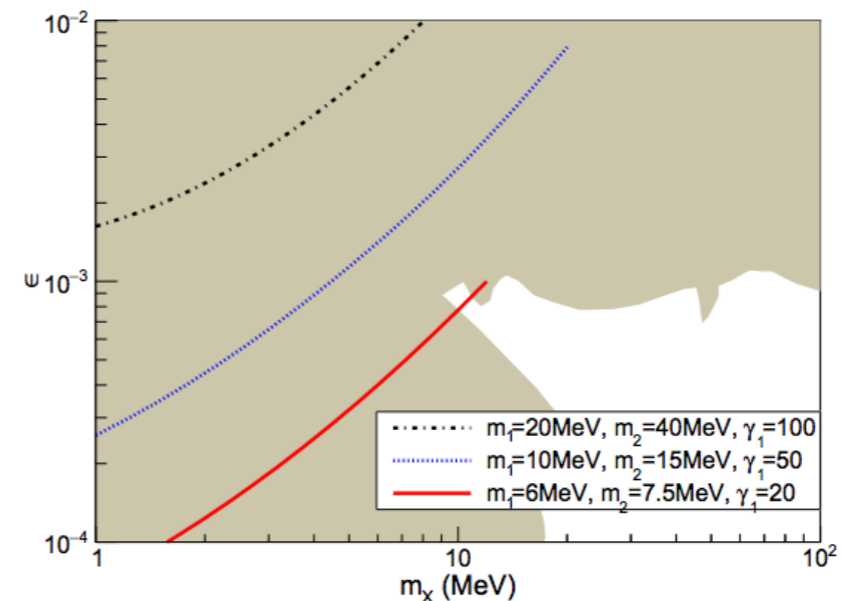
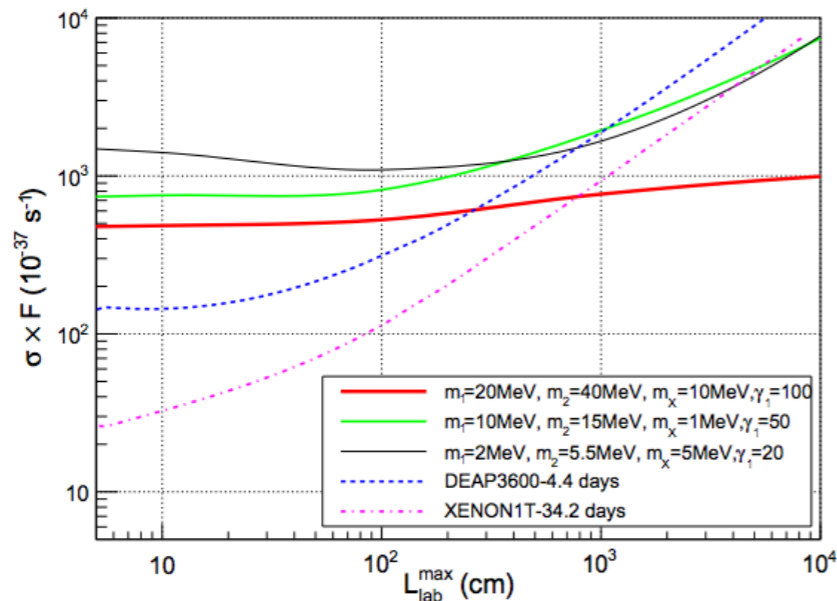
Giudice, Kim, Park, **SS**, PLB 780, 543 (2018)



2200L of liquid scintillator  
(~ 2 ton)

106kg array of 8 ultra-pure NaI(Tl) crystals  
immersed in an active veto detector

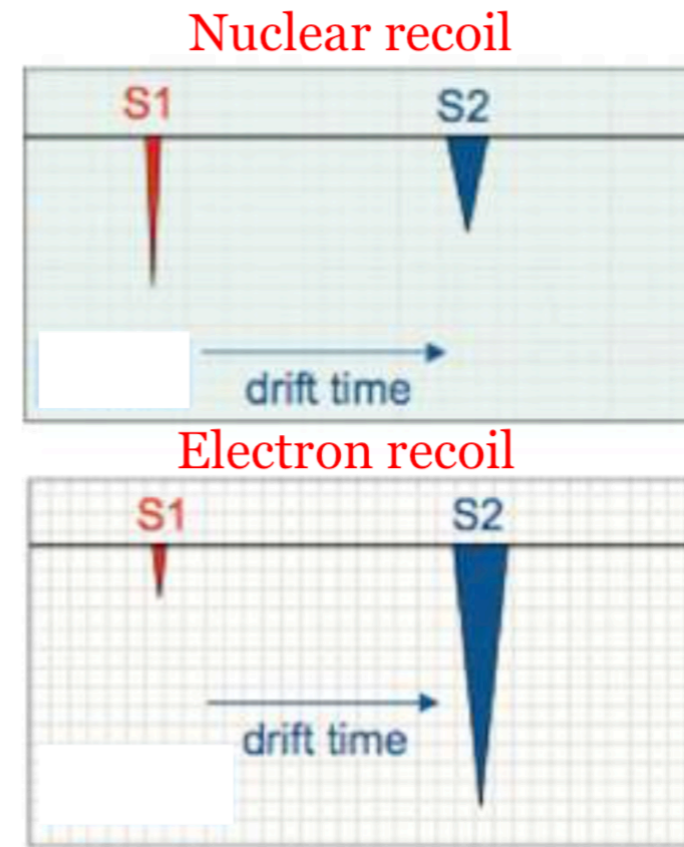
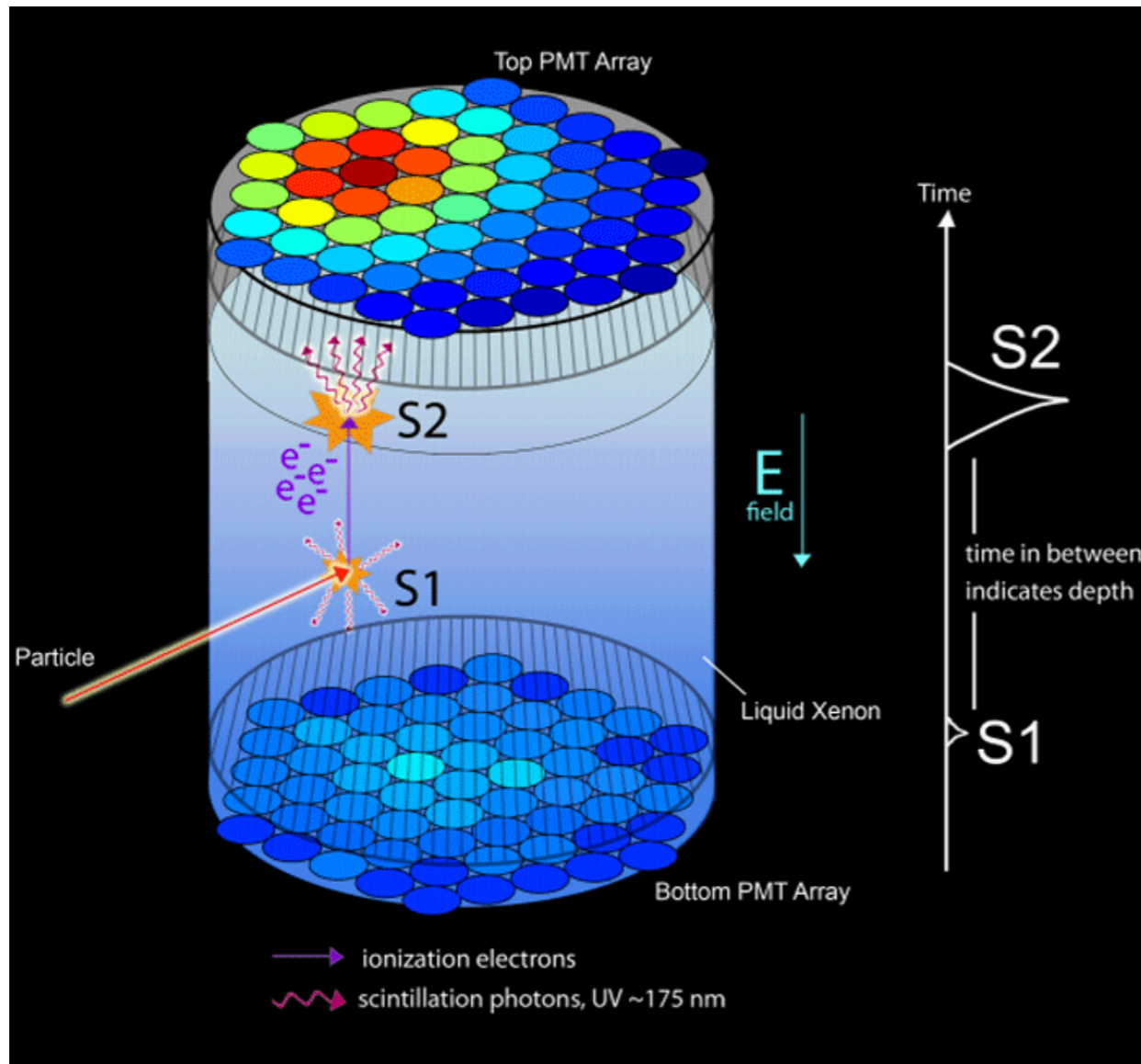
Observed: 21 events, Background expected:  $16.4 \pm 2.1$



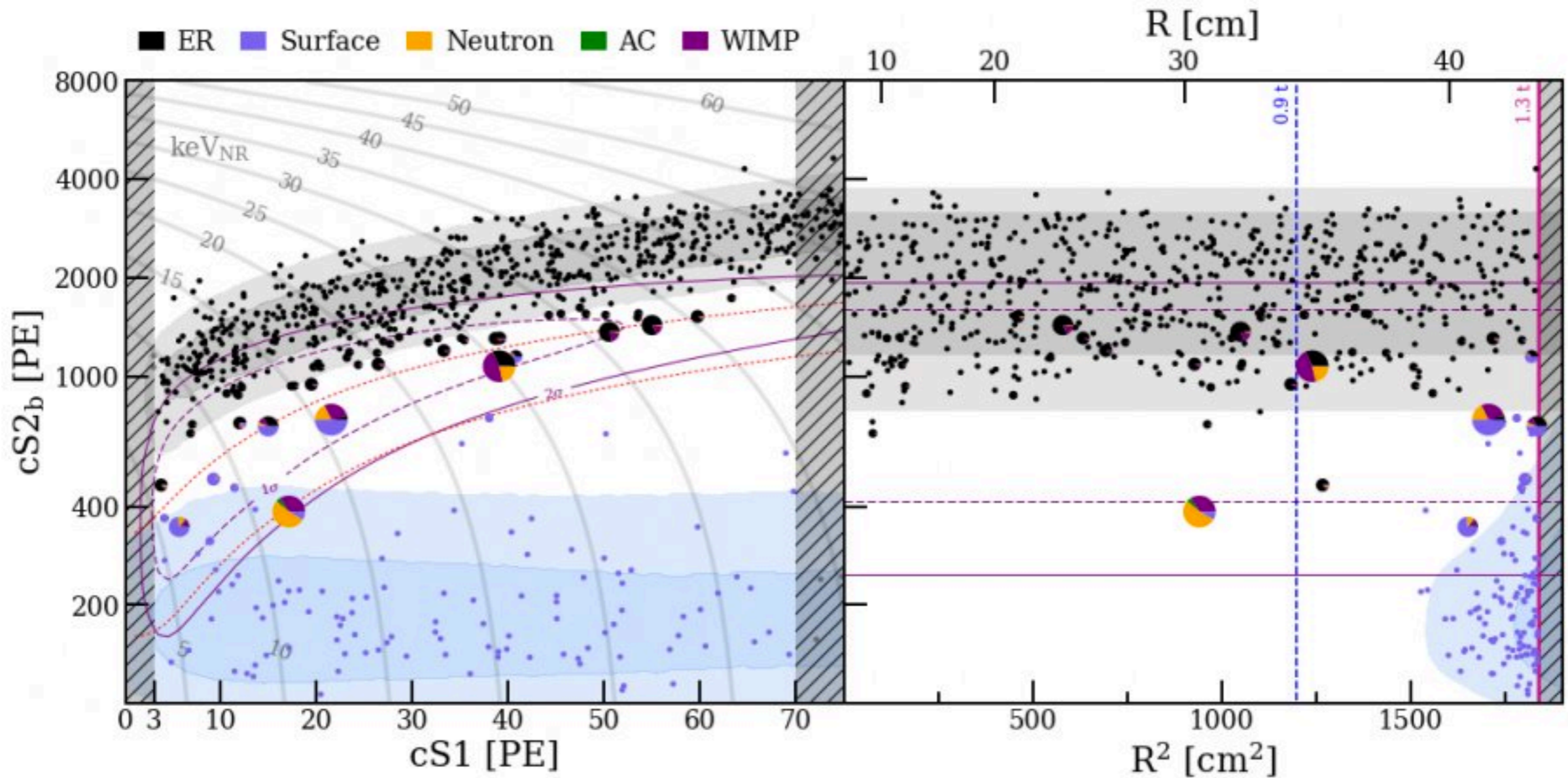


# XENON detectors

XENON1T, NT, LUX-ZEPLIN, PANDA, ...



# XENON detectors

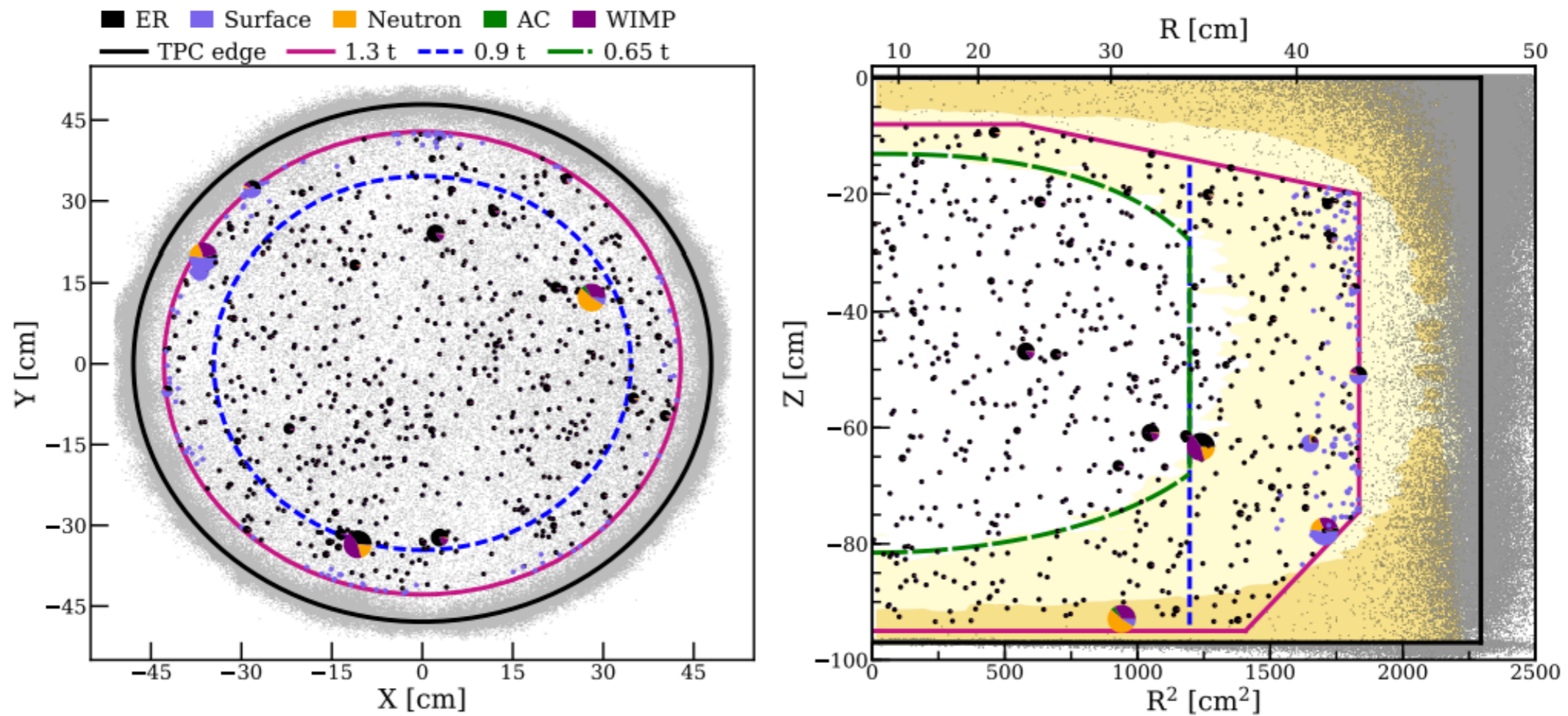


XENON1T, 1805.12562

Neutrons include cosmogenic neutrons and radiogenic neutrons from detector materials.

AC: Artificial Coincidence, an artificial background from detector effects

# XENON detectors

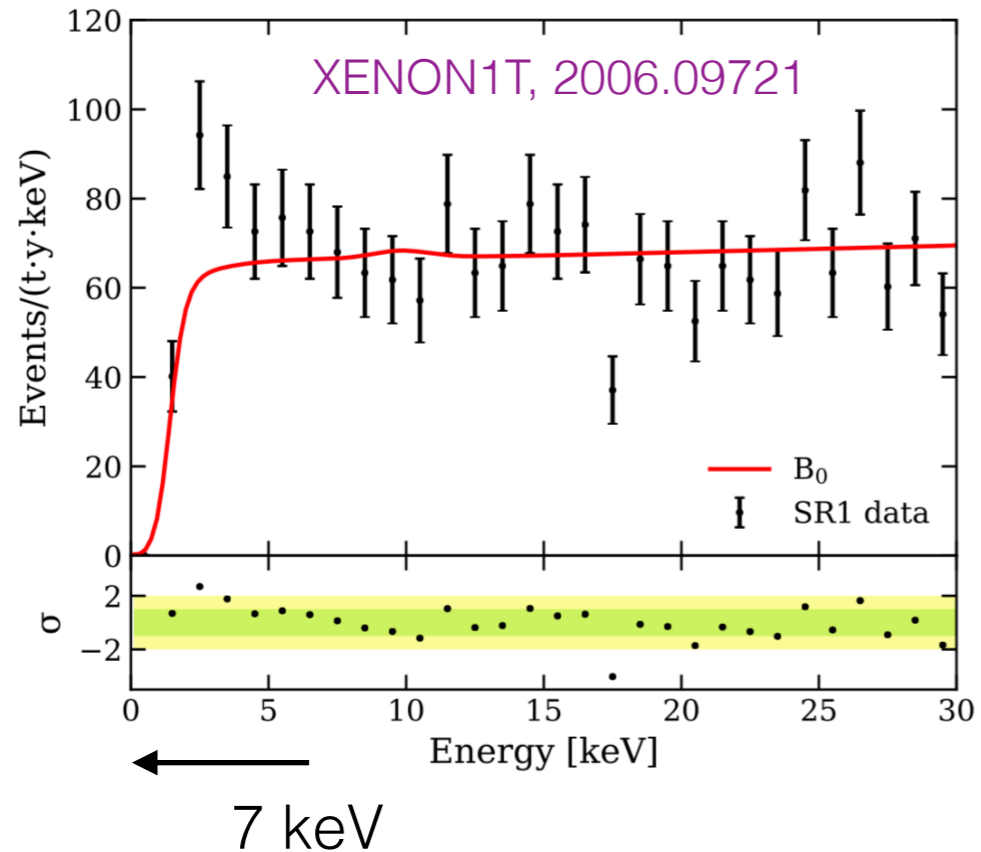


XENON1T, 1805.12562

# XENON1T 2020

0.65 ton·year

$76 \pm 2(\text{stat})$  events exceeding background expectation

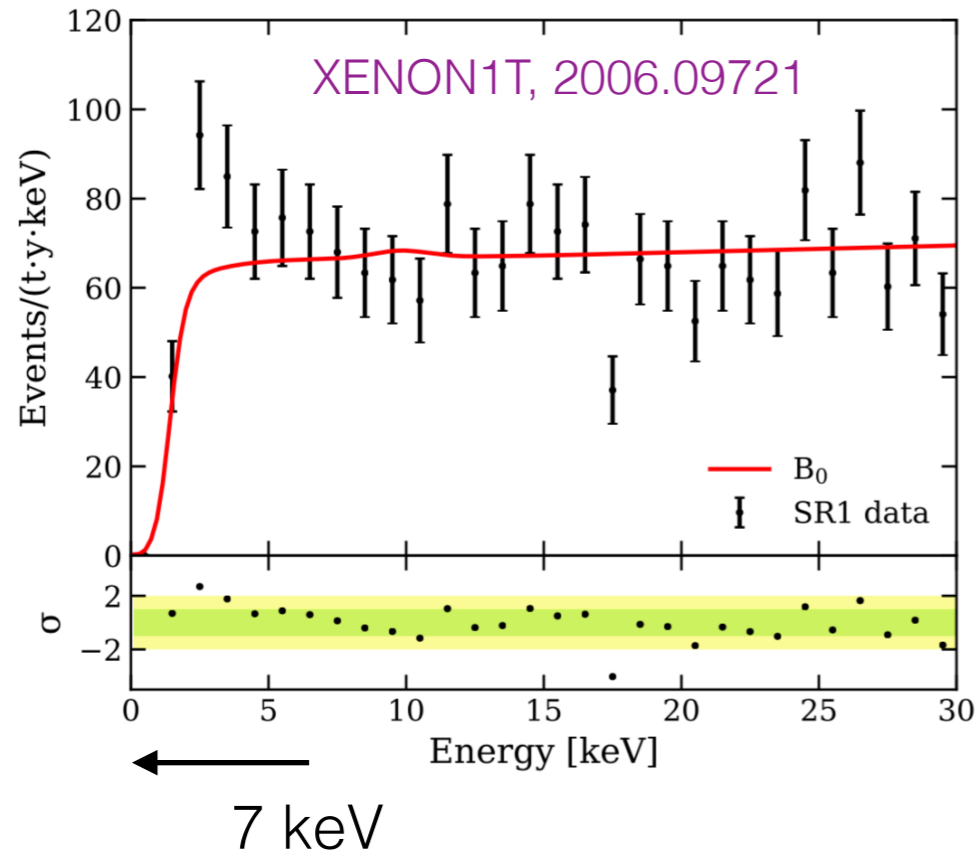


- Background? Tritium, Ar37 decays (most probable?)
- Solar axion, neutrino MDM  $\sim 3\sigma$  ?
- Dark Matter recoil?

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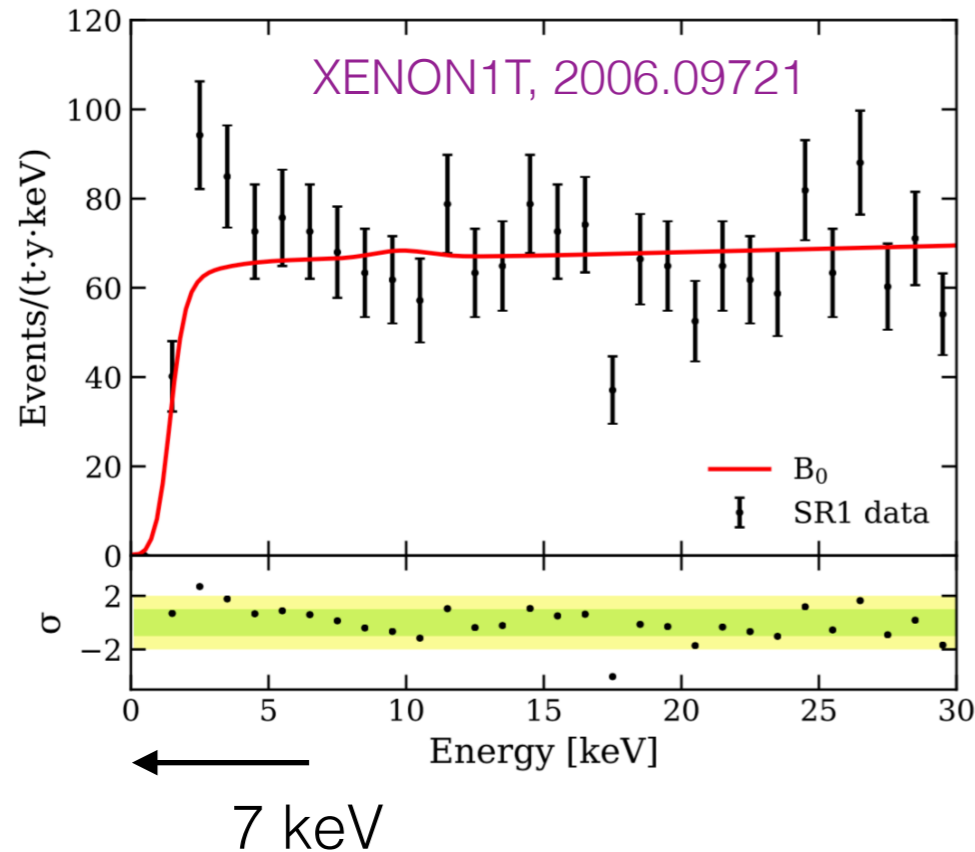
- No nuclear recoil signals
- Too high energy  $e$ -recoil for WIMP or WIMP-like light DM



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The electron target is in a bound state with a typical wavelength size of  $R_{\text{Bohr}} = 1/(\alpha m_e)$

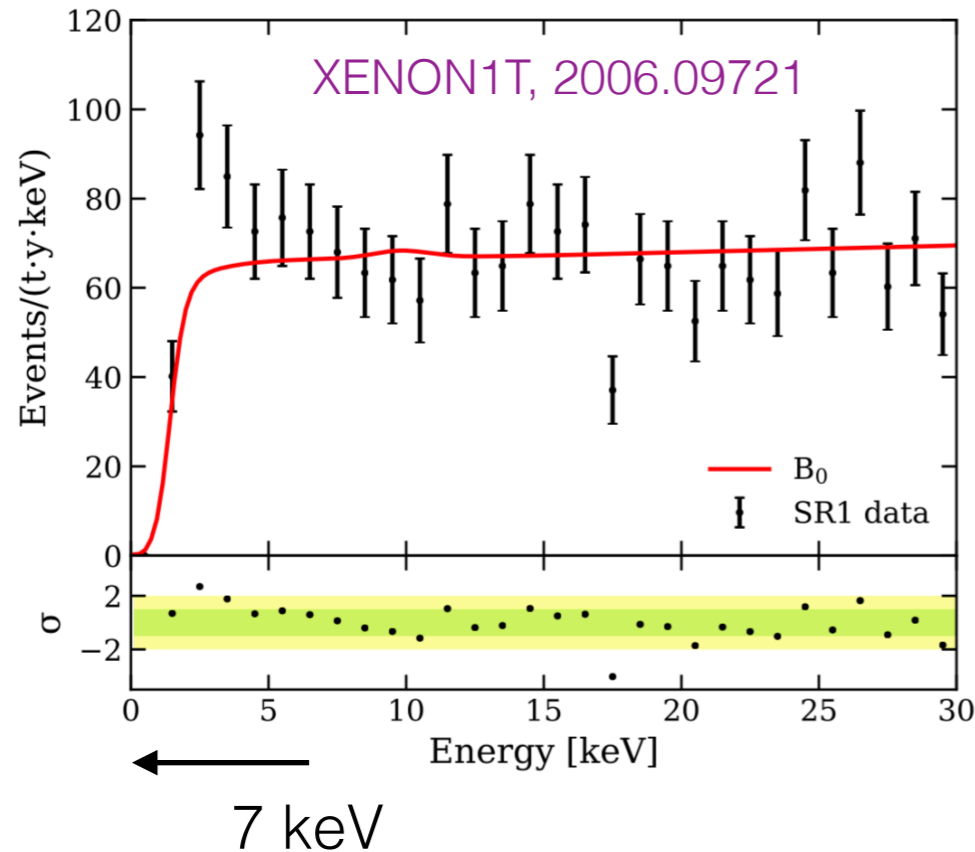
The typical momentum of the electron is  $k_e \simeq 1/R_{\text{Bohr}} = \alpha m_e \longrightarrow v_e \simeq \alpha \lesssim 0.01$

The deposited energy by slowly moving DM ( $v \sim 10^{-3}$ ) is at most  $\mathcal{O}(\text{eV}) \ll \text{keV}$ .

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0.65 ton·year

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**Fast-moving Light DM**

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- Solar axion, neutrino MDM  $\sim 3\sigma$  ?
- Dark Matter recoil?

- No nuclear recoil signals
- Too high energy  $e$ -recoil for WIMP or WIMP-like light DM



The electron target is in a bound state with a typical wavelength size of  $R_{\text{Bohr}} = 1/(\alpha m_e)$

The typical momentum of the electron is  $k_e \simeq 1/R_{\text{Bohr}} = \alpha m_e \longrightarrow v_e \simeq \alpha \lesssim 0.01$

The deposited energy by slowly moving DM ( $v \sim 10^{-3}$ ) is at most  $\mathcal{O}(\text{eV}) \ll \text{keV}$ .

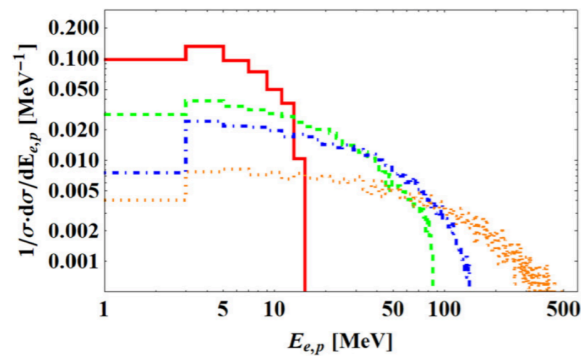
# XENON1T 2020

The first proposal of searching for high energy  $e$ -recoil by fast-moving light DM

## Inelastic Boosted Dark Matter at direct detection experiments

Gian F. Giudice<sup>a,\*</sup>, Doojin Kim<sup>a,\*</sup>, Jong-Chul Park<sup>b,\*</sup>, Seodong Shin<sup>c,d,\*</sup>

*Theoretical Physics Department, CERN, Geneva, Switzerland*  
*Department of Physics, Chungnam National University, Daejeon 34134, Republic of Korea*  
*Enrico Fermi Institute, University of Chicago, Chicago, IL 60637, USA*  
*Department of Physics & IPAP, Yonsei University, Seoul 03722, Republic of Korea*



Giudice, Kim, Park, **SS**, PLB 780, 543 (2018)

	$m_1$	$m_2$	$m_X$	$\gamma_1$	$\epsilon$
ref1 (red solid)	2	5.5	5	20	$4.5 \times 10^{-5}$
ref2 (green dashed)	3	8.5	7	50	$6 \times 10^{-5}$
ref3 (blue dot-dashed)	20	35	11	50	$7 \times 10^{-4}$
ref4 (orange dotted)	20	40	15	100	$6 \times 10^{-4}$

But the reference parameters were different at that time.  $E_R \sim \text{MeV}$

Probably, the U. Chicago group misunderstood that the model cannot explain their preliminary results. :(



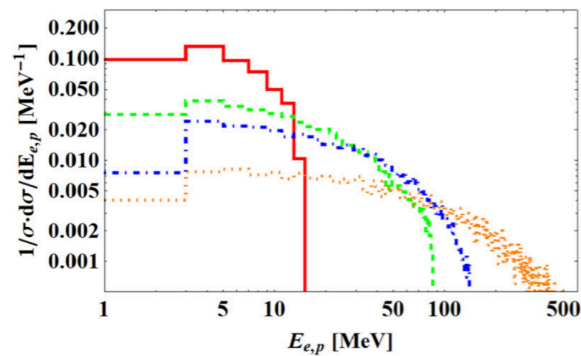
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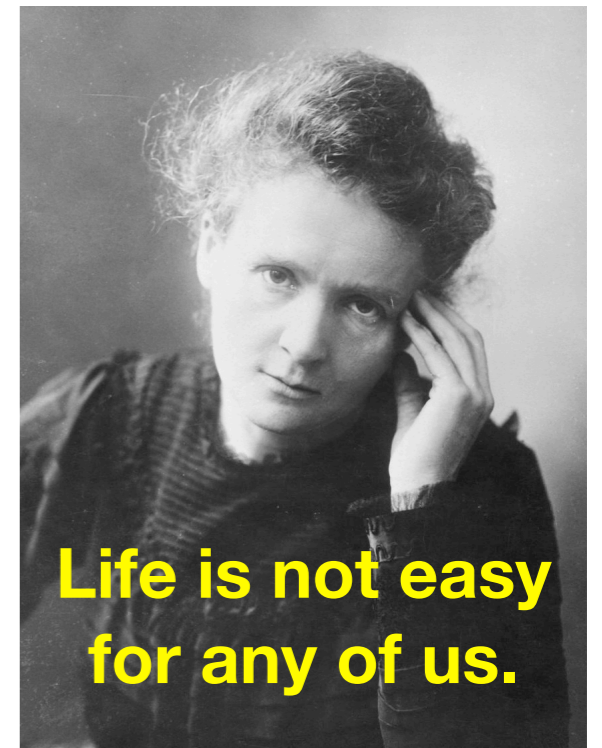
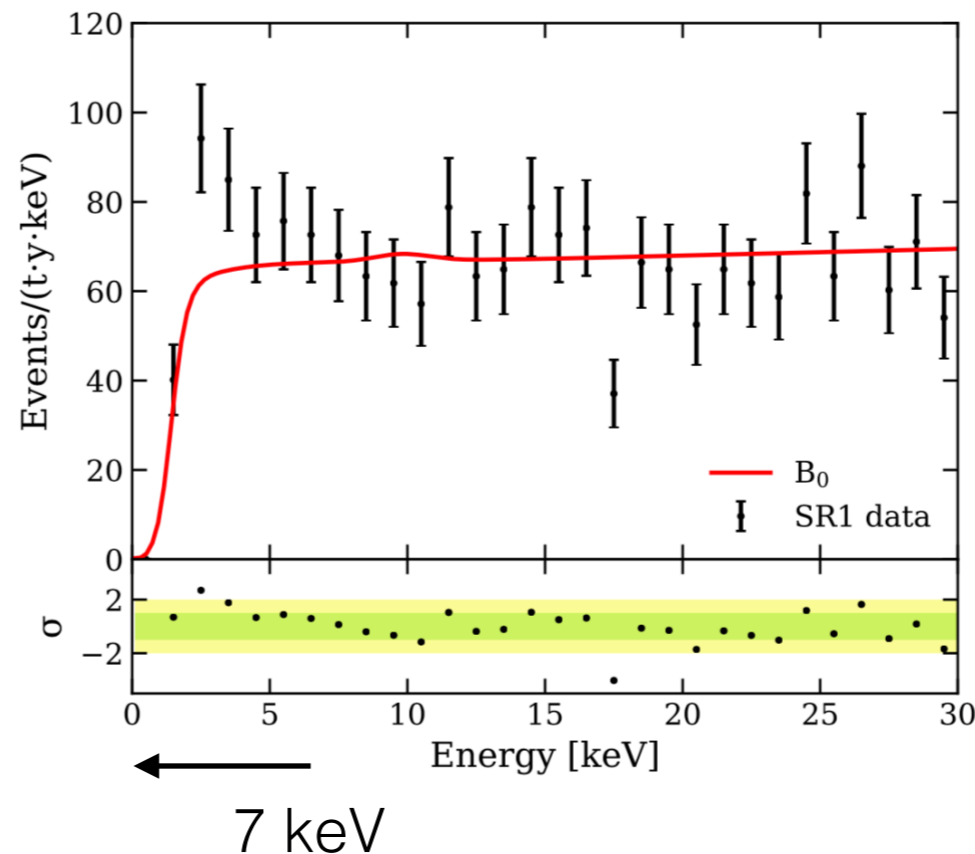
Regardless of the reliability of the excess, further analyses are useful in the sense that

1. XENON-NT can consider our reference models in confirming the excess.
2. we can provide guidelines to future ton-scale experiments, e.g., DarkSide, COSINE-200, in searching for fast-moving DM signals.

# Favored parameter region

Not so simple task (some authors had to update their original papers)

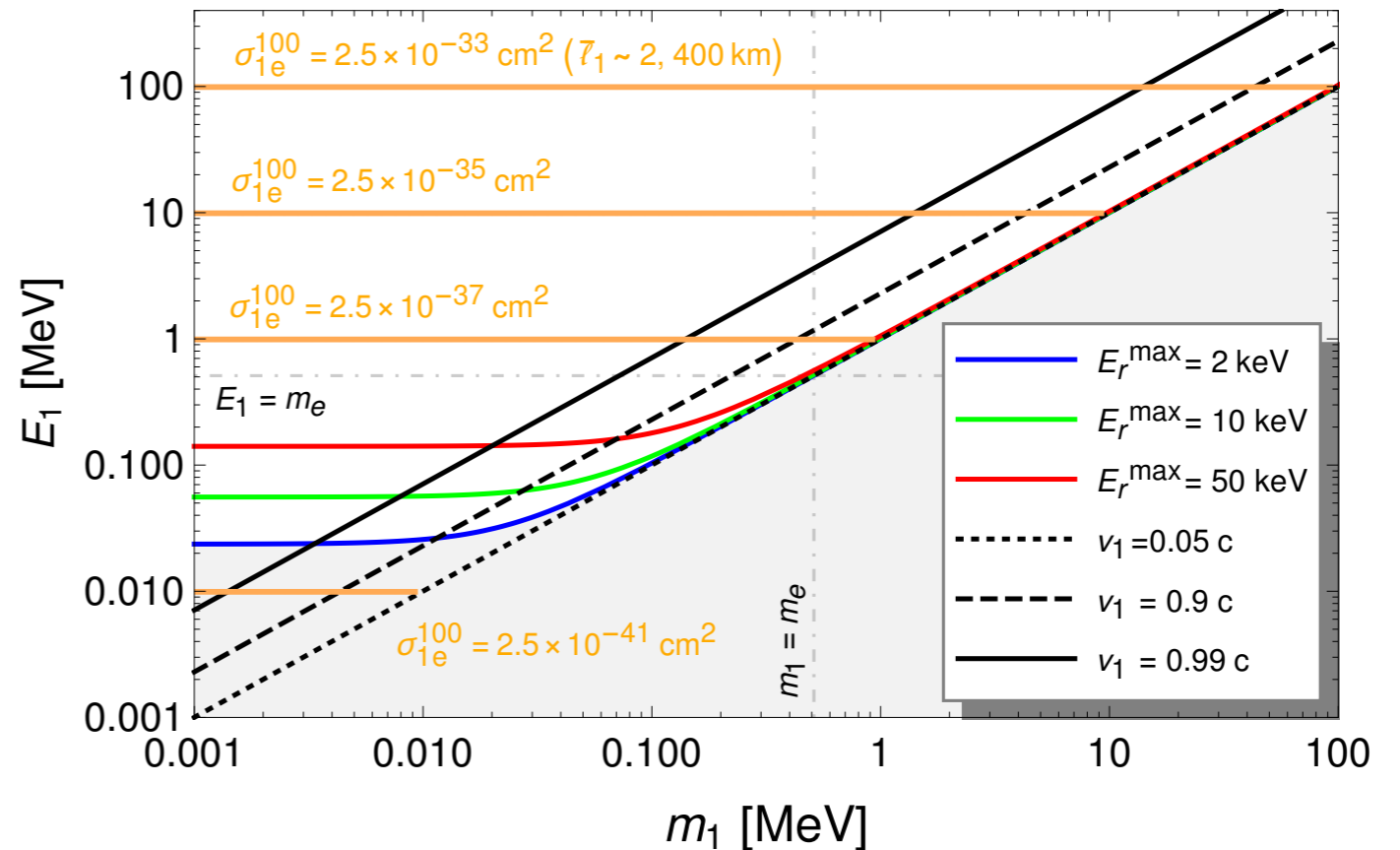
- Large number of events: large cross section with the material of Earth (deflected and loose energy)
- A narrow range of  $2 \text{ keV} \leq E_R \leq 7 \text{ keV}$  is preferred.
- The binding energy of electrons in Xe is not negligible. (for  $E_R \sim \text{keV}$ )



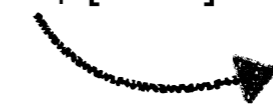
# Favored parameter region

Alhazmi, Kim, Kong, Mohlabeng,  
Park, **SS**, arXiv: 2006.16252

Energy of  
the incoming BDM



$m_1$  [MeV]



Mass of BDM

# Favored parameter region

Alhazmi, Kim, Kong, Mohlabeng,  
Park, **SS**, arXiv: 2006.16252

Energy of  
the incoming BDM

Two-component BDM from GC

$$\mathcal{F}_1 = 1.6 \text{ cm}^{-2} \text{ s}^{-1}$$

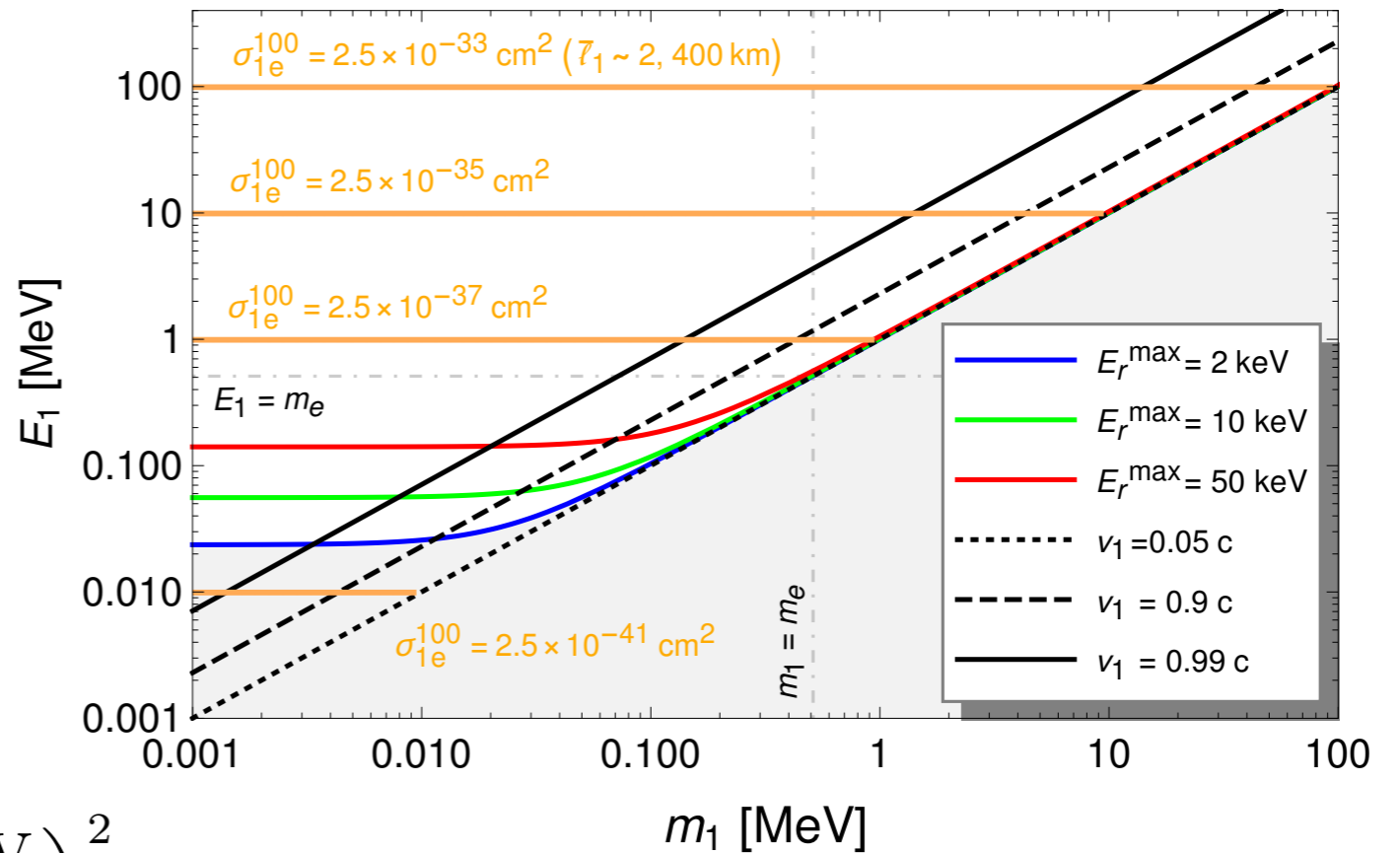
$$\times \left( \frac{\langle \sigma_{0 \rightarrow 1 \nu} \rangle}{5 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}} \right) \left( \frac{10 \text{ MeV}}{m_0} \right)^2$$

$\chi_1 - e$  (fiducial) scattering  
cross section

$$N_{\text{sig}} = \mathcal{F}_1 \sigma_{1e} N_e^{\text{eff}} t_{\text{exp}}$$

Number of effective target electrons  
in the fiducial volume

Total exposure time

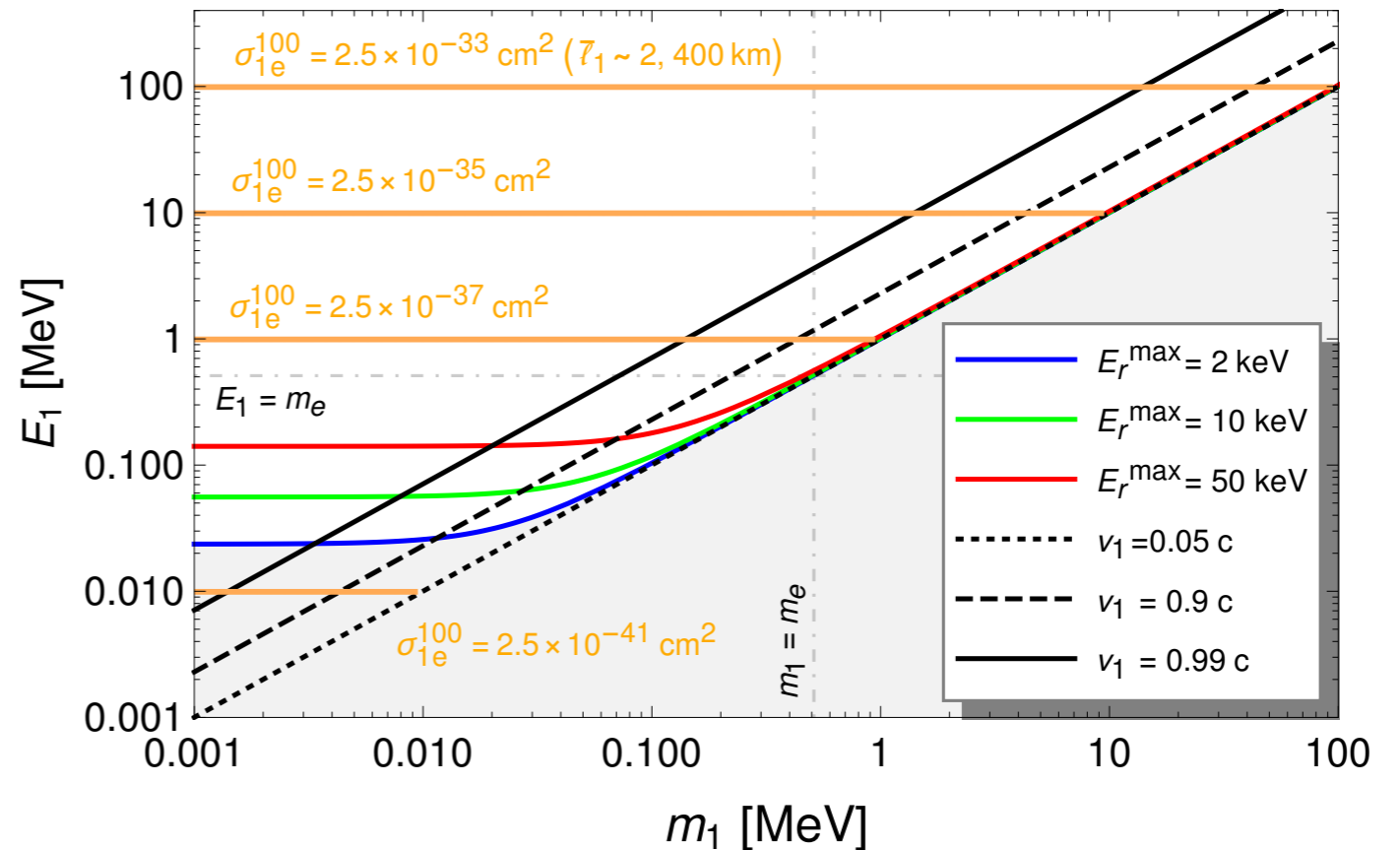


Mass of BDM

# Favored parameter region

Alhazmi, Kim, Kong, Mohlabeng,  
Park, **SS**, arXiv: 2006.16252

Energy of  
the incoming BDM



$$E_r^{\max} = \frac{2m_e p_1^2}{s} = \frac{E_1^2 - m_1^2}{m_1^2 + m_e^2 + 2m_e E_1}$$

Mass of BDM

- The maximum  $E_R$  of electrons scattered by BDM  $\geq 2$  keV (non-shaded). This is **model independently** given as above.
- $E_1 \gtrsim 20$  keV is preferred (depending on  $m_1$ ).

# Favored parameter region

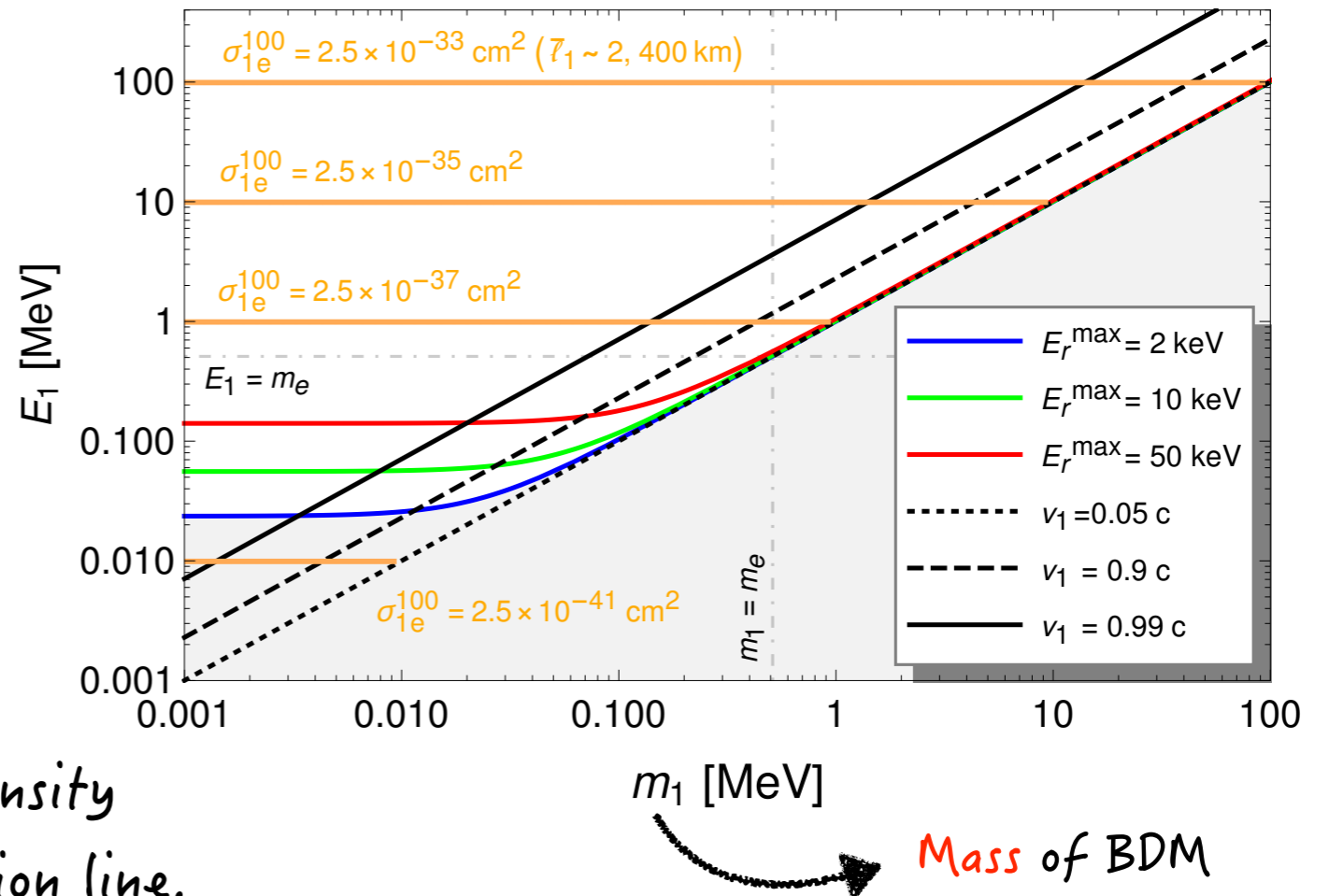
Alhazmi, Kim, Kong, Mohlabeng,  
Park, **SS**, arXiv: 2006.16252

Energy of  
the incoming BDM

$E_1$  [MeV]

$$\bar{l}_1 \sim \frac{1}{\langle n_e \rangle \sigma_{1e}}$$

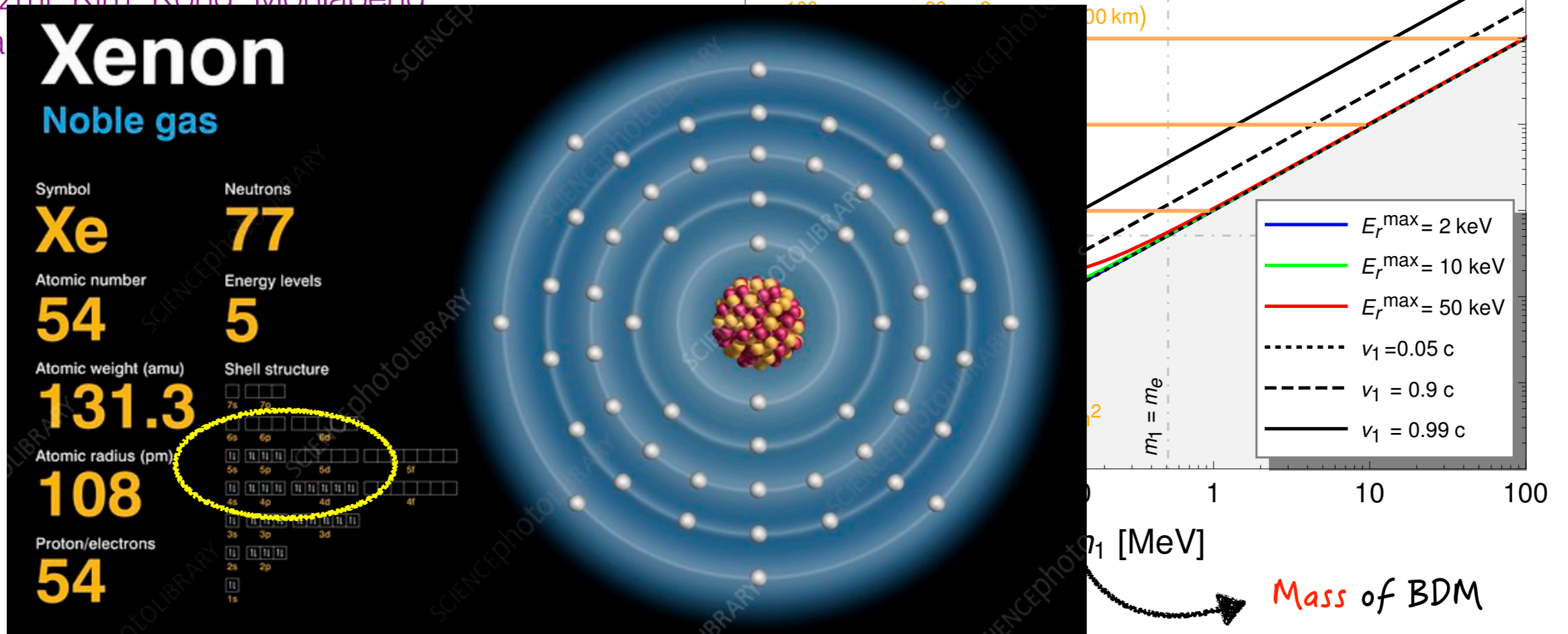
Mean  $e$ -number density  
along the  $\chi_1$  propagation line.



- The values of (fiducial) cross section  $\sigma_{1e}$  giving  $N_{\text{sig}} = 100$  are shown, assuming  $N_e^{\text{eff}} = 26$  (will be discussed in the next slide).
- To avoid the Earth attenuation, the mean free path  $\approx \mathcal{O}(1000 \text{ km})$  is preferred.  
(at least larger than the depth of XENON1T  $\sim 1.6 \text{ km}$ )

# Favored parameter region

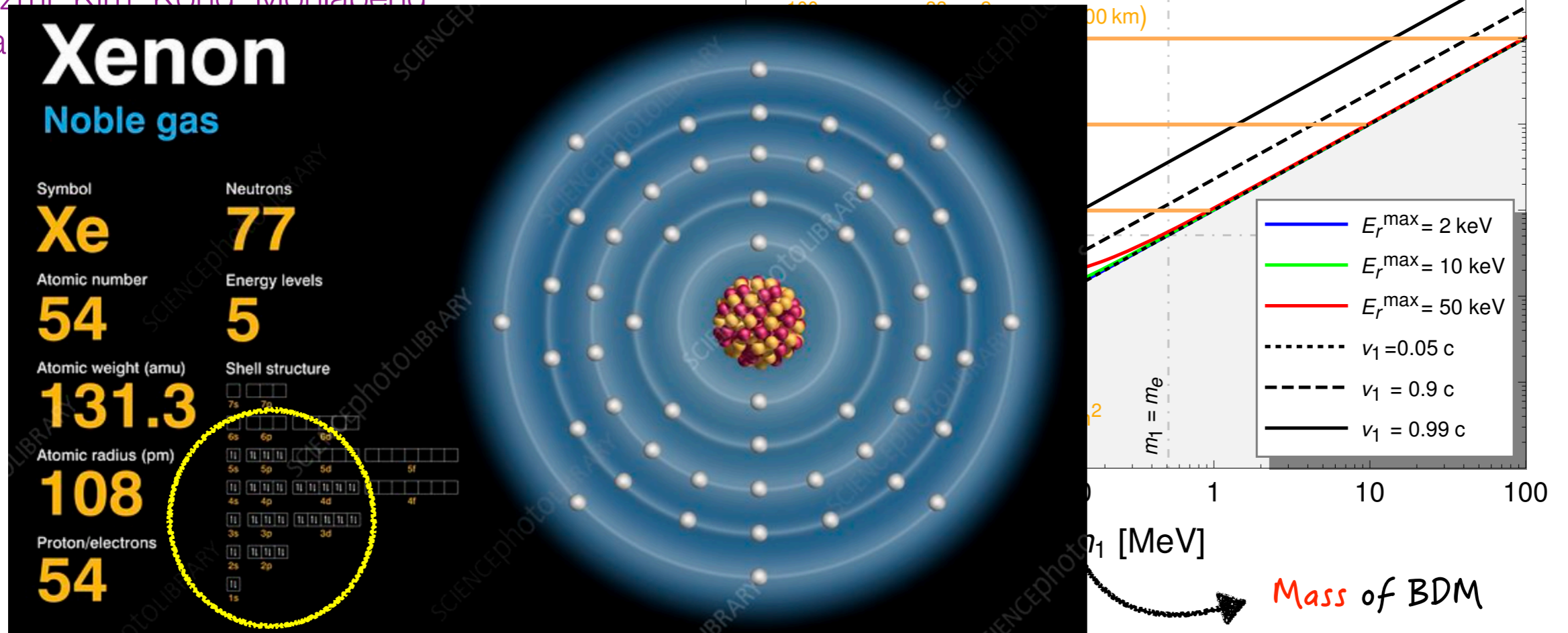
Alhazmi Kim Kong Mohlabeng  
Pa



- The value of  $N_e^{\text{eff}} = 26$  is from naively considering the *O*-shell (outermost, 8) and the *N*-shell (18, binding energy  $\approx 213$  eV  $\sim 10\%$  of the energy resolution).
- In practice, we need to consider ionization form factor from all the shells.  
(Do not significantly change the results but it depends on  $E_1$  and mediators.)

# Favored parameter region

Alhazmi, Kim, Kong, Mohlabeng, Park



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*Preliminary* Work in progress with Alhazmi, Kim, Kong, Mohlabeng, Park, **SS**

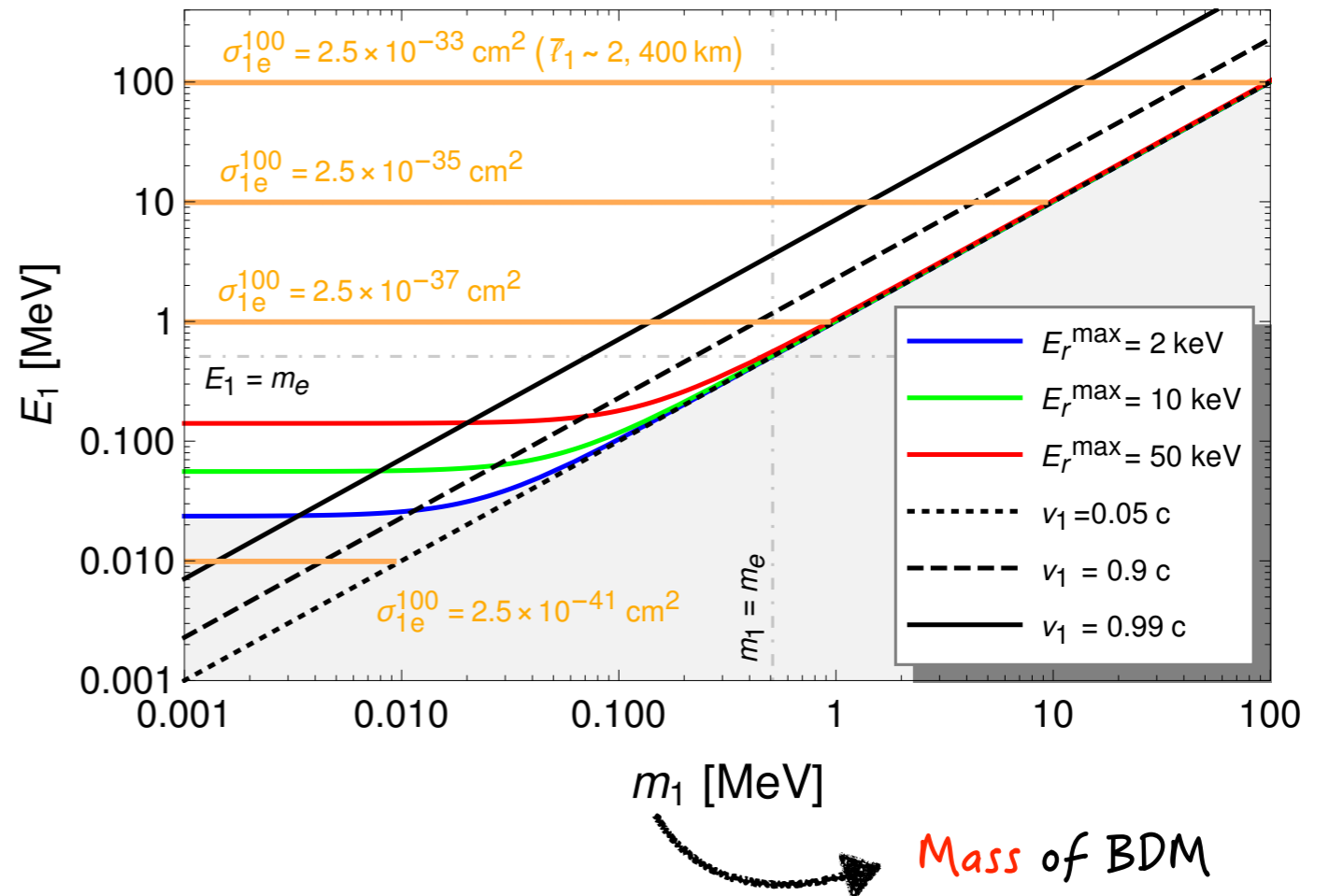
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# Favored parameter region

Alhazmi, Kim, Kong, Mohlabeng,  
Park, **SS**, arXiv: 2006.16252

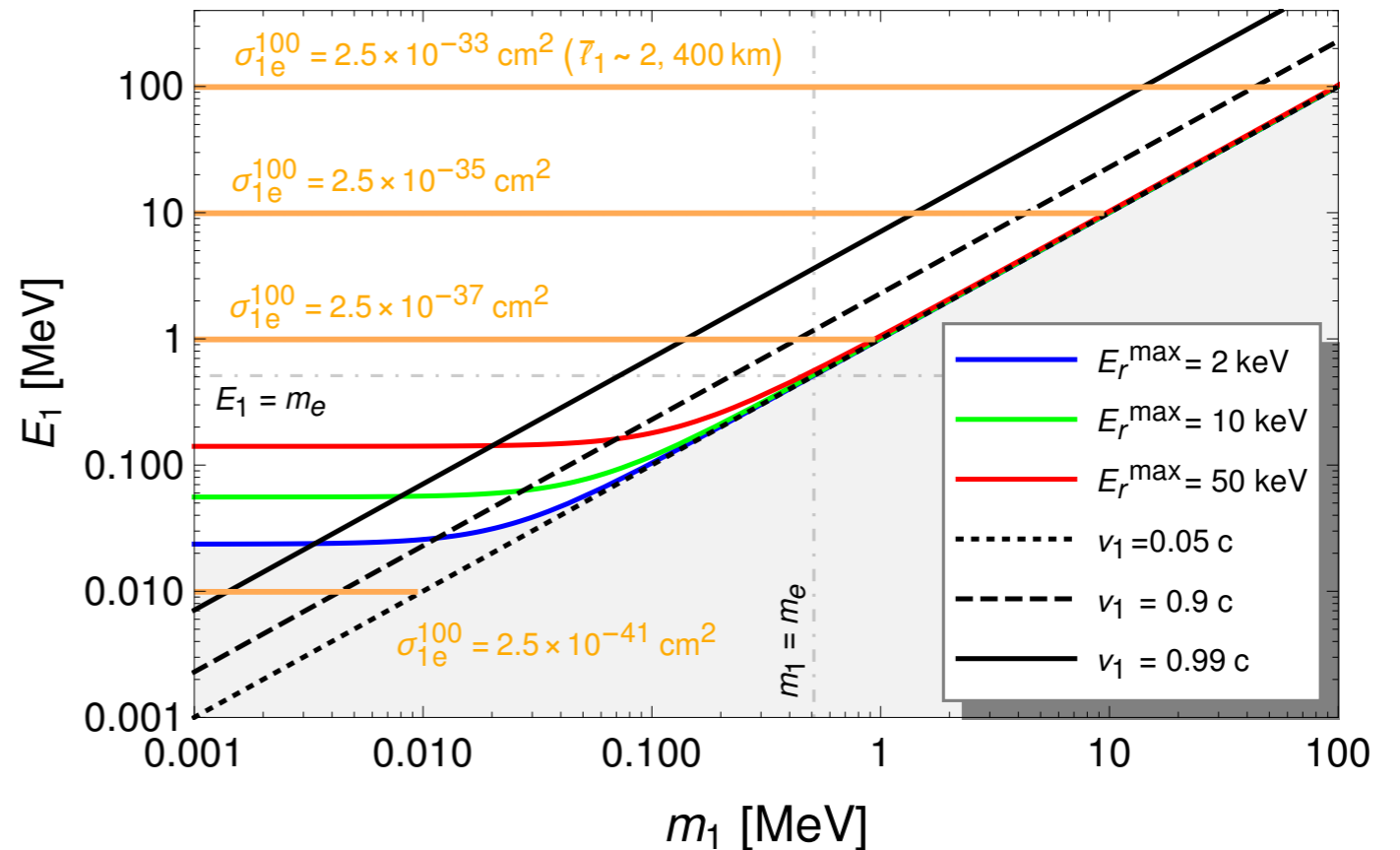
Energy of  
the incoming BDM



- The velocity of BDM,  $v_1$ , can be close to  $c$  in a wide range of parameter space ( $\gg 0.1c$  is also preferred).
- Shade regions and the black lines are model independent while the orange lines are applied for conventional BDM (but readily applicable).

# Dependence on mediators

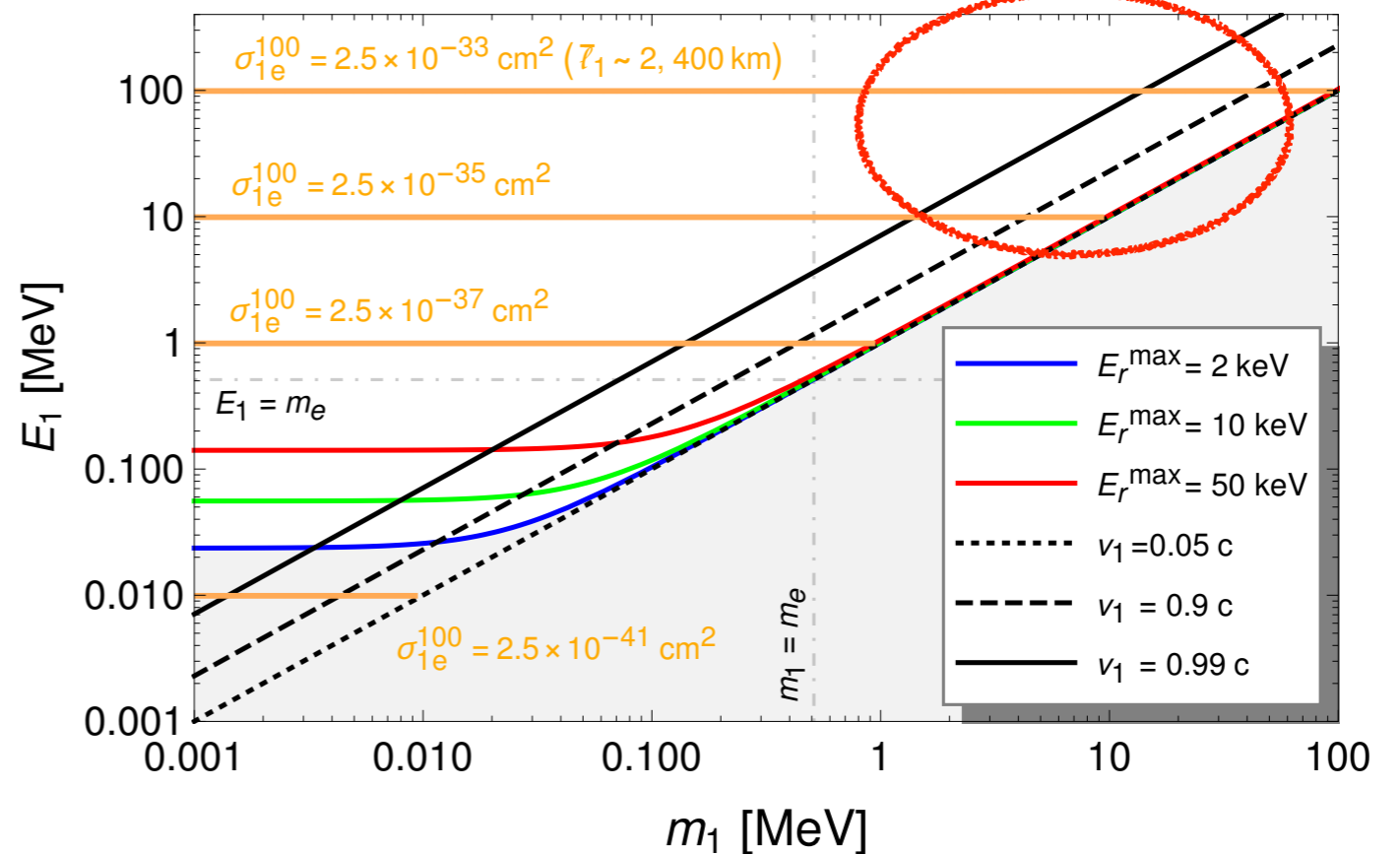
mediator	DM
V	F
V	S
P	F
P	S
S	F
S	S



- The scales of mass and coupling parameters preferred by the excess depend on the type of the mediators.
- We analyze the shape of the spectrum for various types of mediators (vector:V, pseudoscalar: P, scalar: S) and DM (fermion: F, scalar: S).
- Three reference parameter regions are chosen.

# Dependence on mediators

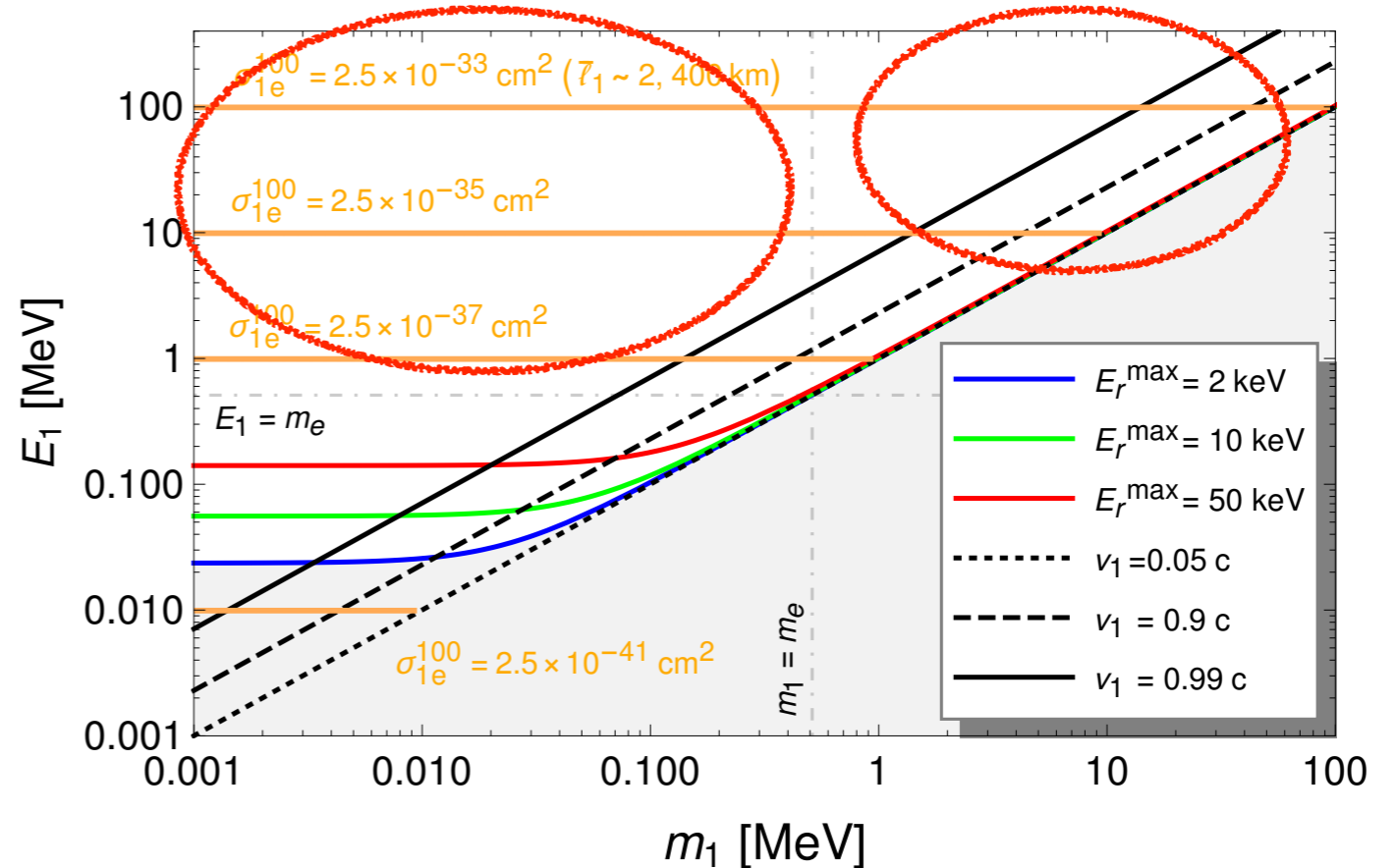
mediator	DM
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  - (i)  $E_1 \approx m_1 \gg m_e, m_i \gg m_e$
  - (ii)  $E_1 \approx m_1 \gg m_e, m_i < m_e$
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# Dependence on mediators

mediator	DM
V	F
V	S
P	F
P	S
S	F
S	S



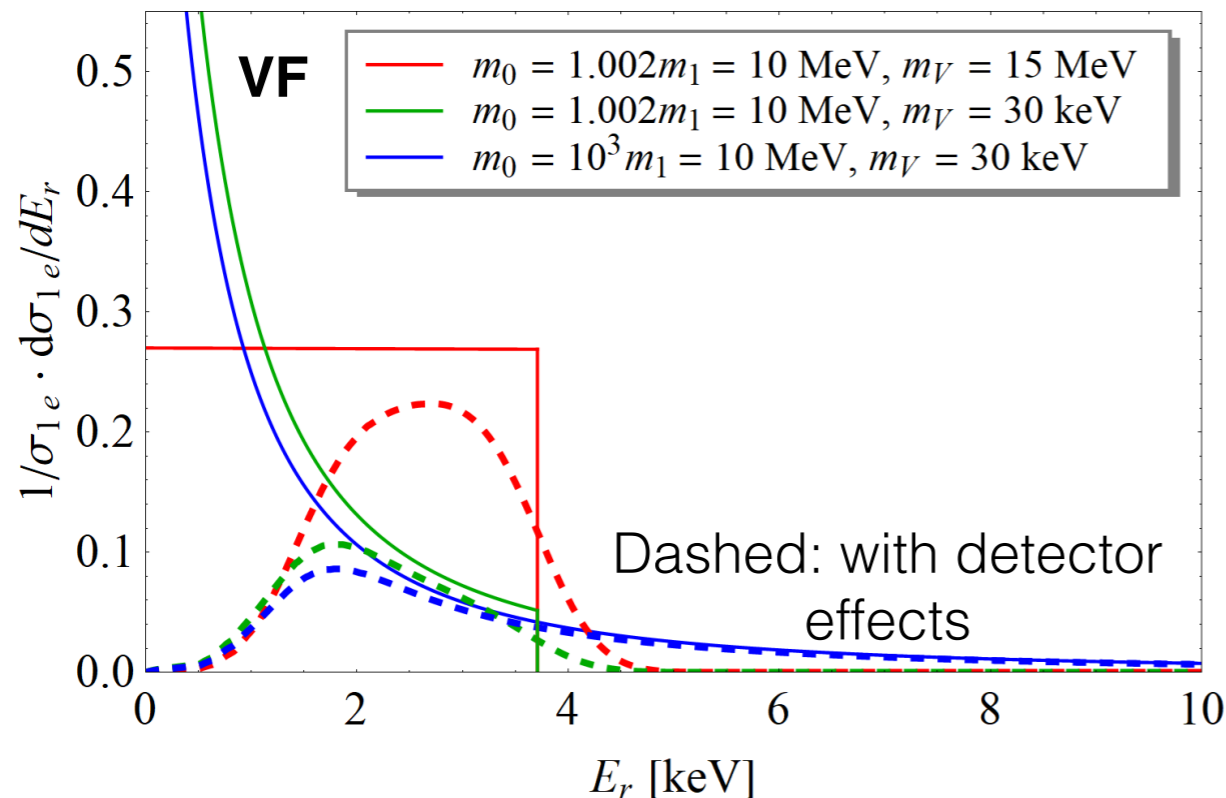
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  - (iii)  $E_1 \gg m_e > m_1, m_i < m_e$
- Three reference parameter regions are chosen.

# Dependence on mediators

$$\frac{d\sigma_{1e}}{dE_r} = \frac{(g_j^i g_e^i)^2}{8\pi\lambda(s, m_e^2, m_1^2)(2m_e E_r + m_i^2)^2} |\bar{\mathcal{A}}|^2 \quad \text{without detector resolution and efficiency}$$

Case	Mediator	Dark matter	$\mathcal{L}_{\text{int}}$	$ \bar{\mathcal{A}} ^2$
VF	$V_\mu$	$\chi_1$	$(g_e^V \bar{e}\gamma^\mu e + g_\chi^V \bar{\chi}_1 \gamma^\mu \chi_1) V_\mu$	$8m_e \{m_e(2E_1^2 - 2E_1 E_r + E_r^2) - (m_e^2 + m_1^2)E_r\}$
VS	$V_\mu$	$\varphi_1$	$(g_e^V \bar{e}\gamma^\mu e + g_\varphi^V \varphi_1^* \partial^\mu \varphi_1 + \text{h.c.}) V_\mu$	$8m_e \{2m_e E_1(E_1 - E_r) - m_1^2 E_r\}$
PF	$a$	$\chi_1$	$(ig_e^a \bar{e}\gamma^5 e + ig_\chi^a \bar{\chi}_1 \gamma^5 \chi_1) a$	$4m_e^2 E_r^2$
PS	$a$	$\varphi_1$	$(ig_e^a \bar{e}\gamma^5 e + ig_\varphi^a m_1 \varphi^* \varphi) a$	$8m_e m_1^2 E_r$
SF	$\phi$	$\chi_1$	$(g_e^\phi \bar{e}e + g_\chi^\phi \bar{\chi}_1 \chi_1) \phi$	$4m_e(E_r + 2m_e)(2m_1^2 + m_e E_r)$
SS	$\phi$	$\varphi_1$	$(g_e^\phi \bar{e}e + g_\varphi^\phi m_1 \varphi^* \varphi) \phi$	$8m_e m_1^2 (E_r + 2m_e)$

$$\text{VF} \quad \text{(i)} \quad \frac{d\sigma_{1e}}{dE_r} \propto \frac{m_e m_1^2}{m_V^4} \quad \text{(ii)} \quad \frac{d\sigma_{1e}}{dE_r} \propto \frac{m_e m_1^2}{(2m_e E_r + m_V^2)^2} \quad \text{(iii)} \quad \frac{d\sigma_{1e}}{dE_r} \propto \frac{m_e E_1^2}{(2m_e E_r + m_V^2)^2}$$

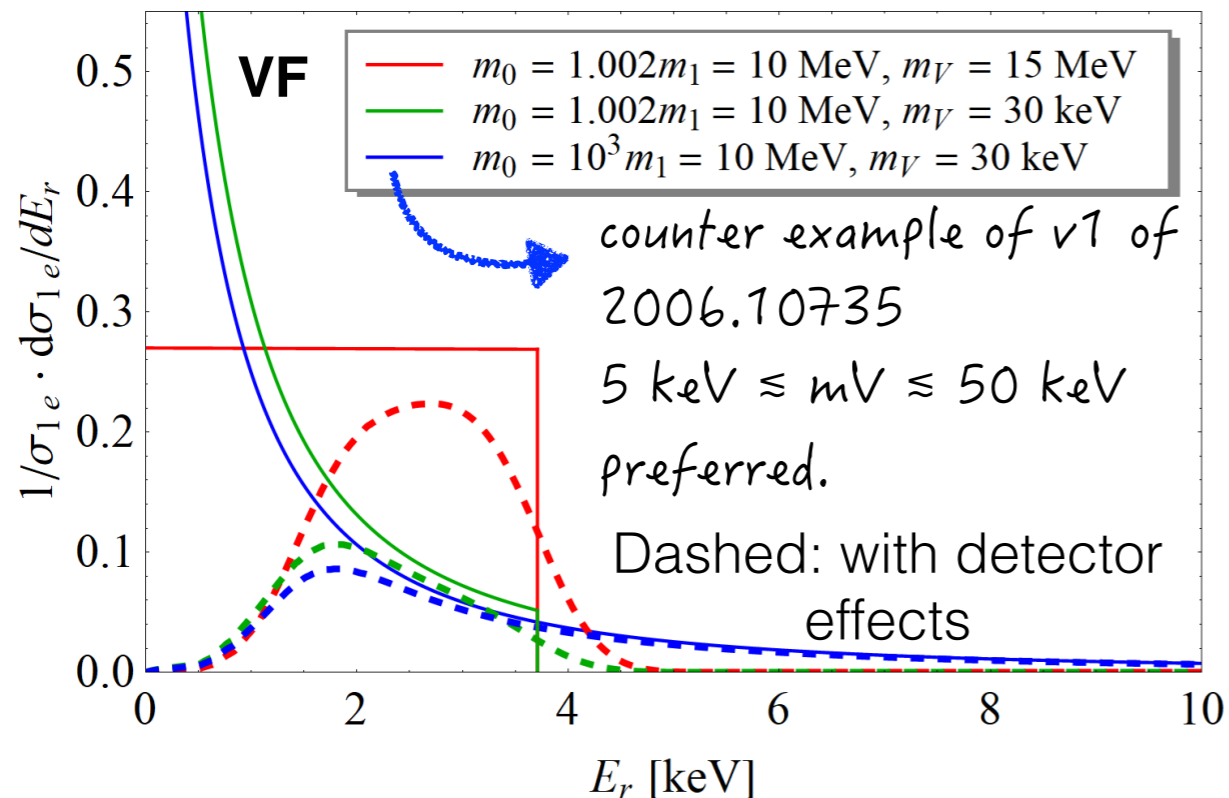


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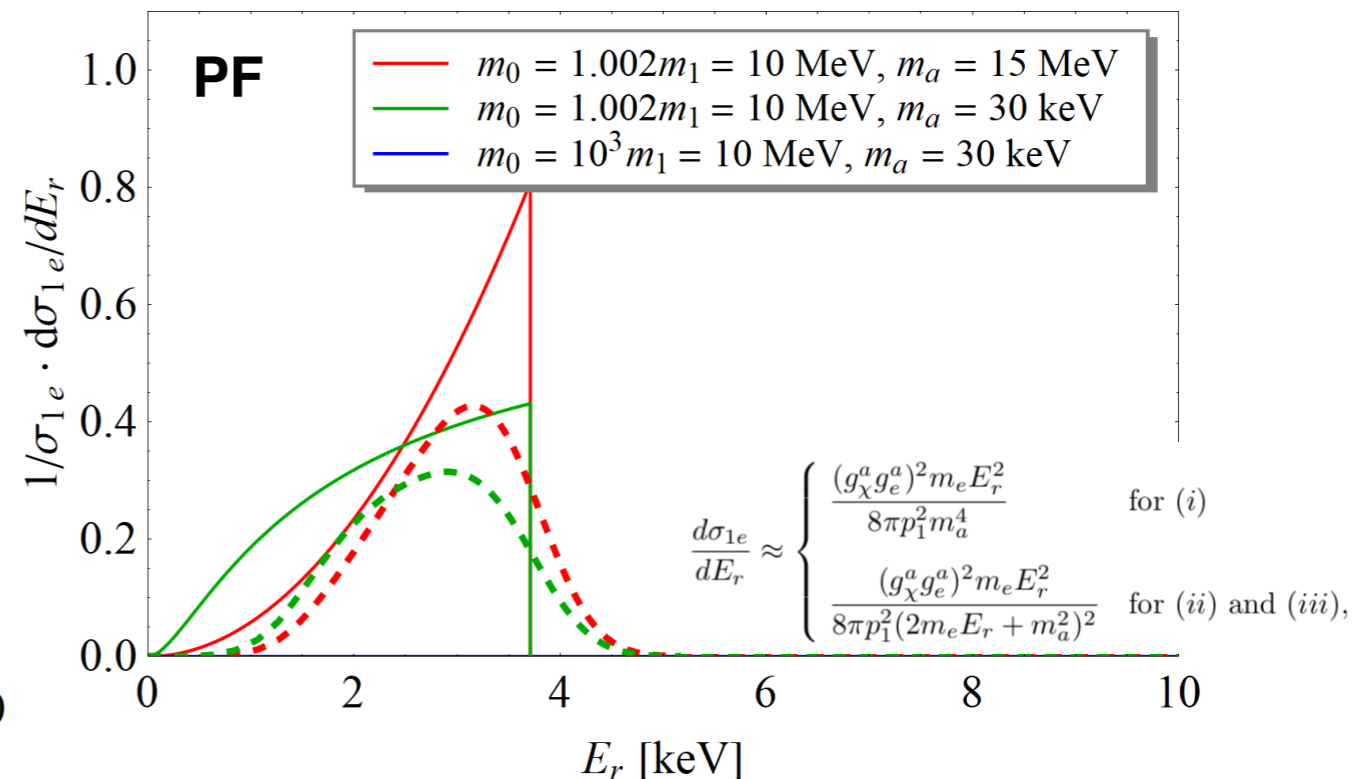
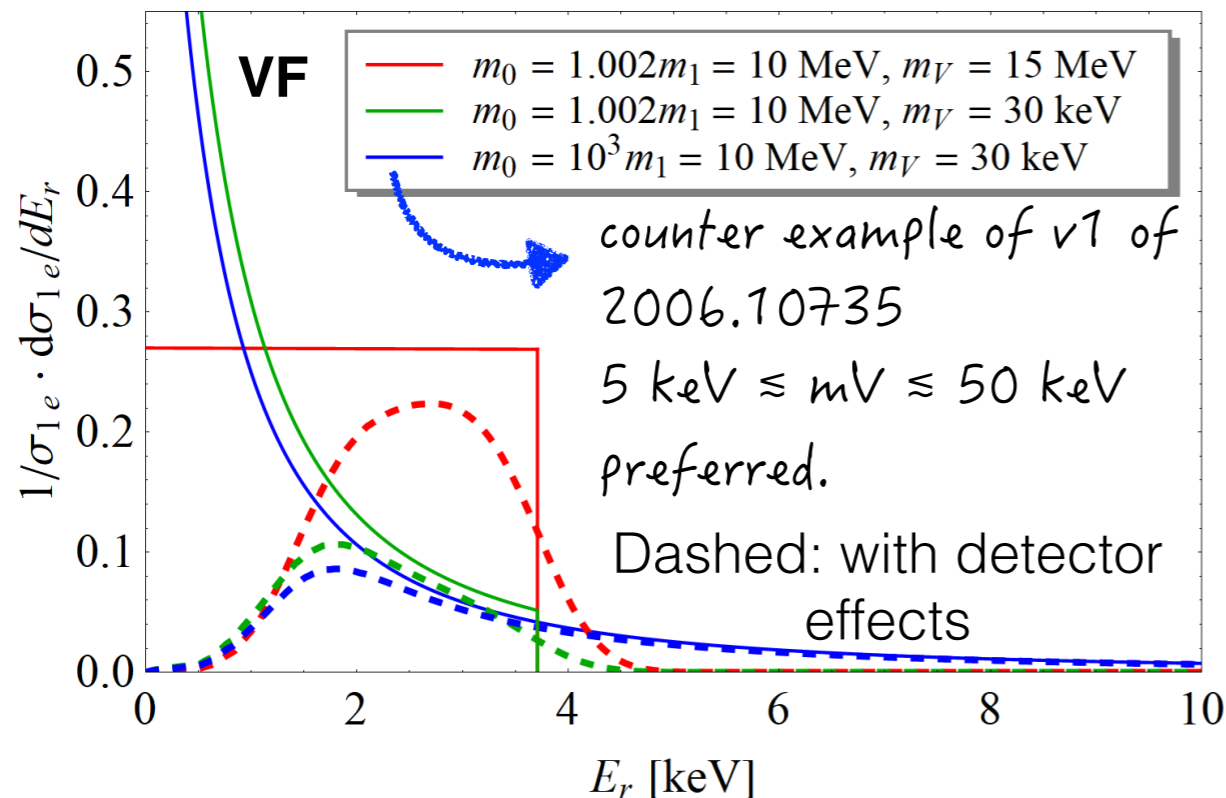


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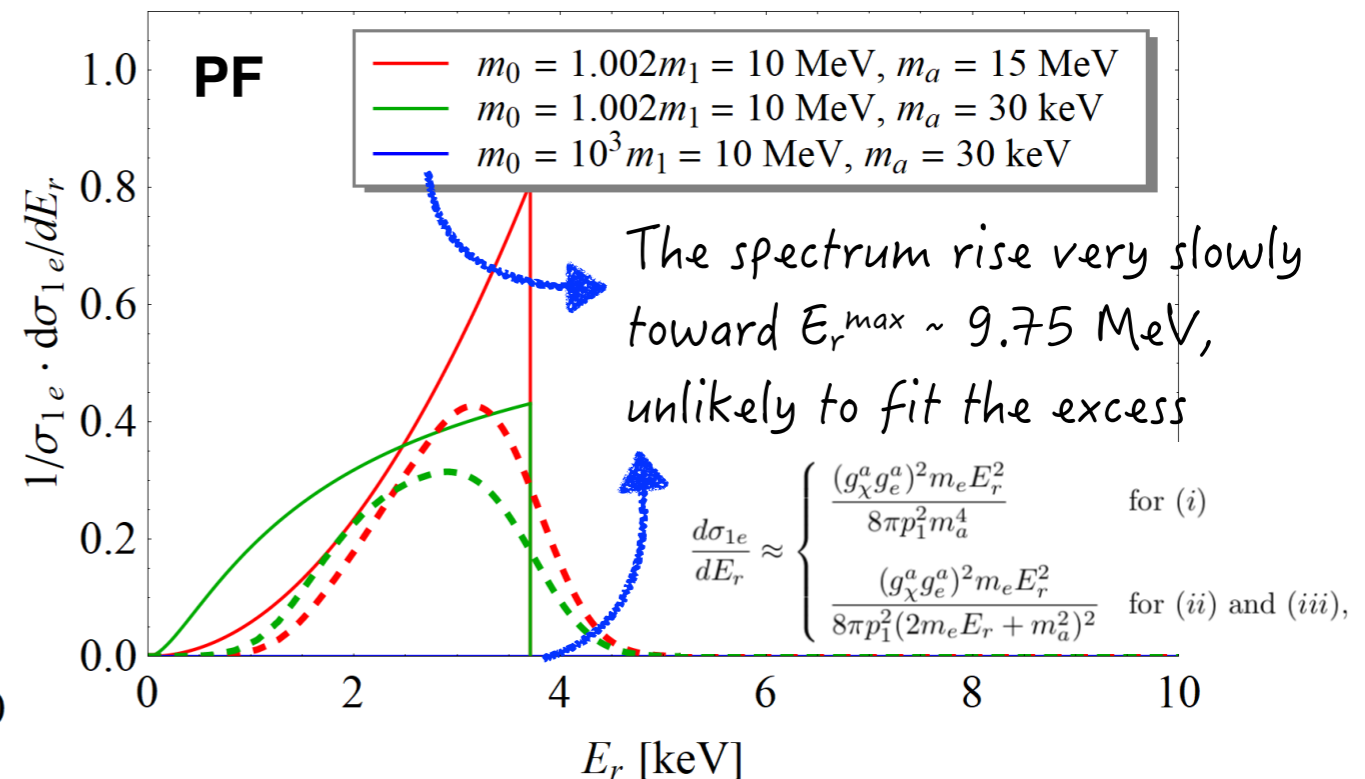
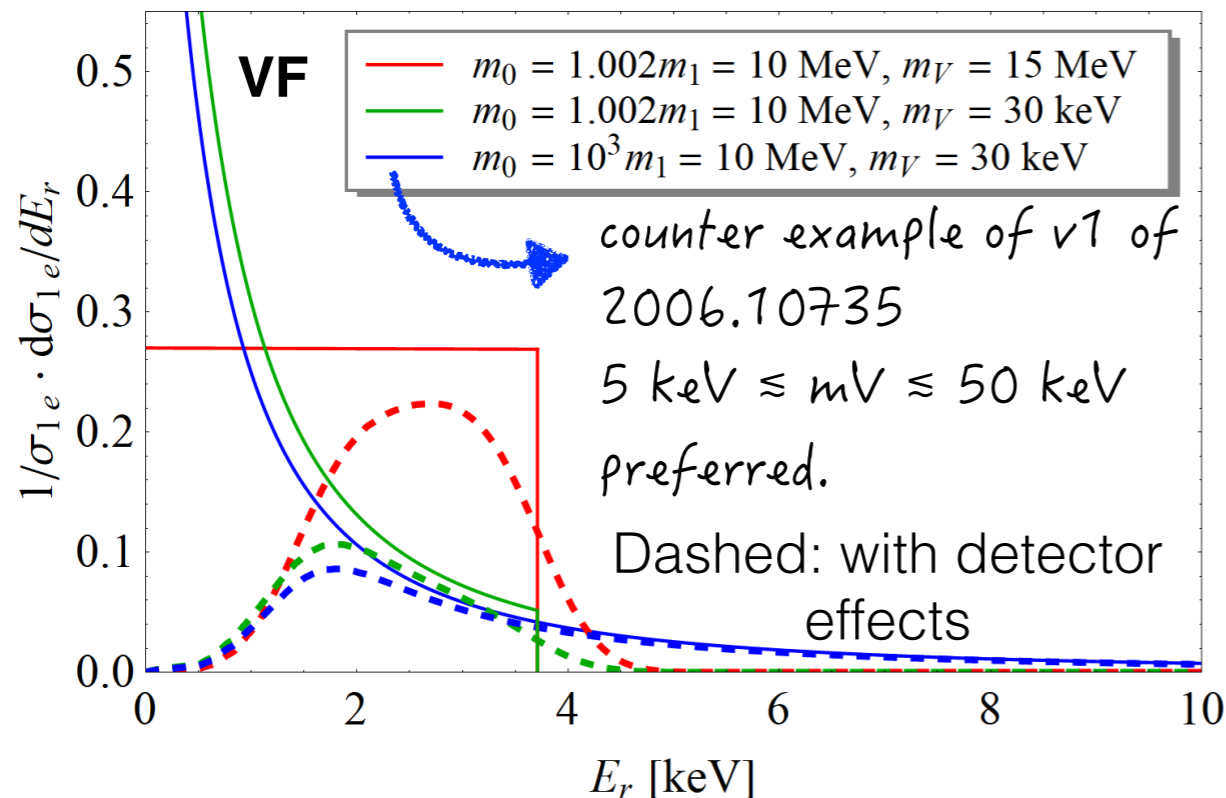


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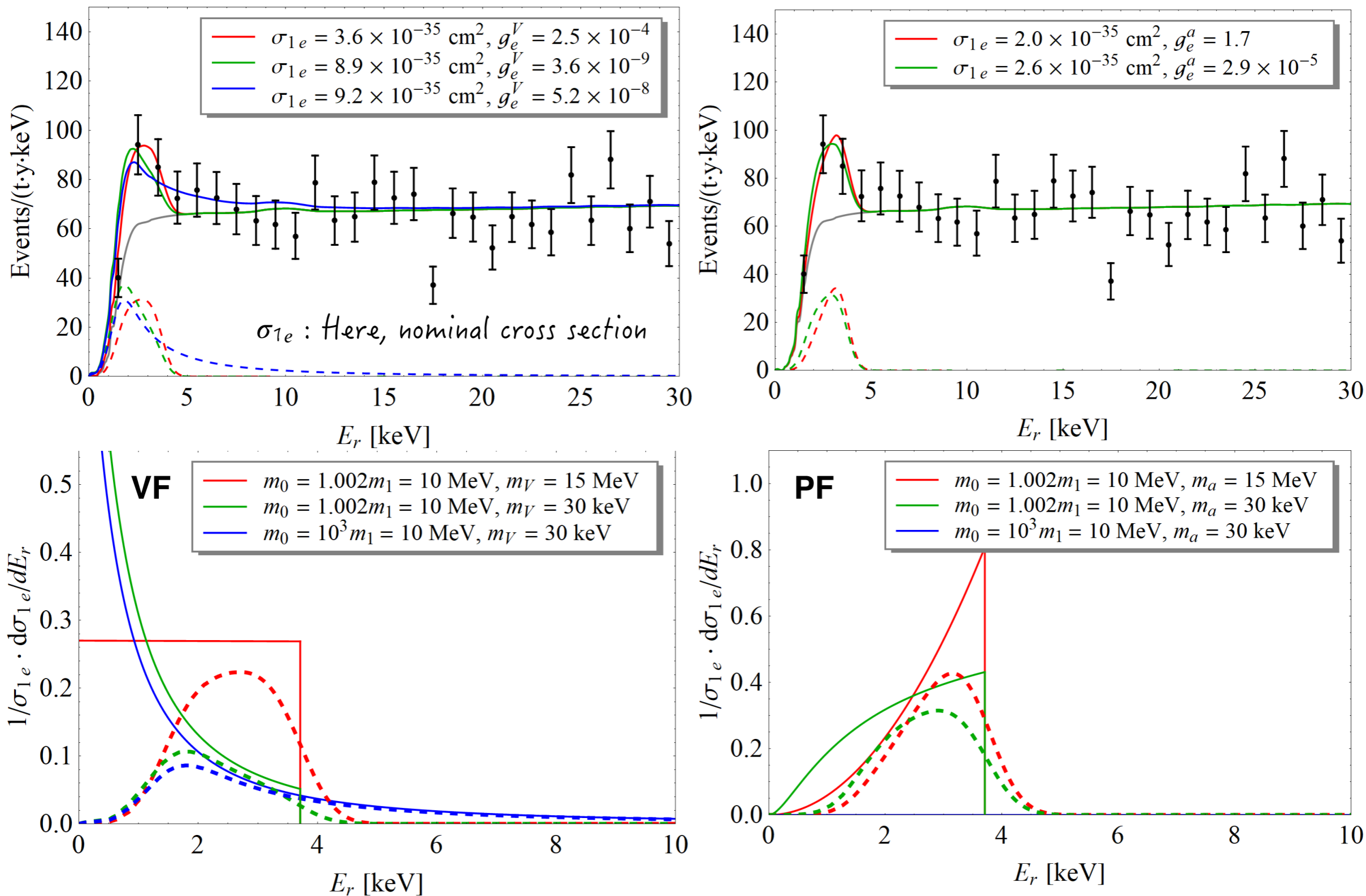
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VF (i)  $\frac{d\sigma_{1e}}{dE_r} \propto \frac{m_e m_1^2}{m_V^4}$  (ii)  $\frac{d\sigma_{1e}}{dE_r} \propto \frac{m_e m_1^2}{(2m_e E_r + m_V^2)^2}$  (iii)  $\frac{d\sigma_{1e}}{dE_r} \propto \frac{m_e E_1^2}{(2m_e E_r + m_V^2)^2}$

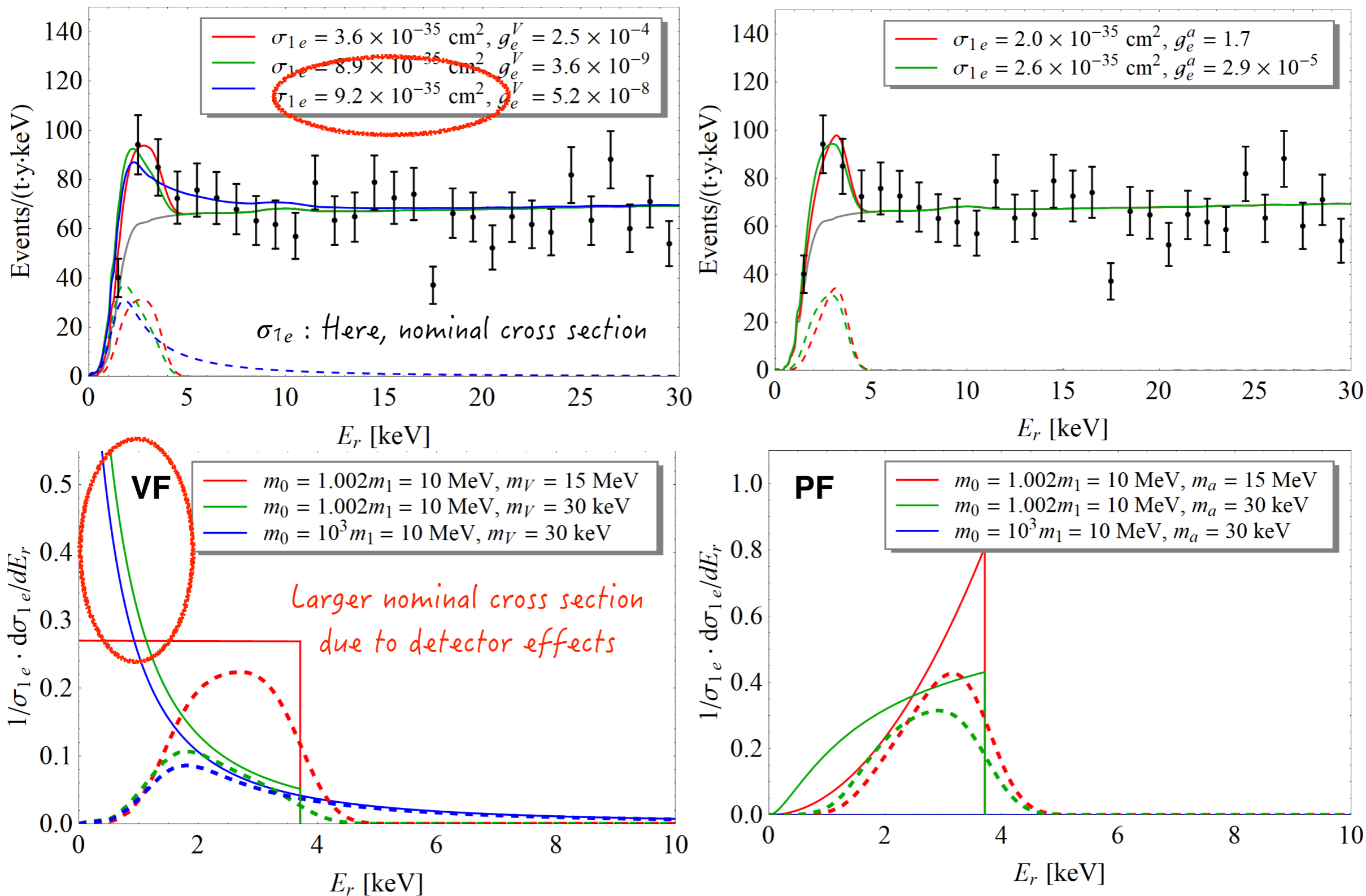




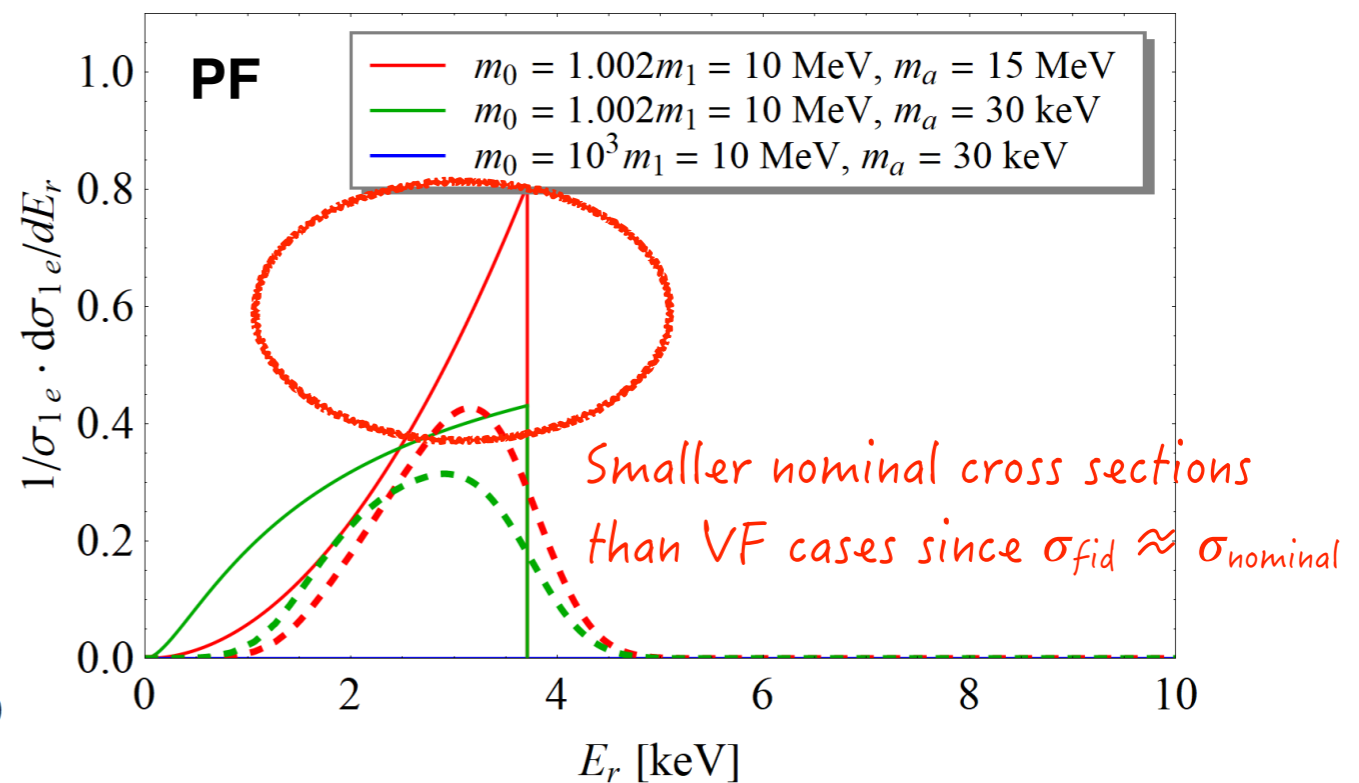
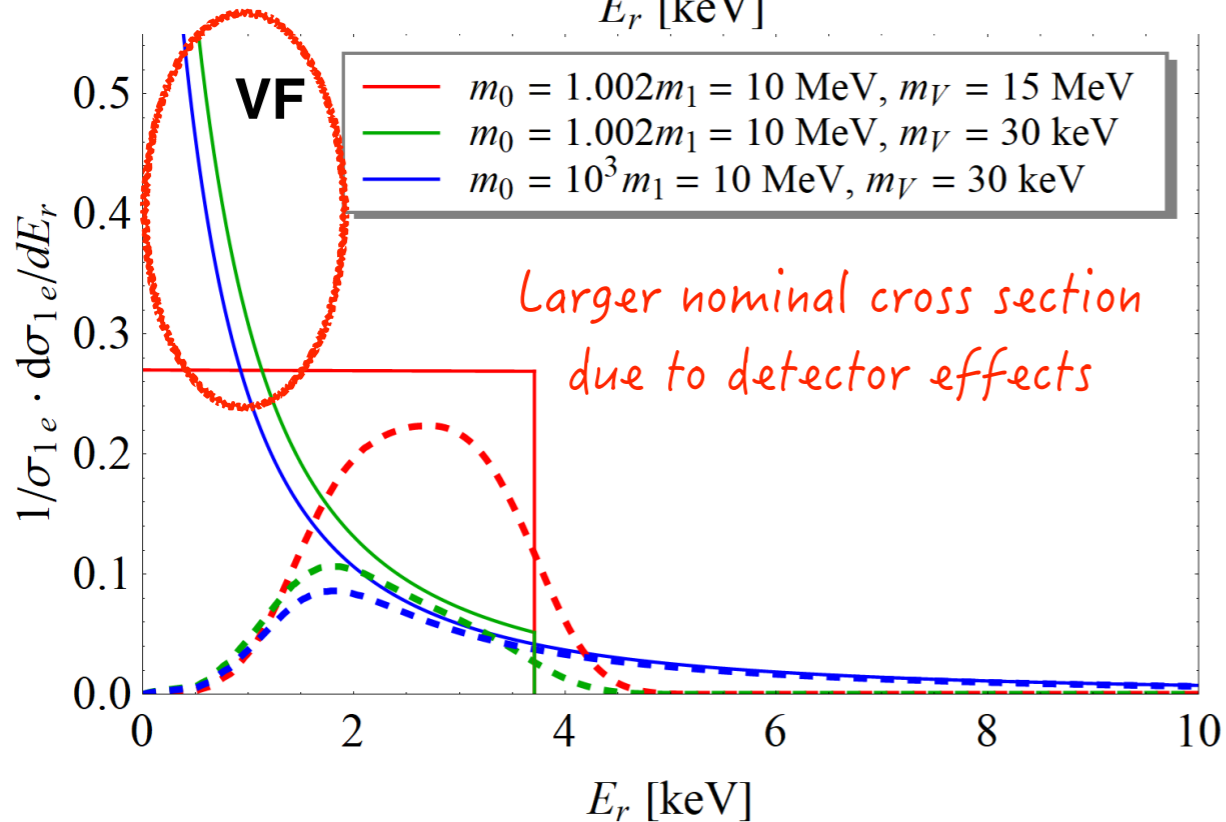
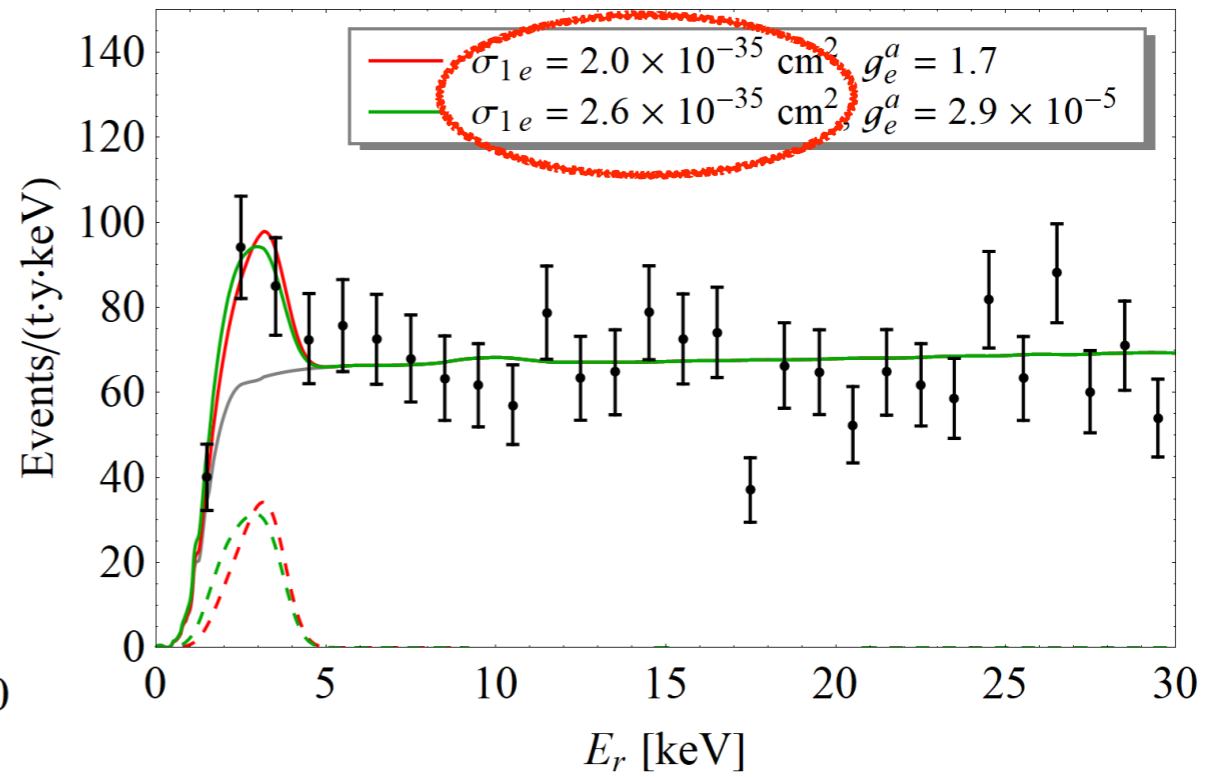
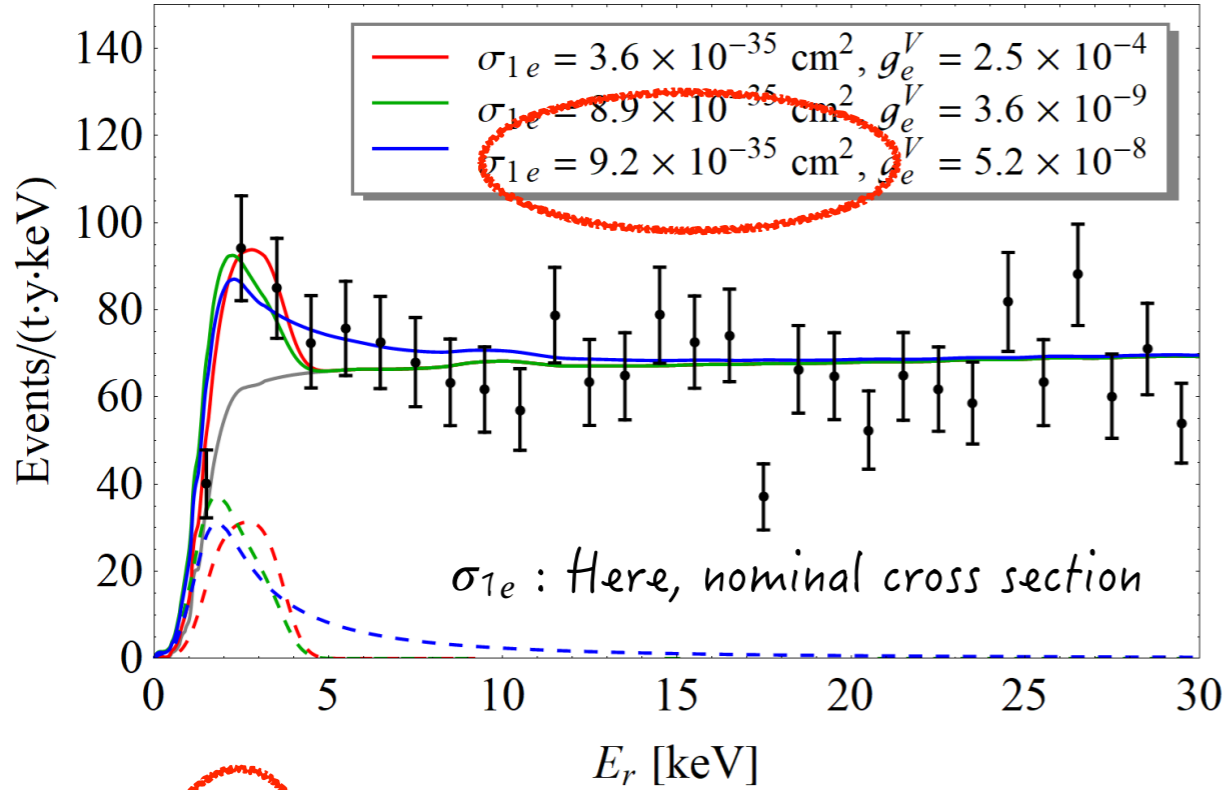
# Fit to the excess



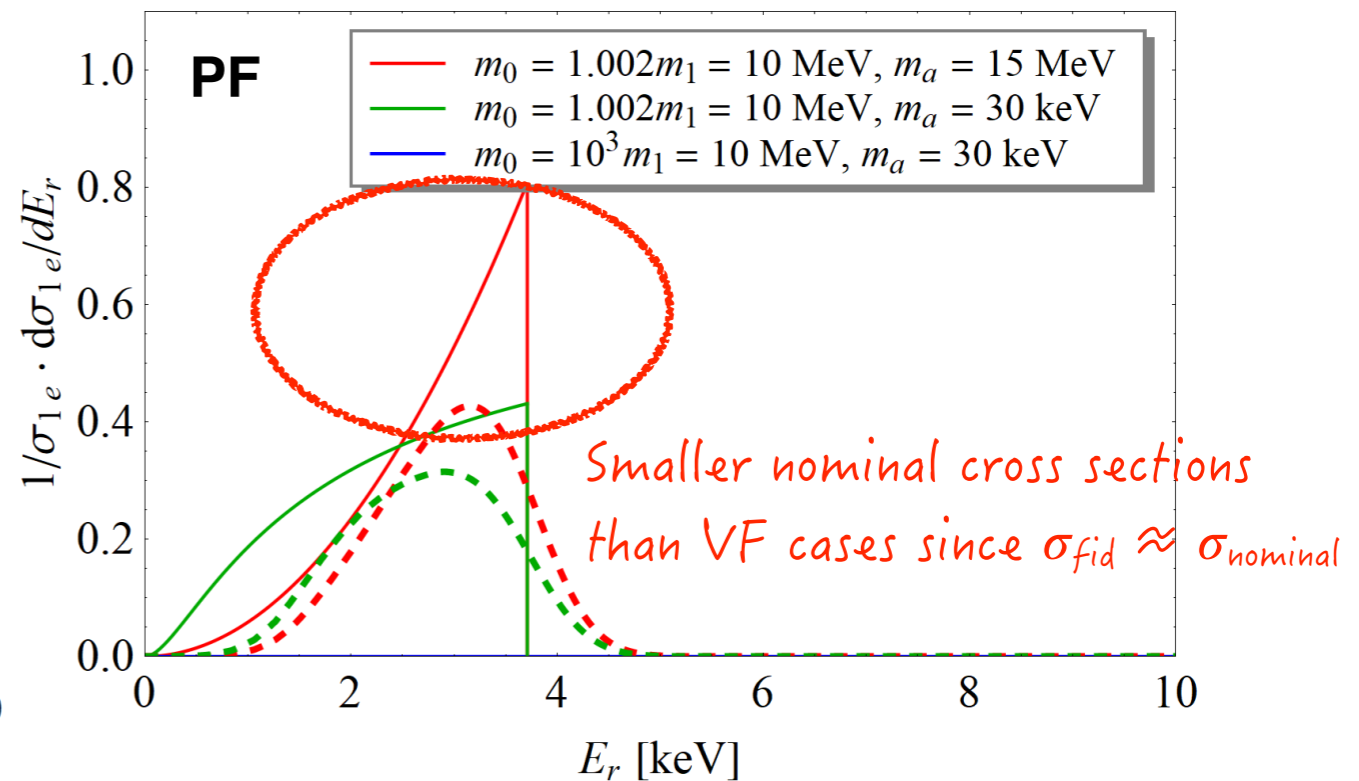
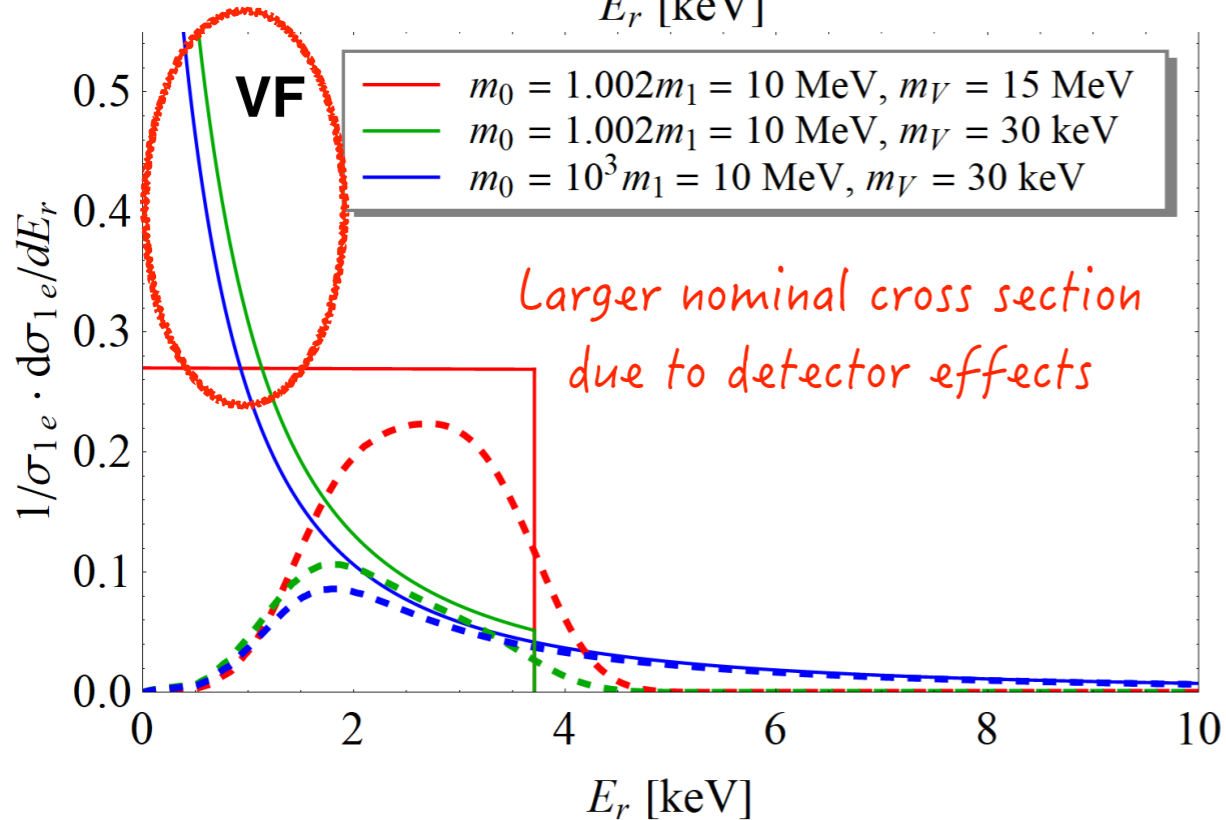
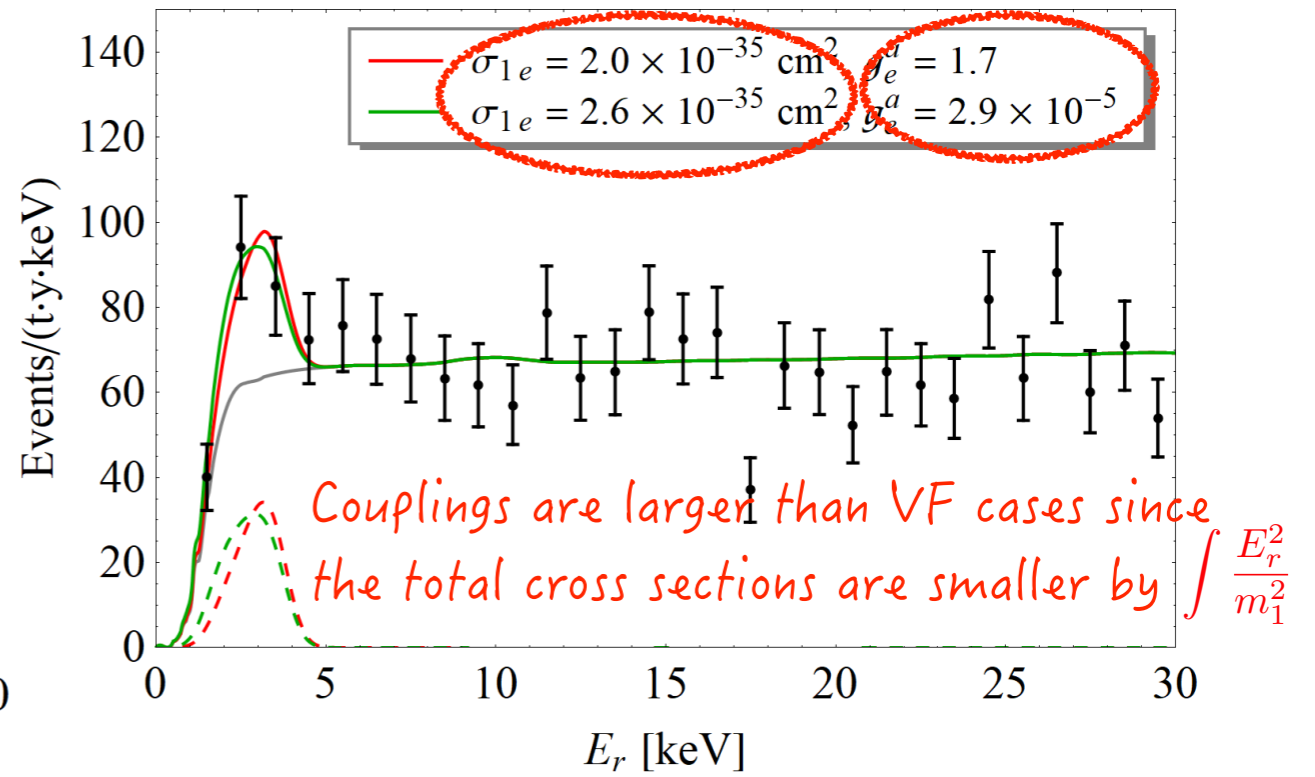
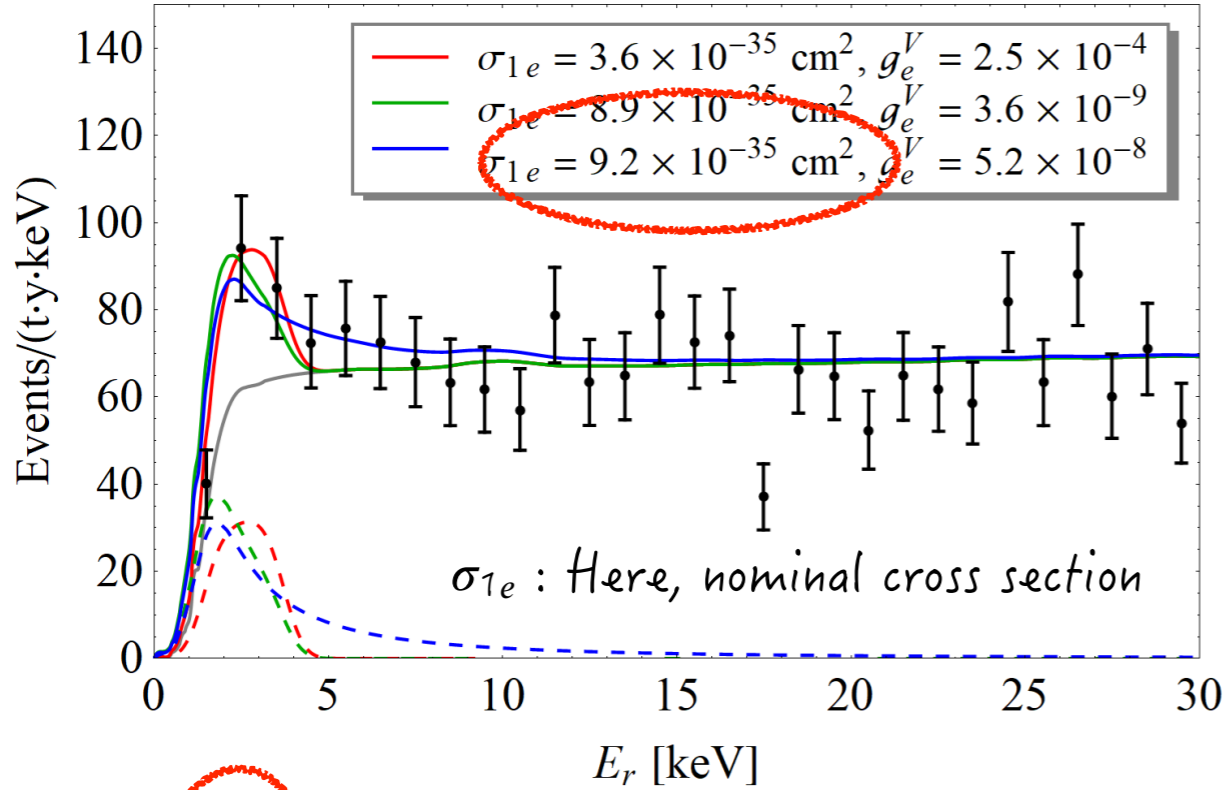
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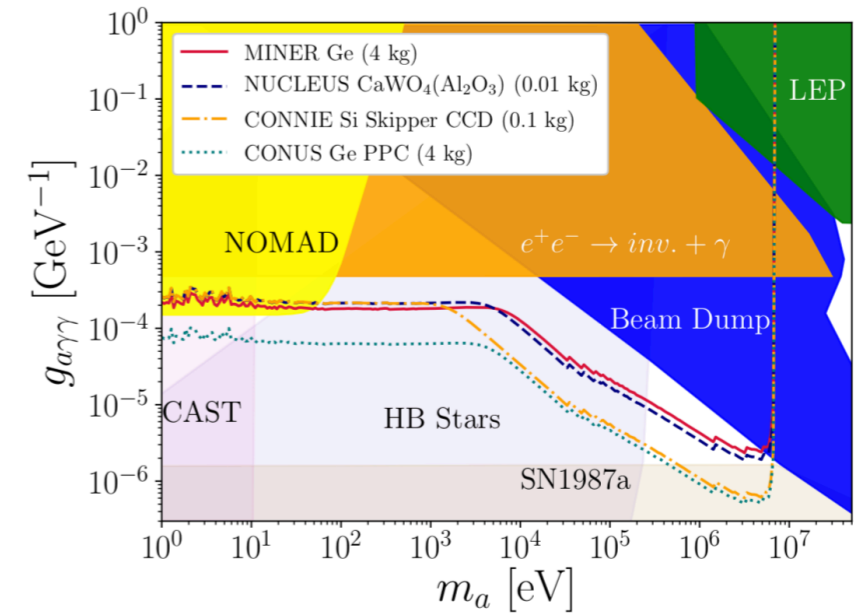
# Fit to the excess

	Region (i)	Region (ii)	Region (iii)
$\gamma_{\text{BDM}}$	$\approx 1$	$\approx 1$	$\gg 1$
VF	✓(flat)	✓(falling)	✓(falling)
VS	✓(flat)	✓(falling)	✓(falling)
PF	✓(rising)	✓(rising)	✗(-)
PS	✓(rising)	✓(rising-and-falling)	✓(rising-and-falling)
SF	✓(flat)	✓(falling)	✓(rising-and-falling)
SS	✓(flat)	✓(falling)	✓(falling)

- ✓: One can find mass spectra to reproduce XENON1T excess and satisfy the conditions of the associated regions.
- ✓: A certain range of mediator mass may not reproduce the XENON1T excess.
- ✗: It is generally hard to find a mass spectrum to explain the excess.

# Further discussions

- Bounds from accelerators, astrophysical and cosmological observations?



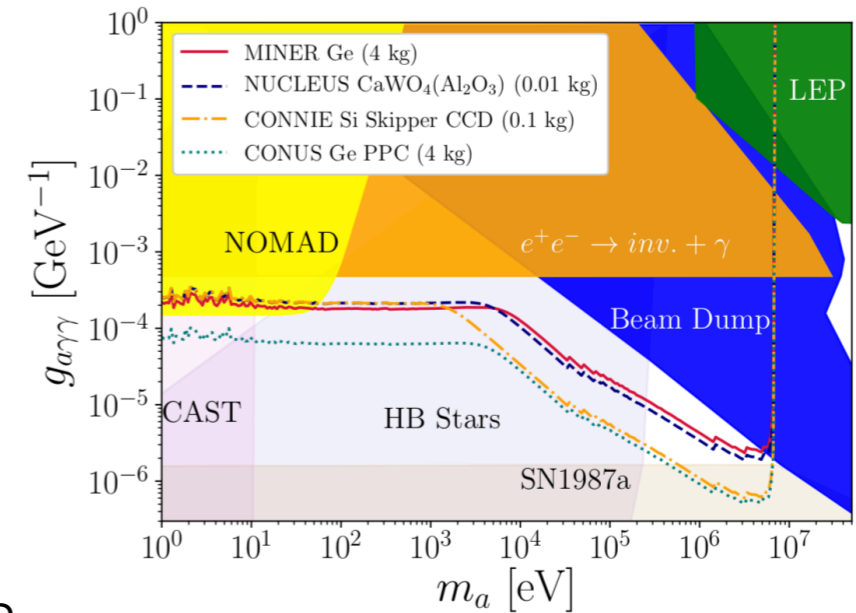
Dent et al., PRL 2020

# Further discussions

- Bounds from accelerators, astrophysical and cosmological observations?

- If the coupling constant and the mass parameter have effective dependence upon environmental conditions of astrophysical objects such as temperature and matter density, the limits can be relaxed by several orders of magnitude.

- Some regions can be probed in future accelerators.



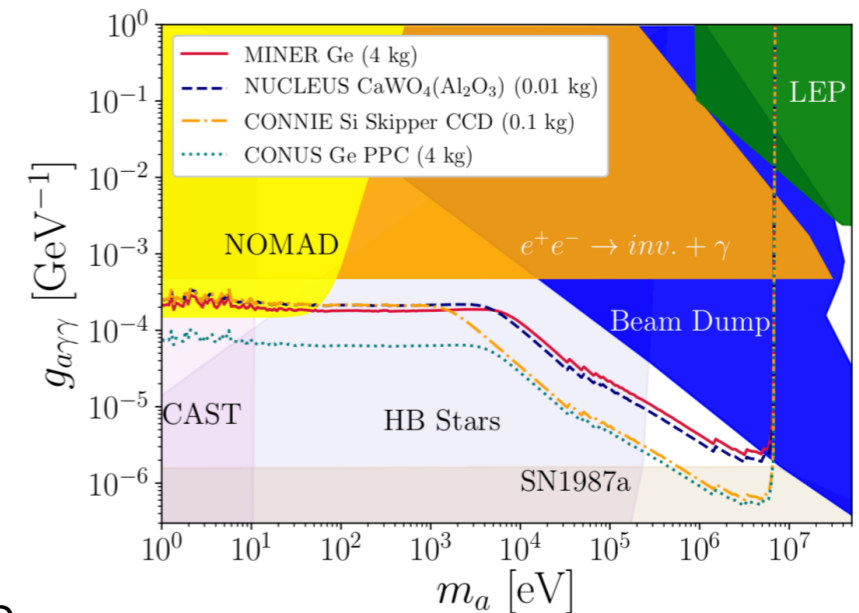
Dent et al., PRL 2020

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Dent et al., PRL 2020

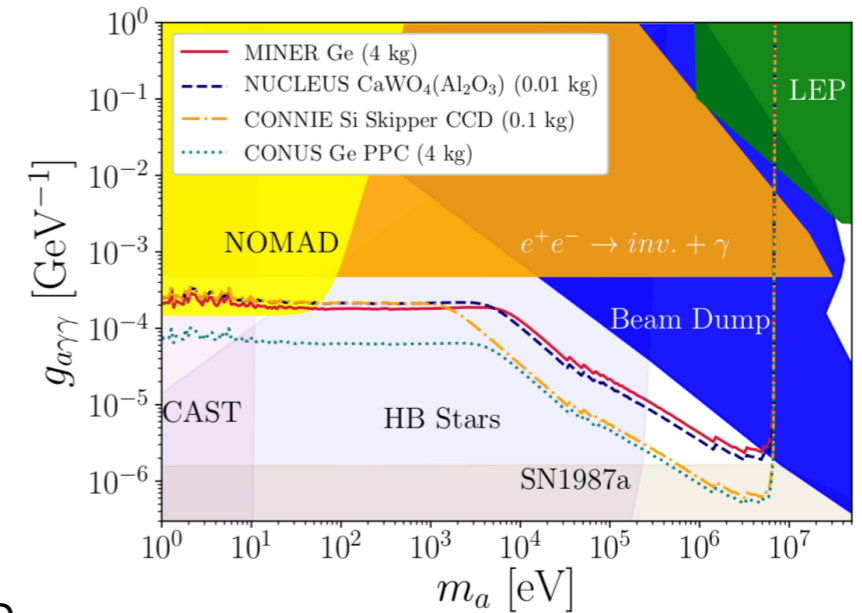
- Nuclear scattering can occur when  $E_1$  increases over  $\mathcal{O}(10 \text{ MeV})$ .  
(reference parameters do not induce nuclear scattering due to kinematics)



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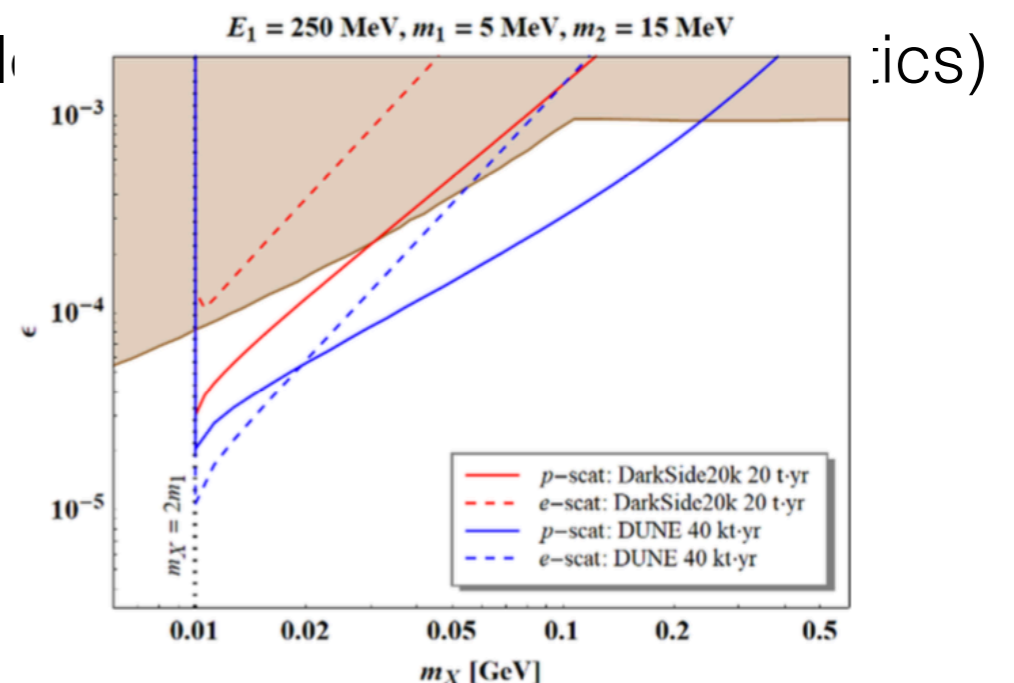
Dent et al., PRL 2020

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- Nuclear scattering can occur when  $E_1$  increases over  $\mathcal{O}(10 \text{ MeV})$ .

(reference parameters do not induce nuclear scattering)

- Complimentary searches are possible.



Kim, Machado, Park, **SS**, JHEP 2007, 057 (2020)

# Conclusions

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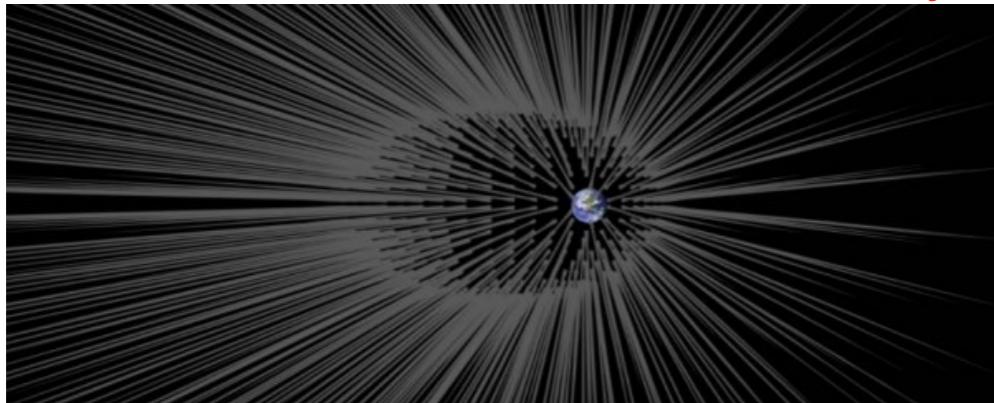
- The XENON1T excess can be a smoking gun signal of Dark World beyond WIMP: fast moving DM, e.g., Boosted Dark Matter.
- The dedicated analyses will be helpful in future studies in XENON-NT, DarkSide, COSINE-200, ...
- Fitting the excess with BDM is nontrivial (large cross section, mean free path, a narrow range of  $E_R$ , binding energy of the electron).
- We found a wide range of BDM parameter region that can fit the excess including  $v \sim c$  and  $m_1 < m_e$  for various mediator masses.
- The scales of mass and couplings preferred by the excess relies on the type of mediator (further studies for more reference models needed).
- Our method is general, and hence readily applicable to the interpretation of observed data in the DM direct detection experiments.

# Energetic dark matter from the Universe

---

~~WIMP wind  $\sim O(100 \text{ km/s})$~~

$\gg O(100 \text{ km/s})$



in promising theories beyond WIMP

- Anti-DM from DM-induced nucleon decay in the Sun

Huang, Zhao, 1312.0011

- Solar reflection: light DM scattered with hot solar nuclei or electrons

An, Pospelov, Pradler, Ritz, 1708.03642

Emken, Kouvaris, Nielsen, 1709.06573

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Bringmann, Pospelov, 1810.10543

Yin, 1809.08610

Ema, Sala, Sato, 1811.00520

Cappiello, Beacom, 1906.11283

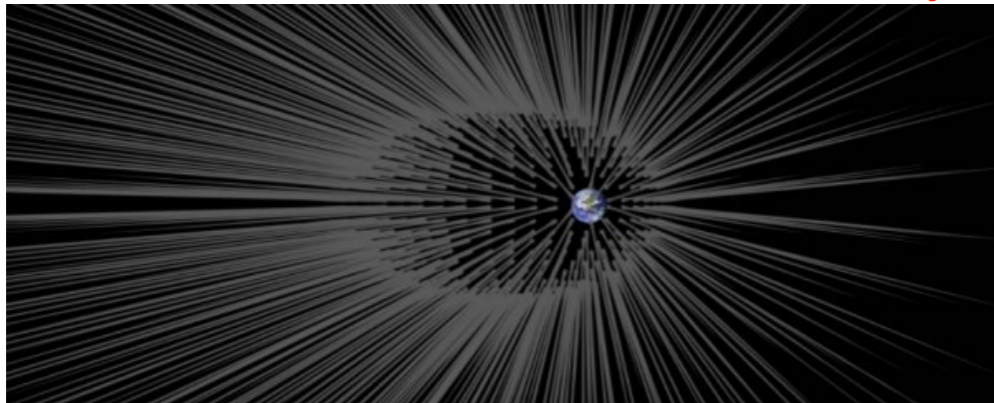
Cappiello, Ng, Beacom, 1810.07705

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- Boosted Dark Matter: DM boosted by the dark sector structure

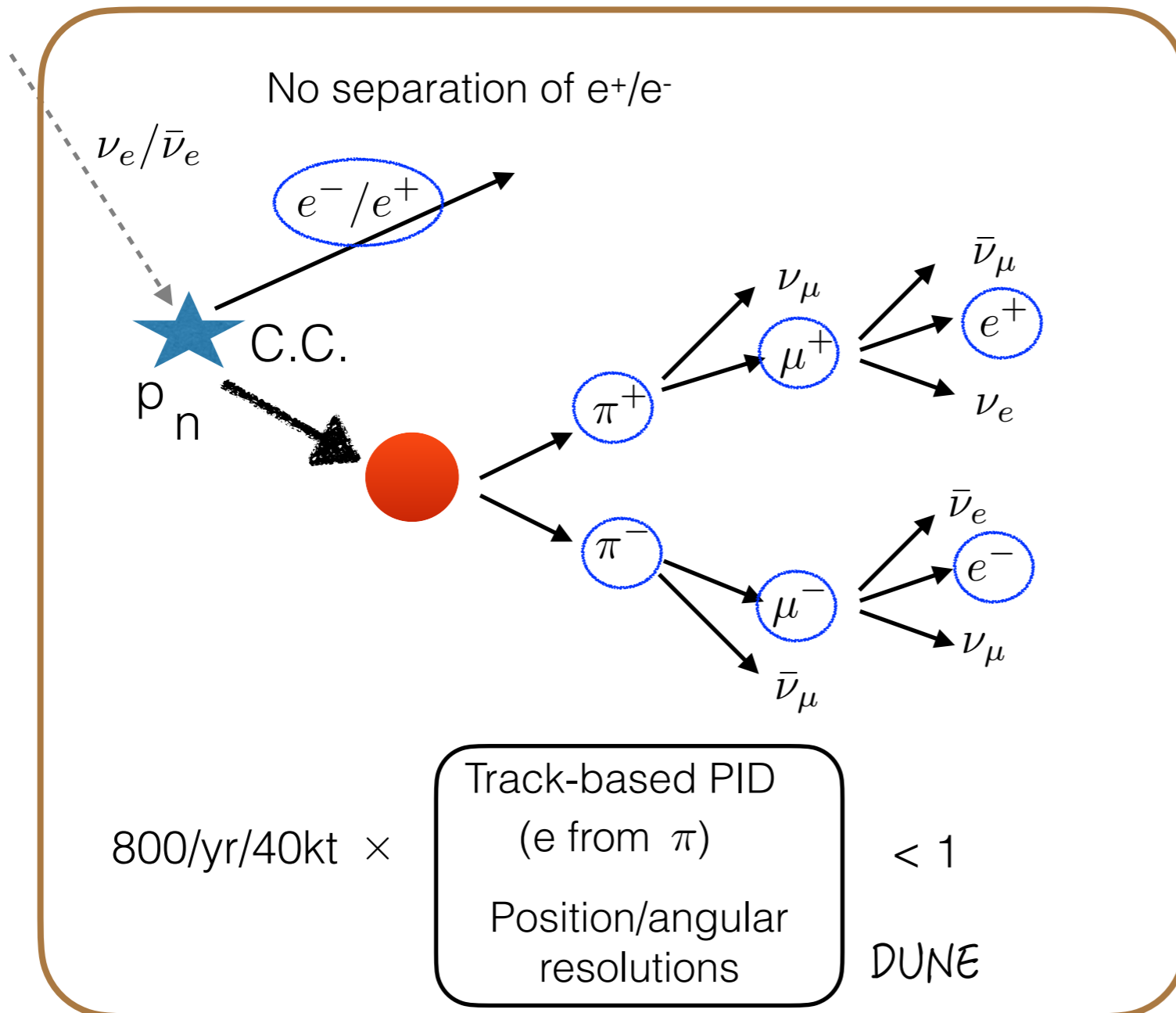
from 2014 (not from scattering with the energetic SM particles)

# Backup: Background candidate

e.g., primary: e-scattering, secondary  $e^+ e^-$  iBDM

De Roeck, Kim, Moghaddam, Park, **SS**,  
Whitehead, 2005.08979

Fiducial volume



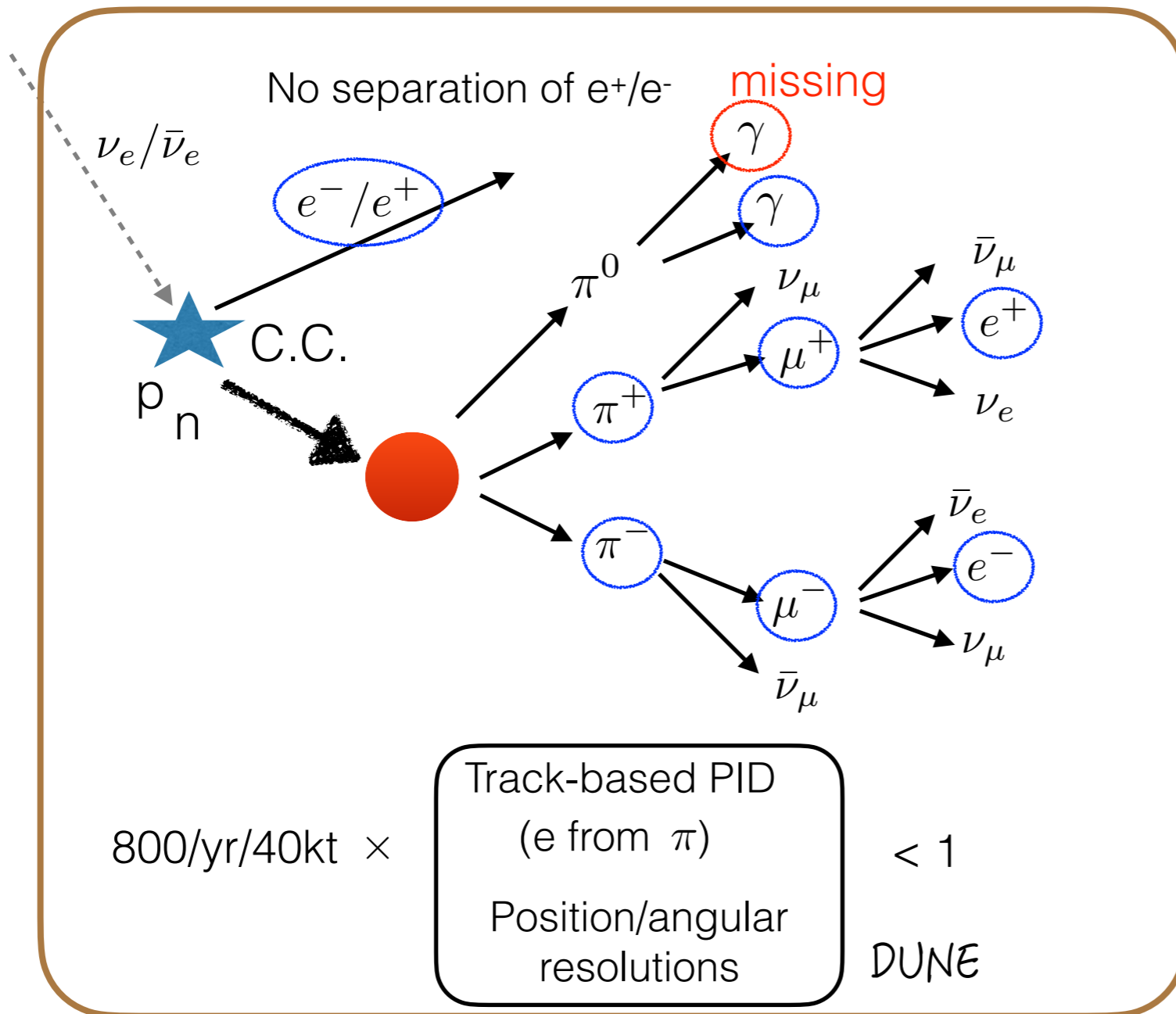
$\chi_1$ : light BDM,  $\chi_2$ : excited state

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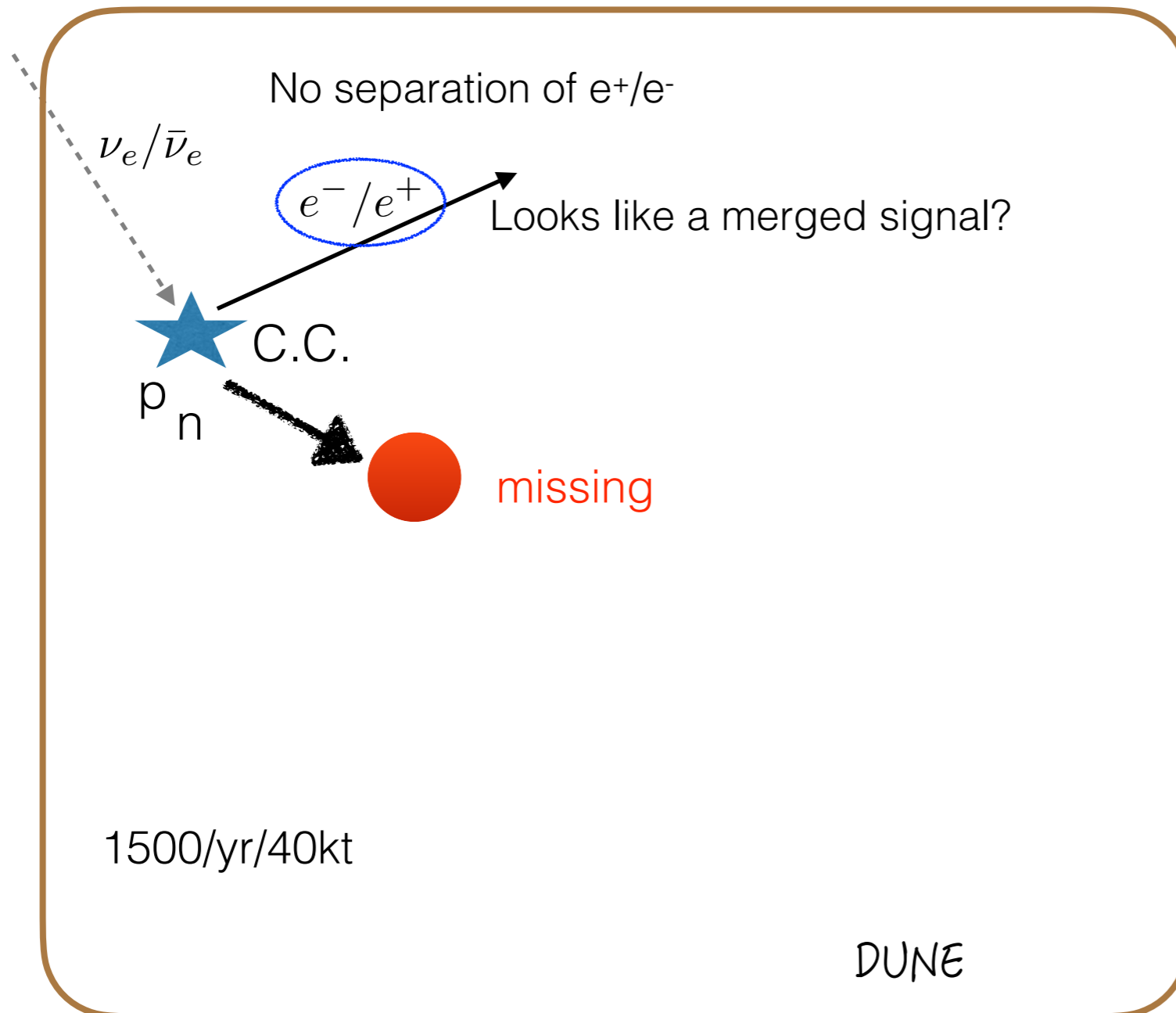
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De Roeck, Kim, Moghaddam, Park, **SS**,  
Whitehead, 2005.08979

Fiducial volume



$\mu$

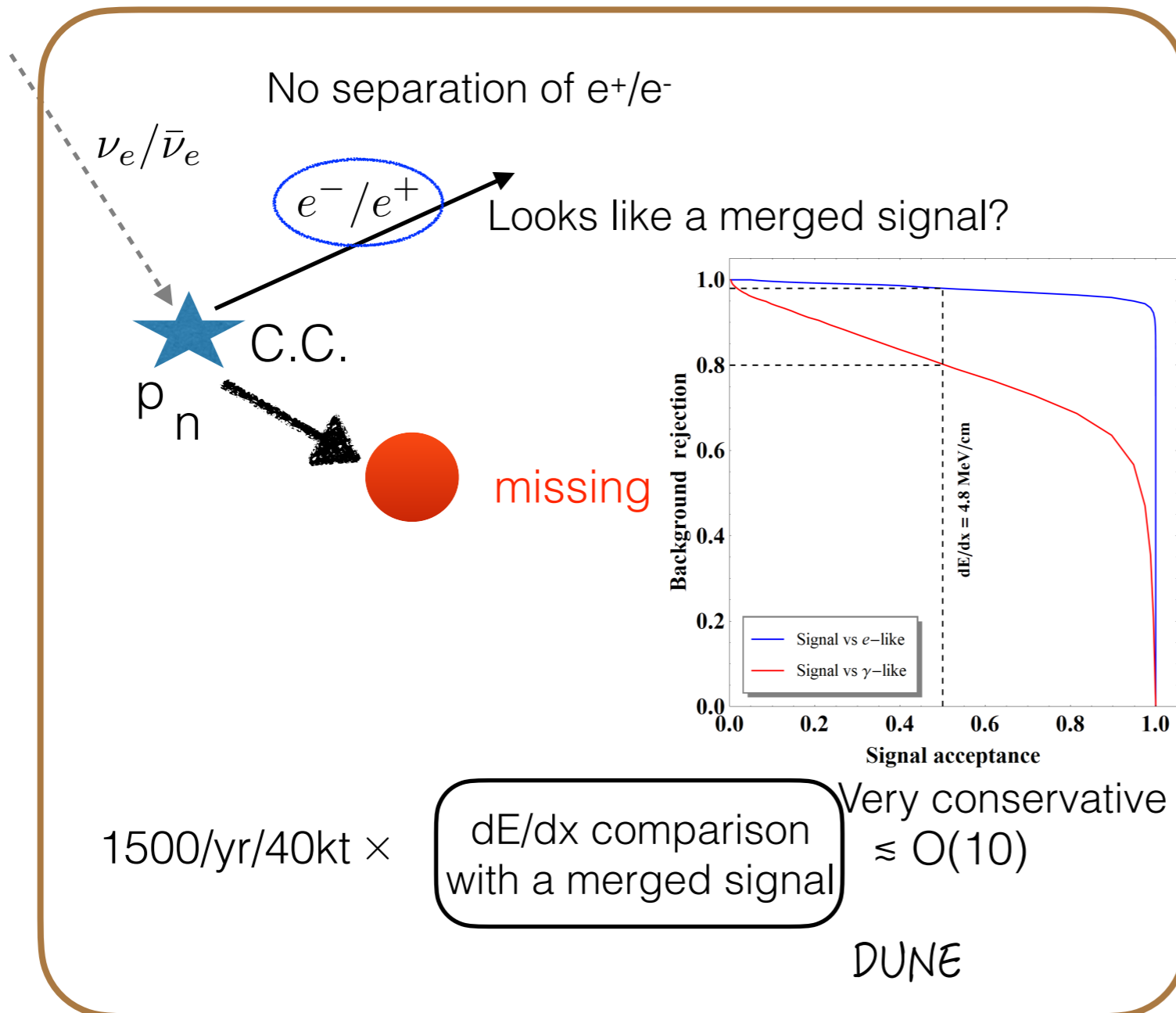
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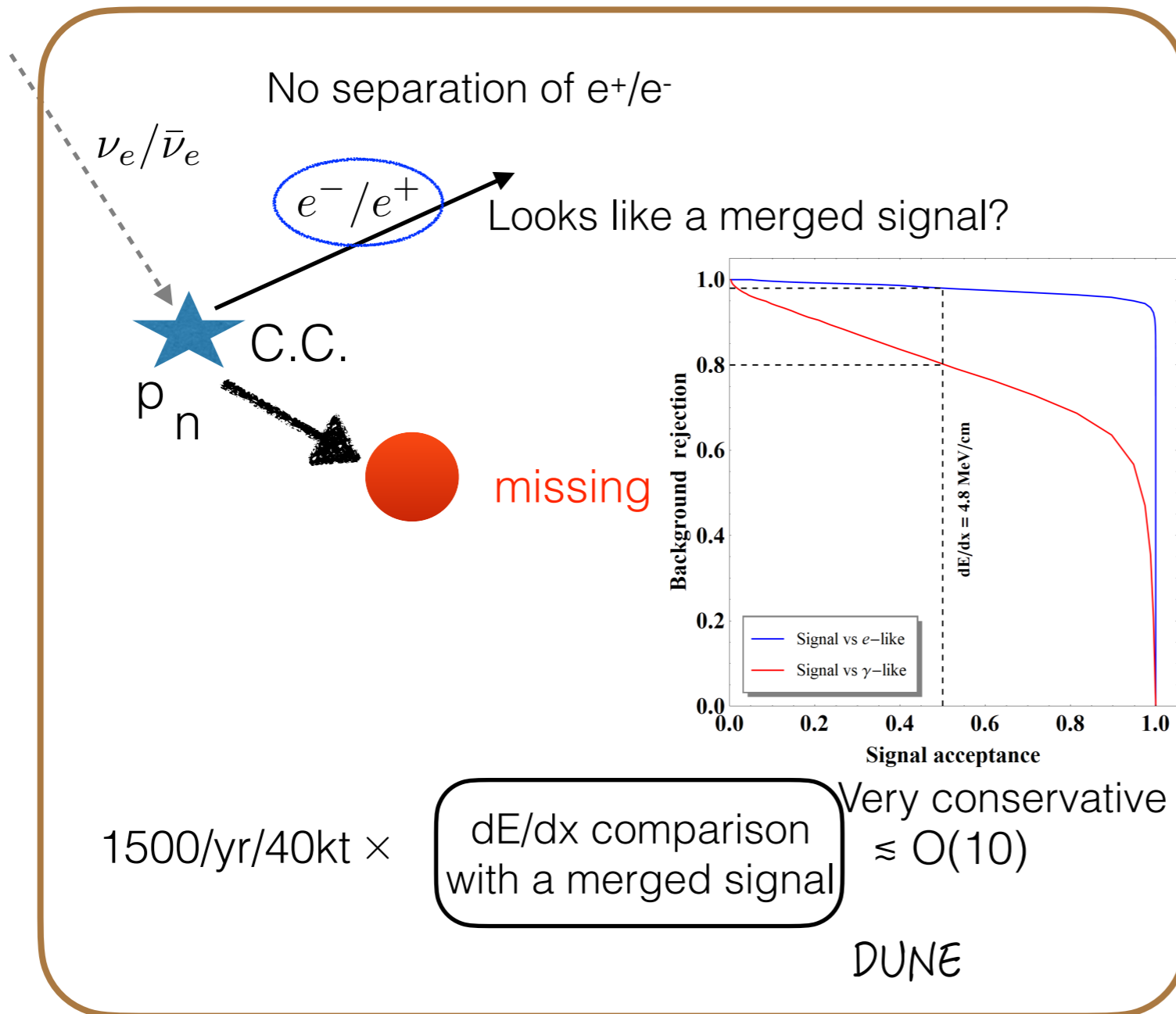


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Fiducial volume



- Other subdominant bkg. negligible
- N.C. events (smaller)
- $\nu_\mu$ : accompanying  $\mu$
- $\nu_\tau$ : too small flux
- Comic-ray: flux & PID
- Zero-bkg. is easily achievable
- (quasi-elastic) proton scattering: less bkg.

$\chi_1$ : light BDM,  $\chi_2$ : excited state