

Glimpses of new physics through neutrinos

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Dark Matter as a Portal to New Physics

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Online

apctp

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asia pacific center for theoretical physics

Plan of the talk

- Basics of oscillation physics and pending issues
- Hint for a light (eV) sterile neutrino
- Non-Standard Neutrino Interaction (NSI)
- Other BSM options (Dark sector?)
- Summary

standard oscillation probability

- $|\nu_\alpha\rangle = \sum_k U_{\alpha k}^* |\nu_k\rangle$ ($\alpha = e, \mu, \tau$) ← *Flavor & Mass related by mixing*

- $H |\nu_k\rangle = E_k |\nu_k\rangle$

← *Solve Schrodinger's eqn.*

- $i \frac{d}{dt} |\nu_k(t)\rangle = H |\nu_k(t)\rangle$

- $P_{\alpha\beta}(L, E) = |\langle \nu_\beta | \nu_\alpha(x) \rangle|^2$
 $= \delta_{\alpha\beta} - 4 \sum_{k>j} \text{Re}[U_{\alpha k}^* U_{\beta j}^* U_{\beta k} U_{\alpha j}] \sin^2 \Delta_{kj}$
 $+ \sum_{k>j} \text{Im}[U_{\alpha k}^* U_{\beta j}^* U_{\beta k} U_{\alpha j}] \sin 2\Delta_{kj},$

(where, $\Delta_{ij} = 1.27 \times \frac{\Delta m_{ij}^2 [\text{eV}^2] \times L [\text{km}]}{E [\text{GeV}]}$)

- Parameters responsible for neutrino oscillation:

$$\theta_{13}, \theta_{12}, \theta_{23}, \Delta m_{21}^2, \Delta m_{31}^2, \delta_{cp}$$

Status of oscillation parameters

Table: de Salas, Forero, Gariazzo, Martinez-Mirave, Mena, Ternes, Tortola, Valle: 2006.11237

Oscillation parameter	Best fit value	3σ range
$\theta_{12}/^\circ$	34.3	[31.4, 37.4]
$\theta_{23}/^\circ$	48.8	[41.6, 51.3]
$\theta_{13}/^\circ$	8.6	[8.2, 8.9]
δ_{13}/π	-0.8	$[-1, 0] \cup [0.8, 1]$
$\Delta m_{21}^2/10^{-5} \text{ eV}^2$	7.5	[6.9, 8.1]
$\Delta m_{31}^2/10^{-3} \text{ eV}^2$	2.6	[2.5, 2.7]

} 3 mixing angles
 } 1 CP phase
 } 2 mass squared differences

$$\begin{aligned}
 P_{\mu e} = & \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2(1-A)\Delta}{1-A} \\
 & + \alpha^2 \sin^2 2\theta_{12} \cos^2 \theta_{23} \frac{\sin^2 A\Delta}{A} \\
 & + \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \frac{\sin(1-A)\Delta}{1-A} \frac{\sin A\Delta}{A} \cos(\delta + \Delta)
 \end{aligned}$$

where,

$$A = \frac{2\sqrt{2}EG_F n_E}{\Delta m_{31}^2}$$

$$\Delta = \frac{\Delta m_{31}^2 L}{4E}$$

$$\alpha = \frac{\Delta m_{21}^2}{\Delta m_{31}^2}$$

Leptonic CP violation?

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Is the CP phase non-zero?
could help explain baryon asymmetry

$$\begin{aligned}
 P_{\mu e} = & \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2(1-A)\Delta}{1-A} \\
 & + \alpha^2 \sin^2 2\theta_{12} \cos^2 \theta_{23} \frac{\sin^2 A\Delta}{A} \\
 & + \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \frac{\sin(1-A)\Delta}{1-A} \frac{\sin A\Delta}{A} \cos(\delta + \Delta)
 \end{aligned}$$

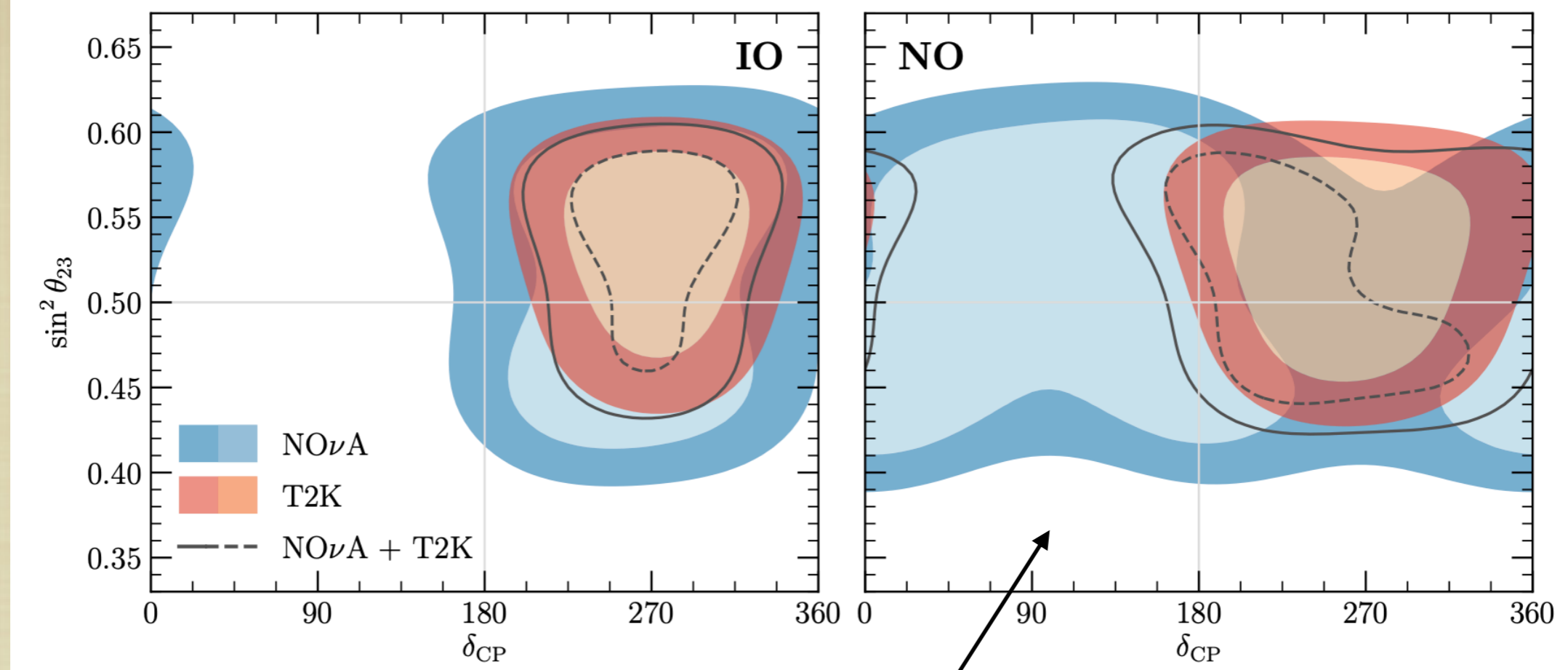
where,

$$A = \frac{2\sqrt{2}EG_F n_E}{\Delta m_{31}^2}$$

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Leptonic CP violation (Tension between NO ν A & T2K)

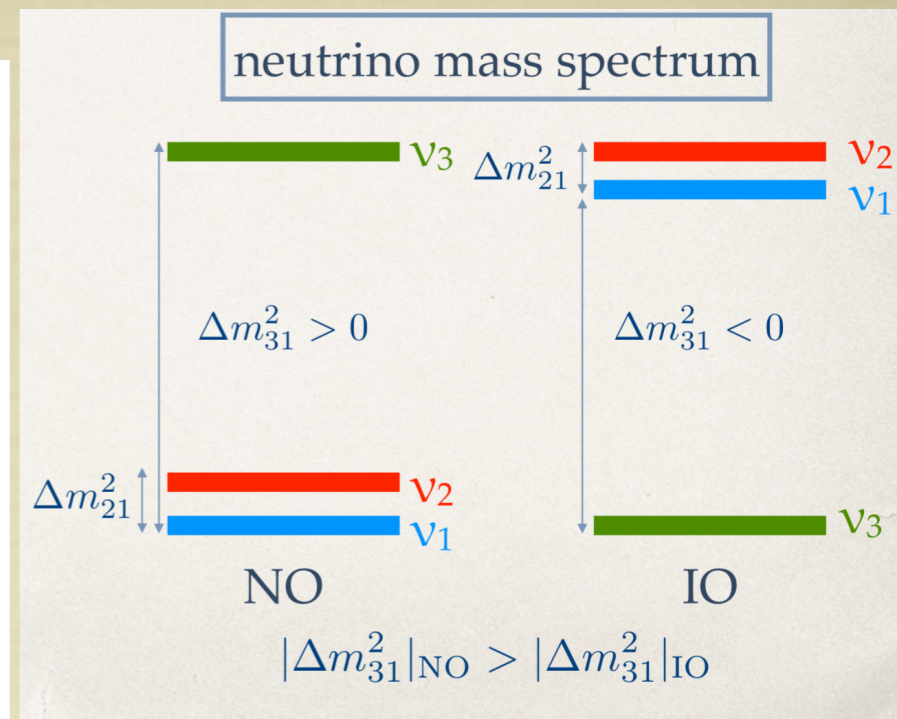


Tension () between allowed regions in $\delta_{13} - \sin^2 \theta_{23}$ plane by NO ν A and T2K (for NO)

Mass ordering ambiguity

Table: de Salas, Forero, Gariazzo, Martinez-Mirave, Mena, Ternes, Tortola, Valle: 2006.11237

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What is the sign of Δm_{31}^2 ?

$$\begin{aligned}
 P_{\mu e} = & \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2(1-A)\Delta}{1-A} \\
 & + \alpha^2 \sin^2 2\theta_{12} \cos^2 \theta_{23} \frac{\sin^2 A\Delta}{A} \\
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where,

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$$\Delta = \frac{\Delta m_{31}^2 L}{4E}$$

$$\alpha = \frac{\Delta m_{21}^2}{\Delta m_{31}^2}$$

Mass ordering ambiguity (Tension between NOvA & T2K)

Experiment(s)	$\Delta\chi^2_{(\text{NO,IO})}$
T2K	+1.2
NOvA	+0.13
SK18/SK20	+3.4/+3.2
T2K + NOvA	-1.8
T2K + SK18	+5.7
NOvA + SK18	+3.6
T2K+NOvA+SK18	+2.2

Tension happens
because they prefer
different value of δ_{13}

Octant degeneracy

Table: de Salas, Forero, Gariazzo, Martinez-Mirave, Mena, Ternes, Tortola, Valle: 2006.11237

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$\theta_{23} > \pi/4$ or $\theta_{23} < \pi/4$?

$$\begin{aligned}
 P_{\mu e} = & \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2(1-A)\Delta}{1-A} \\
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 \end{aligned}$$

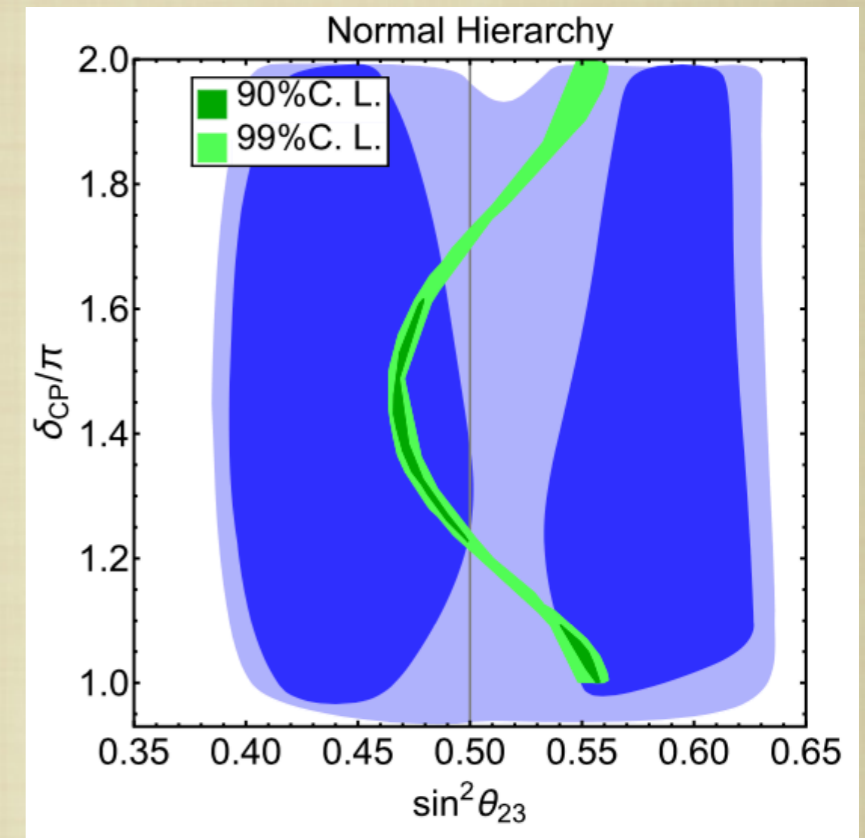
$$\sin 2\theta_{23} = \sin 2(\pi/2 - \theta_{23})$$

Octant degeneracy

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Chatterjee, MM, Pasquini, Valle (2017)

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 \end{aligned}$$

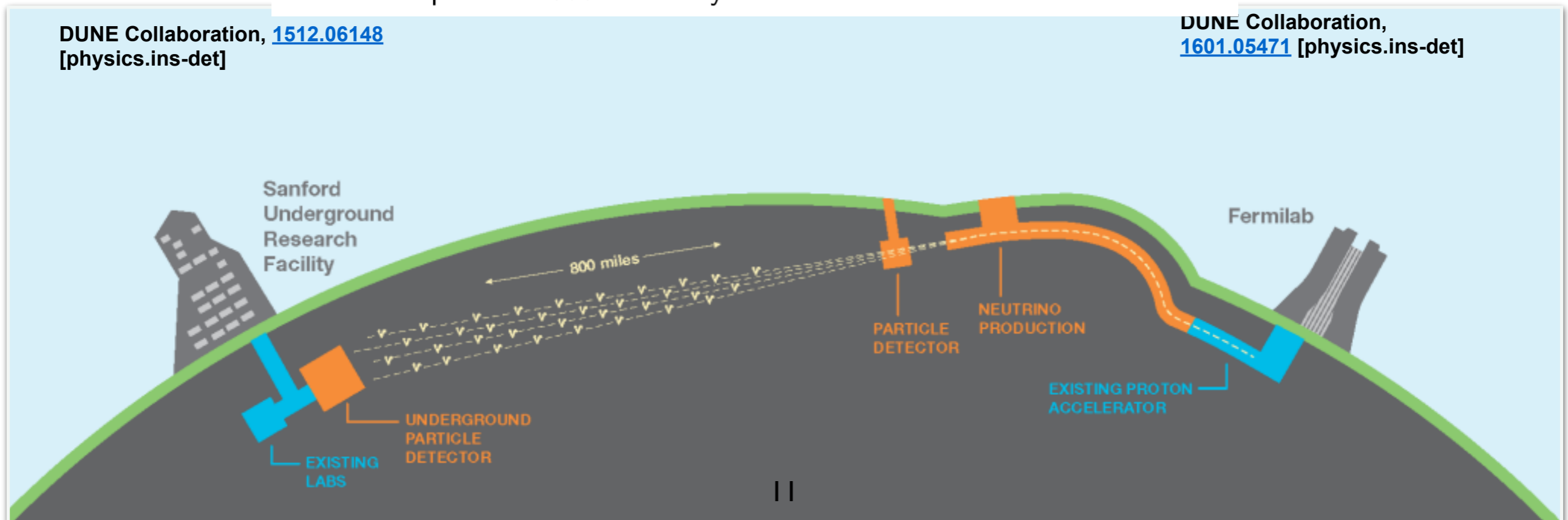
$$\sin 2\theta_{23} = \sin 2(\pi/2 - \theta_{23})$$

Deep Underground Neutrino Experiment (DUNE)

- A proposed long baseline experiment (the erstwhile LBNE) with 1300 km baseline
- likely to have a 40 kt FD with 3.5 yrs. of ν and 3.5 yrs. of $\bar{\nu}$ run.
- The incident ν_μ beam is generated by 80 GeV proton beam delivered at 1.07 MW with a POT of 1.47×10^{21}
- Total exposure : 300 kt-MW-yr.

DUNE Collaboration, [1512.06148](#)
[physics.ins-det]

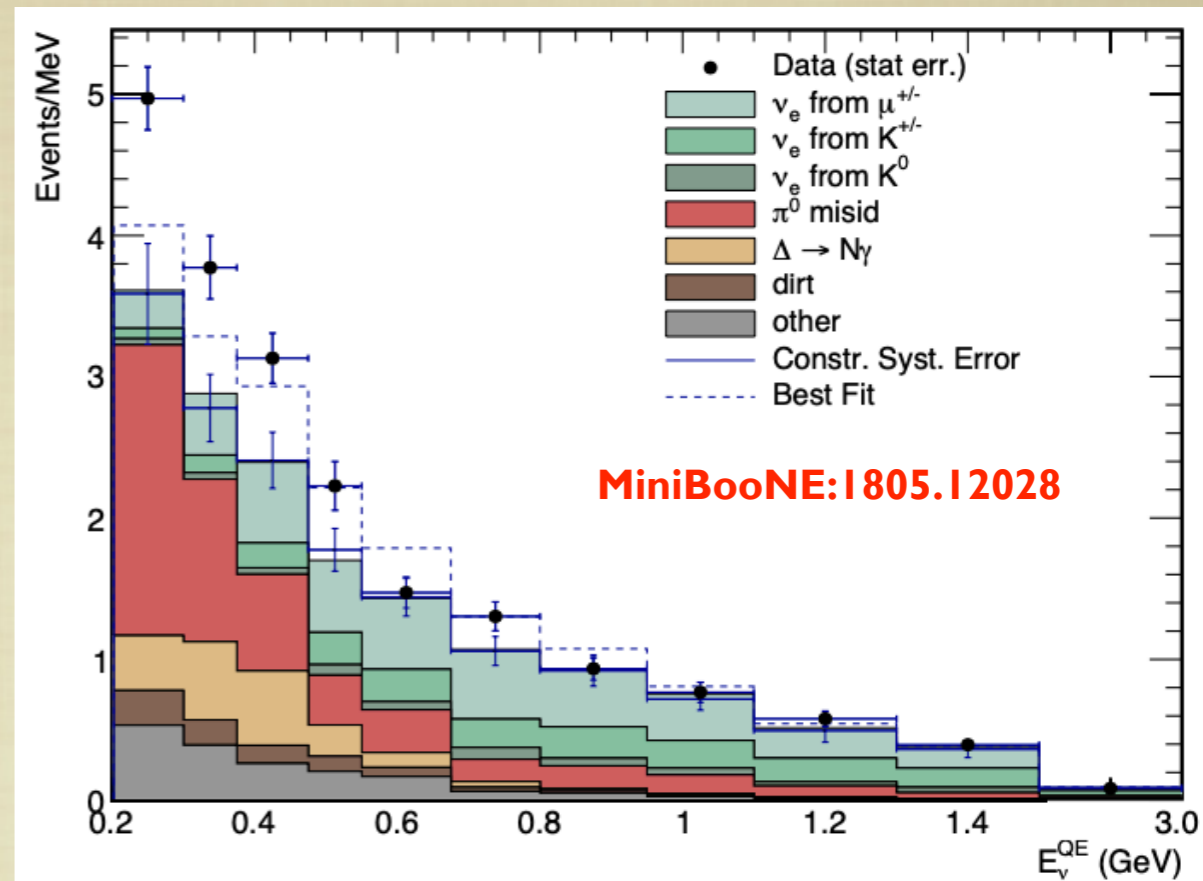
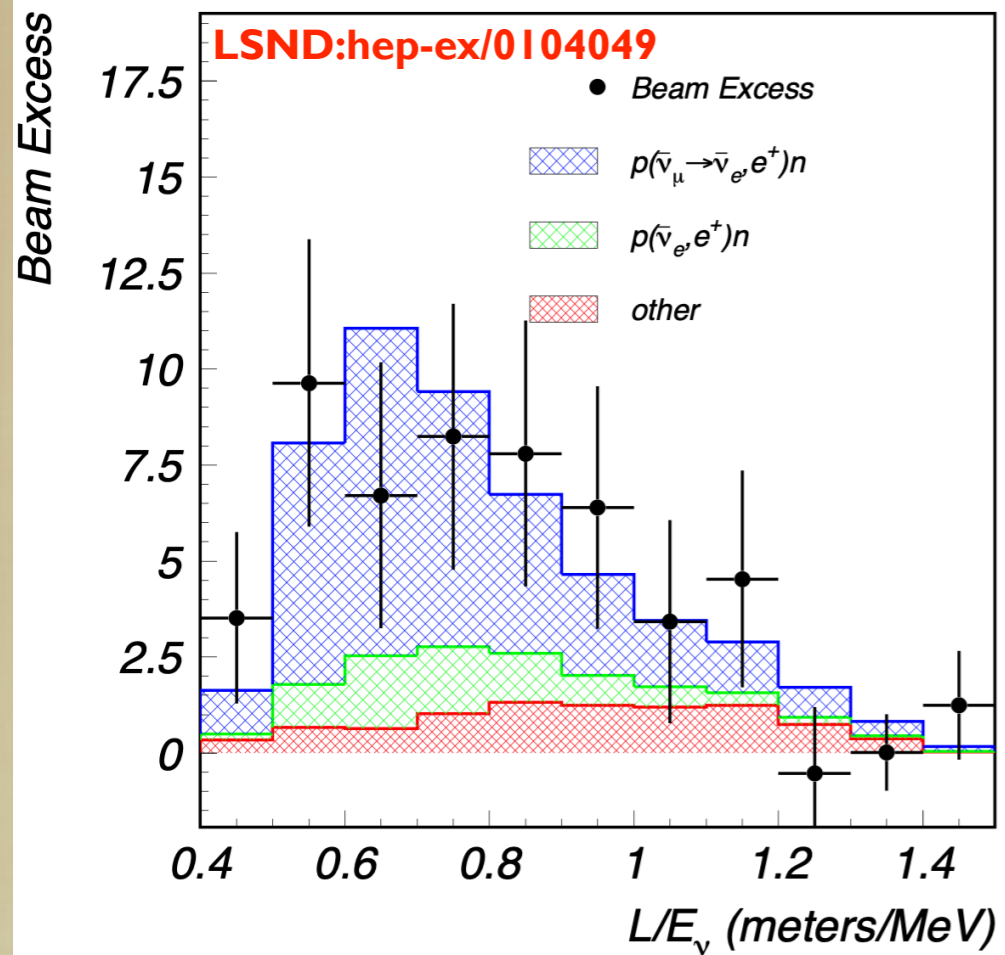
DUNE Collaboration,
[1601.05471](#) [physics.ins-det]



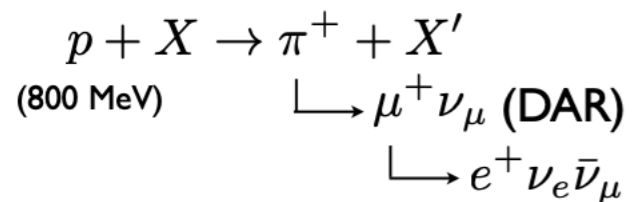
Plan of the talk

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Hint for new physics?



VERY intense proton beam



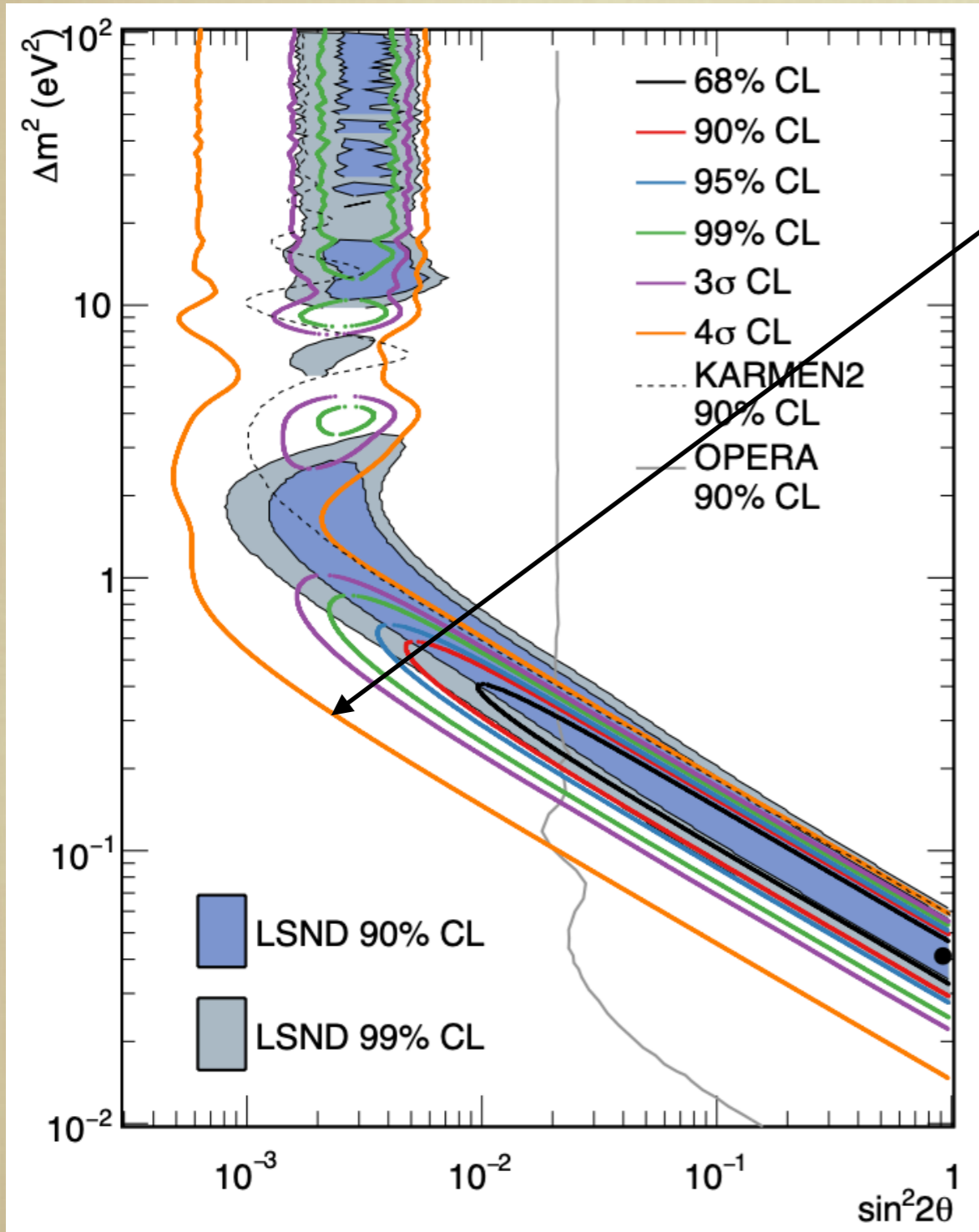
Detection signature:
IBD ($\bar{\nu}_e p \rightarrow e^+ n$)

LSND detected more $\bar{\nu}_e$
than expected: 87.9 ± 23.2 events
(**3.8 σ excess**)

Measures $\nu_\mu \rightarrow \nu_e$ osc.

Event excess: 381.2 ± 85.2 (4.5 σ)

Hint for new physics?



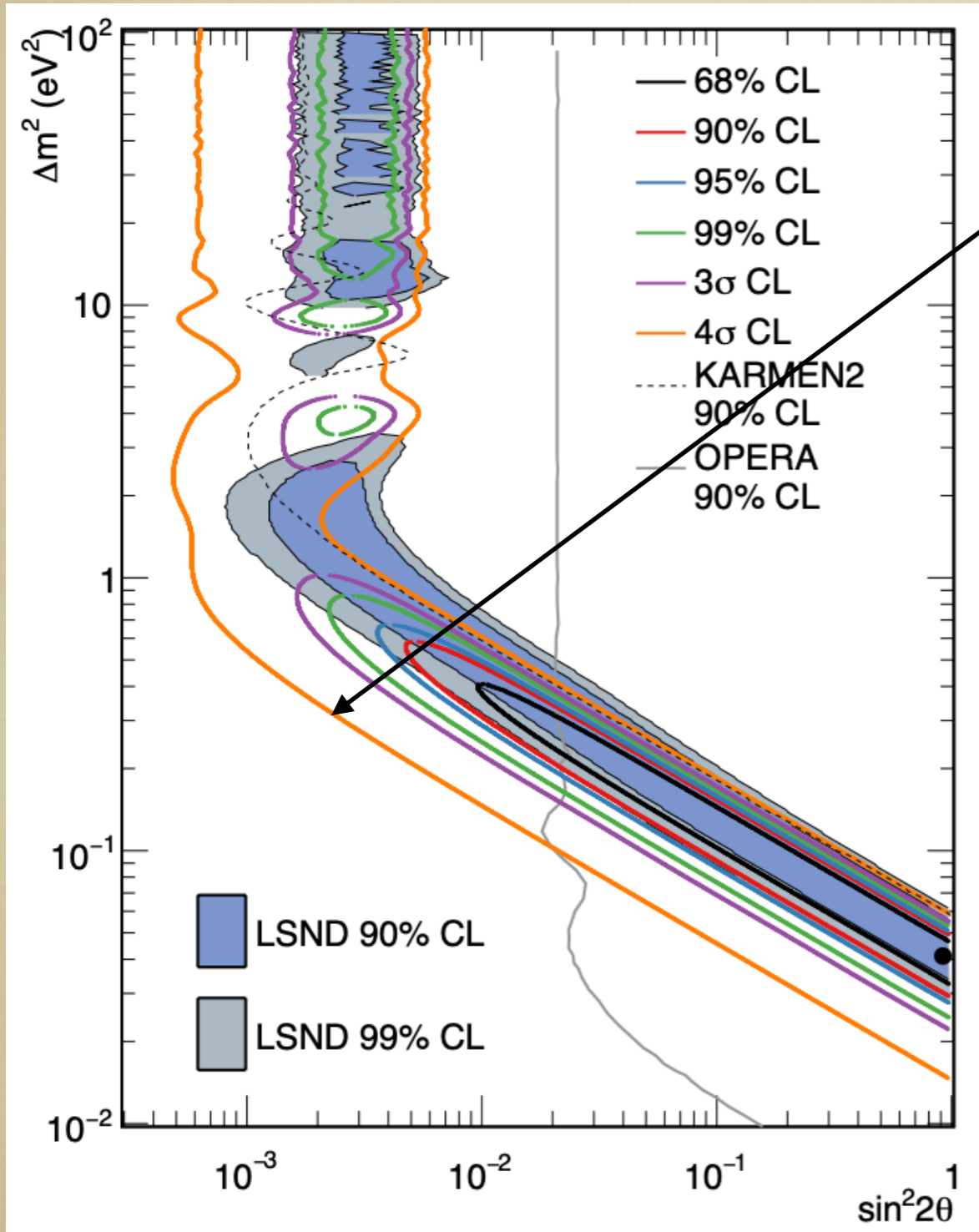
4σ by MiniBooNE

MB + LSND ~ 6σ excess

• Light sterile ν

👍 DANSS, Neutrino-4
👎 Icecube, Cosmology

Hint for new physics?




4 σ by MiniBooNE

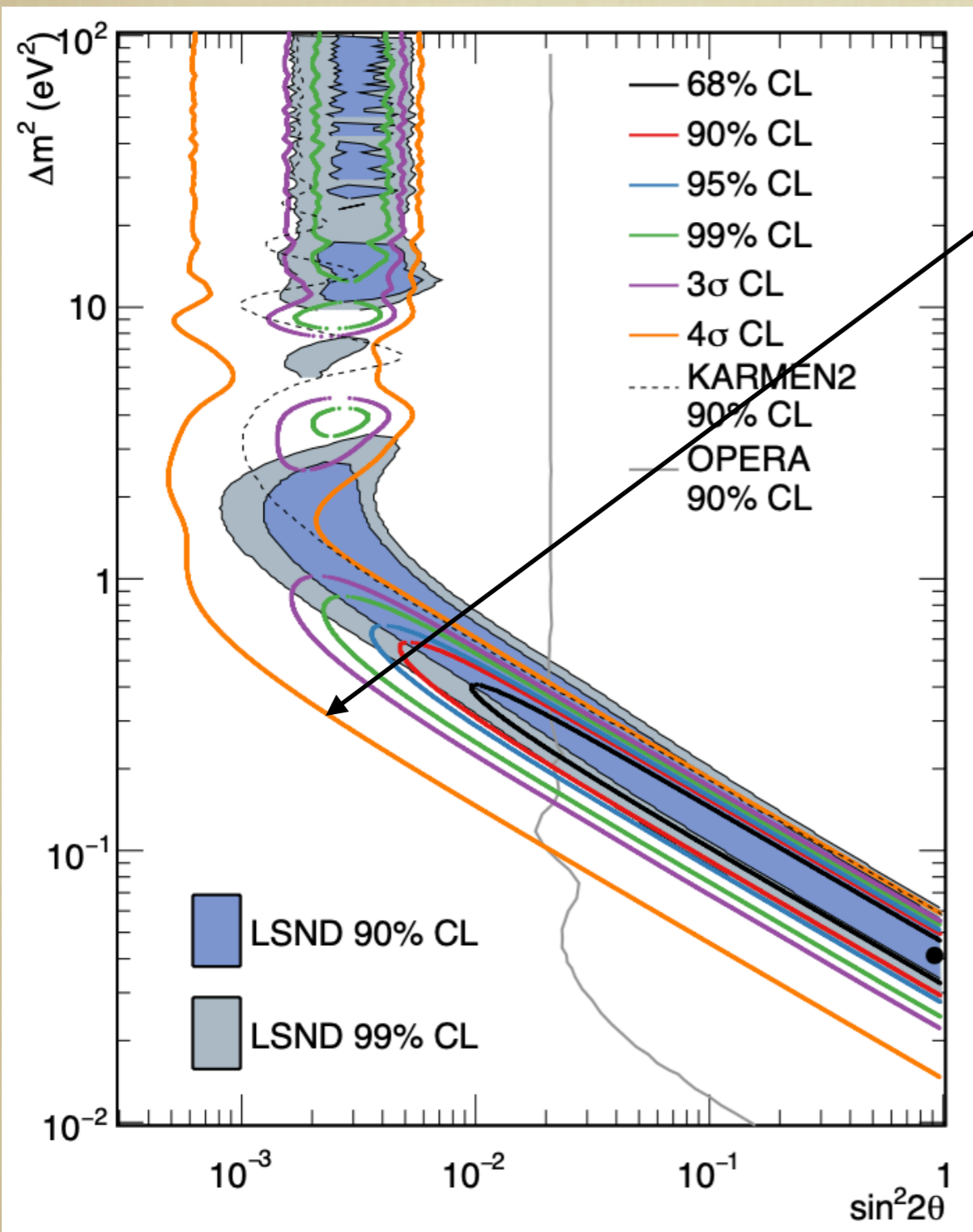
MB + LSND \sim 6 σ excess

• Light sterile ν

Stay tuned for Fermilab SBN


 Icecube, Cosmology

Hint for new physics?



4σ by MiniBooNE

MB + LSND ~ 6σ excess

Stay tuned for Fermilab SBN

• Light sterile ν

• $\nu_{\mu} \rightarrow \nu_{\tau}$, Neutrino-4
 Icecube, Cosmology

• Decay of heavy sterile ν

Dentler, Esteban, Kopp, Machado: 1911.01427

Ballett, Pascoli, Ross-Lonergan: 1808.02915

• The Dark Sector

Bertuzzo, Jana, Machado, Funchal: 1807.09877

Abdallah, Gandhi, Roy: 2006.01948

• ??

Light sterile neutrino (3+1) & oscillation phenomenology

Gandhi, Kayser, MM, Prakash : JHEP (2015)

Agarwalla, Chatterjee, Palazzo : JHEP (2016)

Giunti : NPB (2016)

Coloma, Forero, Parke : JHEP (2018)

.....

3+1 case :basics

$$P_{\alpha\beta} = \delta_{\alpha\beta} - \sum_{k>j} \text{Re}[U_{\alpha k}^* U_{\beta j}^* U_{\alpha j} U_{\beta k}] \sin^2 \Delta_{kj} - \sum_{k>j} \text{Im}[U_{\alpha k}^* U_{\beta j}^* U_{\alpha j} U_{\beta k}] \sin 2\Delta_{kj}$$

$$\text{where } \Delta_{ij} = 1.27 \times \frac{\Delta m_{ij}^2 [\text{eV}^2] \times L [\text{km}]}{E [\text{GeV}]}$$

- 3 more mixing angles and 2 additional CP phases

- Unitarity in 4x4 sector : $\sum_{k=1,2,3,4} U_{\alpha k} U_{\beta k}^* = \delta_{\alpha\beta}$

- $\Delta m_{41}^2 \sim 1 \text{ eV}^2, \Delta m_{21}^2 \approx 7.5 \times 10^{-5} \text{ eV}^2, \Delta m_{31}^2 \approx 2.5 \times 10^{-3} \text{ eV}^2$
 $\Delta m_{41}^2 \approx \Delta m_{42}^2 \approx \Delta m_{43}^2 \gg \Delta m_{31}^2 \gg \Delta m_{21}^2$

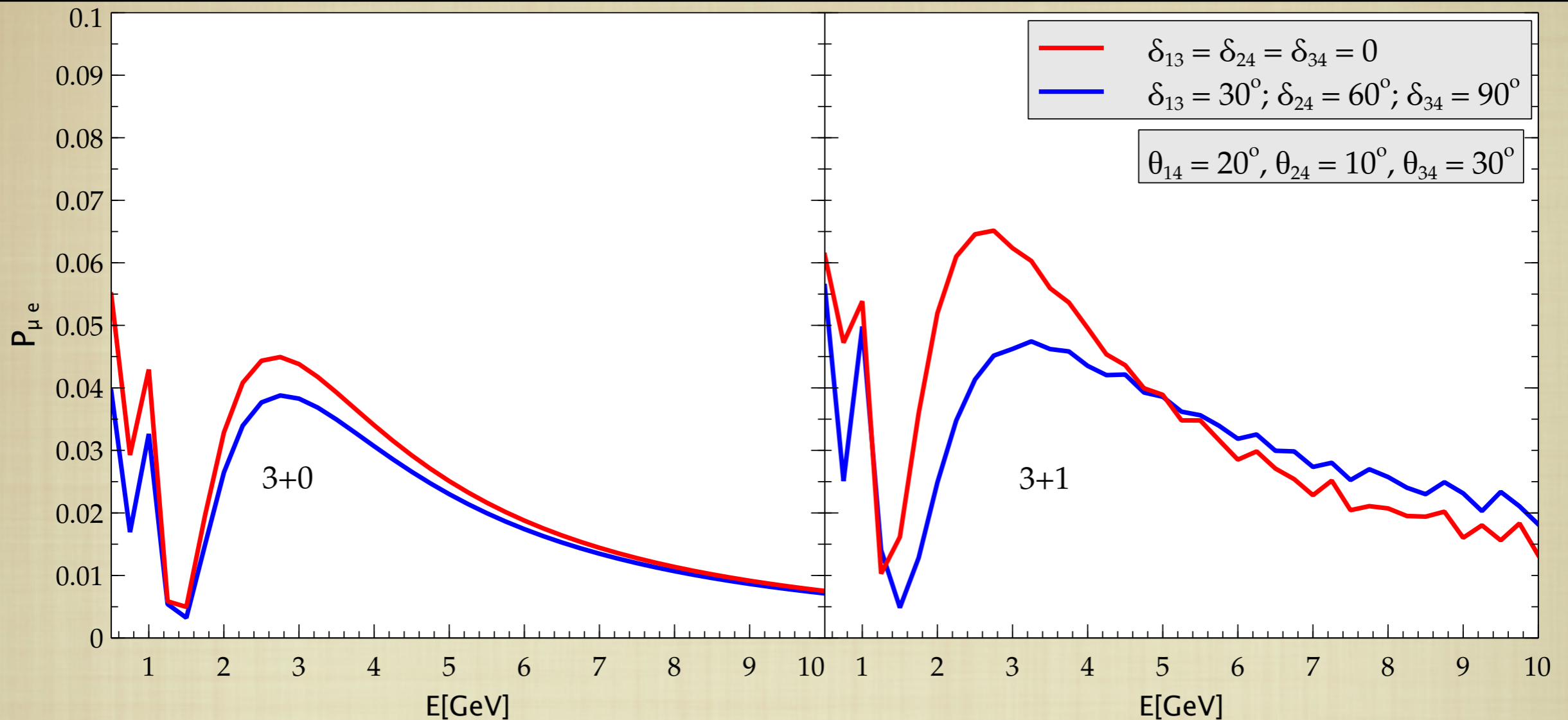
$$P_{\alpha\beta} \approx P_{\alpha\beta}^{3+0} + |U_{\alpha 4}|^2 |U_{\beta 4}|^2 \sin^2 \Delta_{41} \quad (\alpha \neq \beta)$$

Excess

$$P_{\alpha\alpha} \approx P_{\alpha\alpha}^{3+0} - |U_{\alpha 4}|^2 (1 - |U_{\alpha 4}|^2) \sin^2 \Delta_{41}$$

Dip

3+1 case: Effect on probability



- Significant modification of vacuum probability for both CP conserving and CP violating case

Gandhi, Kayser, MM, Prakash (2015)

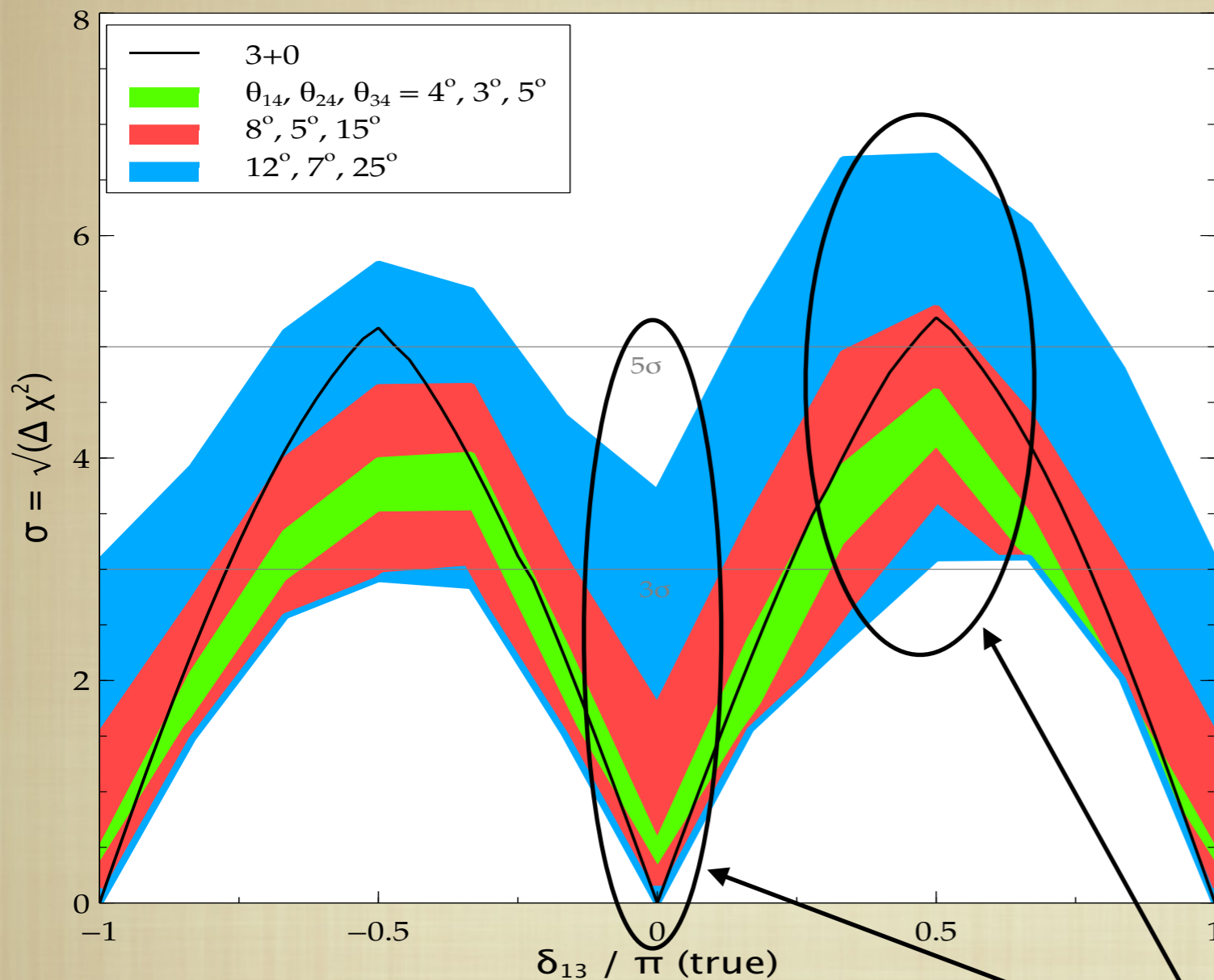
3+1 case :Global analysis

$$|U_{e4}|^2 \lesssim 0.1, \quad |U_{\mu 4}|^2 \lesssim 0.01, \quad |U_{\tau 4}|^2 \lesssim 0.17$$

$$\theta_{14} \lesssim 18.4^\circ, \quad \theta_{24} \lesssim 6.05^\circ, \quad \theta_{34} \lesssim 25.6^\circ$$

Phases are unconstrained

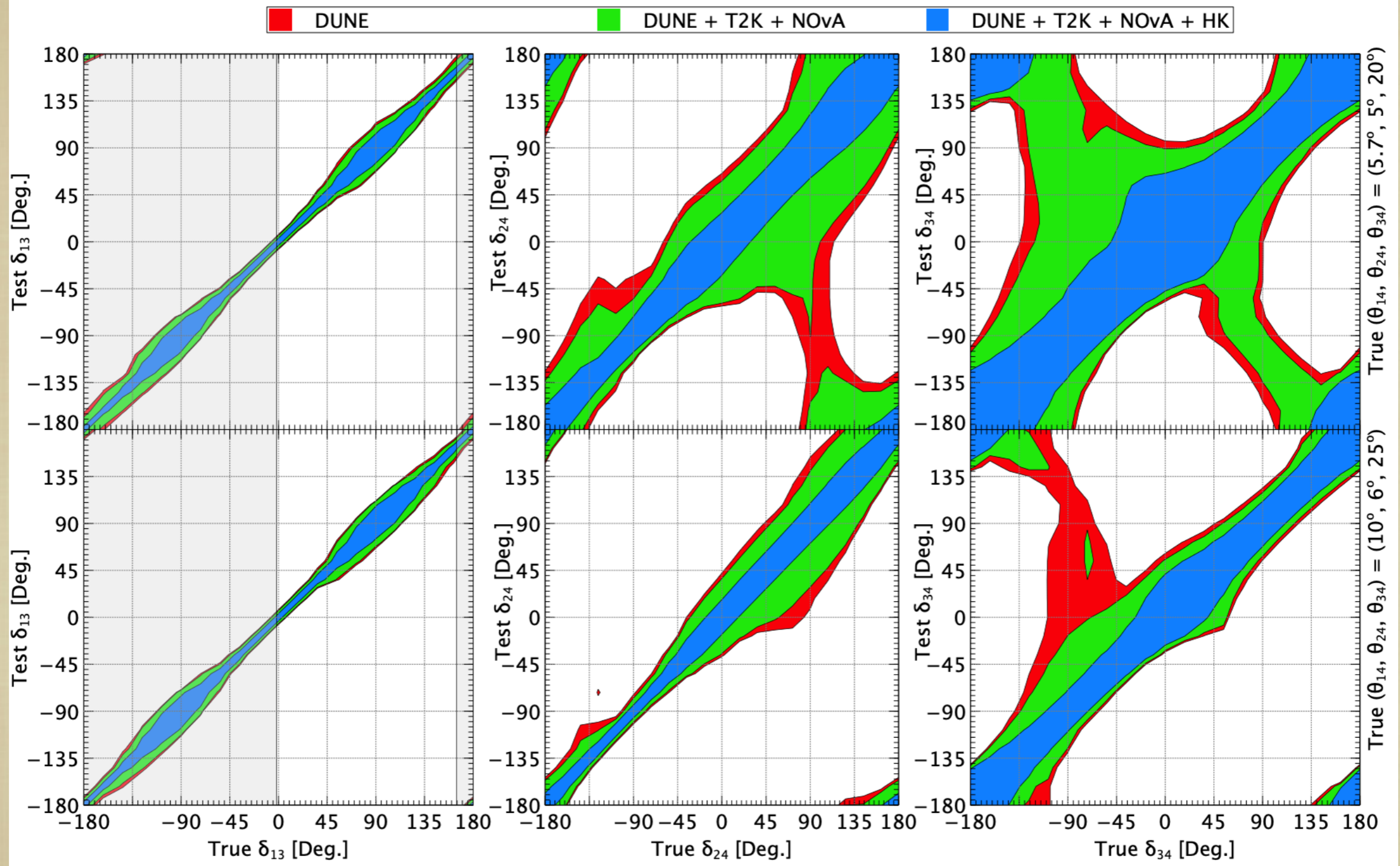
3+1 case: Impact on CPV sensitivity



- Fitted with, $\delta_{13}, \delta_{24}, \delta_{34} = 0$ or π
- The band shows the variation in data for $\delta_{24}^{true}, \delta_{34}^{true} \in [-\pi, \pi]$

CP conservation/ violation?

3+1 case: Exploring the phases



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Non-standard interaction (NSI)

MM, Chatterjee, Mehta : JPG (2016)

de Gouvea, Kelly : NPB (2016)

Liao, Marfatia, Whisnant : PRD (2016)

Deepthi, Goswami, Nath : NPB (2018)

Farzan, Tortola : Review (2018)

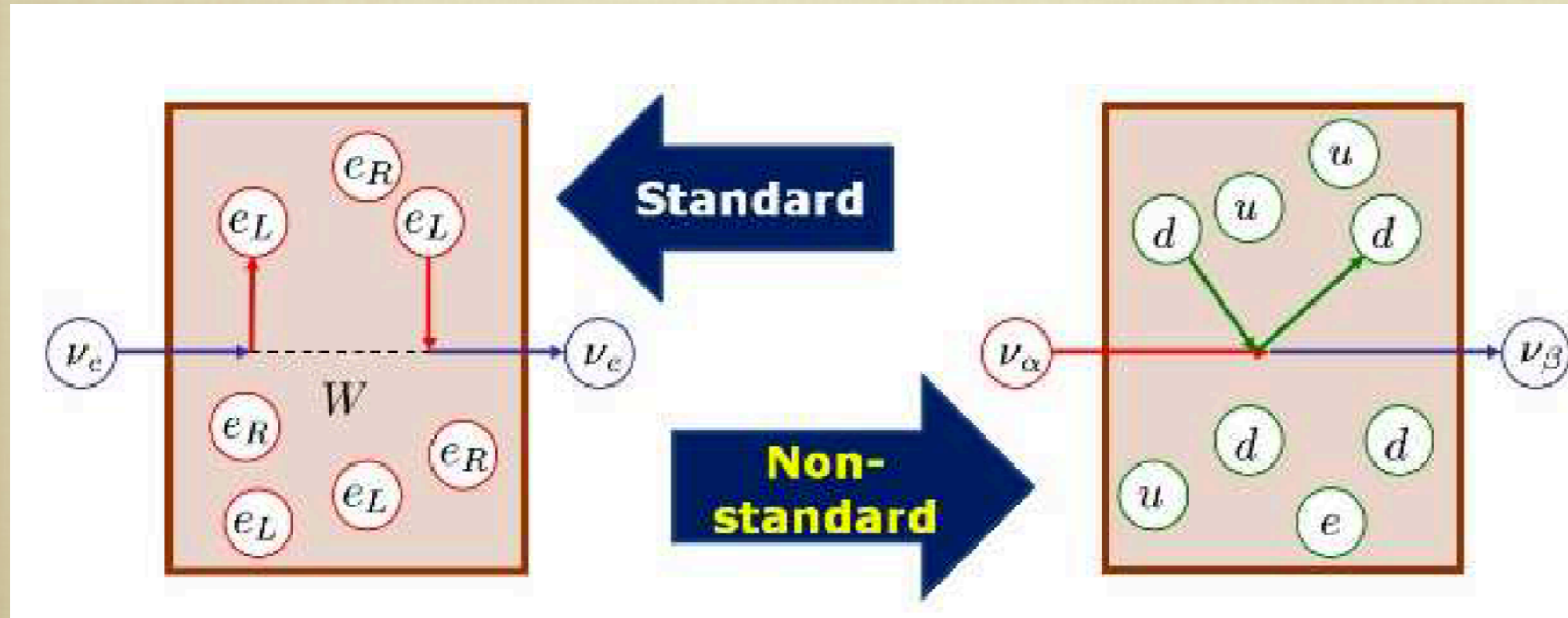
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NSI theory background

Wolfenstein (1978), Barger et al. (1991), Grossman (1995), Ohlsson (2012)...

$$\mathcal{L}_{\text{CC-NSI}} = -2\sqrt{2}G_F \epsilon_{\alpha\beta}^{ff'X} (\bar{\nu}_\alpha \gamma^\mu P_L \ell_\beta) (\bar{f}' \gamma_\mu P_X f), \quad f \neq f' \in (u, d)$$

$$\mathcal{L}_{\text{NC-NSI}} = -2\sqrt{2}G_F \epsilon_{\alpha\beta}^{fX} (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) (\bar{f} \gamma_\mu P_X f), \quad f \in (e, u, d)$$



Theory background

Wolfenstein (1978), Barger et al. (1991), Grossman (1995), Ohlsson (2012)...

$$\epsilon_{\alpha\beta} \equiv \epsilon_{\alpha\beta}^{eV} + \frac{N_u}{N_e} \epsilon_{\alpha\beta}^{uV} + \frac{N_d}{N_e} \epsilon_{\alpha\beta}^{dV}$$

1. Osc. expts. (DUNE, T2HK, NOvA, INO)
2. Scattering expts. (COHERENT..)

Inside sun: $\frac{N_u}{N_e} \approx 2 \frac{N_d}{N_e} \approx 1$

Inside earth: $\frac{N_u}{N_e} \approx \frac{N_d}{N_e} \approx 3$

$$H = H_{vac} + H_{mat} + H_{NSI}$$

$$i \frac{d}{dt} \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \frac{1}{2E} \left[U \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^\dagger + A \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix} \right] \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

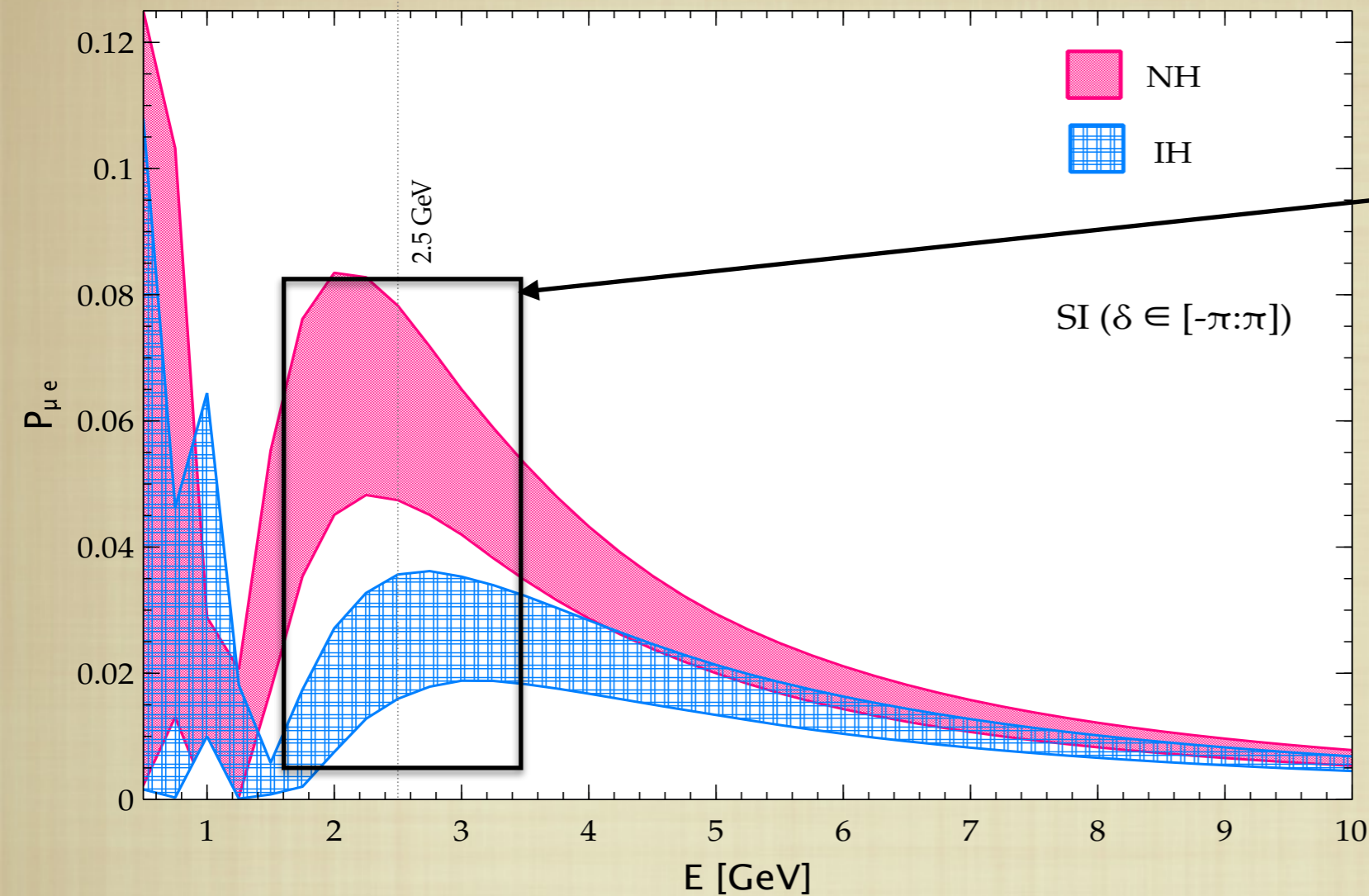
NSI at long baseline experiment

- NSI amplitudes with their new CP phases can potentially spoil CPV, MH and octant sensitivities
- Most relevant for $\nu_\mu \rightarrow \nu_e$: $\varepsilon_{e\mu}, \varepsilon_{e\tau}, \varepsilon_{ee}$

$$\nu_\mu \rightarrow \nu_\mu : \varepsilon_{\mu\mu}, \varepsilon_{\mu\tau}$$

$$i \frac{d}{dt} \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \frac{1}{2E} \left[U \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^\dagger + A \begin{pmatrix} 1 + \varepsilon_{ee} & \varepsilon_{e\mu} & \varepsilon_{e\tau} \\ \varepsilon_{e\mu}^* & \varepsilon_{\mu\mu} & \varepsilon_{\mu\tau} \\ \varepsilon_{e\tau}^* & \varepsilon_{\mu\tau}^* & \varepsilon_{\tau\tau} \end{pmatrix} \right] \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

NSI at long baseline experiment (MH)



high sensitivity for SI

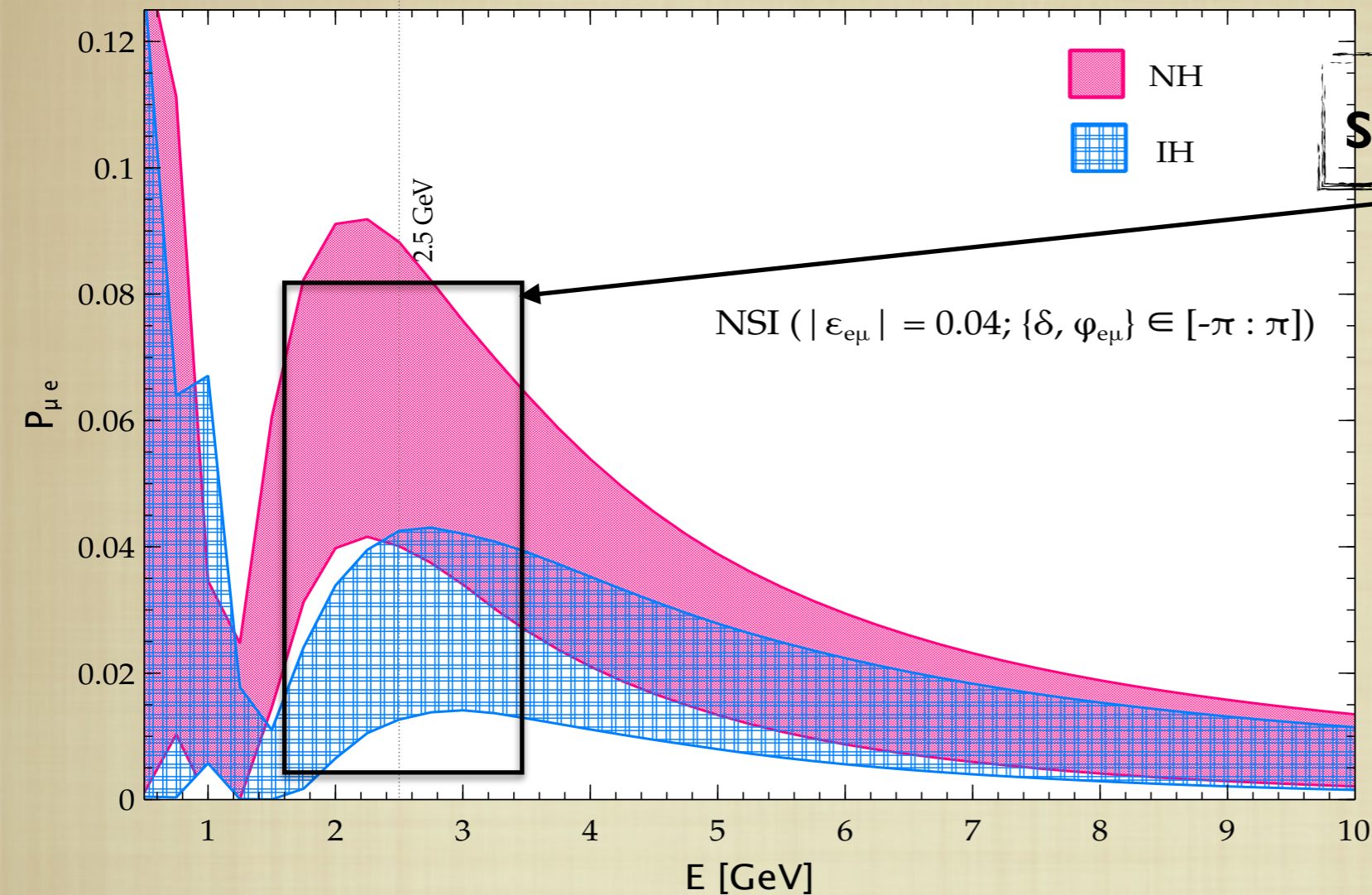
changes signs for IH

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 \end{aligned}$$

$$\begin{aligned}
 A &= \frac{2\sqrt{2}EG_F n_E}{\Delta m_{31}^2} \\
 \Delta &= \frac{\Delta m_{31}^2 L}{4E} \\
 \alpha &= \frac{\Delta m_{21}^2}{\Delta m_{31}^2}
 \end{aligned}$$

MM, Mehta (2016)

NSI at long baseline experiment (MH)



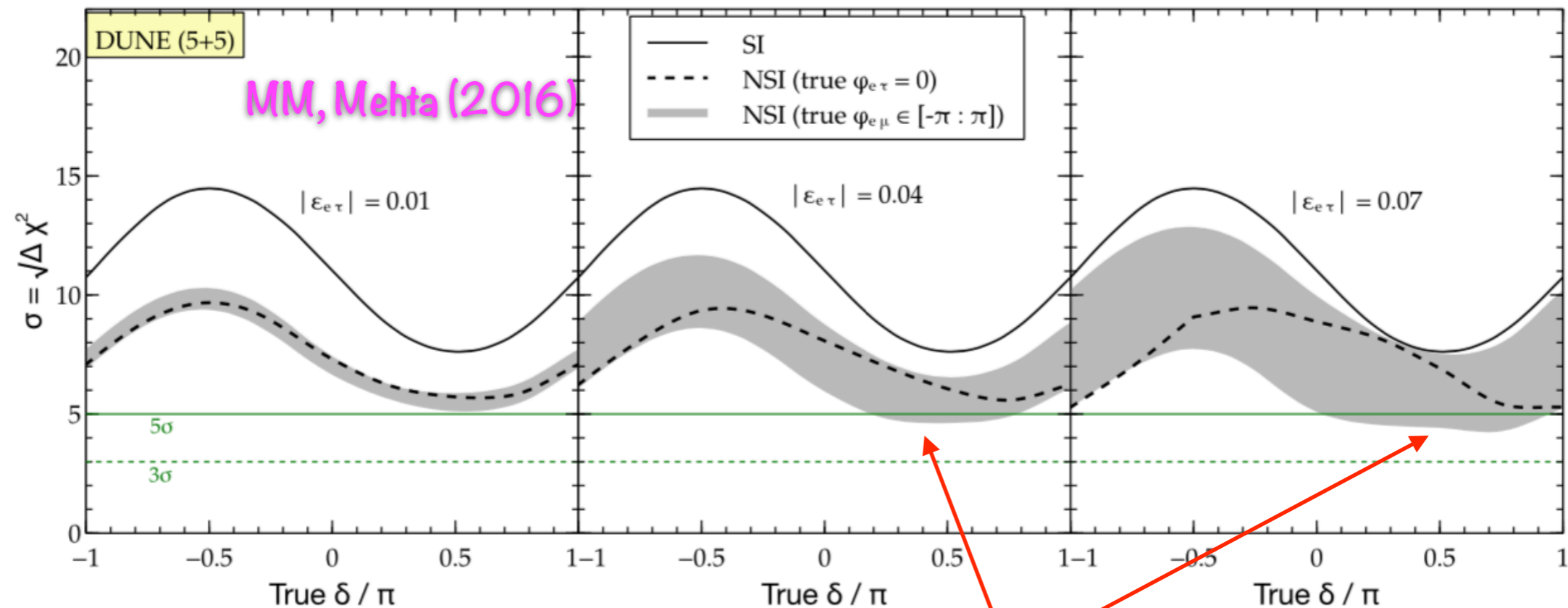
sensitivity decreases for NSI

changes signs for IH

$$\begin{aligned}
 P_{\mu e} = & \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2(1-A)\Delta}{1-A} \\
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NSI at long baseline experiment (MH)

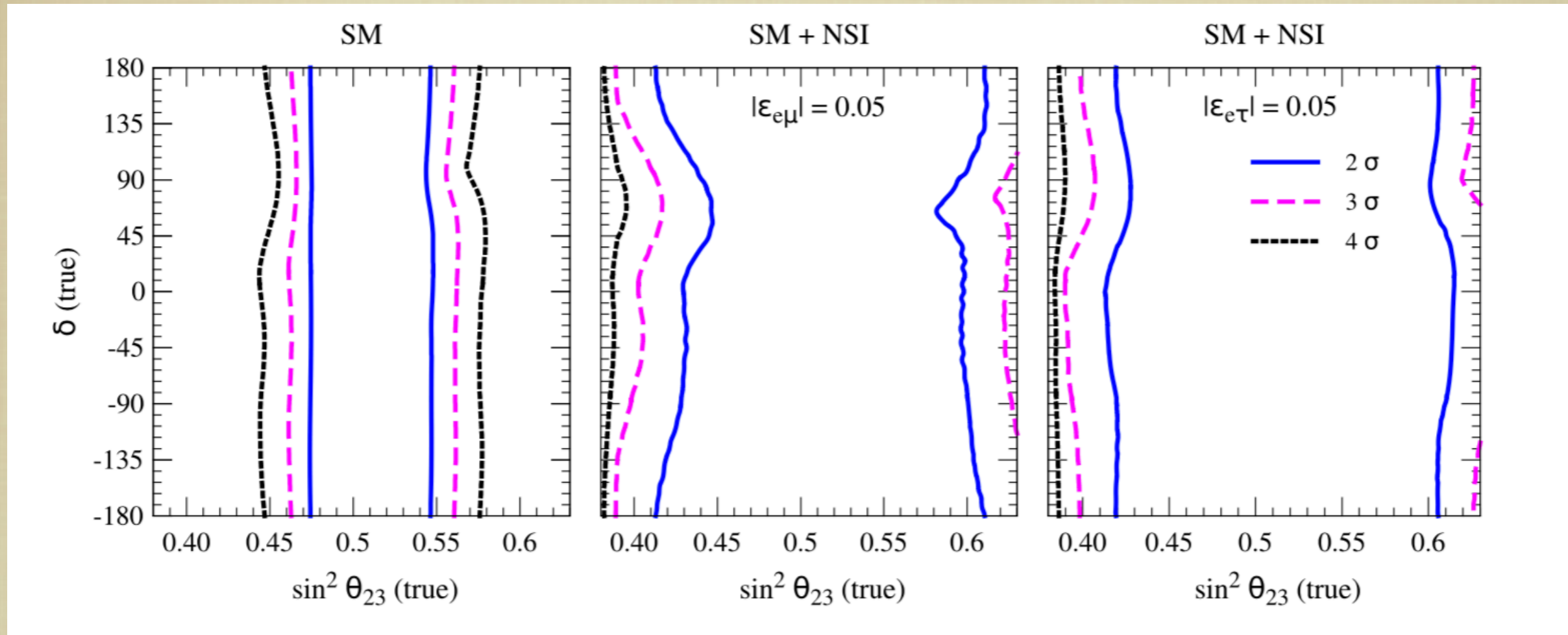


MH sensitivity < 5sigma

MH sensitivity at DUNE

Also see, Soumya, Mohanta (2017);
Deepthi, Goswami, Nath (2017)

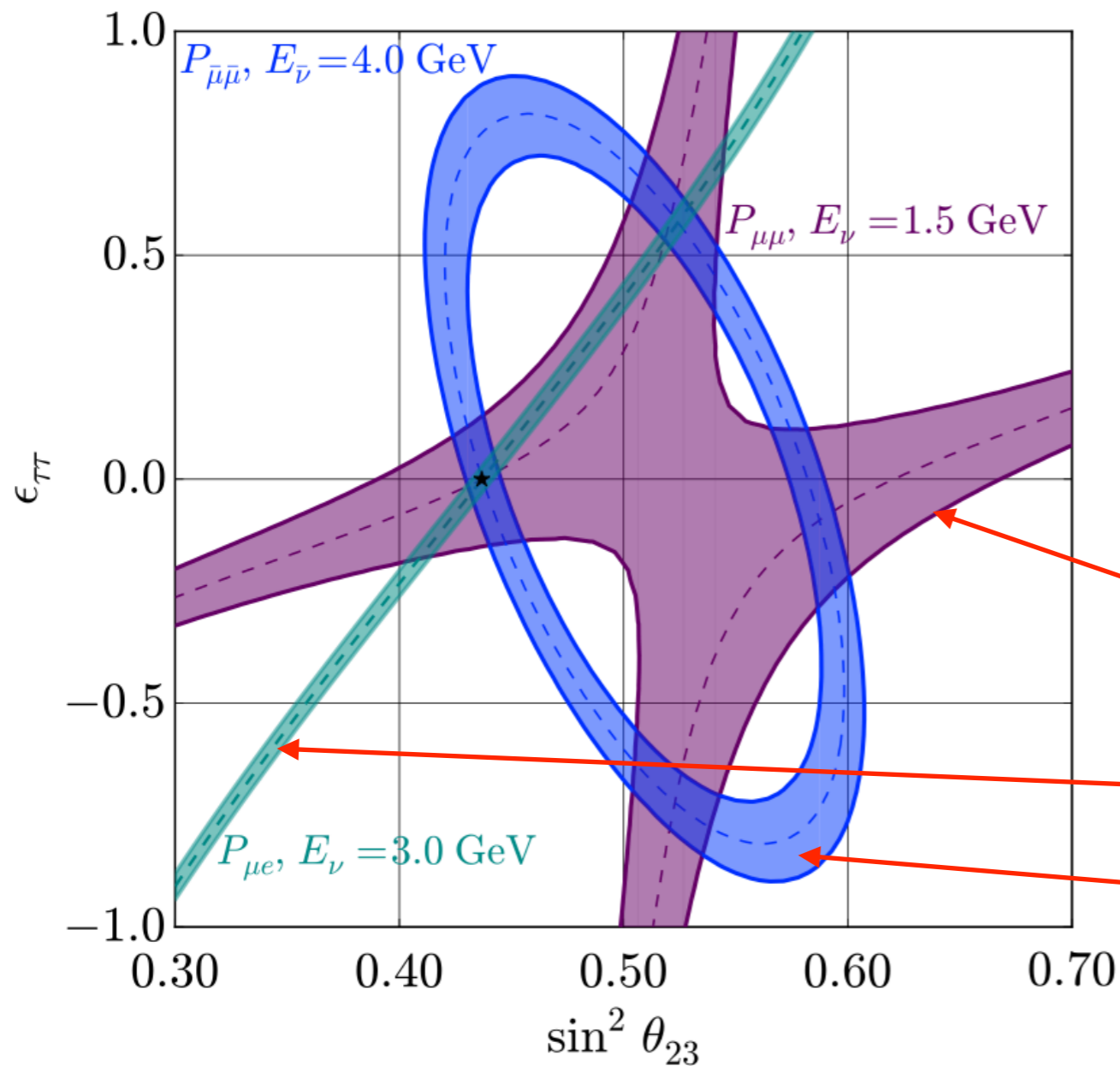
NSI at long baseline experiment (Octant)



Agarwalla, Chatterjee, Palazzo (2016)

$\theta_{23} - \delta$ exclusion region at DUNE

NSI at long baseline experiment (SI-NSI degeneracy)



De Gouvea, Kelly (2016)

Degeneracy in:

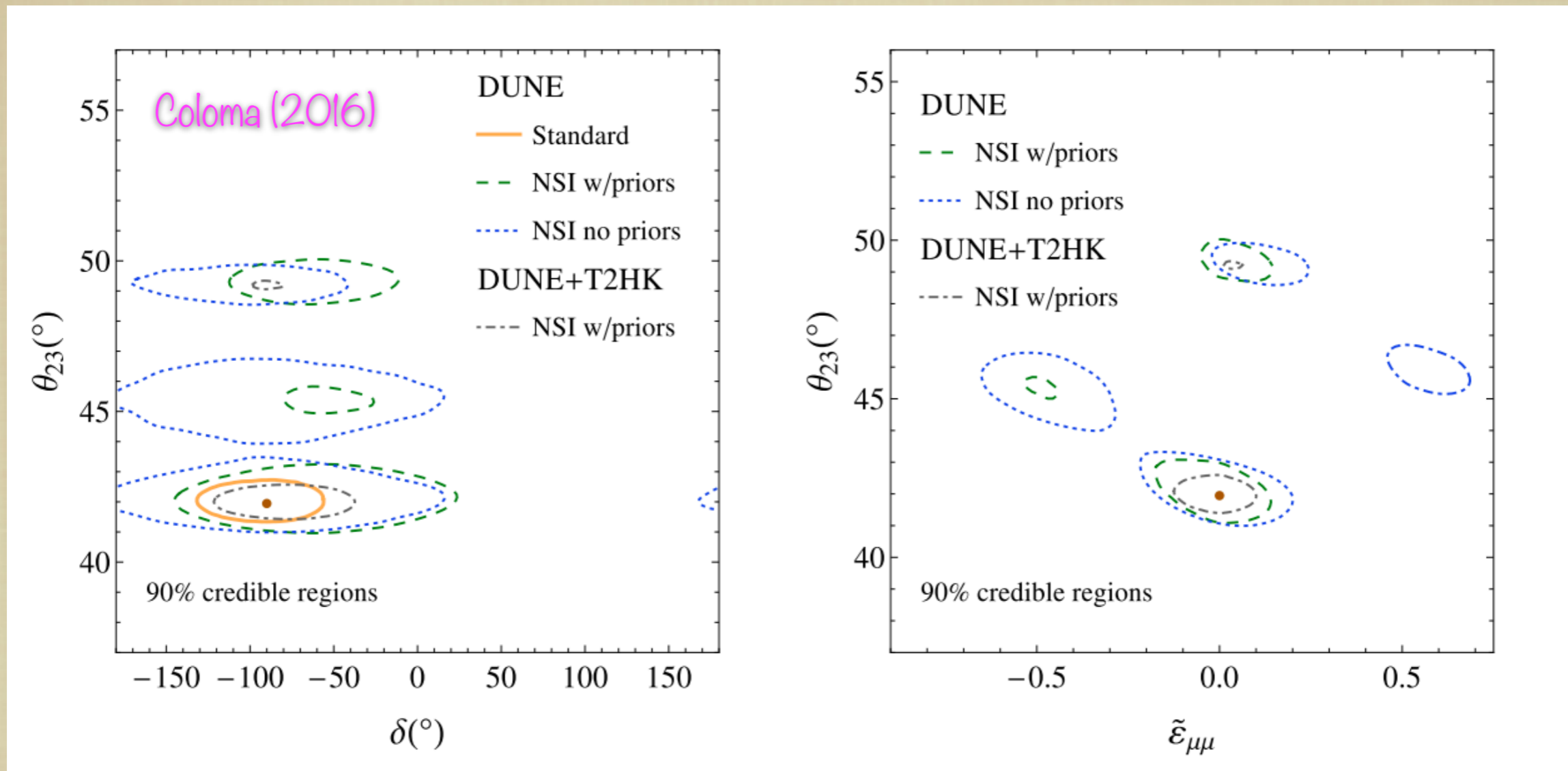
$P_{\mu\mu}$

$P_{\mu e}$

$\bar{P}_{\mu\mu}$

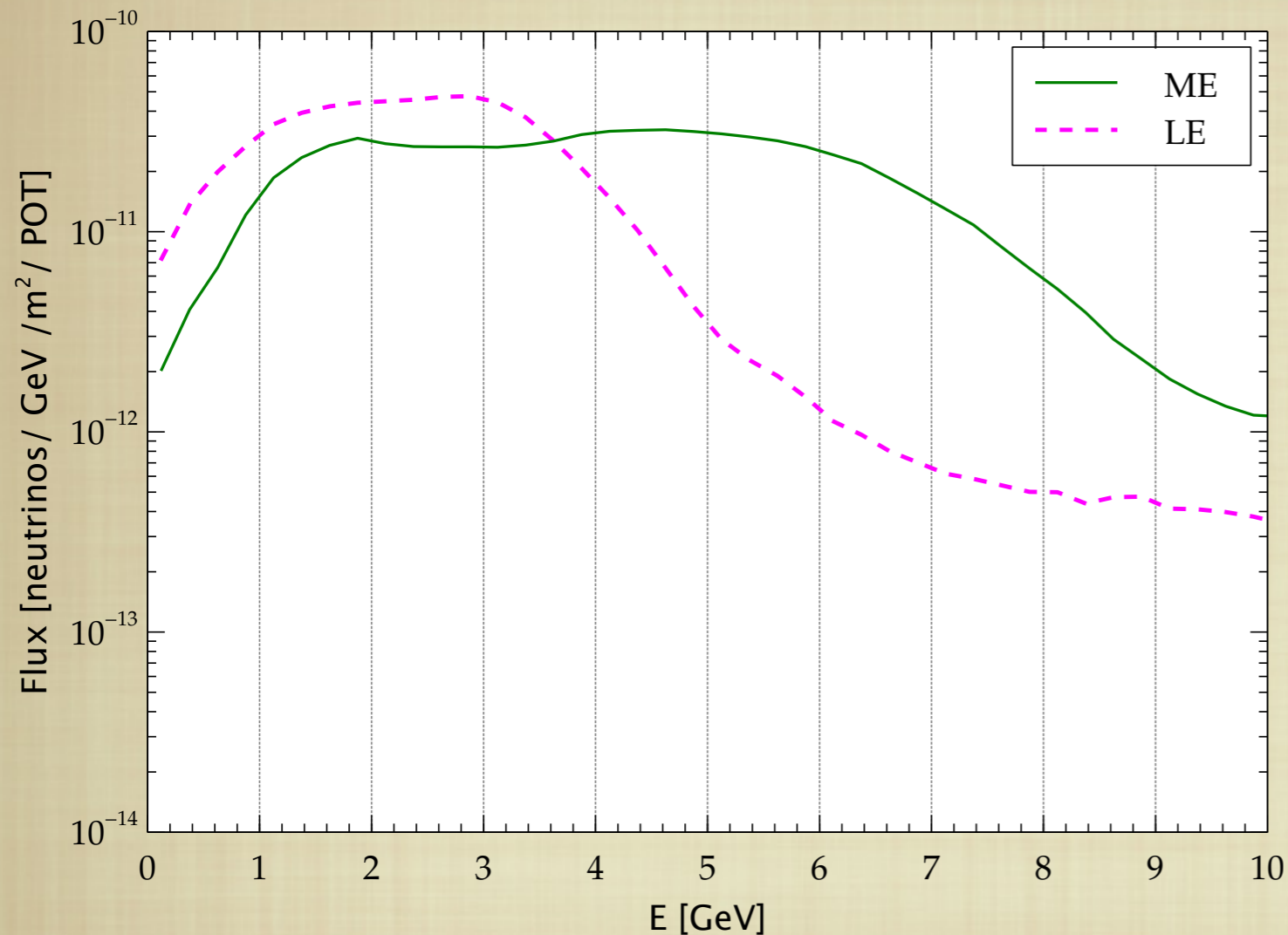
Also see [Liao, Marfatia, Whisnant \(2017\)](#) for degeneracy in the χ^2 level

SI-NSI degeneracy and priors



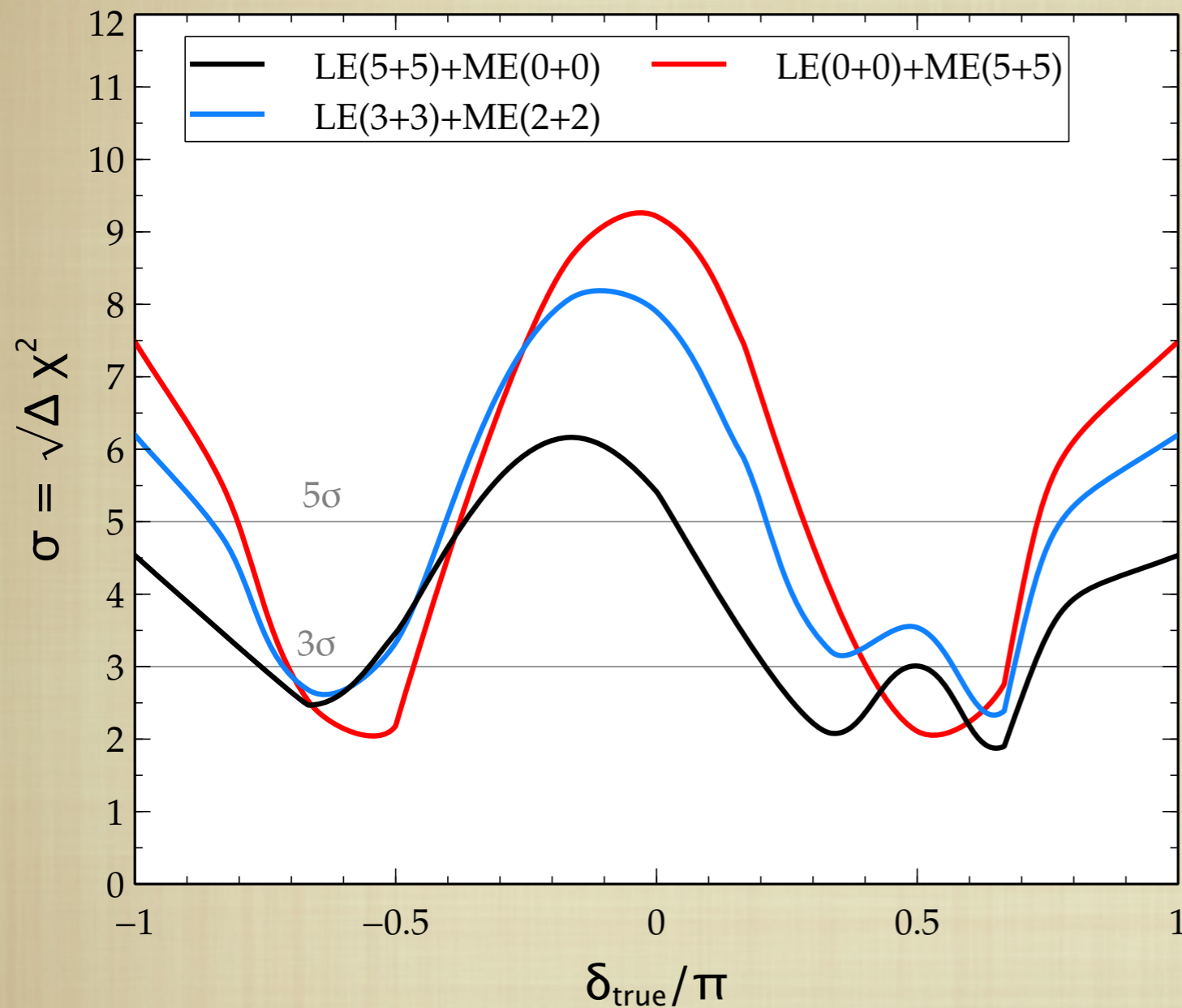
- Using priors and adding T2HK data help to reduce degeneracy

Can we distinguish NSI from SI?



Standard (LE) flux falls quickly!
 \implies Exploit a medium energy (ME) tuned flux

Can we distinguish NSI from SI?



data simulated assuming

$$\varepsilon_{e\mu} = 0.05$$

Significant improvement of
separability

NSI Global analysis (complementarity of experiments)

NSI WITH QUARKS

ϵ_{ee}^{dL}	[-0.3, 0.3]	CHARM
ϵ_{ee}^{dR}	[-0.6, 0.5]	CHARM
ϵ_{ee}^{dV}	[0.030, 0.55]	Oscillation data + COHERENT
ϵ_{ee}^{uV}	[0.028, 0.60]	Oscillation data + COHERENT
$\epsilon_{\mu\mu}^{dV}$	[-0.042, 0.042]	Atmospheric + accelerator
$\epsilon_{\mu\mu}^{uV}$	[-0.044, 0.044]	Atmospheric + accelerator
$\epsilon_{\mu\mu}^{dA}$	[-0.072, 0.057]	Atmospheric + accelerator
$\epsilon_{\mu\mu}^{uA}$	[-0.094, 0.14]	Atmospheric + accelerator
$\epsilon_{\tau\tau}^{dV}$	[-0.075, 0.33]	Oscillation data + COHERENT
$\epsilon_{\tau\tau}^{uV}$	[-0.09, 0.38]	Oscillation data + COHERENT
$\epsilon_{\tau\tau}^{qV}$	[-0.037, 0.037]	Atmospheric

NSI WITH ELECTRONS

ϵ_{ee}^{eL}	[-0.021, 0.052]	Solar + KamLAND
ϵ_{ee}^{eR}	[-0.07, 0.08]	TEXONO
$\epsilon_{\mu\mu}^{eL}, \epsilon_{\mu\mu}^{eR}$	[-0.03, 0.03]	Reactor + accelerator
$\epsilon_{\tau\tau}^{eL}$	[-0.12, 0.06]	Solar + KamLAND
$\epsilon_{\tau\tau}^{eR}$	[-0.98, 0.23]	Solar + KamLAND and Borexino
	[-0.25, 0.43]	Reactor + accelerator
$\epsilon_{\tau\tau}^{eV}$	[-0.11, 0.11]	Atmospheric

Biggio, Blennow, Fernandez-Martinez (2009);
Gonzalez-Garcia, Maltoni, Salvado;
Tortola, Farzan (2018)

Diagonal NSI

NSI Global analysis (complementarity of experiments)

NSI WITH QUARKS

$\epsilon_{e\mu}^{qL}$	$[-0.023, 0.023]$	Accelerator
$\epsilon_{e\mu}^{qR}$	$[-0.036, 0.036]$	Accelerator
$\epsilon_{e\mu}^{uV}$	$[-0.073, 0.044]$	Oscillation data + COHERENT
$\epsilon_{e\mu}^{dV}$	$[-0.07, 0.04]$	Oscillation data + COHERENT
$\epsilon_{e\tau}^{qL}, \epsilon_{e\tau}^{qR}$	$[-0.5, 0.5]$	CHARM
$\epsilon_{e\tau}^{uV}$	$[-0.15, 0.13]$	Oscillation data + COHERENT
$\epsilon_{e\tau}^{dV}$	$[-0.13, 0.12]$	Oscillation data + COHERENT
$\epsilon_{\mu\tau}^{qL}$	$[-0.023, 0.023]$	Accelerator
$\epsilon_{\mu\tau}^{qR}$	$[-0.036, 0.036]$	Accelerator
$\epsilon_{\mu\tau}^{qV}$	$[-0.006, 0.0054]$	IceCube
$\epsilon_{\mu\tau}^{qA}$	$[-0.039, 0.039]$	Atmospheric + accelerator

NSI WITH ELECTRONS

$\epsilon_{e\mu}^{eL}, \epsilon_{e\mu}^{eR}$	$[-0.13, 0.13]$	Reactor + accelerator
$\epsilon_{e\tau}^{eL}$	$[-0.33, 0.33]$	Reactor + accelerator
$\epsilon_{e\tau}^{eR}$	$[-0.28, -0.05] \& [0.05, 0.28]$	Reactor + accelerator
	$[-0.19, 0.19]$	TEXONO
$\epsilon_{\mu\tau}^{eL}, \epsilon_{\mu\tau}^{eR}$	$[-0.10, 0.10]$	Reactor + accelerator
$\epsilon_{\mu\tau}^{eV}$	$[-0.018, 0.016]$	IceCube

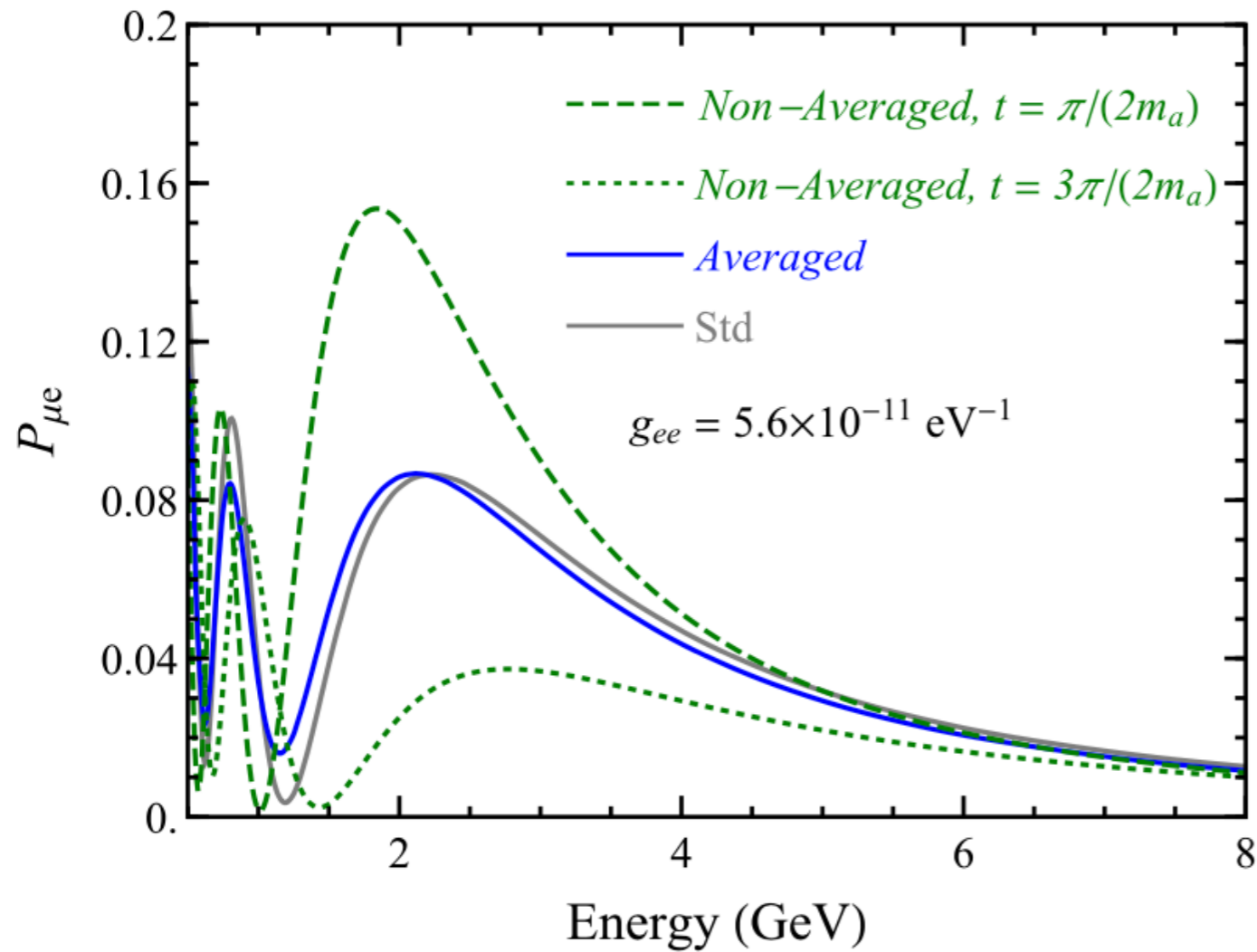
Biggio, Blennow, Fernandez-Martinez (2009);
Gonzalez-Garcia, Maltoni, Salvado;
Tortola, Farzan (2018)

Off-diagonal NSI

Plan of the talk

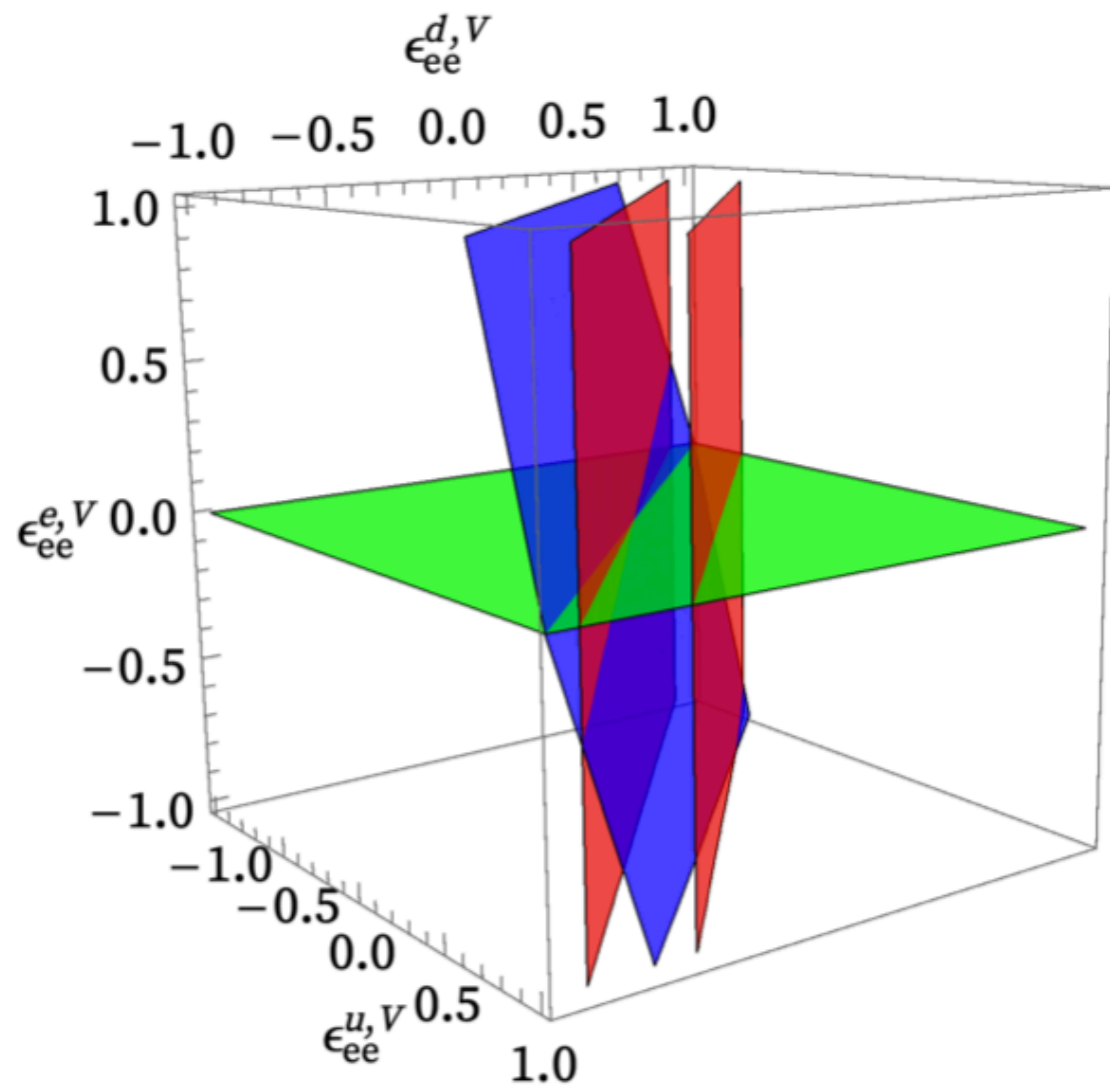
- Basics of oscillation physics and pending issues
- Hint for a light (eV) sterile neutrino
- Non-Standard Neutrino Interaction (NSI)
- Other BSM options (Dark sector?)
- Summary

Neutrinophilic ALP



$$H(t) = \frac{1}{2E_\nu} U \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^\dagger + \begin{pmatrix} V + \xi_{ee}(t) & \xi_{e\mu}(t) & \xi_{e\tau}(t) \\ \xi_{e\mu}^*(t) & \xi_{\mu\mu}(t) & \xi_{\mu\tau}(t) \\ \xi_{e\tau}^*(t) & \xi_{\mu\tau}^*(t) & \xi_{\tau\tau}(t) \end{pmatrix}$$

NSI: Complementarity of detectors



CE ν NS, E ν NS, Oscillation

$$\epsilon_{\alpha\beta} = \epsilon_{\alpha\beta}^{e,V} + 3\epsilon_{\alpha\beta}^{u,V} + 3\epsilon_{\alpha\beta}^{d,V}$$

Dutta, Lang, Liao, Sinha, Strigari, Thompson (2020)

Non-unitarity

- Neutral Heavy Leptons (NHL) arise naturally as an extension of SM \implies Non-unitary mixing matrix
- Can induce Lepton Flavour Violation
- Modification of CP sensitivity due to more unknown phases
-

Antusch, Biggio, Fernandez-Martinez, Gavela, Lopez-Pavon : JHEP (2006)

Forero, Morisi, Tortola, Valle : JHEP (2011)

Blennow, Coloma, Hernandez-Garcia, Lopez-Pavon : JHEP (2017)

Summary

- Outstanding issues: CPV, mass ordering, octant

- Tension between T2K/NOvA
- Stay tuned for DUNE/T2HK...

- Hint for new physics (MiniBoone, LSND..)

- MB Excess $\gtrsim 4\sigma$
- Wait for SBN
- Compatible with eV sterile ν !
- DM sector? Decay of heavy ν ?.....

- Sterile neutrinos, NSI,

- Spoils std. param. measurements
- Needs more constraints from osc. (SBN, Neutrino-4, atmos expt....) & nosc. data (CE ν NS)

- Other possibilities (DM sector?...)

- Exploit the complementarity of DM and ν detectors

ধন্যবাদ!

Thank you!

감사합니다

3+1 case: Impact on CPV sensitivity

$$\Delta\chi^2(true) \simeq \text{Min}_{u,s} \left[\sum_{bins} \frac{\{N_{data}(true) - N_{fit}(u, s)\}^2}{N_{data}(true)} + \sum \frac{s_i^2}{\sigma_i^2} \right] \quad \begin{array}{l} u : \text{osc. parameters} \\ s : \text{systematics} \end{array}$$

- GLoBES simulation \implies DUNE data for all true CP phase(s) $\in [-\pi, \pi]$
- Fitted with CP phase(s) = $0, \pi$
- Marginalisation over $u(= \{ \theta_{23}, \theta_{14}, \theta_{24}, \theta_{34}, | \Delta m_{31}^2 | \})$ in the allowed ranges and also over the various systematics(s)
- Small $\Delta\chi^2 \implies$ Better fit with data;
large $\Delta\chi^2 \implies$ fit is less compatible with data

Theory background (backup)

- $|\nu_\alpha\rangle = \sum_k U_{\alpha k}^* |\nu_k\rangle$ ($\alpha = e, \nu, \tau$)
- $H |\nu_k\rangle = E_k |\nu_k\rangle$
- $i \frac{d}{dt} |\nu_k(t)\rangle = H |\nu_k(t)\rangle$
- $P_{\alpha\beta}(L, E) = |\langle \nu_\beta | \nu_\alpha(x) \rangle|^2$
 $= \delta_{\alpha\beta} - 4 \sum_{k>j} \text{Re}[U_{\alpha k}^* U_{\beta j}^* U_{\beta k} U_{\alpha j}] \sin^2 \Delta_{kj}$
 $+ \sum_{k>j} \text{Im}[U_{\alpha k}^* U_{\beta j}^* U_{\beta k} U_{\alpha j}] \sin 2\Delta_{kj},$

(where, $\Delta_{ij} = 1.27 \times \frac{\Delta m_{ij}^2 [\text{eV}^2] \times L [\text{km}]}{E [\text{GeV}]}$)

Theory background (backup)

Kosteletzky et al. (2012), Mavromatos et al.

$$\begin{aligned}\hat{Q}_{AB} &= \sum_I \hat{Q}_{AB}^I \gamma_I \\ &= \hat{S}_{AB} + i\hat{P}_{AB}\gamma_5 + \hat{V}_{AB}^\mu \gamma_\mu + \hat{A}_{AB}^\mu \gamma_5 \gamma_\mu + \frac{1}{2} \hat{T}_{AB}^{\mu\nu} \sigma_{\mu\nu},\end{aligned}$$

$$\gamma^\nu p_\nu \delta_{AB} - M_{AB} + \hat{Q}_{AB} = \hat{\Gamma}_{AB}^\nu p_\nu - \hat{M}_{AB},$$

$$\begin{aligned}\hat{\Gamma}_{AB}^\nu &= \gamma^\nu \delta_{AB} + \hat{c}_{AB}^{\mu\nu} \gamma_\mu + \hat{d}_{AB}^{\mu\nu} \gamma_5 \gamma_\mu \\ &\quad + \hat{e}_{AB}^\nu + i\hat{f}_{AB}^\nu \gamma_5 + \frac{1}{2} \hat{g}_{AB}^{\kappa\lambda\nu} \sigma_{\kappa\lambda}, \\ \hat{M}_{AB} &= m_{AB} + im_{5AB} \gamma_5 + \hat{m}_{AB} + i\hat{m}_{5AB} \gamma_5 \\ &\quad + \hat{a}_{AB}^\mu \gamma_\mu + \hat{b}_{AB}^\mu \gamma_5 \gamma_\mu + \frac{1}{2} \hat{H}_{AB}^{\mu\nu} \sigma_{\mu\nu}.\end{aligned}$$

Theory background (backup)

Kostelecky et al. (2012), Mavromatos et al.

- The effective hamiltonian can be decomposed into 3x3 blocks:

$$Q = S + iP\gamma_5 + \mathcal{V}^\alpha \gamma_\alpha + \mathcal{A}^\alpha \gamma_5 \gamma_\alpha + \frac{1}{2} \mathcal{T}^{\alpha\beta} \sigma_{\alpha\beta}$$

Induce effective Hamiltonian for neutrino mixing

LIV: Theory background (backup)

Kostelecky et al. (2012), Mavromatos et al.

- The effective hamiltonian can be decomposed into 3x3 blocks:

$$(h_{eff})_{ab} = E\delta_{ab} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \frac{1}{2E} \begin{pmatrix} (m^2)_{ab} & 0 \\ 0 & (m^2)_{ab}^* \end{pmatrix} + \frac{1}{E} \begin{pmatrix} [(a_L)^\mu p_\mu - (c_L)^{\mu\nu} p_\mu p_\nu]_{ab} & 0 \\ 0 & [- (a_L)^\mu p_\mu - (c_L)^{\mu\nu} p_\mu p_\nu]_{ab}^* \end{pmatrix}$$

CPT odd LIV parameters: $((a_L)_{ab}^\alpha)^* = - (a_R)_{\bar{a}\bar{b}}^\alpha$

CPT even LIV parameters: $((c_L)_{ab}^{\alpha\beta})^* = (c_R)_{\bar{a}\bar{b}}^{\alpha\beta}$

Theory background (backup)

Kostelecky et al. (2012), Mavromatos et al.

- The effective hamiltonian can be decomposed into 3x3 blocks:

$\nu - \nu$ mixing \rightarrow

$$h_{ab} = E\delta_{ab} + \frac{m_{ab}^2}{2E} + (a_L)_{ab}^\alpha p_\alpha - (c_L)_{ab}^{\alpha\beta} p_\alpha p_\beta E$$

$\bar{\nu} - \bar{\nu}$ mixing \rightarrow

$$h_{\bar{a}\bar{b}} = E\delta_{\bar{a}\bar{b}} + \frac{m_{\bar{a}\bar{b}}^2}{2E} + (a_R)_{\bar{a}\bar{b}}^\alpha p_\alpha - (c_R)_{\bar{a}\bar{b}}^{\alpha\beta} p_\alpha p_\beta E$$

$\nu - \bar{\nu}$ mixing \rightarrow

$$h_{a\bar{b}} = i\sqrt{2}(\epsilon)_\alpha (H_{a\bar{b}}^\alpha - g_{a\bar{b}}^{\alpha\beta} p_\beta E)$$

Backup

$$\begin{aligned}\Delta P_{\mu e}(\varepsilon_{e\mu}) &= P_{\mu e}^{NSI}(\varepsilon_{e\mu}) - P_{\mu e}^{SI} \\ &\approx -4A\Delta \sin \Delta |\varepsilon_{e\mu}| s_{13} s_{2(23)} c_{23} D_1^{e\mu} \sin(\delta + \varphi_{e\mu} - \gamma_1^{e\mu})\end{aligned}$$

&

$$\Delta P_{\mu e}(\varepsilon_{e\tau}) \approx 4A\Delta \sin \Delta |\varepsilon_{e\tau}| s_{13} s_{2(23)} s_{23} D_1^{e\tau} \sin(\delta + \varphi_{e\tau} + \gamma_1^{e\tau})$$

where,

$$\begin{aligned}D_1^{e\mu} &= \left[\sin^2 \Delta + \left(\tan^2 \theta_{23} \frac{\sin \Delta}{\Delta} + \cos \Delta \right)^2 \right]^{1/2} & \gamma_1^{e\mu} &= \tan^{-1} \left(\frac{\tan^2 \theta_{23}}{\Delta} + \cot \Delta \right) \\ D_1^{e\tau} &= \left[\sin^2 \Delta + \left(\frac{\sin \Delta}{\Delta} - \cos \Delta \right)^2 \right]^{1/2}; & \gamma_1^{e\tau} &= \tan^{-1} \left(\frac{1}{\Delta} - \cot \Delta \right)\end{aligned}$$

3+1 case :basics (backup)

$$U_{\text{PMNS}}^{3+1} = V(\theta_{34}, \delta_{34})V(\theta_{24}, \delta_{24})R(\theta_{14})R(\theta_{23})V(\theta_{13}, \delta_{13})R(\theta_{12})$$

where,

$$V(\theta_{24}, \delta_{24}) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos \theta_{24} & 0 & e^{-i\delta_{24}} \sin \theta_{24} \\ 0 & 0 & 1 & 0 \\ 0 & -e^{i\delta_{24}} \sin \theta_{24} & 0 & \cos \theta_{24} \end{pmatrix} \text{ etc.}$$

- $\theta_{14} \in [0^\circ, 11^\circ]$ [DayaBay: PRL 113(2014) 141802]
 $\theta_{24} \in [0^\circ, 7^\circ]$ [IceCube: PRL 117(2016) 7, 071801; Ben Jones' talk (2016)]
 $\theta_{34} \in [0^\circ, 26^\circ]$ [MINOS: PRL 107(2011) 1, 011802]
- $\delta_{13} \in [-\pi, \pi], \delta_{24} \in [-\pi, \pi], \delta_{34} \in [-\pi, \pi]$

- Can the new phases and mixing angles play roles to affect CPV studies?

Background and motivation: CPT violation?

- Observation of P violation in weak interactions
C.N. Yang & T. D. Lee (Nobel 1957)
V-A theory: Weak interaction acts on LH particles and RH antiparticles
Marshak & Sudarshan; Feynman & Gell-Mann
- Although rare, CP violation was also observed
Kobayashi & Maskawa (Nobel 2008)
- What about CPT violation?
Can be related to , for e.g., Non-unitarity, Lorentz Invariance Violation (LIV)