

Probing dark matter minihalos with extremely magnified extragalactic stars

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Dark Matter as a Portal to New Physics

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(QCD) Axion

The strong CP problem: no evidence for CP violation from QCD 't Hooft 1976

A famous solution: Peccei-Quinn symmetry breaking

Peccei & Quinn 1977; Weinberg 1978; Wilczek 1978

The **axion** particle:

A very light scalar particle;

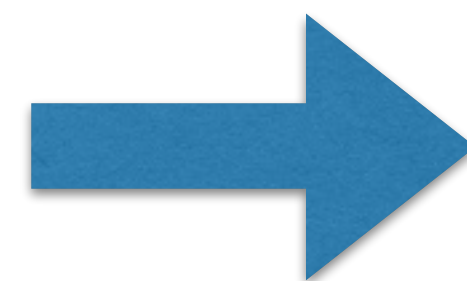
Weakly interacting with standard model particles;

Dynamically drive CP violating term to zero

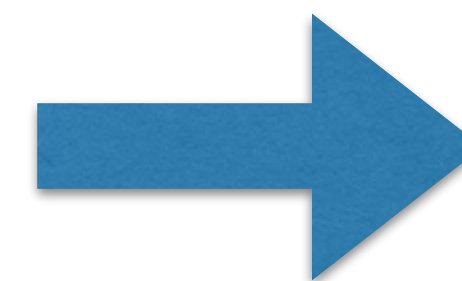
Bonus: axion particles can be the **Cold Dark Matter** (non-thermal relic) $10^{-11} \text{ eV} < m_a < 1 \text{ eV}$

(Not “fuzzy” DM)

PQ symmetry breaking to
incoherent horizon patches
(topological defects)

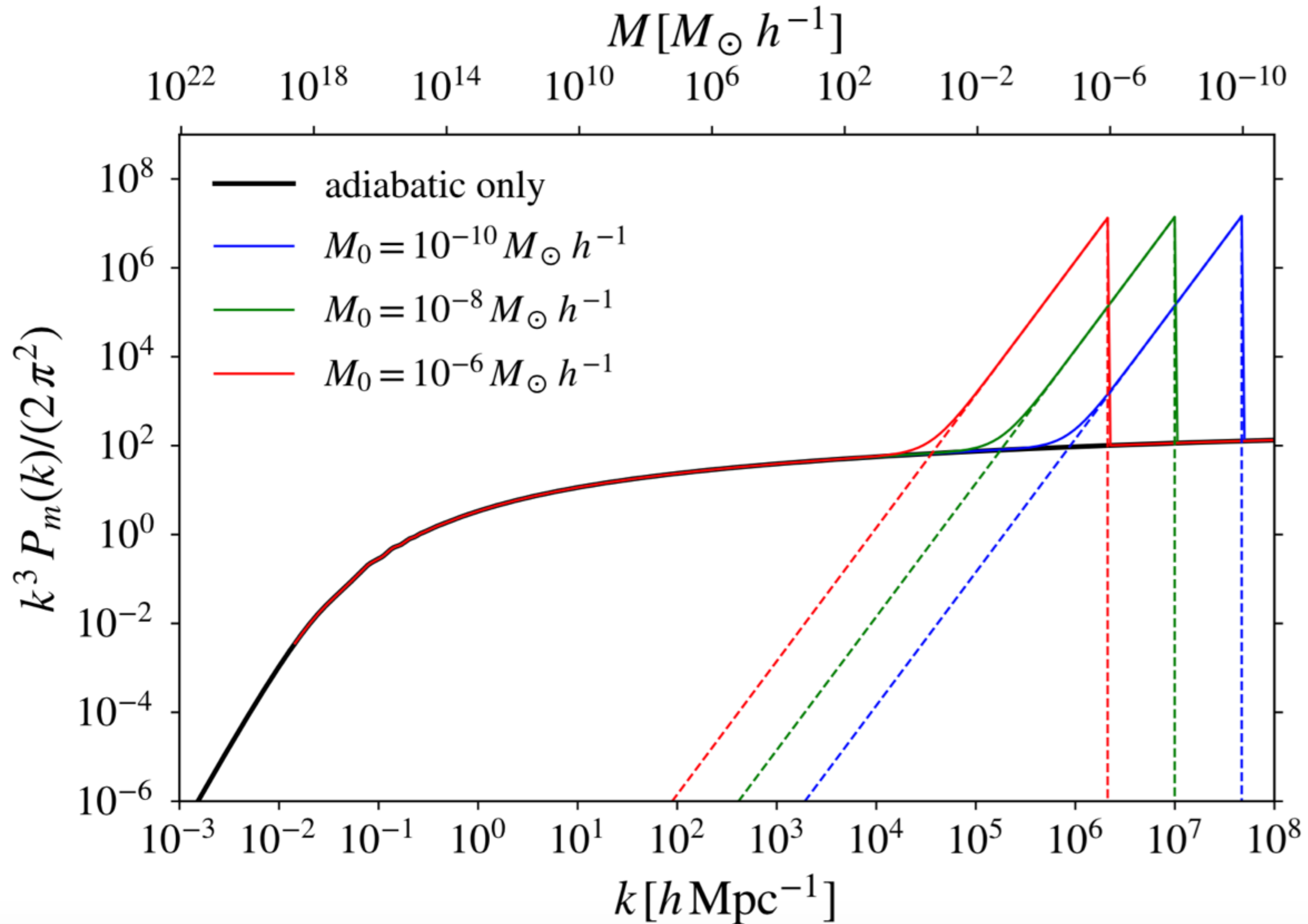


non-negligible mass as
the universe cools



start oscillation at t_0 ;
contribute to CDM

Axion Cosmology: Isocurvature Fluctuations



Initial density fluctuations:

Adiabatic modes from inflation

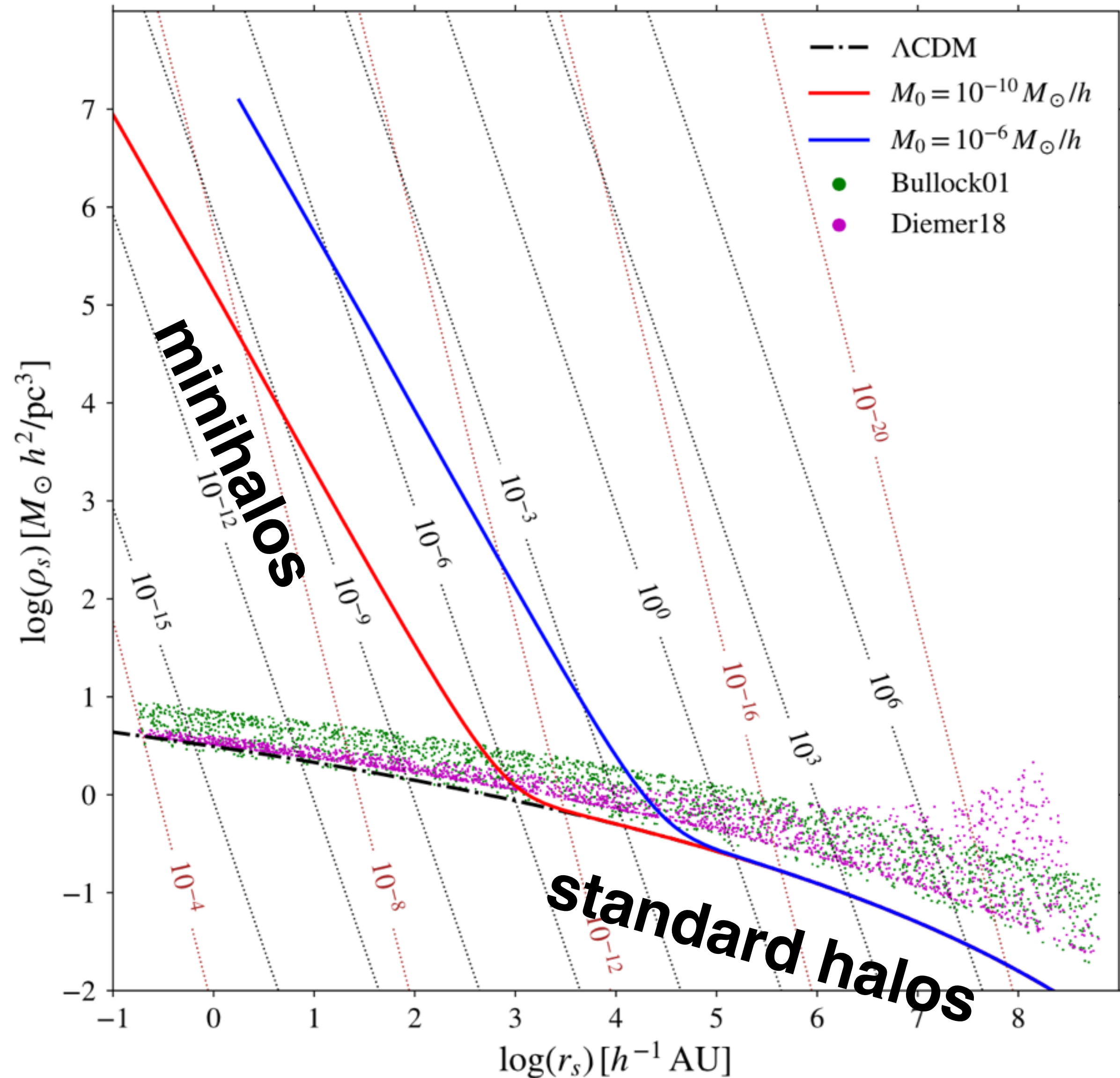
Isocurvature modes from topological defects

$$M_0 = \frac{4\pi}{3} \left(\frac{\pi}{k_0} \right)^3 \bar{\rho}_{m0}$$

(probable values as implied by calculation of $m_a(T)$)

Hogan & Rees 1988
 Kolb & Tkachev 1994

Axion Cosmology: Minihalos



NFW density profile

Navarro, Frenk & White 96'; 97'

$$\rho(r) = \frac{\rho_s}{\left(\frac{r}{r_s}\right) \left(1 + \frac{r}{r_s}\right)^2}$$

Large halos same as in standard cosmology.

Small halos collapse earlier in axion cosmology ($10 < z < 4000$).

They are **denser and more compact** than in standard cosmology.

Axion Minihalos

Xiao, Williams & McQuinn 2101.04177

N-body simulation of axion minihalo formation and hierarchical assembly

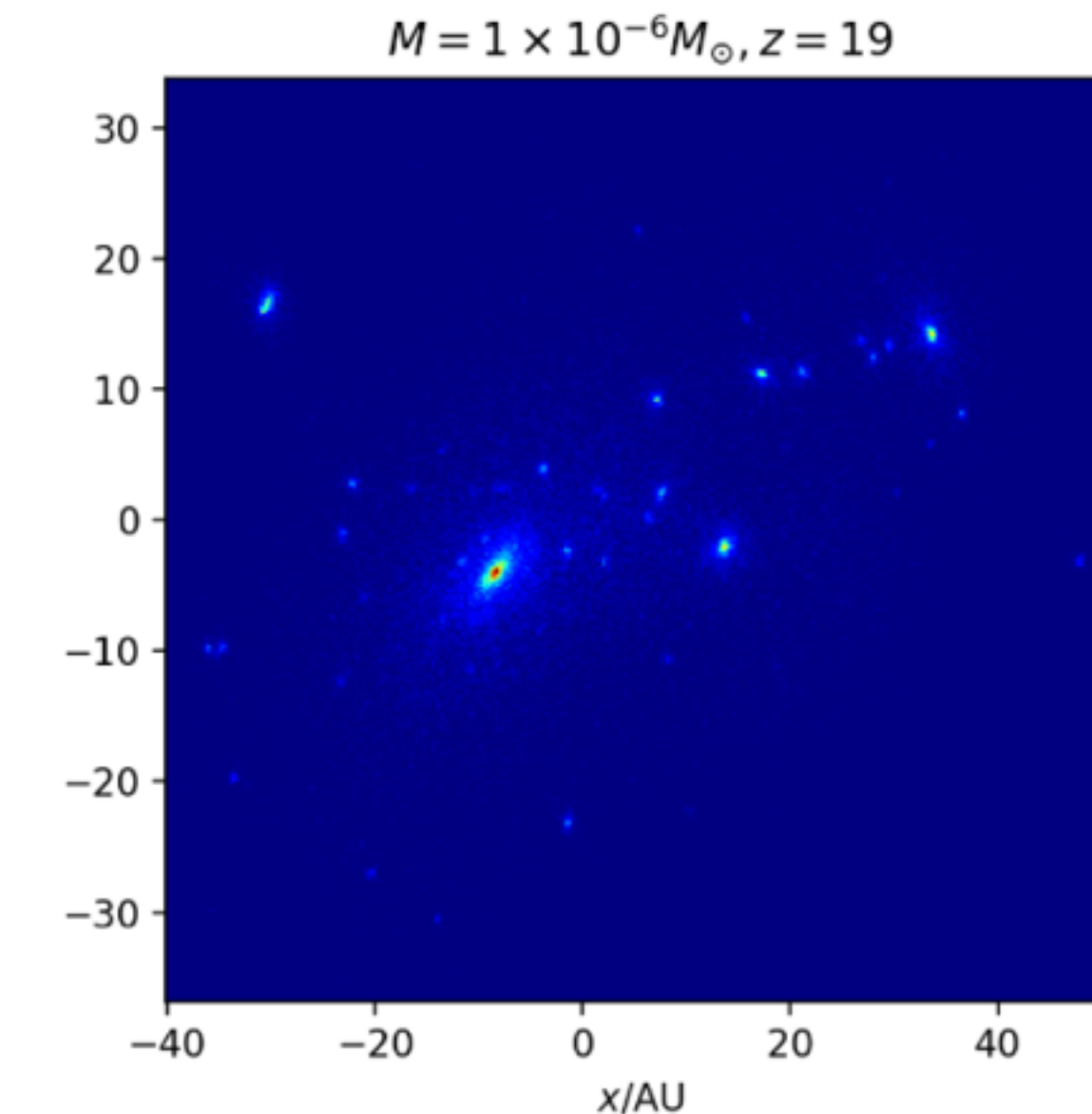
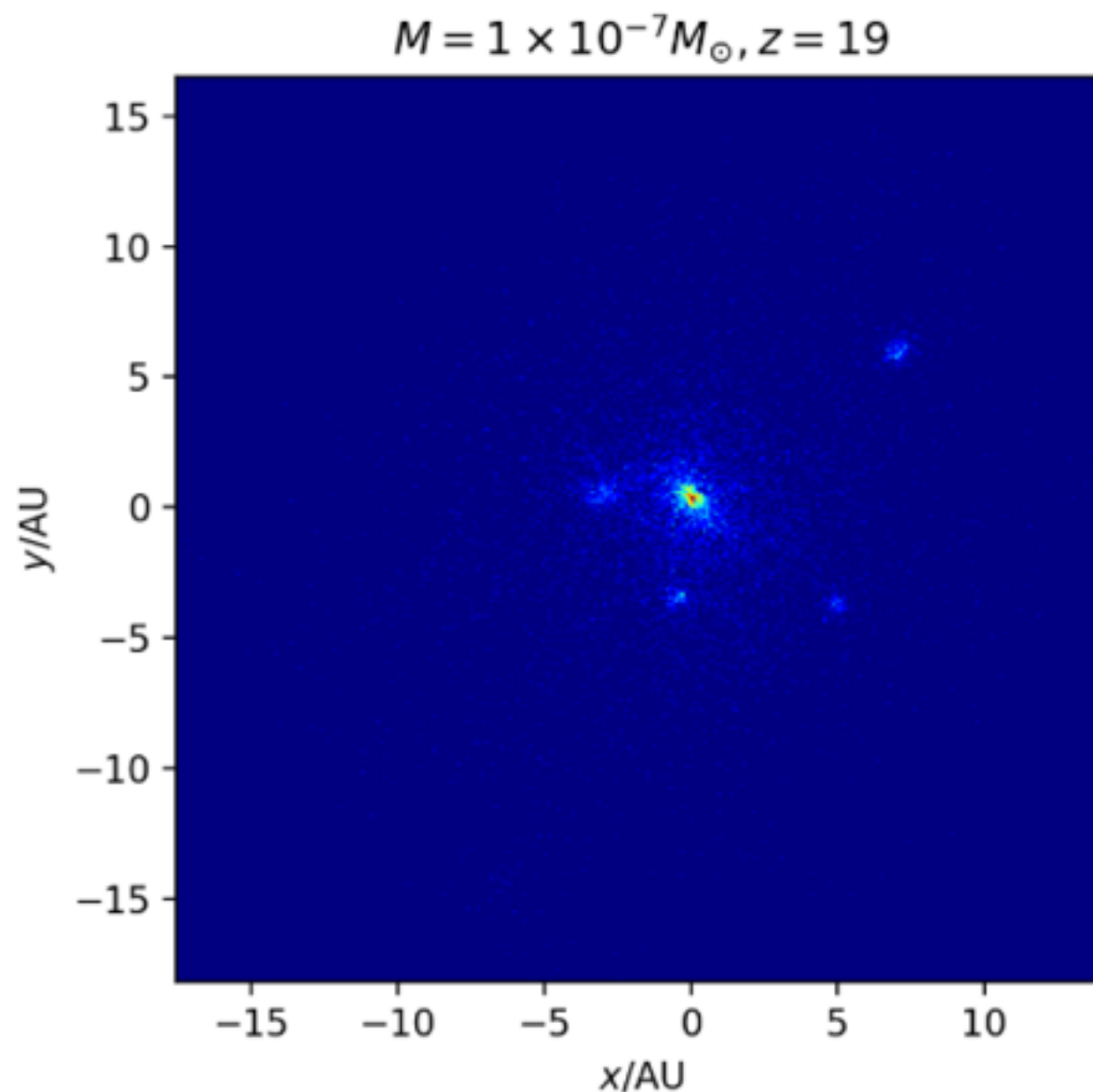
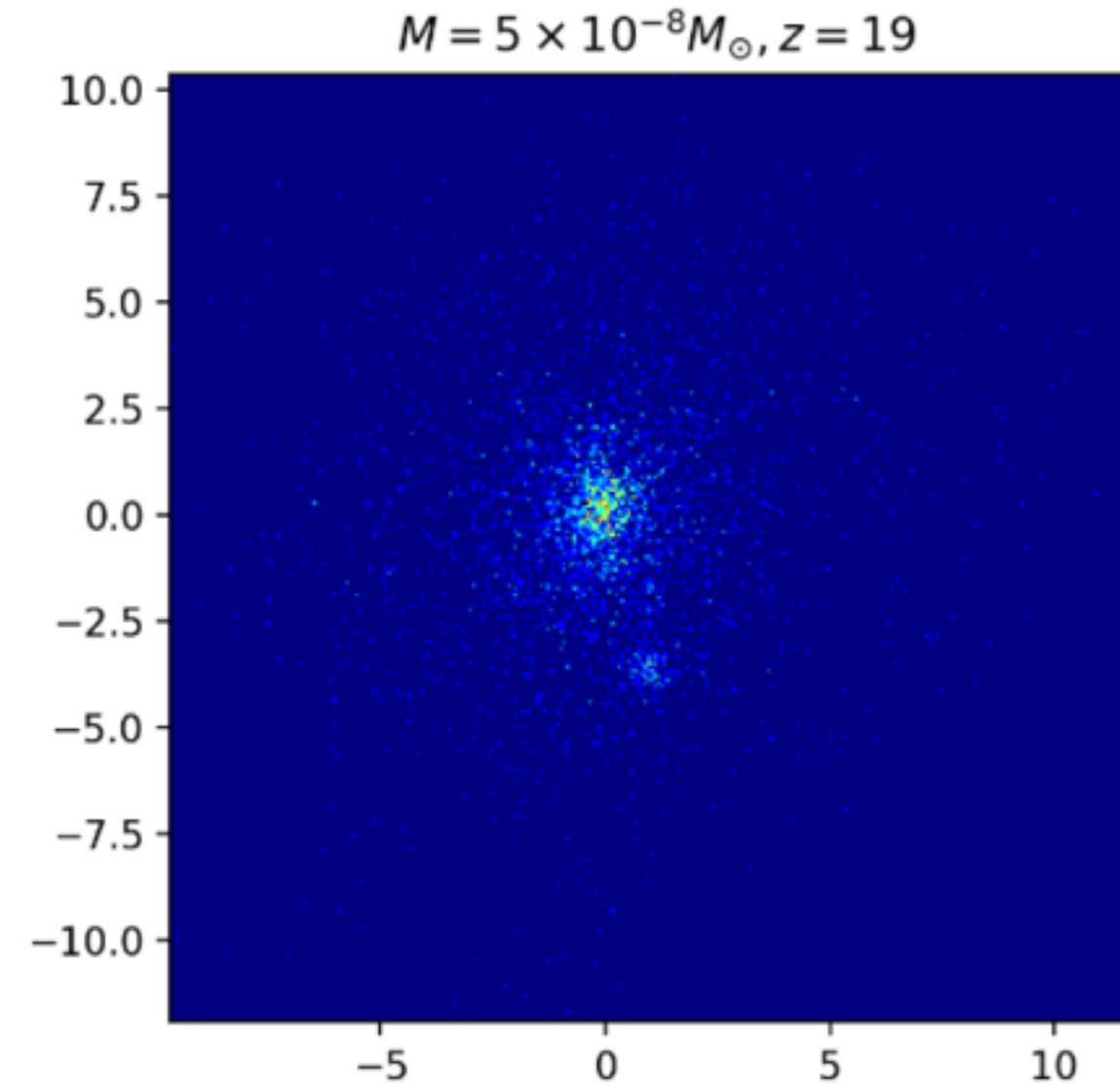
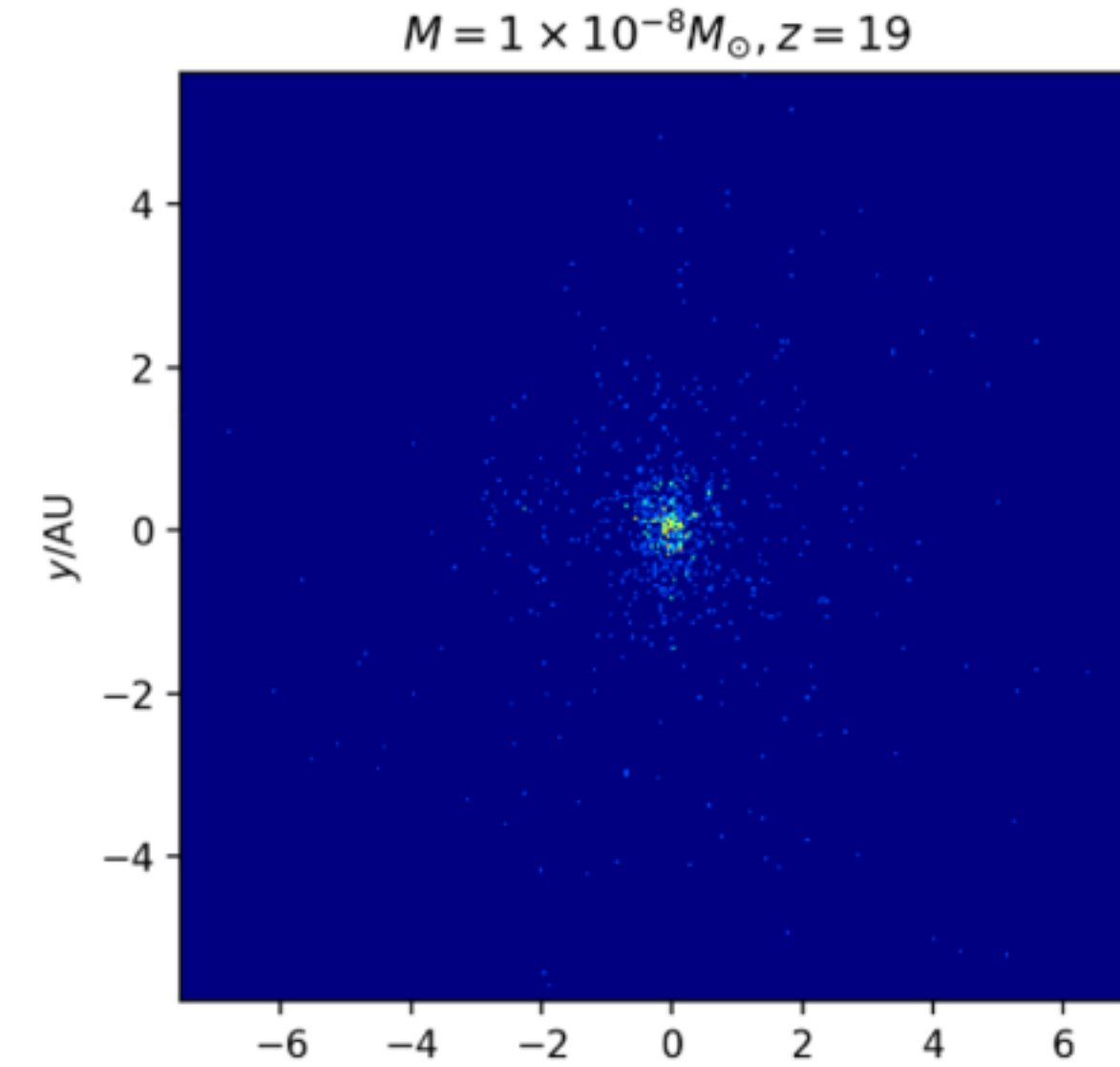
From white-noise initial density fluctuations

NFW profile is a very good fit

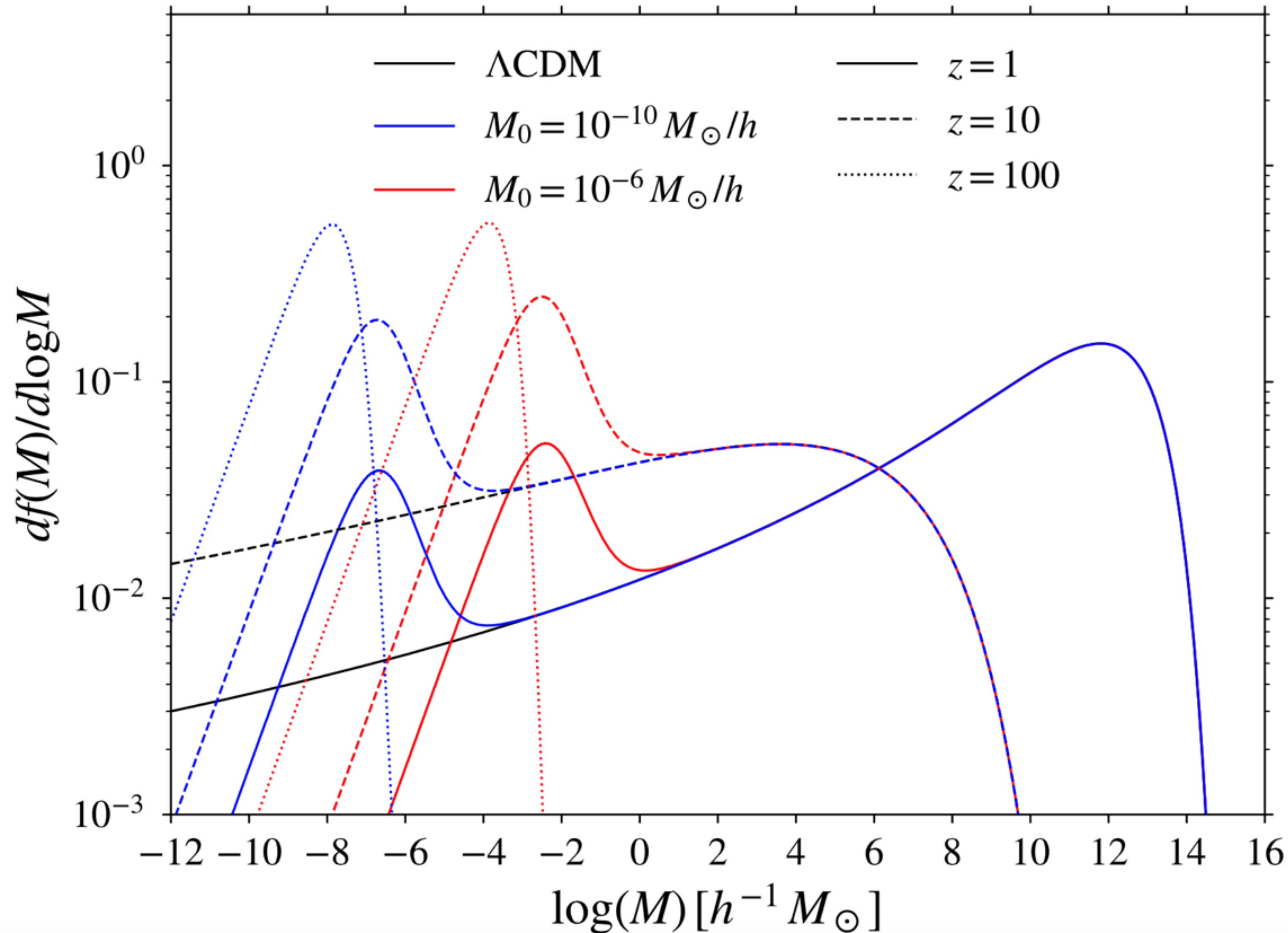
Minihalo concentration and scale radius

$$c(z) \equiv \frac{r_{\text{vir}}}{r_s} = \frac{1.4 \times 10^4}{(1+z)\sqrt{M/(A_{\text{osc}}M_0)}}$$

$$r_s(M) \approx 3.7 \times 10^{-3} h^{-1} \text{pc} \left(\frac{A_{\text{osc}}M_0}{10^{-11}M_\odot/h} \right)^{-1/2} \times \left(\frac{M}{10^{-6}M_\odot/h} \right)^{5/6},$$



Hierarchical Structure Assembly



Press-Schechter halo mass function

Press & Schechter 1974

At $z > 100$:

DM locked into minihalos

Minihalos assemble from bottom up

At $z \lesssim 10$:

“Standard” CDM halos form

Many minihalo assimilated into standard halos

Even as late as $z \sim 1$:

Many field minihalos in isolation

Collapse of cluster-sized halos:

Field minihalos fall into them

Surviving Dynamic Disruption

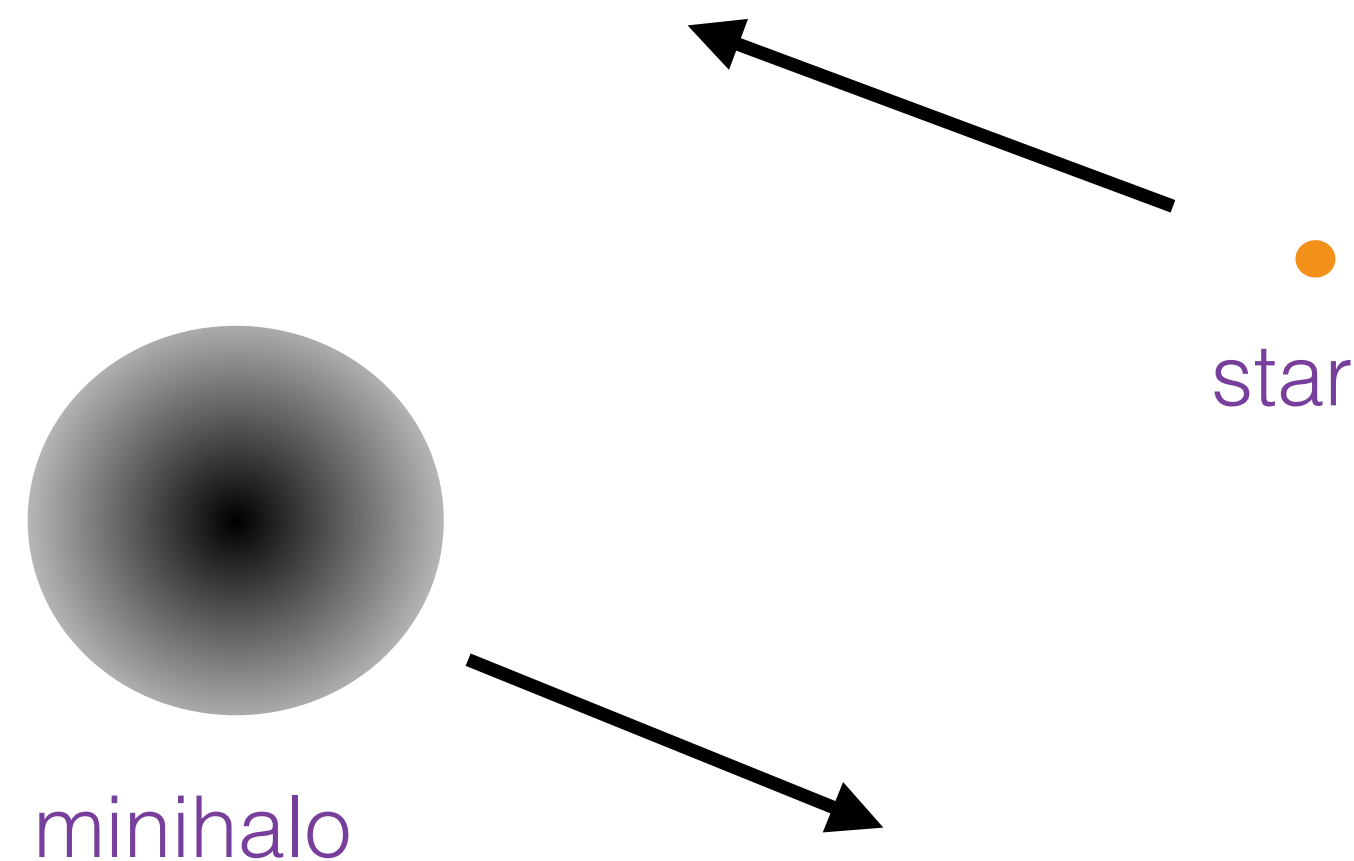
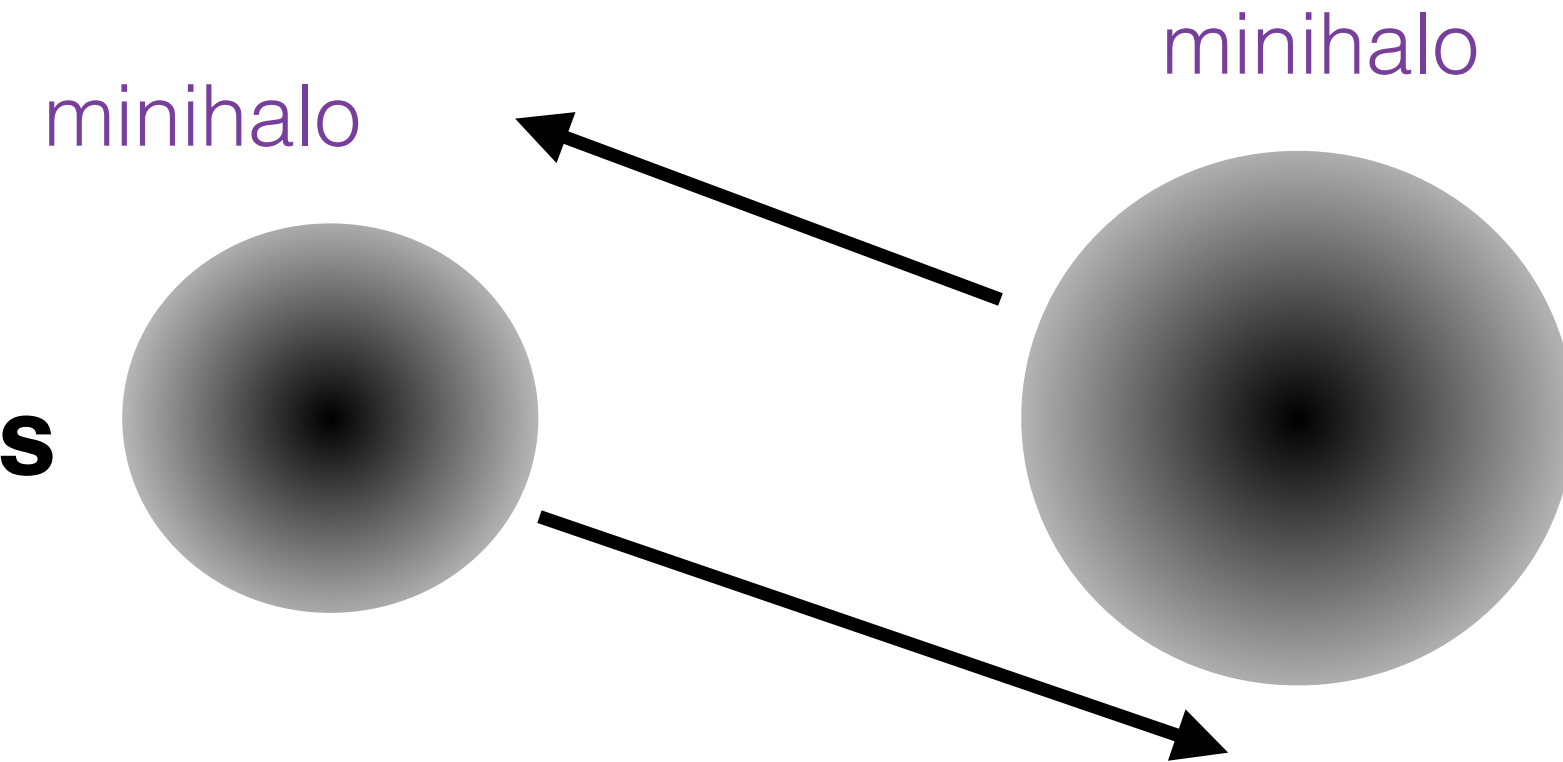
1. Spiral into central dense regions by dynamic friction

Because minihalo have tiny masses, spiral-in timescales are **way too long**



2. Mutual High-speed encounters between minihalos:

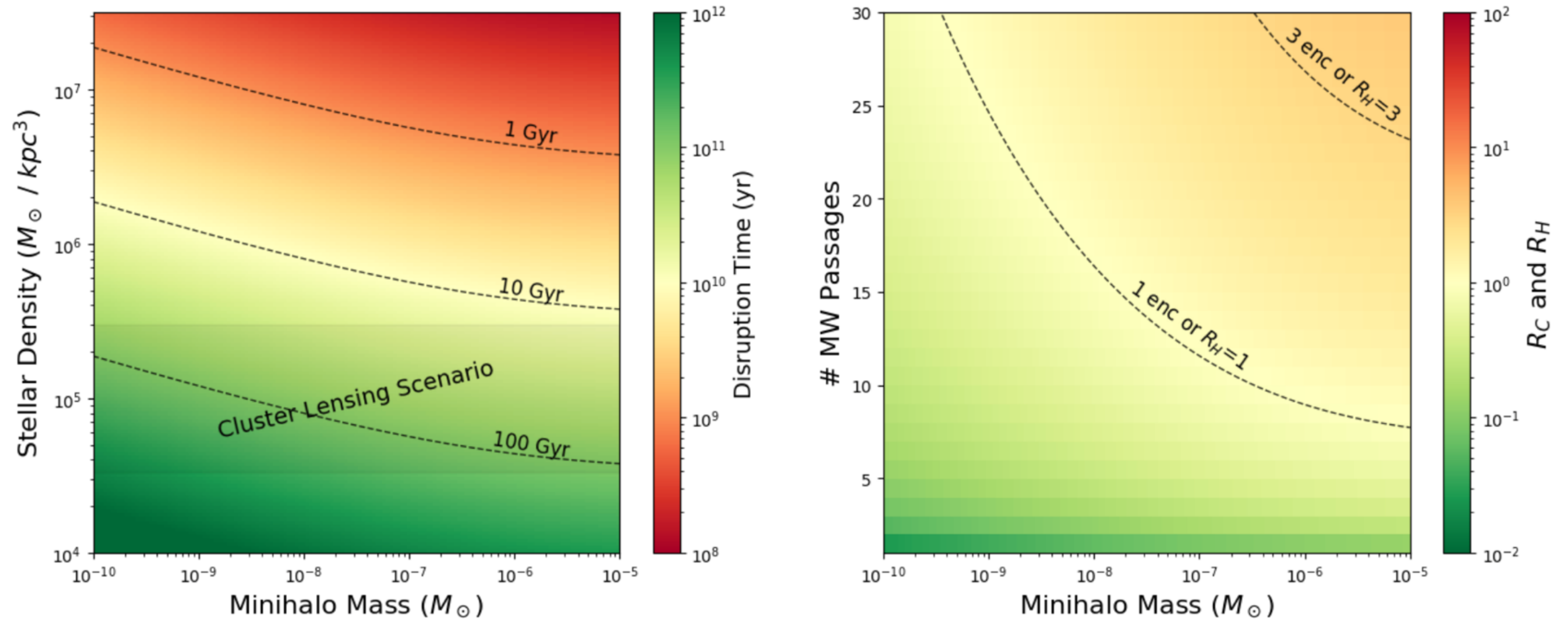
Repeated tidal heating due to **head-on collision between similar mass**
Timescale for disruption much longer than the age of the universe



3. High-speed encounters with passing stars

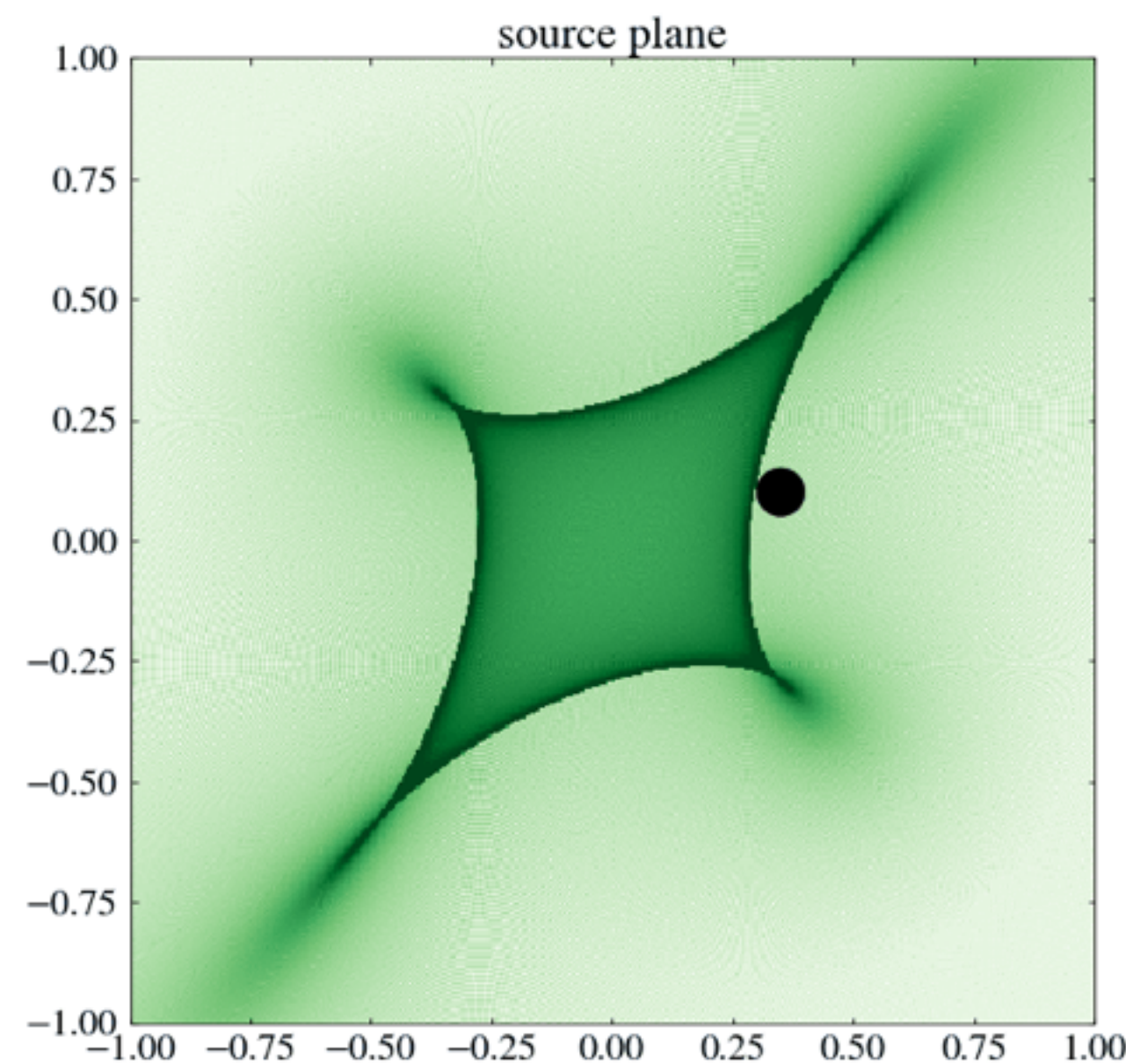
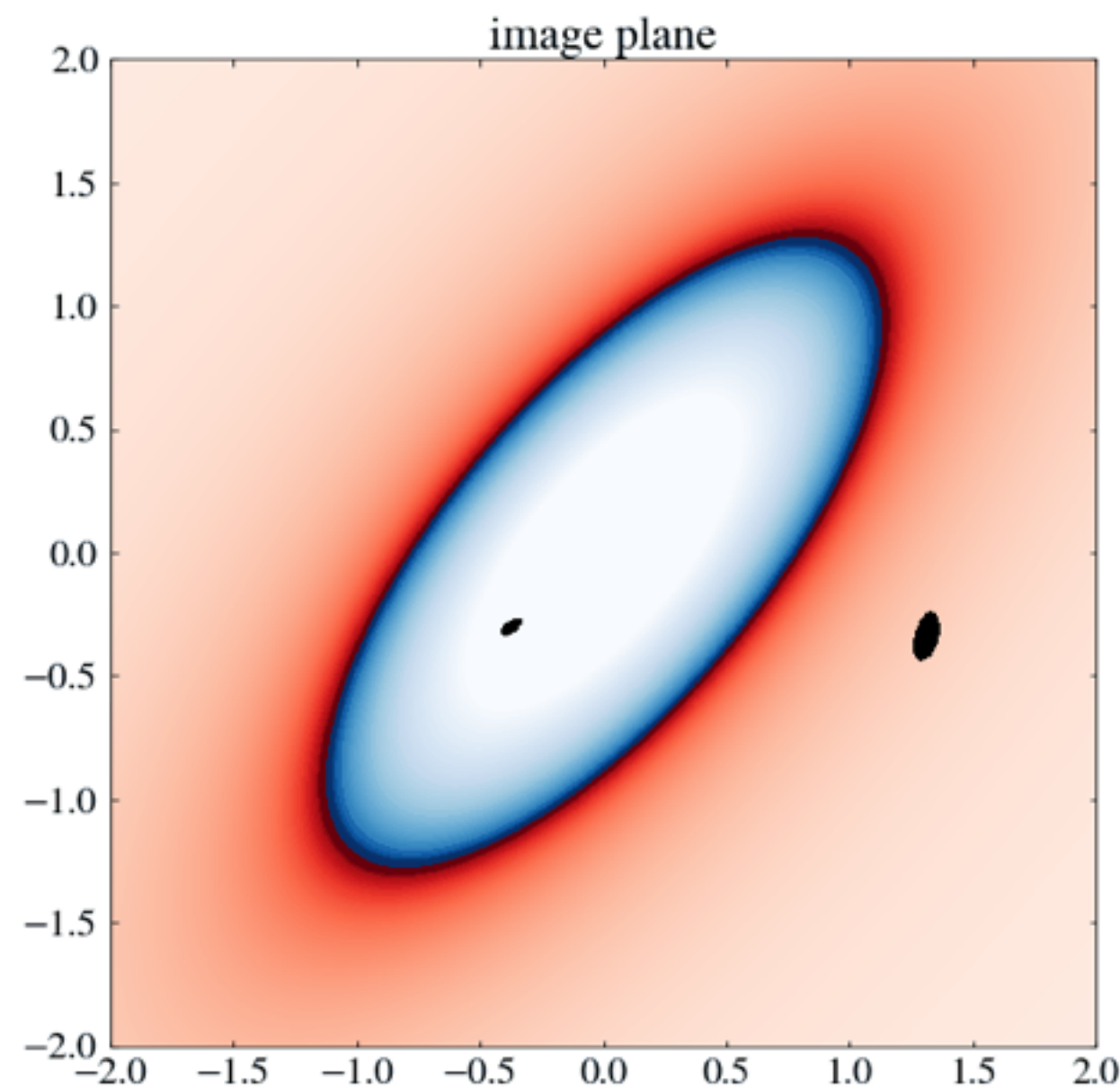
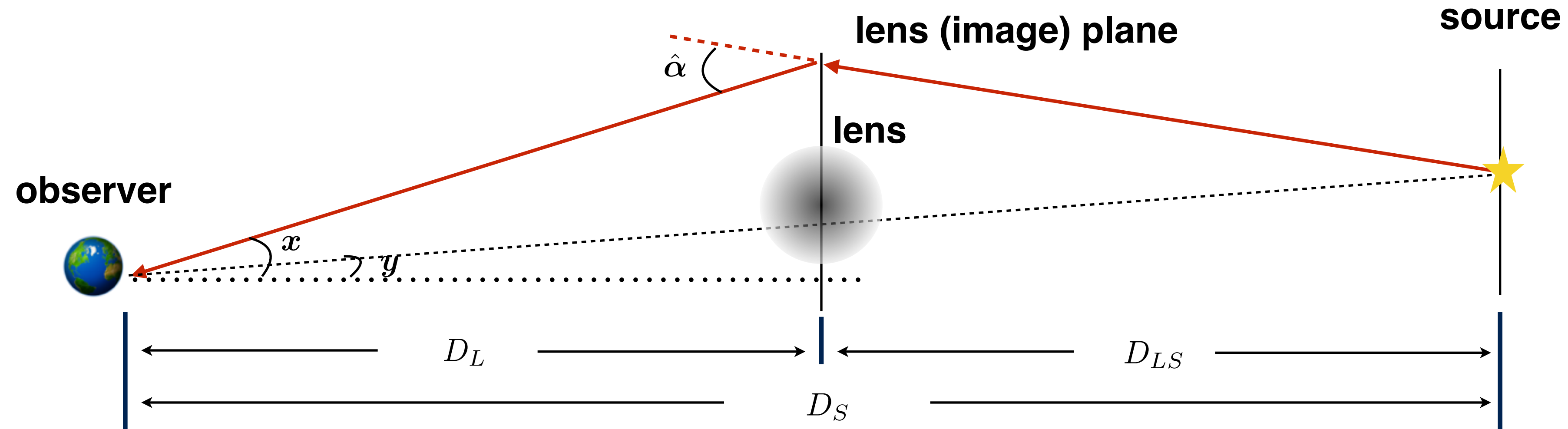
Single flyby disruption within some critical impact parameter
In the intracluster environment, timescale for disruption still a bit longer than the age of the universe

Dynamic disruption in MW and intracluster space



Xiao, Williams & McQuinn 2101.04177

Geometric Description of Lensing



ray equation $\mathbf{y} = \mathbf{x} - \frac{D_{LS}}{D_S} \hat{\alpha}(\mathbf{x}) = \mathbf{x} - \boldsymbol{\alpha}(\mathbf{x})$

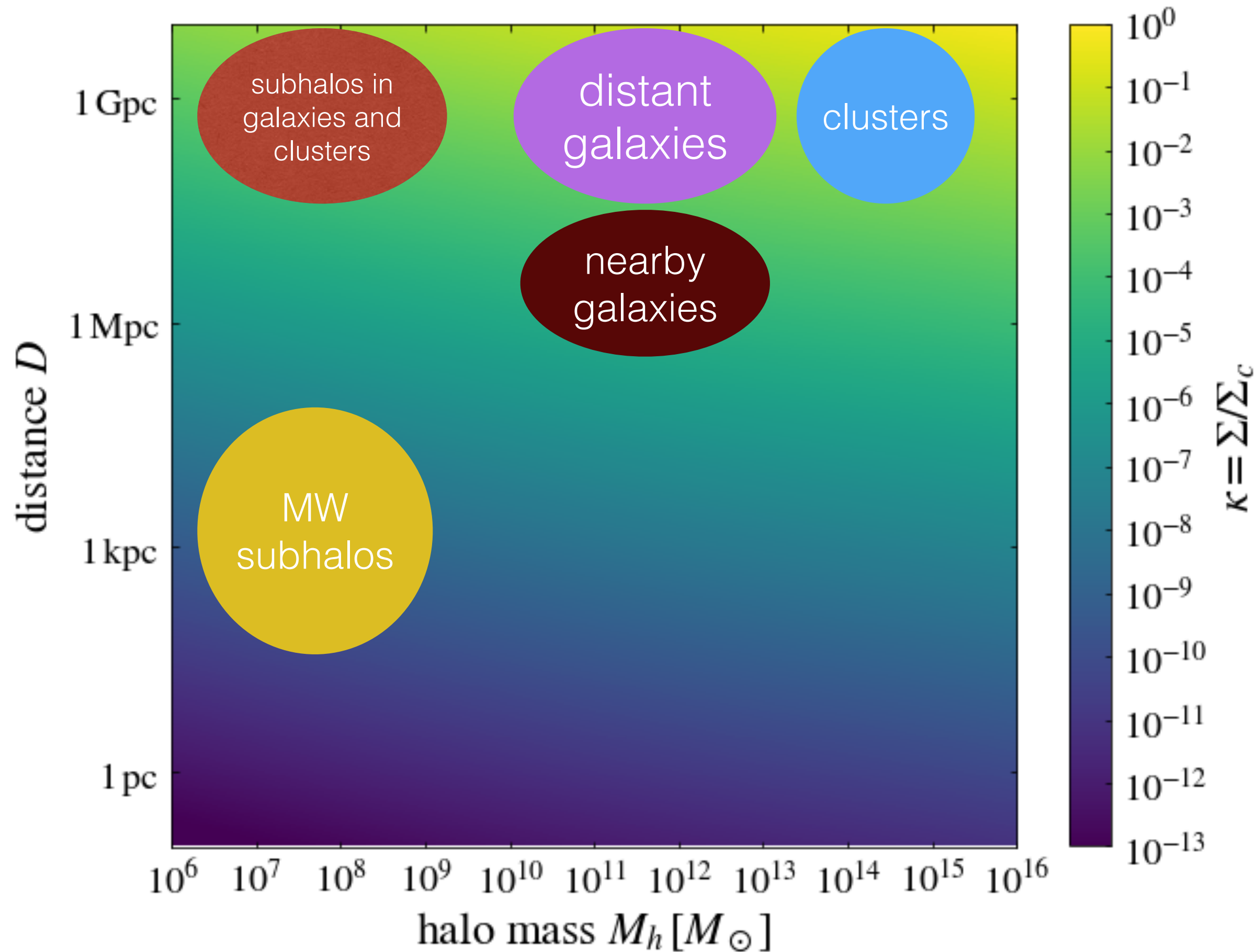
convergence and shear $\frac{\partial \mathbf{y}}{\partial \mathbf{x}} = \begin{bmatrix} 1 - \kappa - \gamma_1 & -\gamma_2 \\ -\gamma_2 & 1 - \kappa + \gamma_1 \end{bmatrix}$

flux magnification $\mu = 1 / [(1 - \kappa)^2 - (\gamma_1^2 + \gamma_2^2)]$

Problem Detecting Extended Lenses

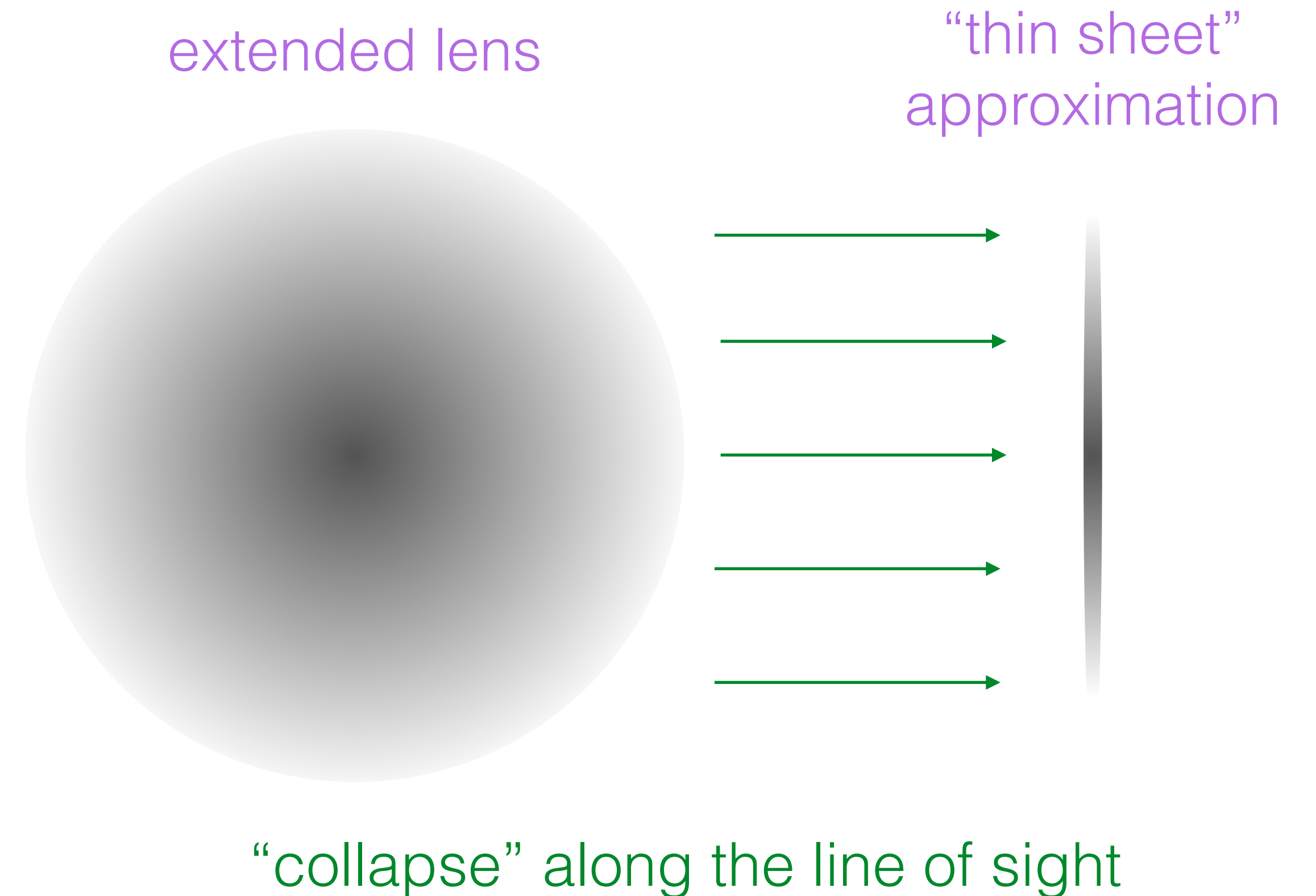
Lensing convergence

$$\kappa = \Sigma / \Sigma_c$$

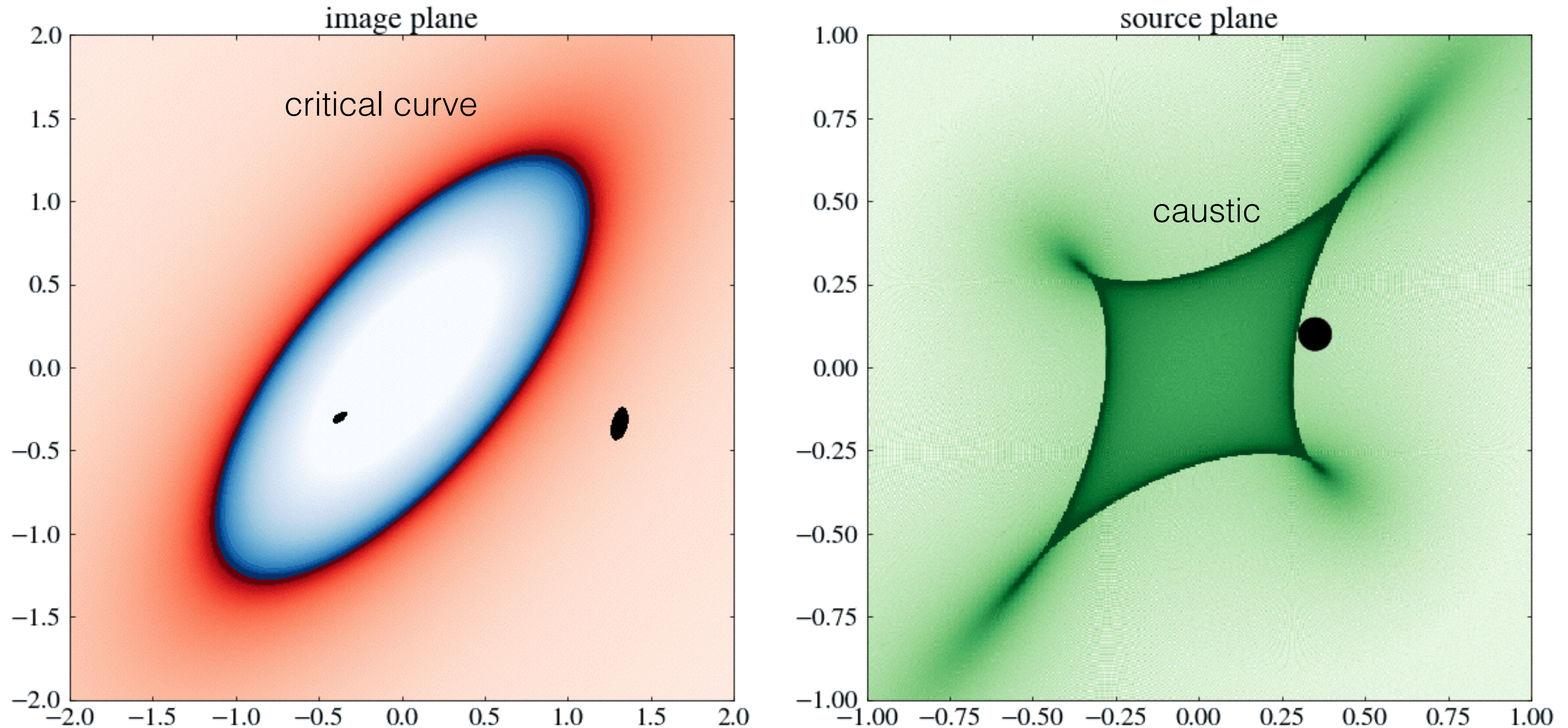


Critical lens surface density

$$\Sigma_c \sim \frac{1}{4\pi G D}$$

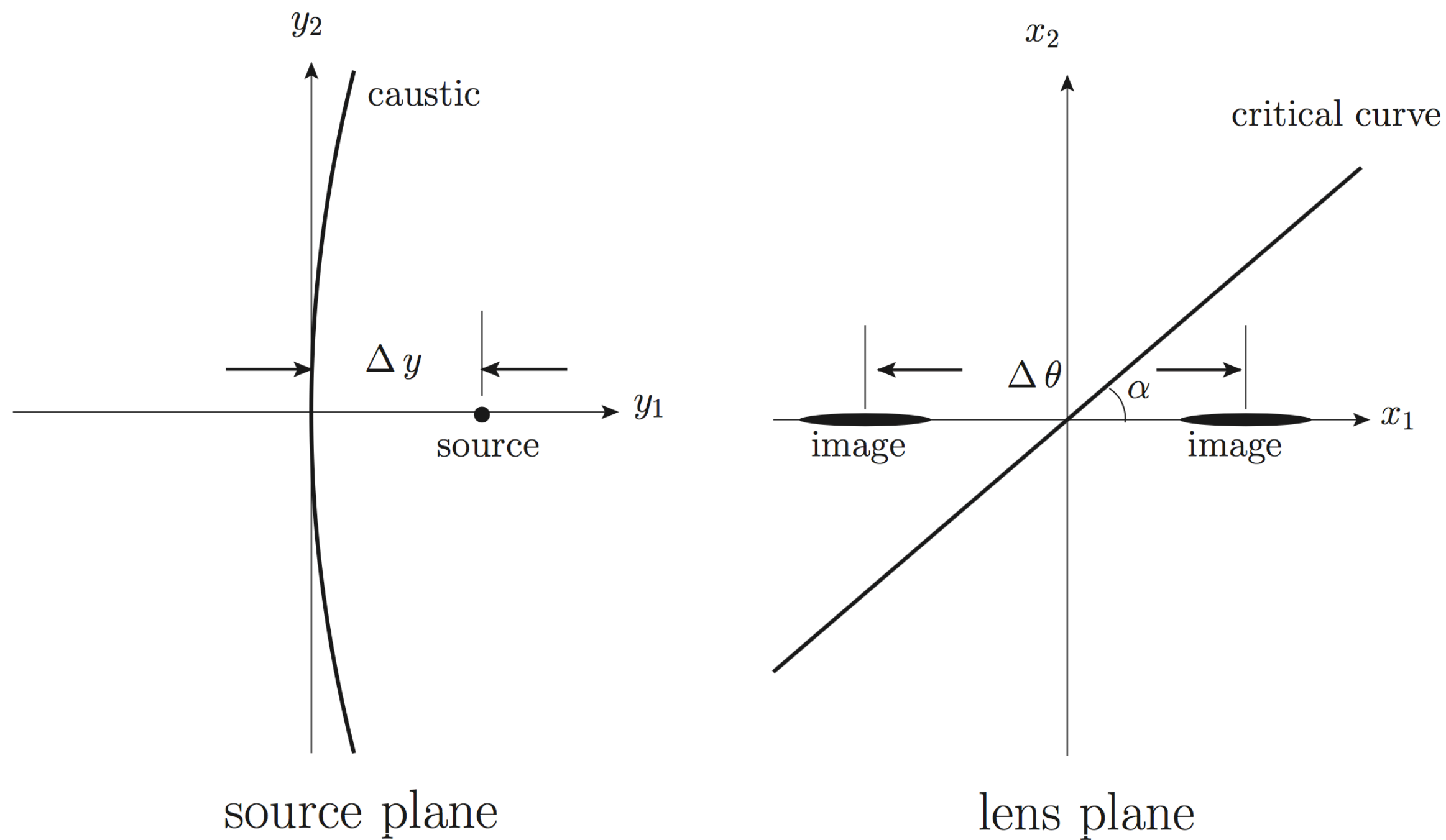


Caustic and Critical Curve



Discontinuous change in the number and configuration of images.
Exhibit **universal behaviors** governed by Catastrophe Theory

Caustic Transit of Compact Sources



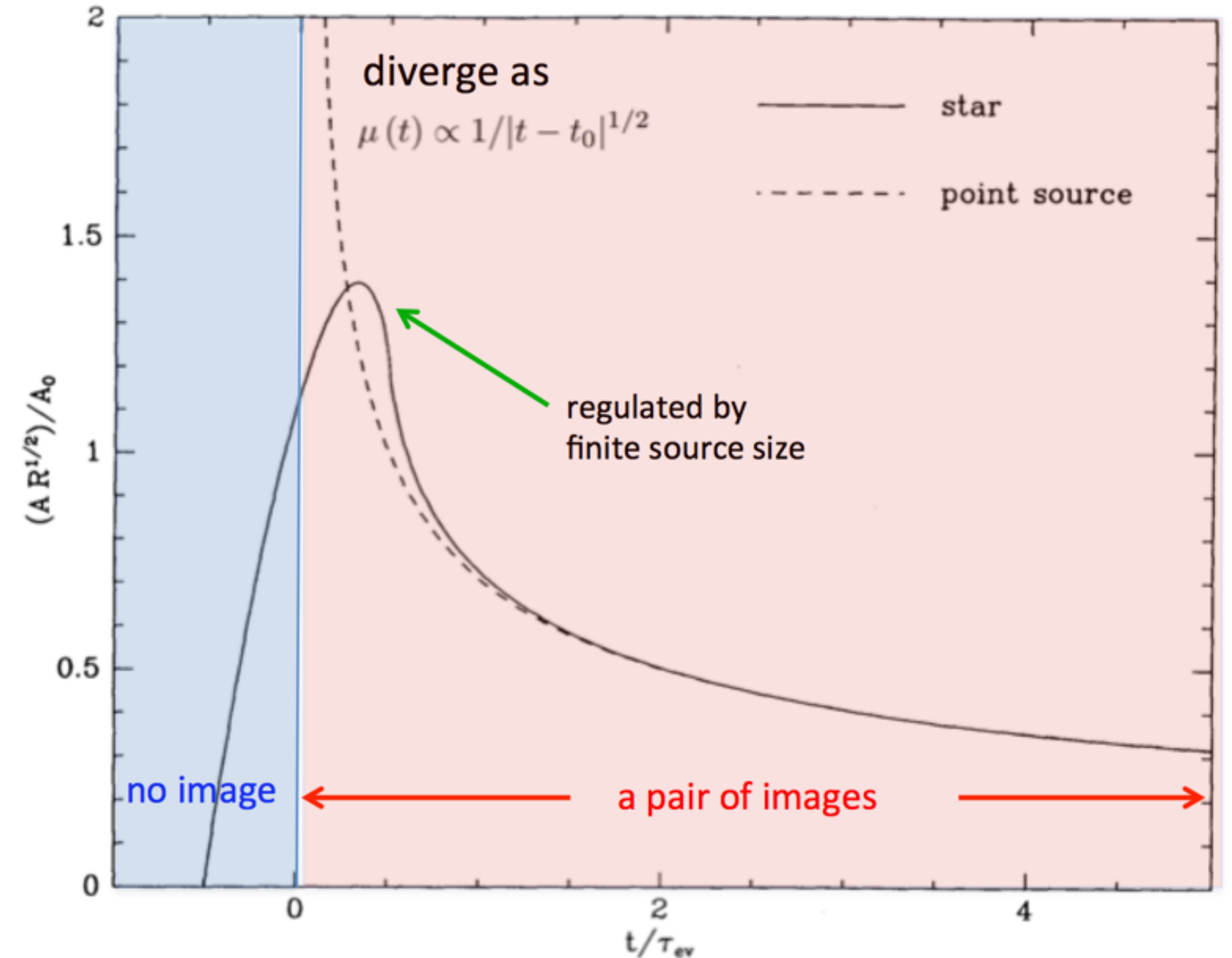
$$\mu \sim \theta_C / \Delta\theta$$

$$\Delta\theta \sim 10 - 100 \text{ mas}$$

$$\theta_C \sim 10''$$

$$\mu \sim \frac{\theta_C}{\Delta\theta} \sim \text{few} \times 10^{2-3}$$

Miralda-Escudé 91'



finite source size

$$\mu_{pk} \propto R_S^{1/2}$$

wave diffraction

$$\mu_{pk} \propto \lambda^{-1/3}$$

Extremely Magnified Stars

Miralda-Escudé 1991

Kelley++ 2017 **MACS J1149 LS1**

Chen++ 2019

Kaurov++ 2019

MACS J0416

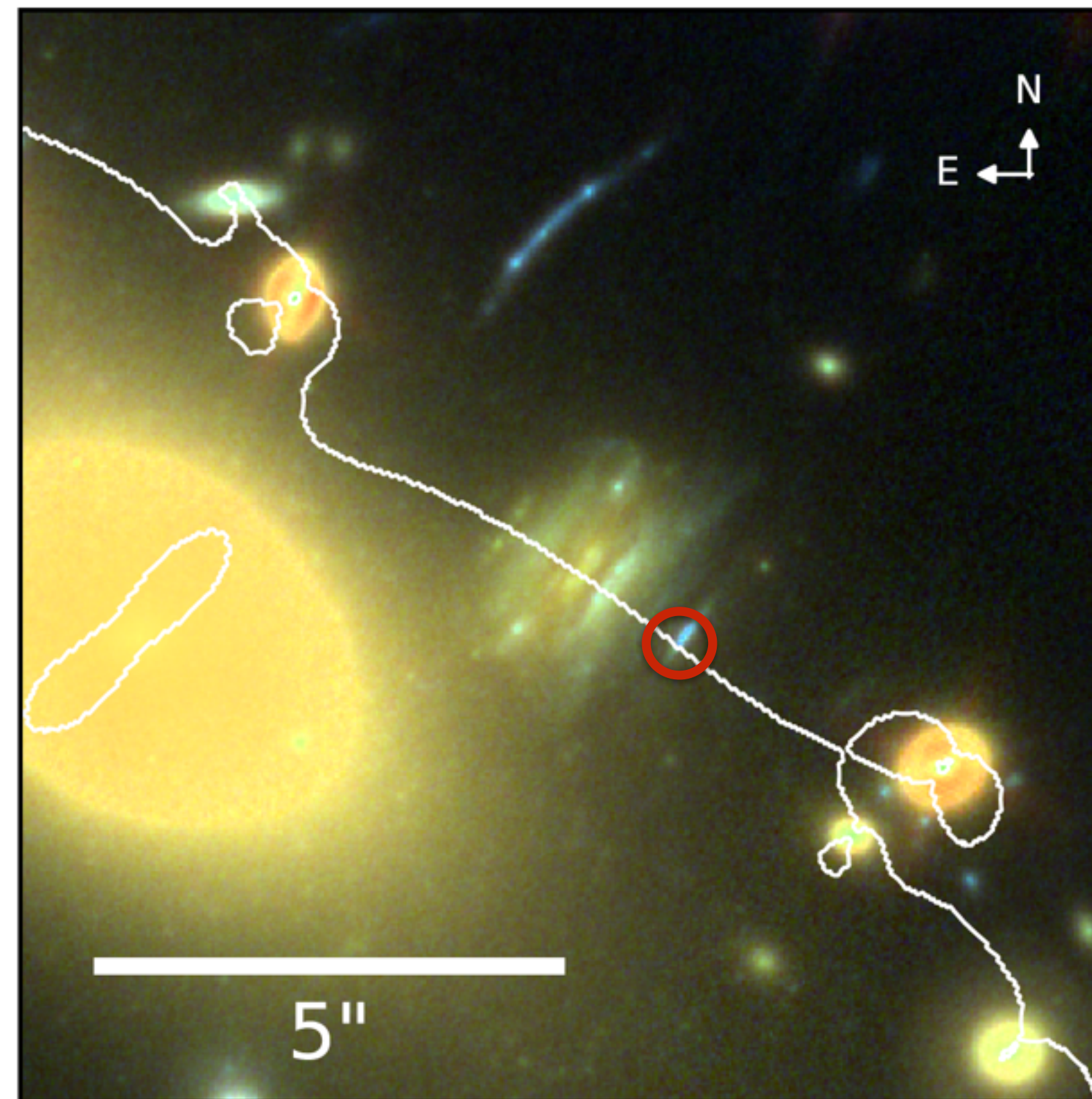
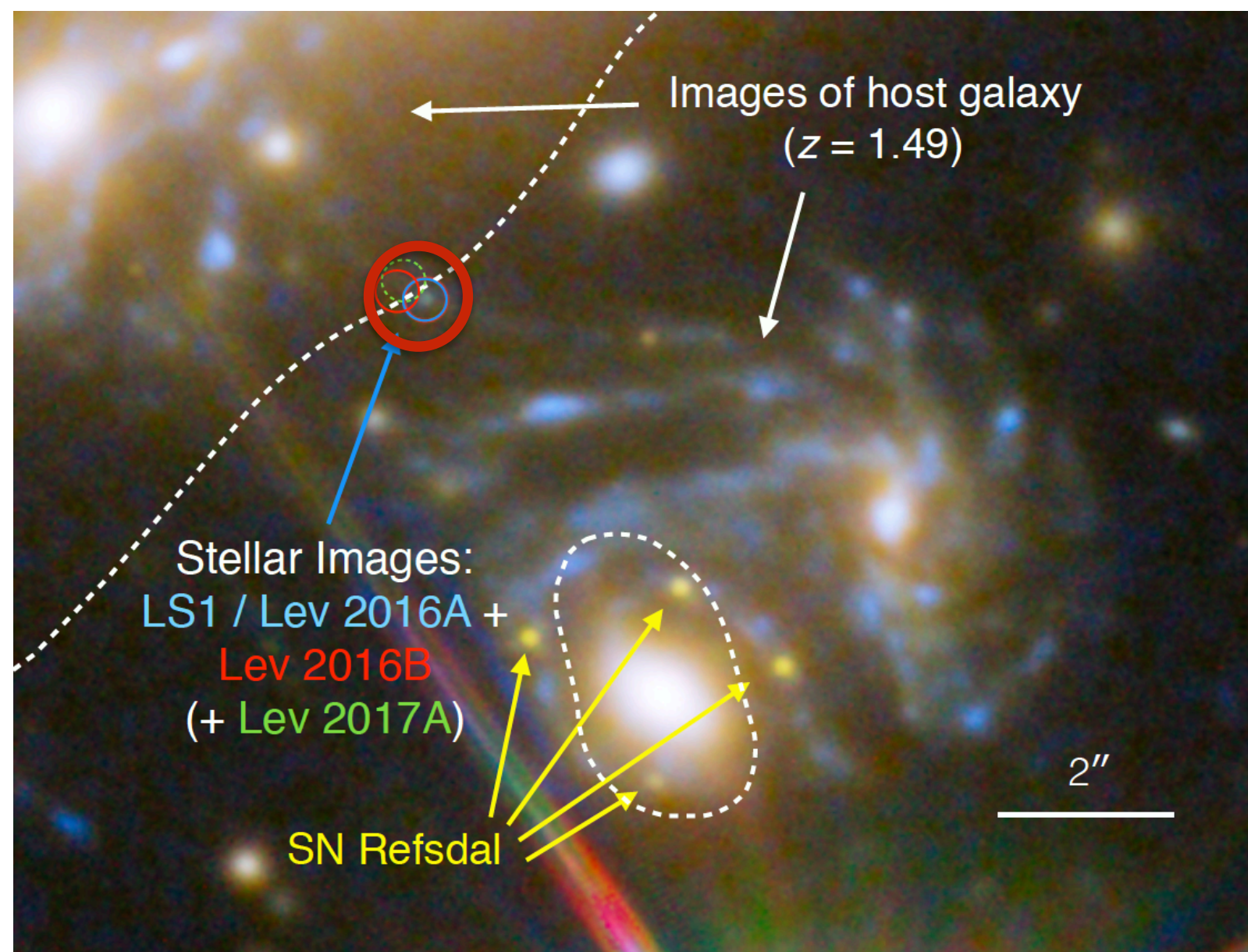
Supergiant stars residing
in $z \sim 1$ lensed arcs

Persistently magnified by
 $O(100)$ fold. Reach AB
mag = 28 — 29

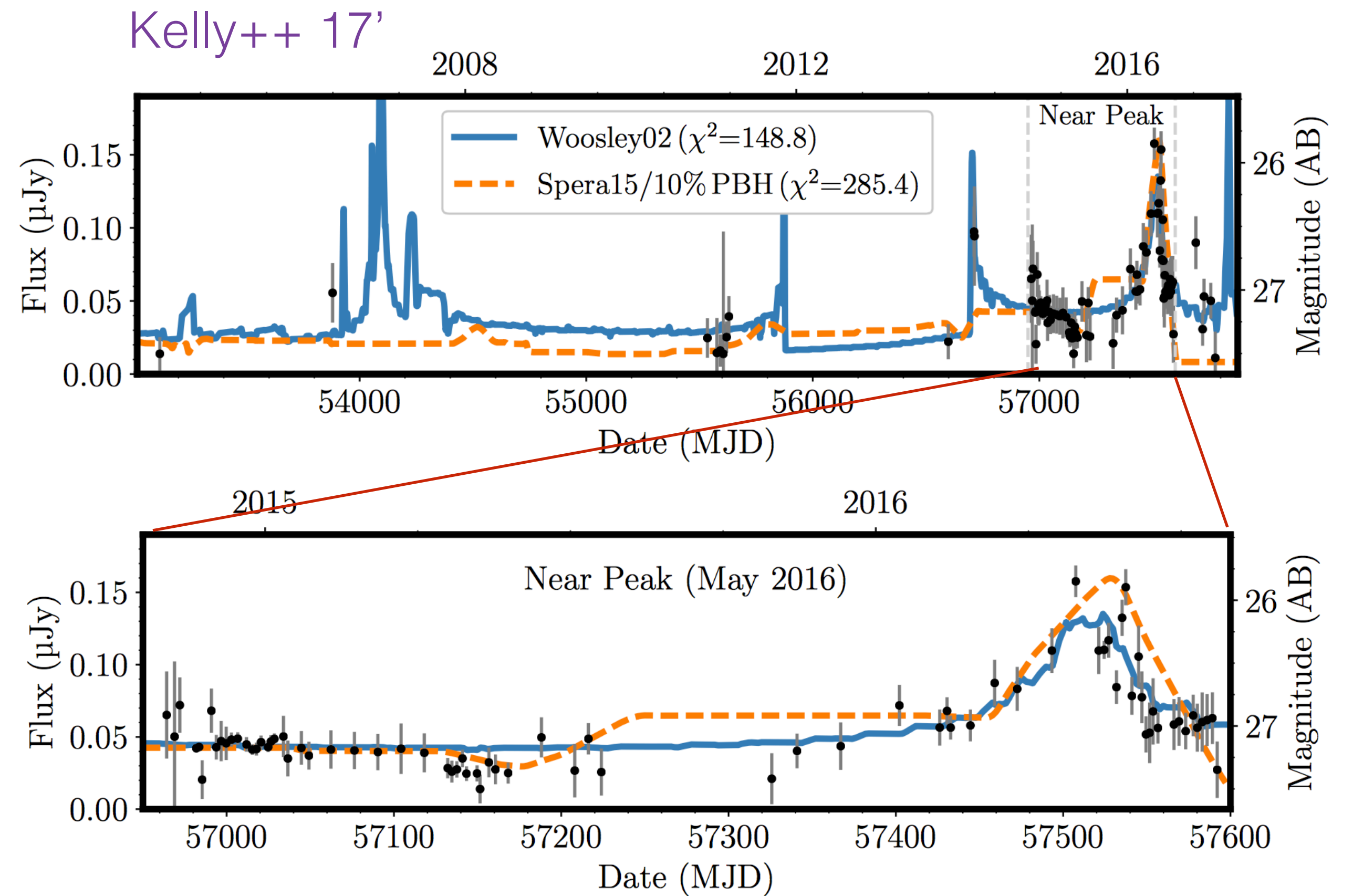
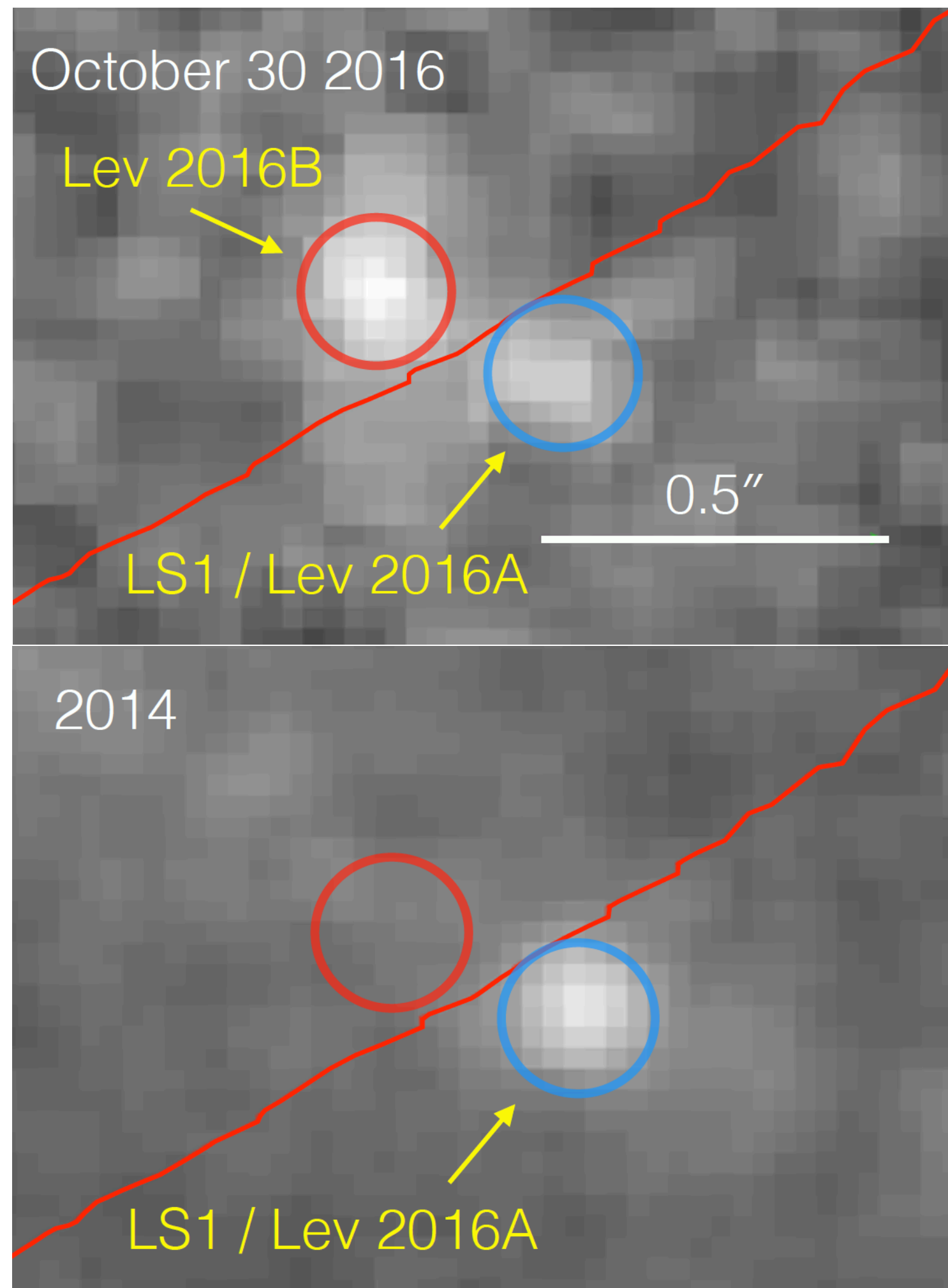
Elongated macro image
pair; separated by
telescope

Phenomenology very
sensitive to lumpiness in
the lens

Dai++ 2018



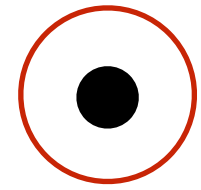
Caustic Transient Variability



A pair of images resolved.
Source intermittently brightens

“Microlensing” Near Critical Curve

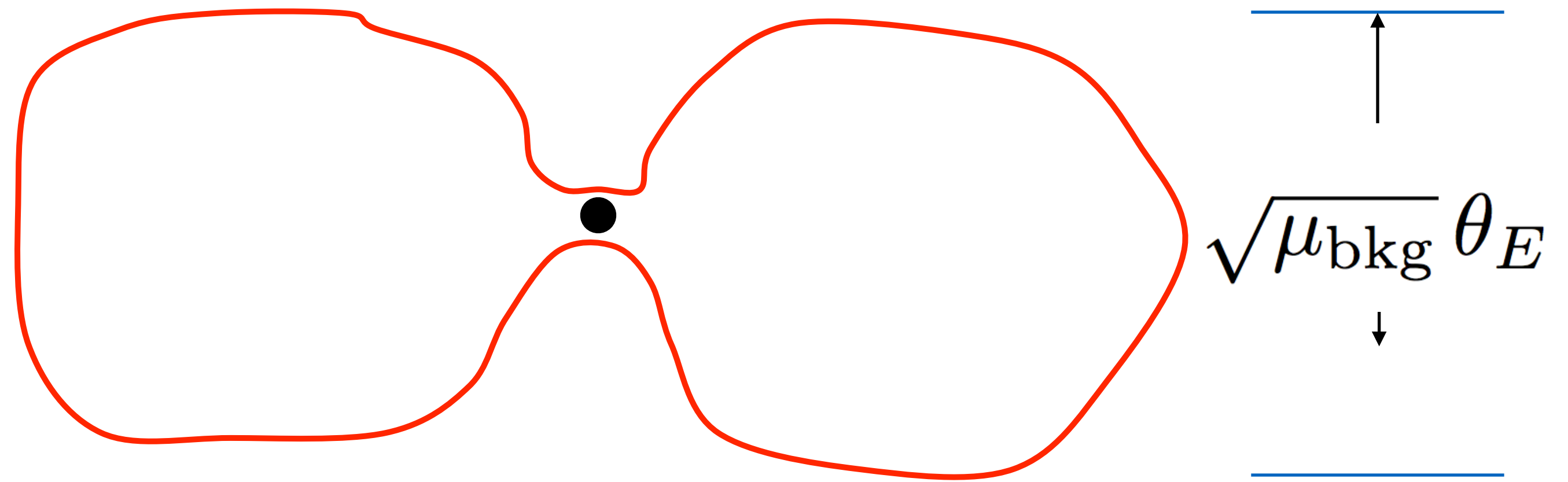
An isolated microlens



$$\theta_E = \left[\frac{4GM}{c^2} \frac{D_{LS}}{D_L D_S} \right]^{1/2} \sim 10^{-6} \text{ arcsec}$$

Covering factor of “Einstein disks” is minuscule for isolated micro-lenses

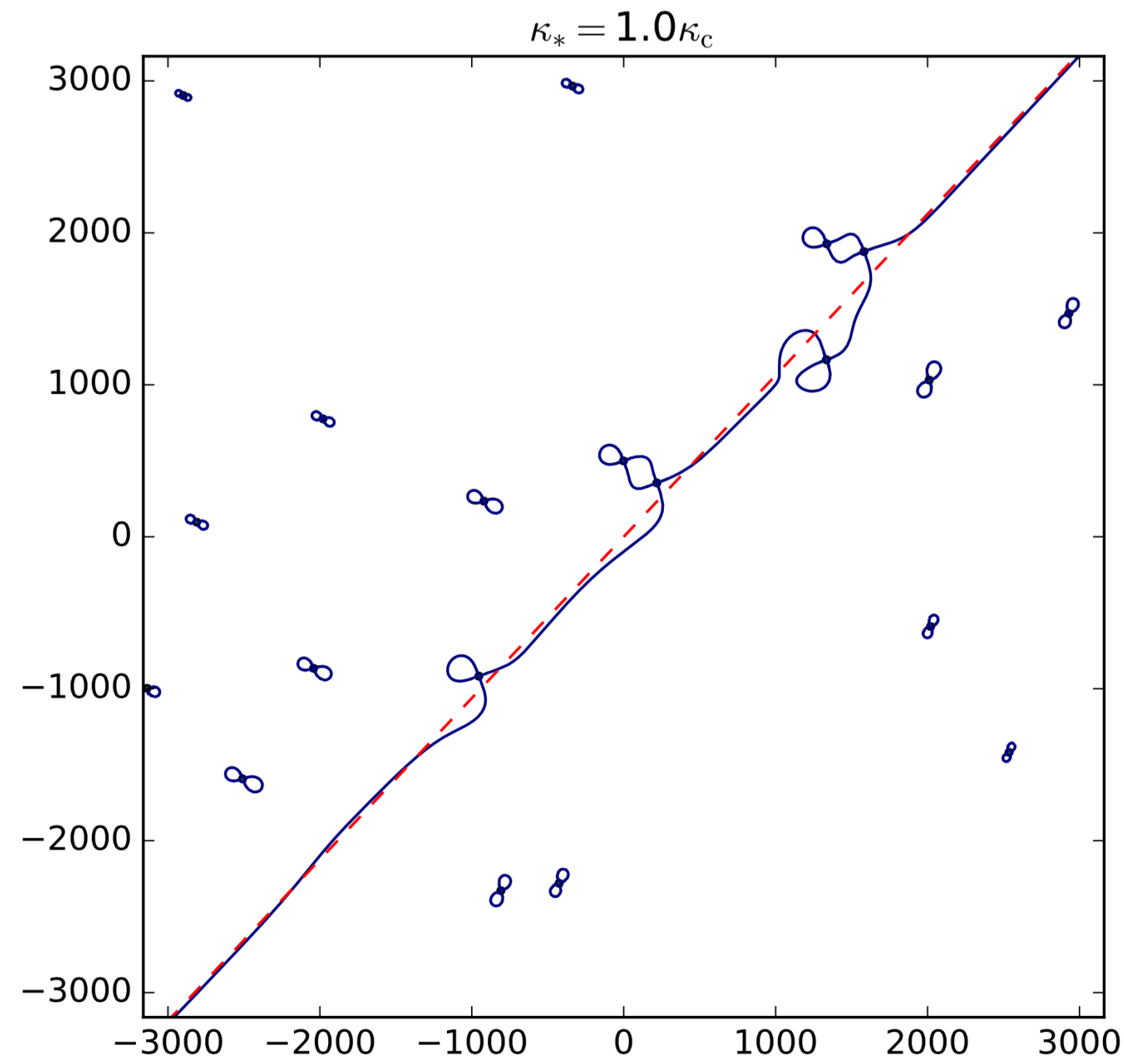
An microlens (in projection) near the critical curve



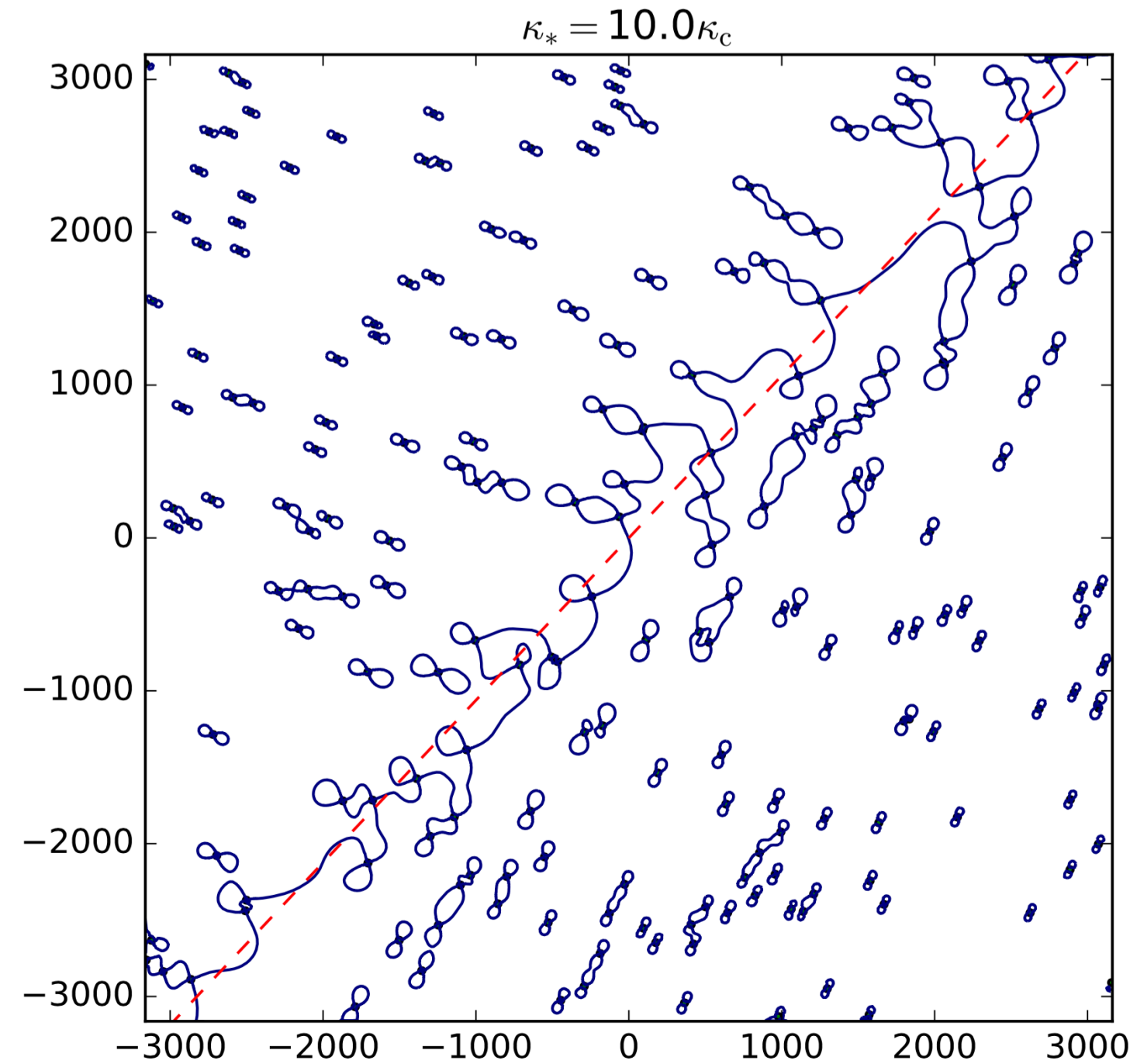
Microlensing cross section enhanced by a huge factor $\sim 10^{2-3}$

Effects of Many Micro-lenses

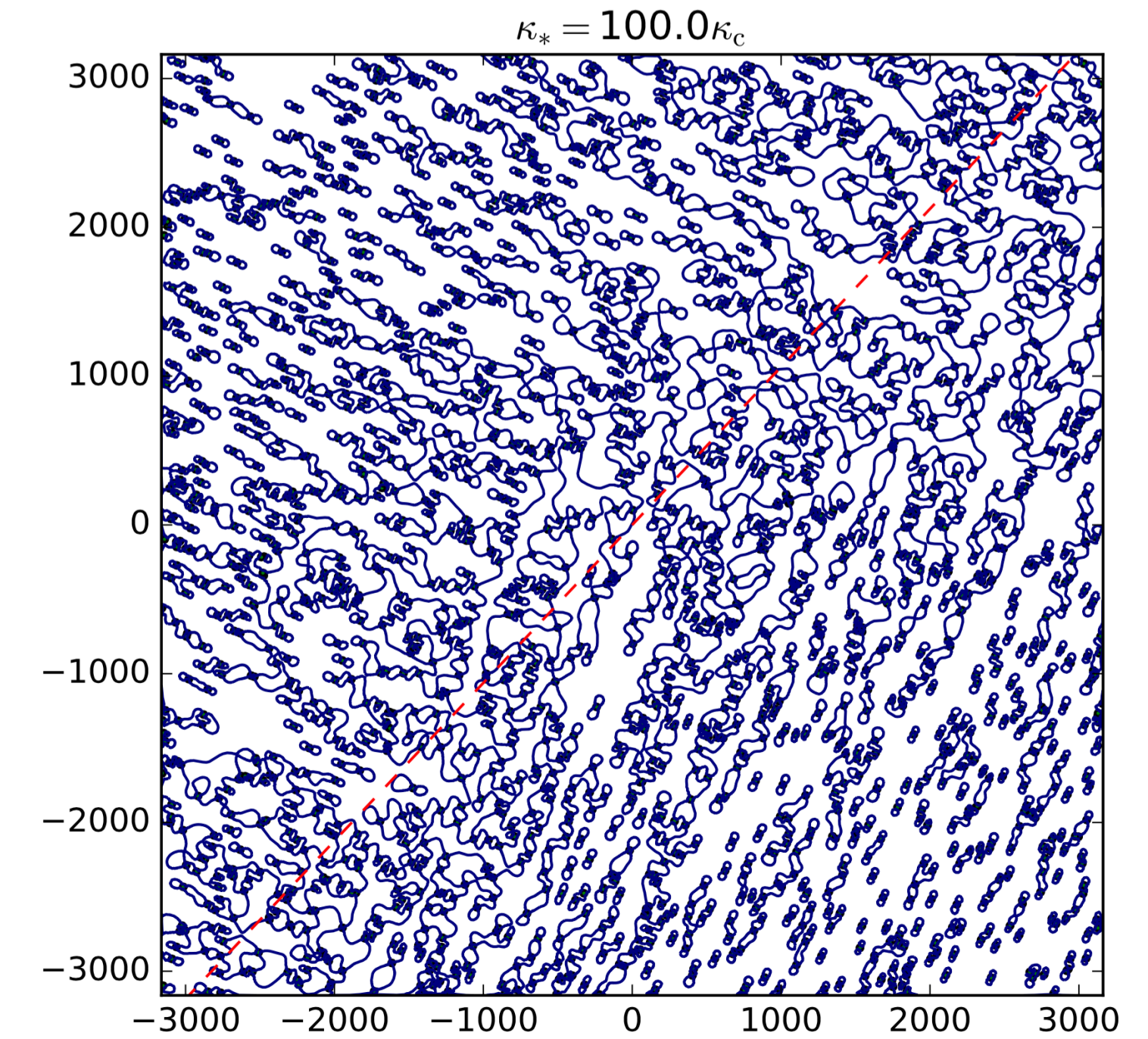
Venumadhav, LD & Miralda-Escudé 17'
Diego++ 17'
Oguri++ 18'



occasionally perturbed

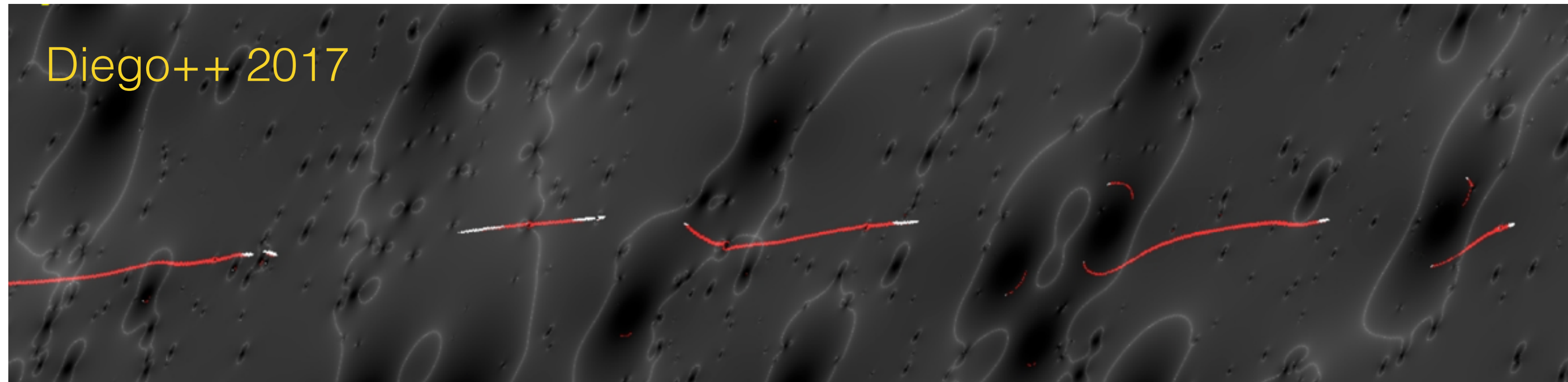


completely disrupted



corrugated network

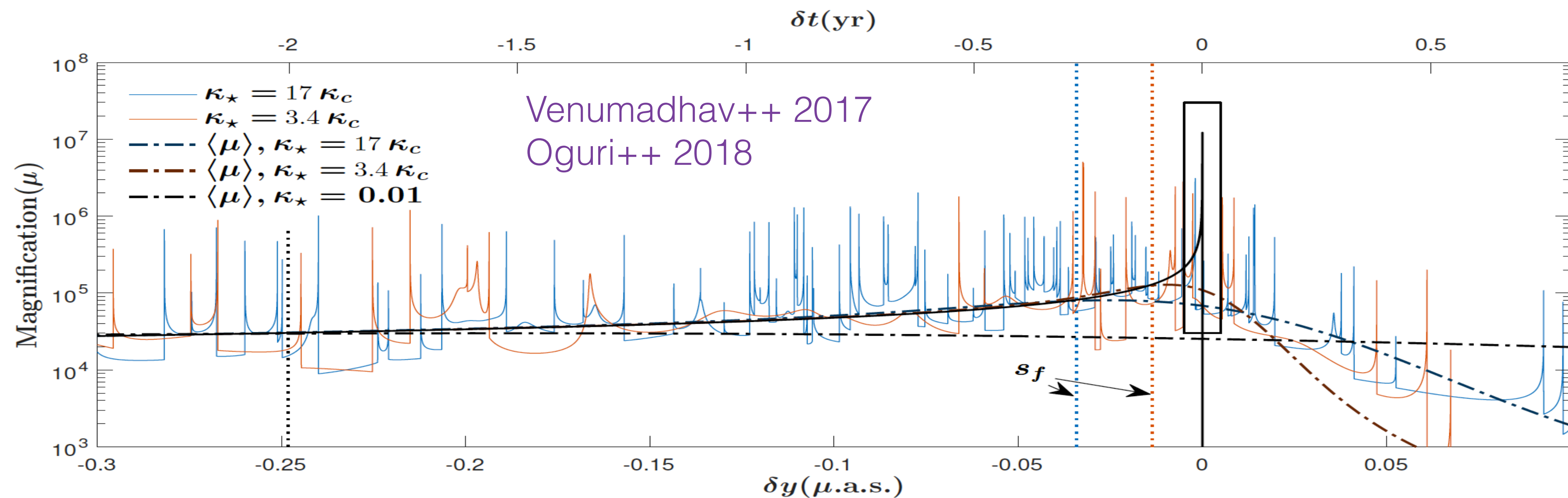
Intermittent Microlensing by Intracluster Stars



frequent microlensing
peaks \sim few times per yr

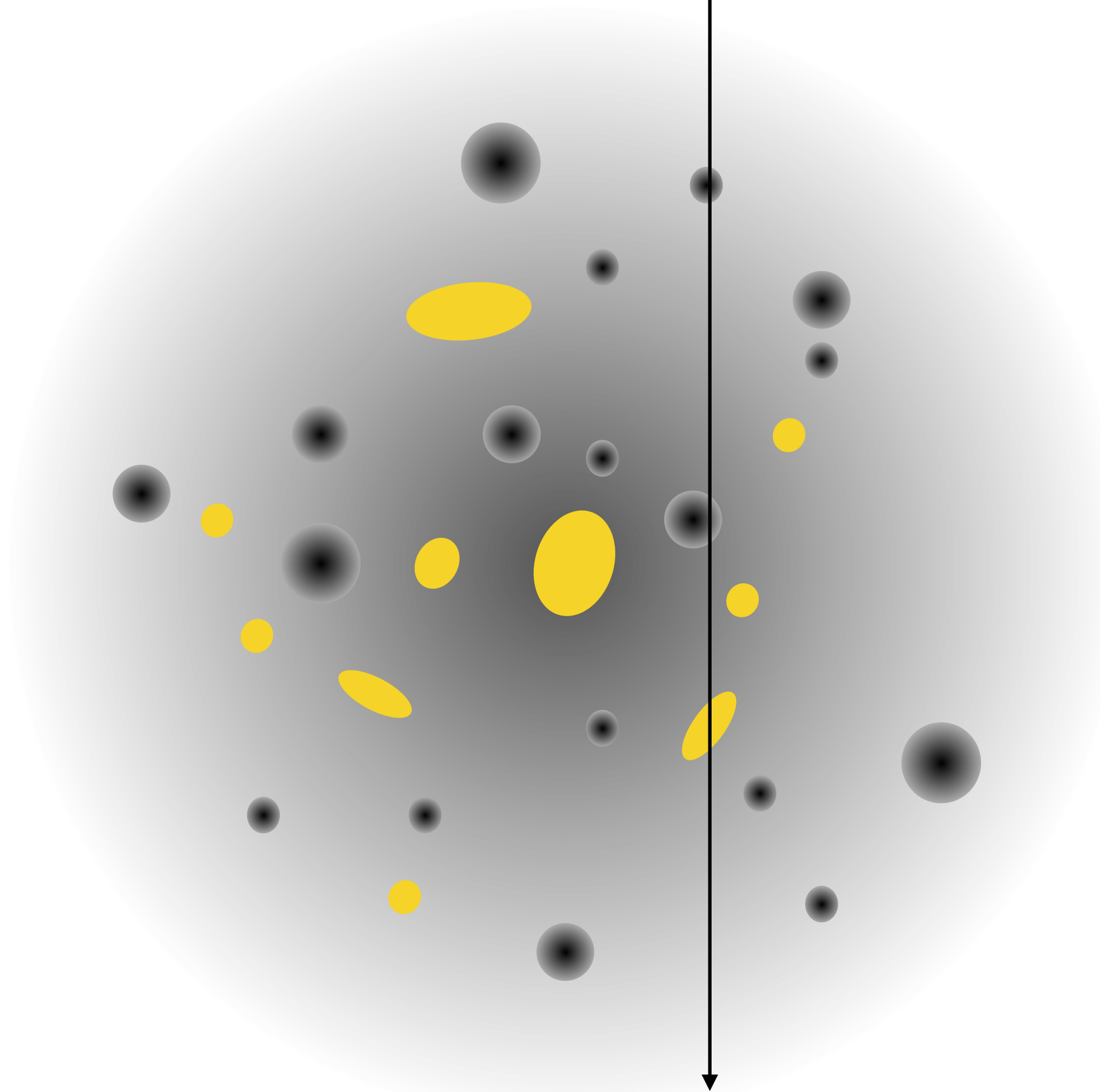
$$\mu_{\text{peak}} \sim 10^4 (R/10 R_{\odot})^{-1/2}$$

cross corrugated
network for $\sim \mathbf{10^4 \text{ yrs}}$!!



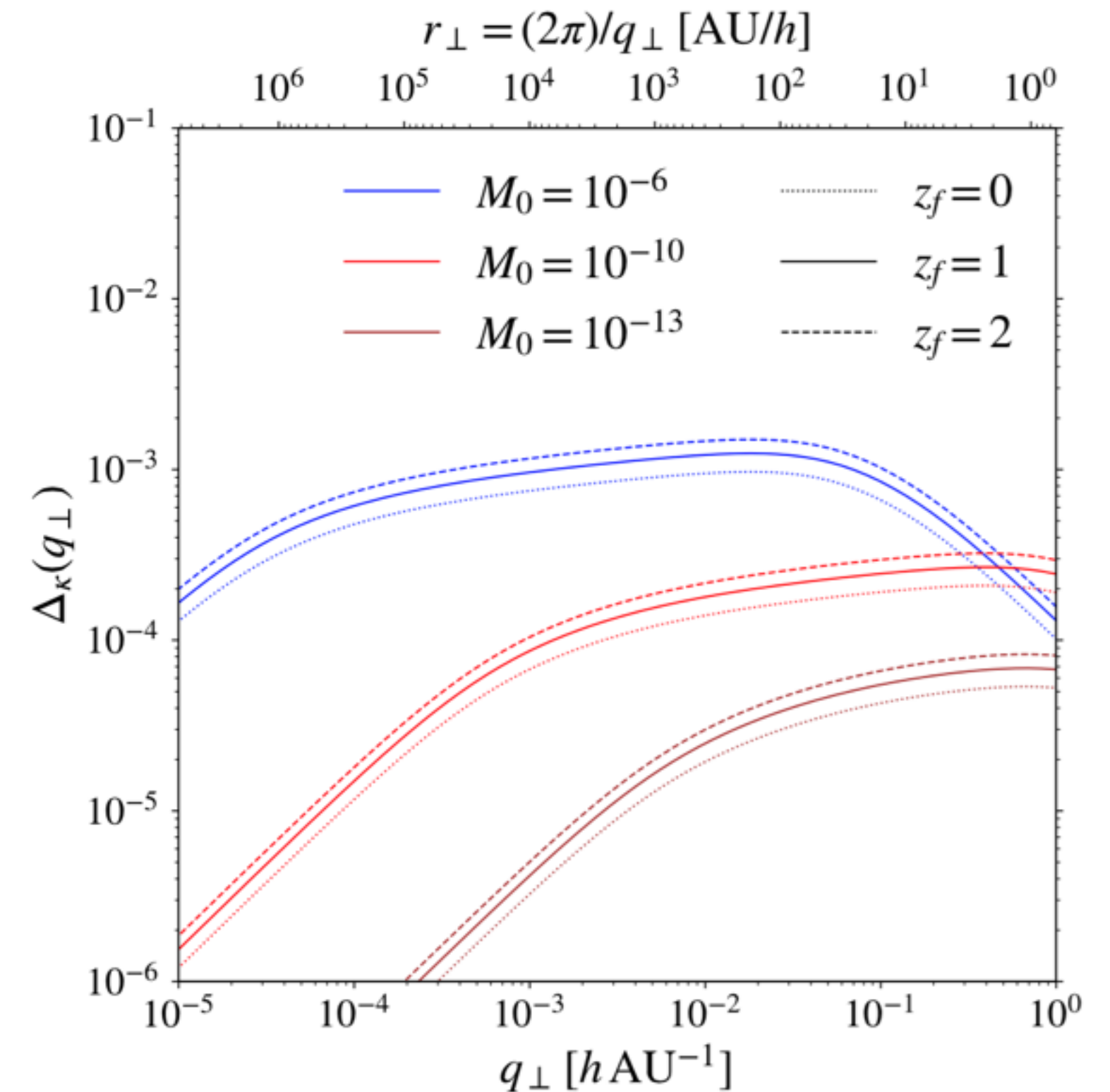
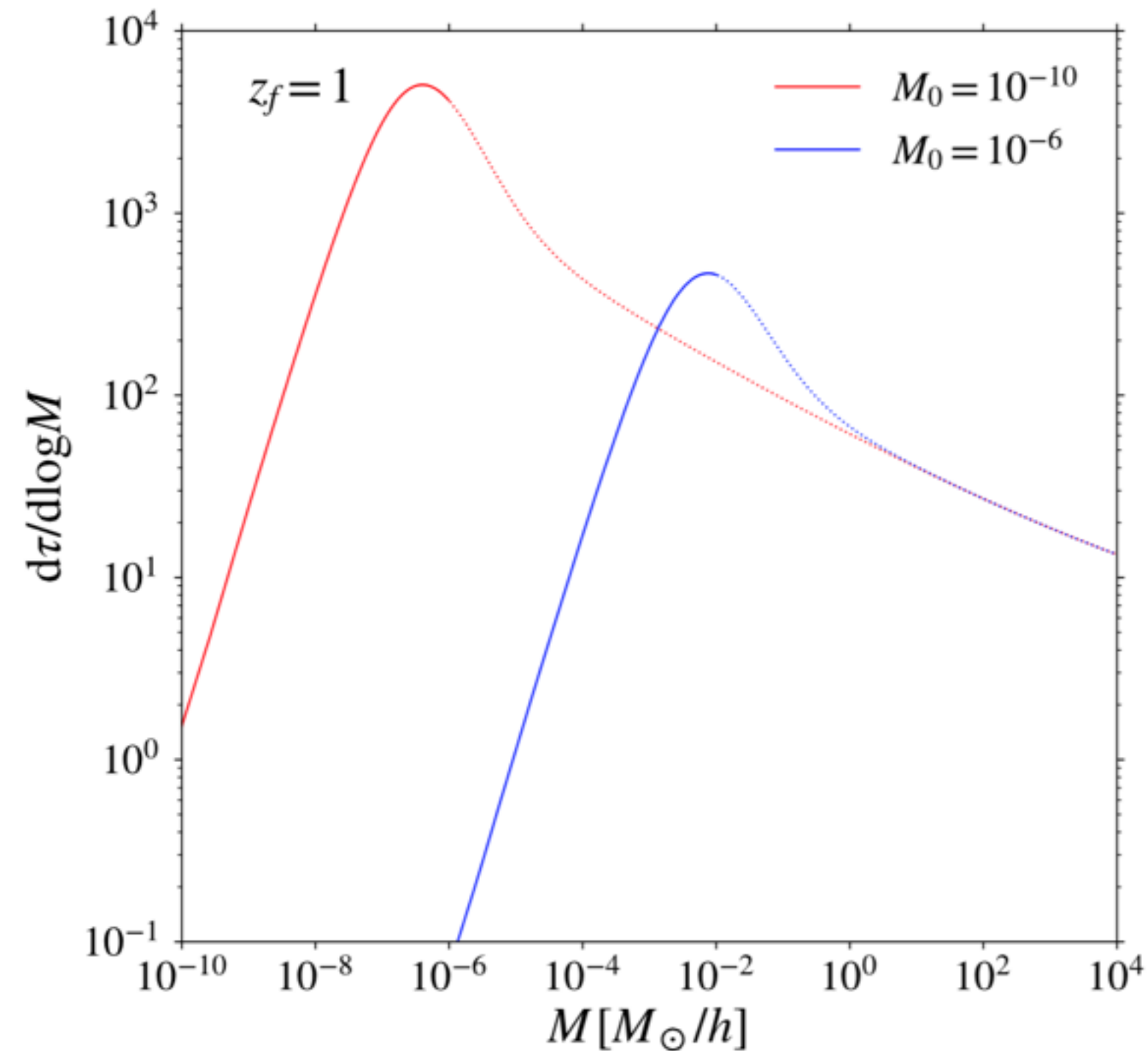
Many Minihalos Stacked Along LoS

line of sight toward the highly magnified star



galaxy cluster

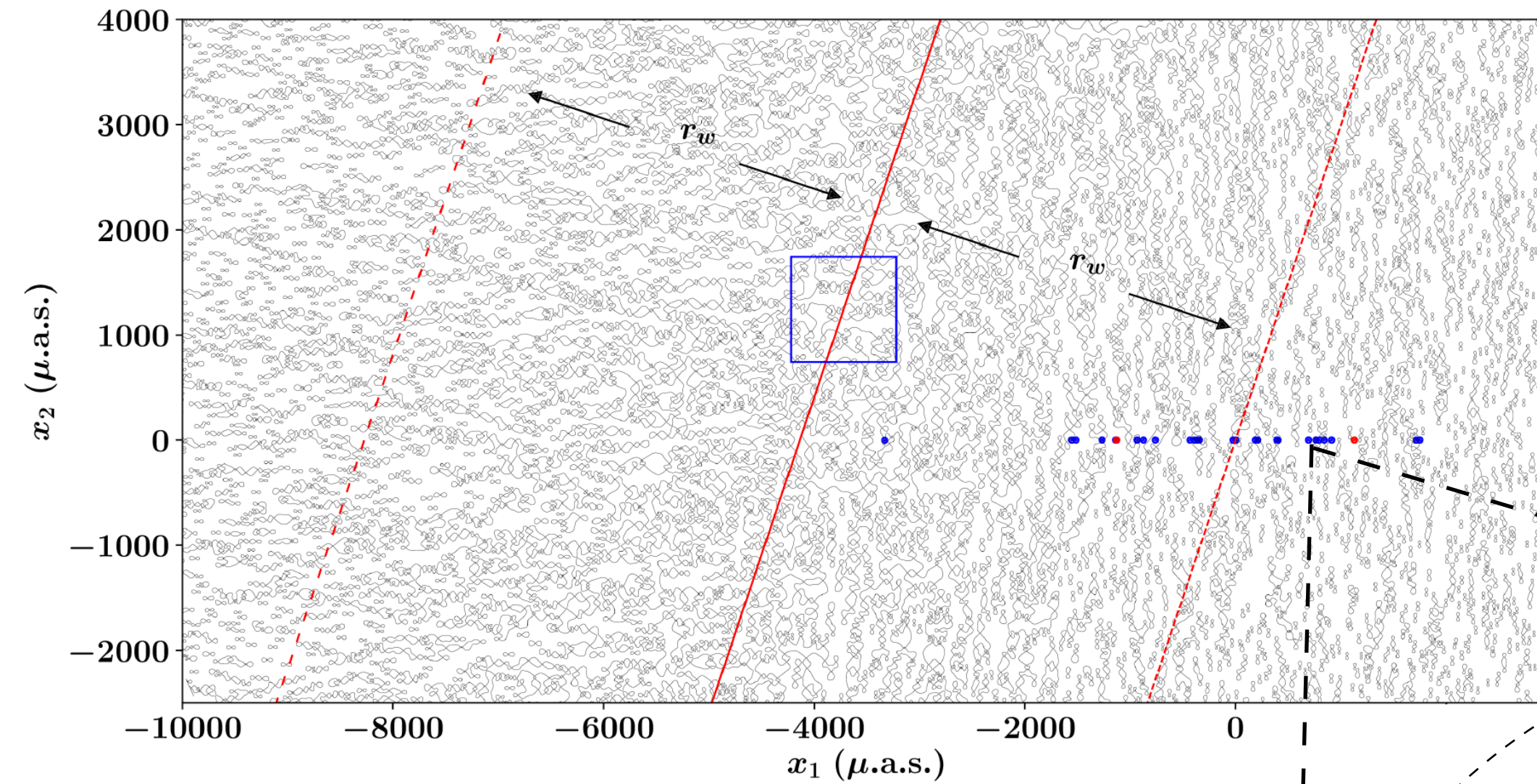
(minihalos too small and too numerous to be shown here)



volume occupation factor $\ll 1$
area covering factor $\gg 1$

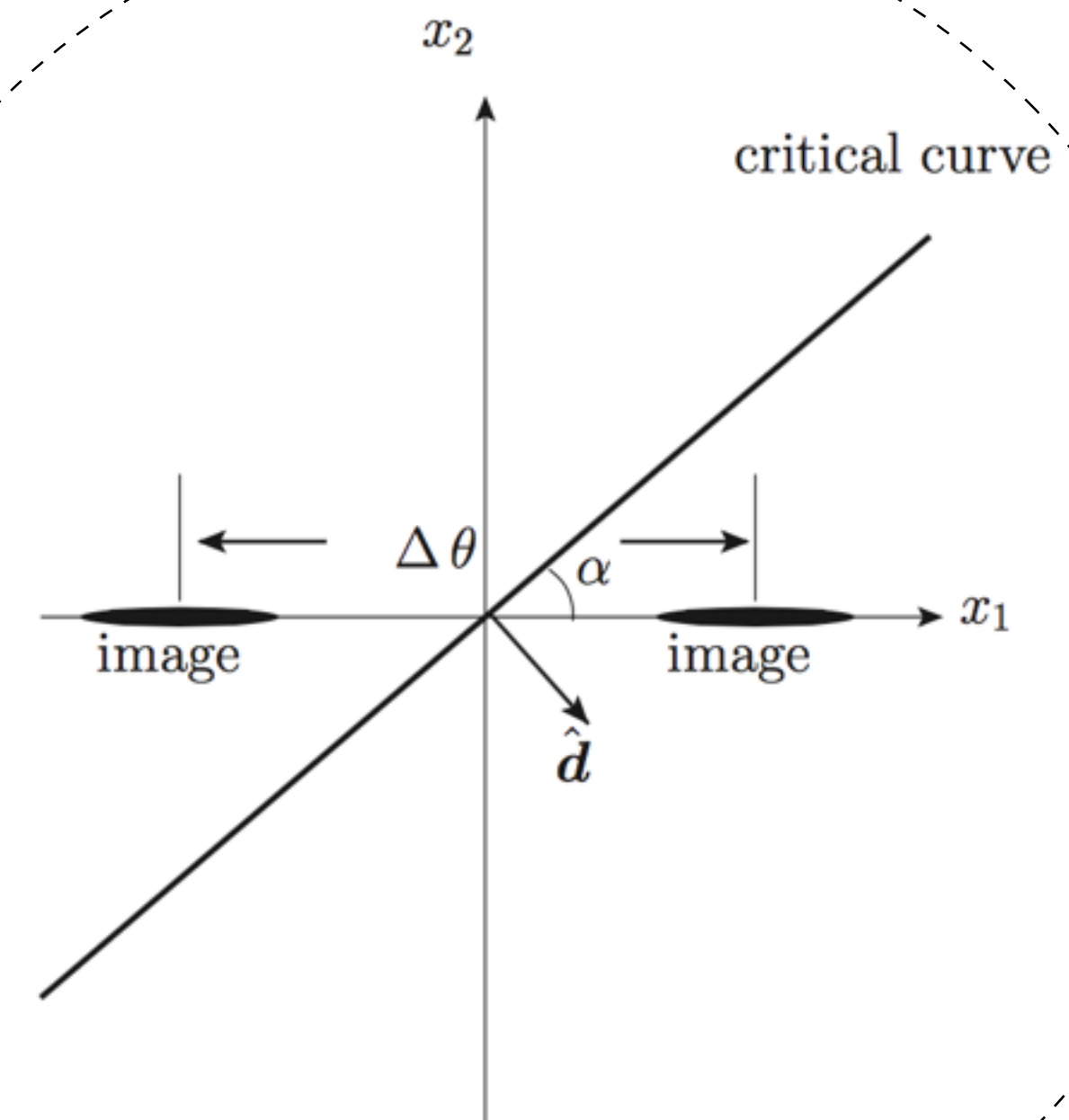
(Gaussianized) surface density fluctuations
 $\sim 10^{-4} - 10^{-3}$ on scales $\sim 10^2 - 10^4$ AU

Minuscule Fluctuations Can Matter ?!

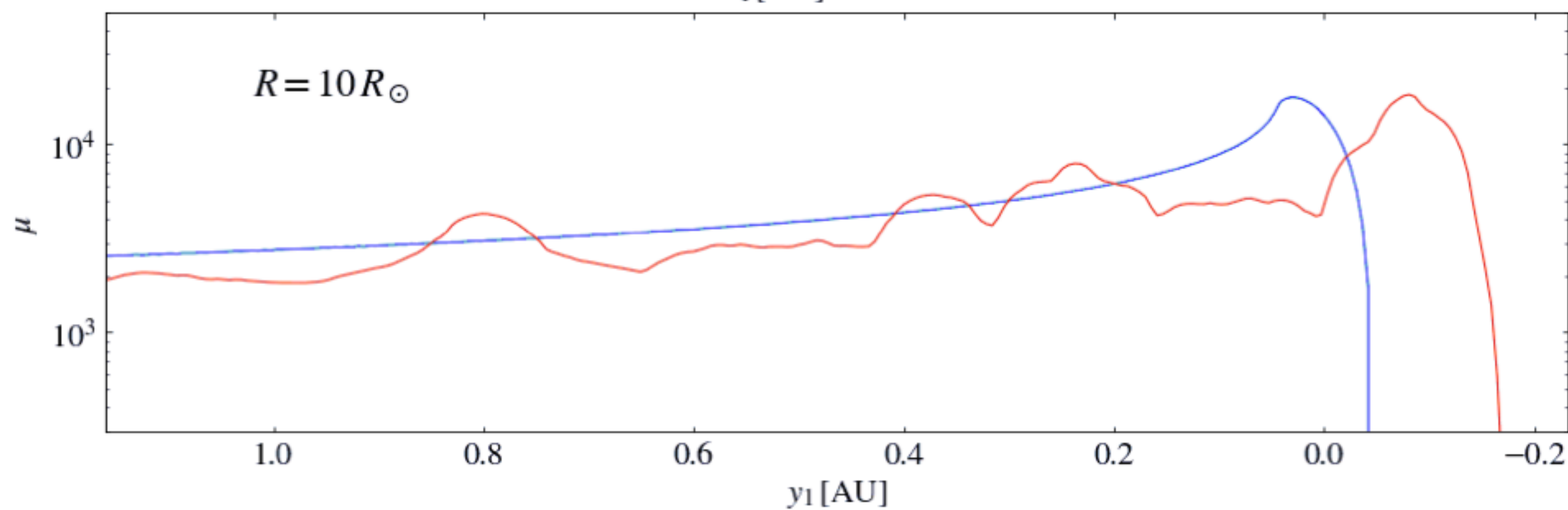
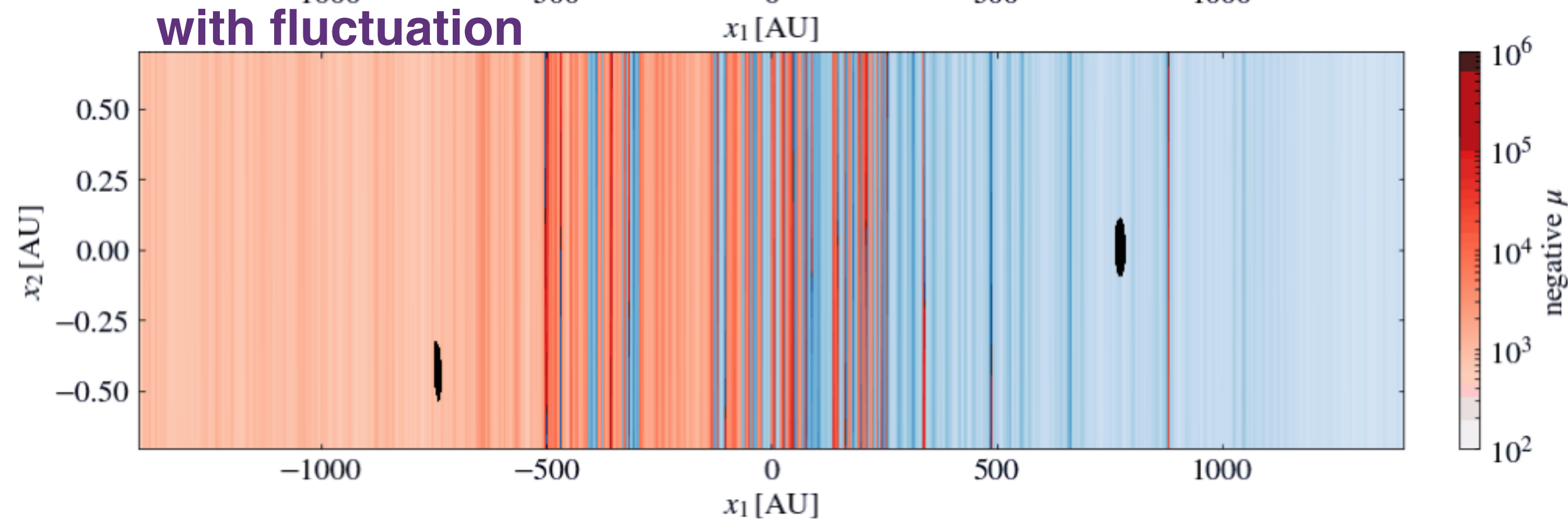
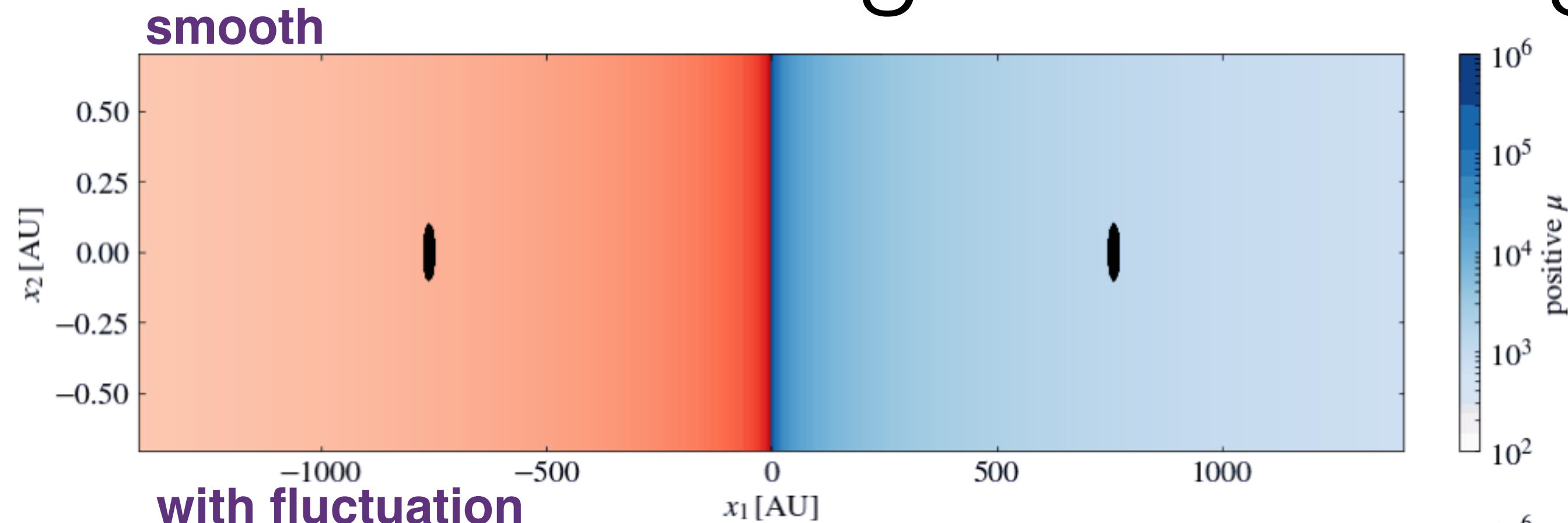


Around the times of microlensing peaks, two micro images are magnified up to $\sim 10^3\text{--}10^4$

If surface density varies on scales close to the size of the (elongated) micro images, fractional changes on the order of **the inverse magnification** alter the flux by **order unity**!



Light Curve Irregularities



An example of a microlensing peaking event

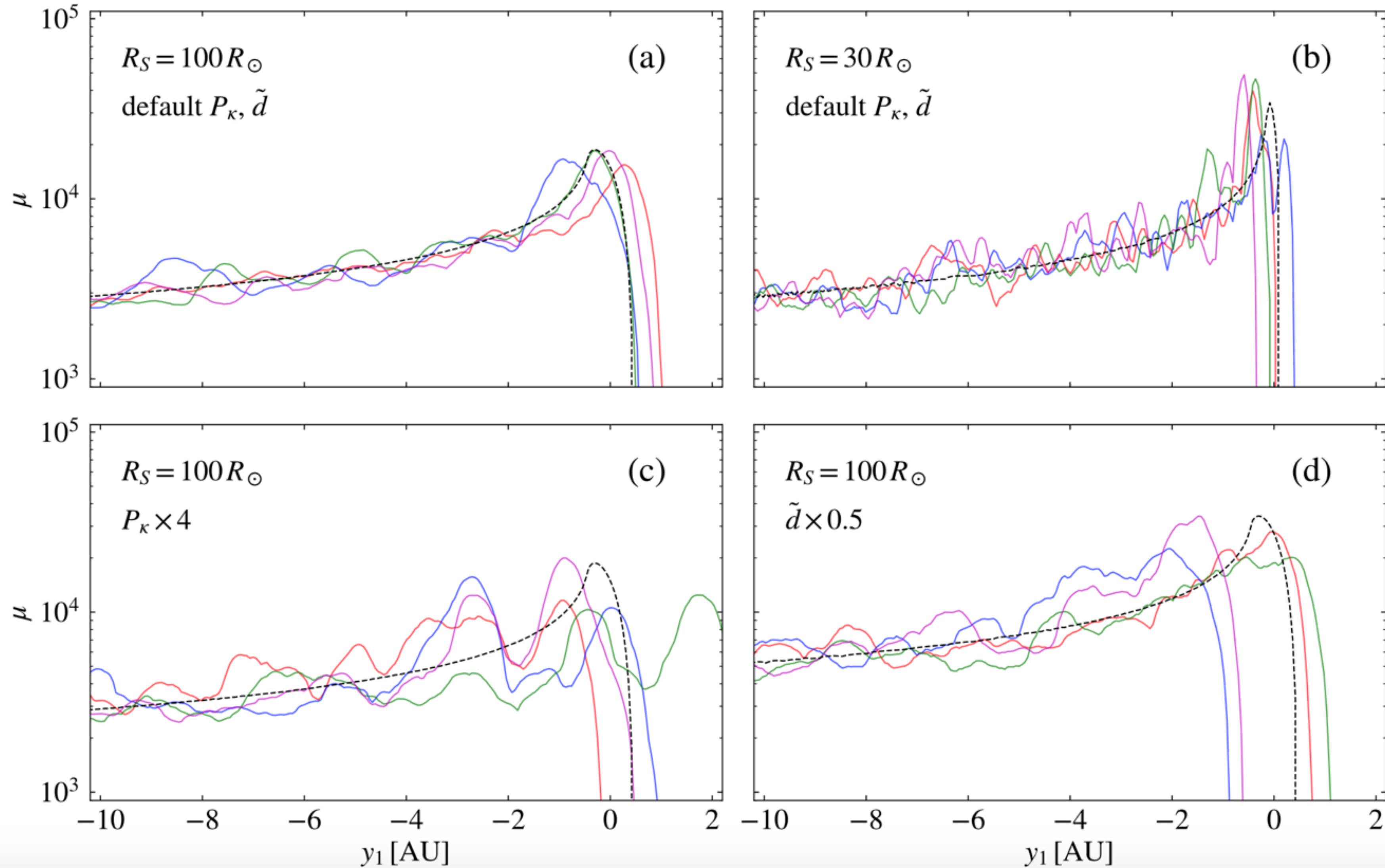
Parameters inferred for **MACS J1149 LS1**

$$\text{time to traverse 1 AU} = 6 \text{ d} \left(\frac{v_t}{300 \text{ km/sec}} \right)$$

Conclusion:

We should follow up microlensing peak events with dedicated observations !

Light Curve Irregularities



Dependence on:
stellar size, convergence
fluctuation,
and peak magnification.

Summary

- The axion particle is a viable candidate for CDM. Additional gravitationally bound structures can form on very small scales if Peccei-Quinn phase transition occurs after inflation. However, gravitational lensing detection of individual mini-halos seems infeasible.
- High abundance of mini halos are expected in intracluster space and are expected to survive the dynamic environment there. They collectively produce subtle surface density fluctuations on minuscule scales.
- Highly magnified extragalactic stars have been observed near the lensing caustic behind galaxy cluster lenses. They undergo frequency microlensing events due to intracluster stars.
- Order unity irregularities in the light curve can be produced by the minihalo population, whenever a caustic crossing star is temporarily magnified by a tremendous factor ($\sim 10^4$)