



Fermionic dark matter on galaxy scales

observations, theory and predictions

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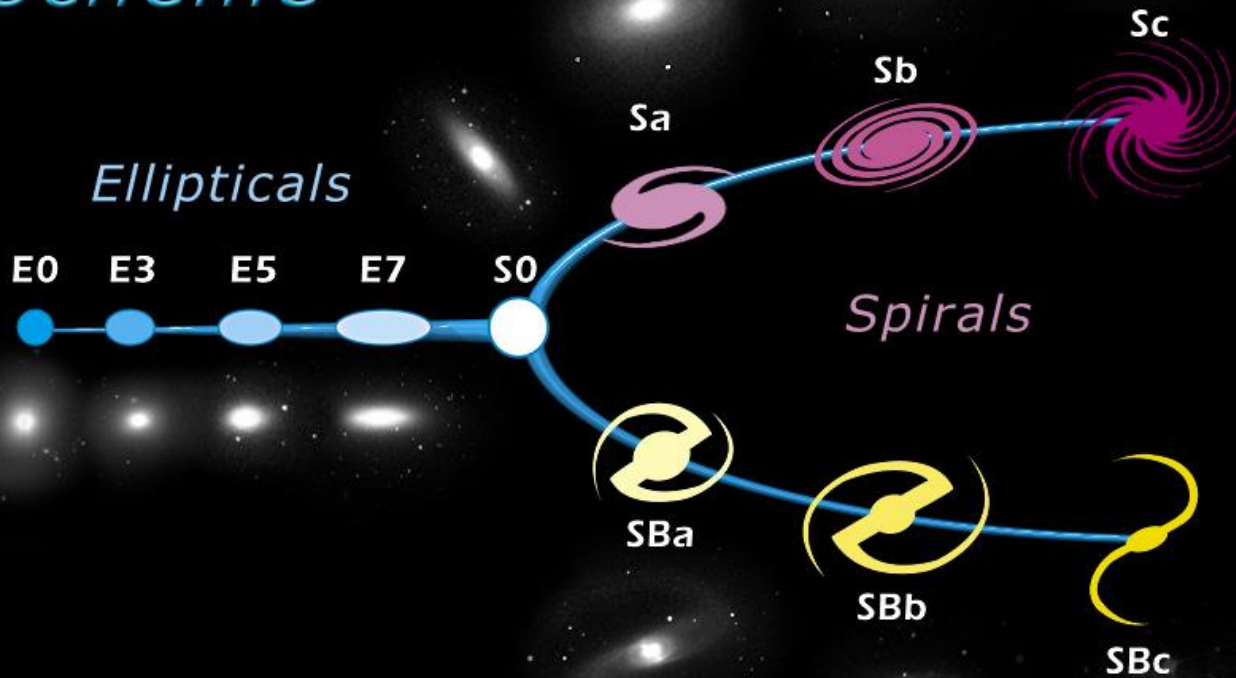


- Short introduction
dark matter and galaxy structures
- Fermionic Dark Matter
Ruffini-Argüelles-Rueda (RAR)
- Results
predictions for dwarfs to ellipticals
- Conclusion

Galaxy Morphology



Edwin Hubble's Classification Scheme



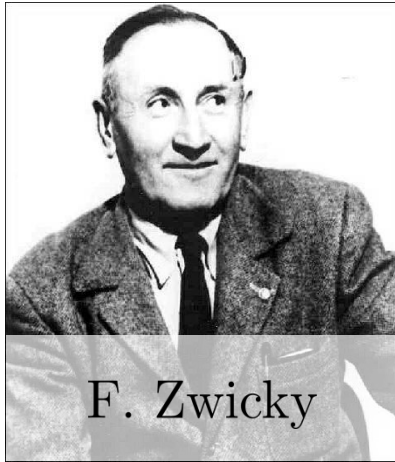
More than baryonic matter!?



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Flat rotation curve

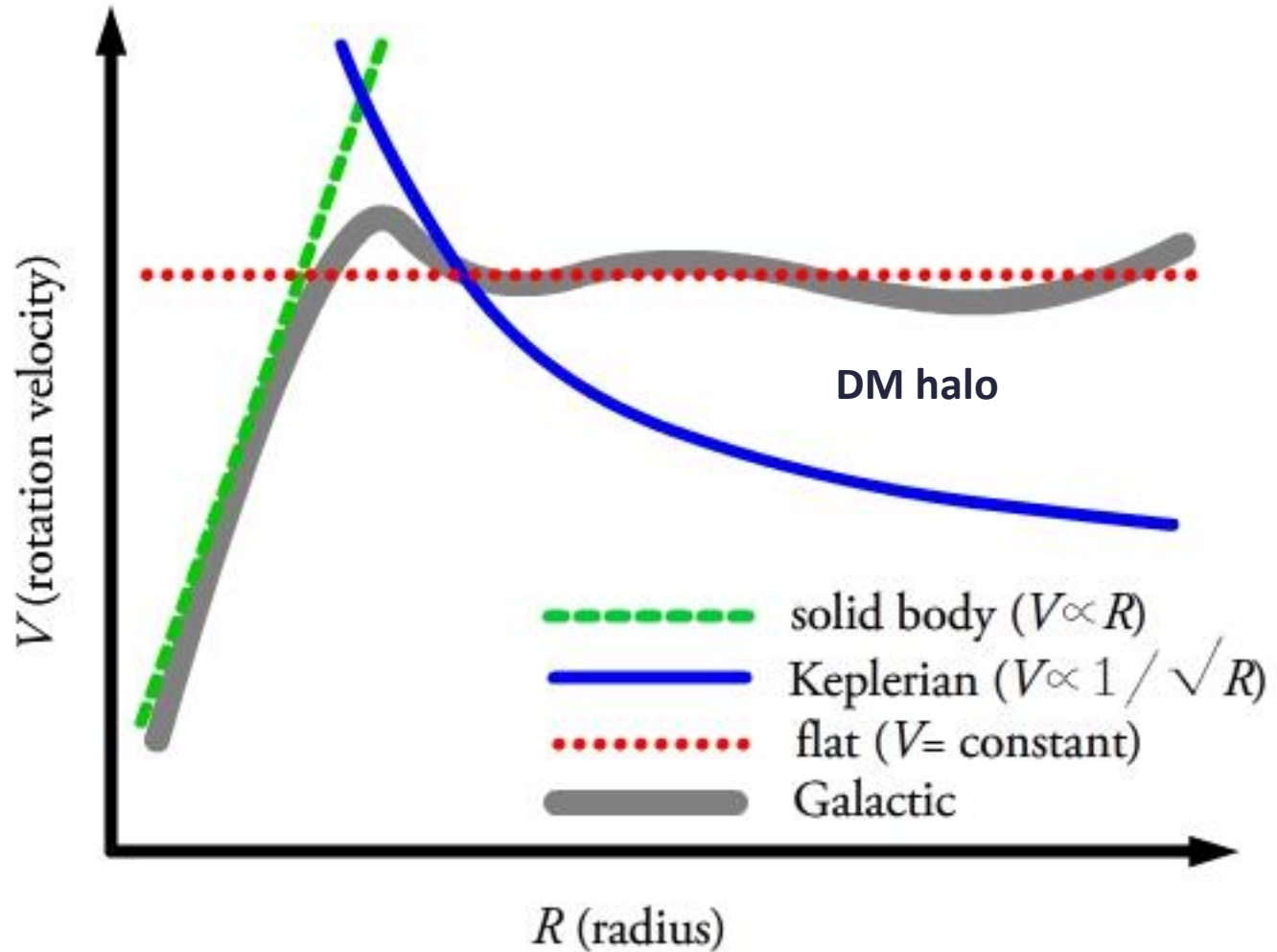
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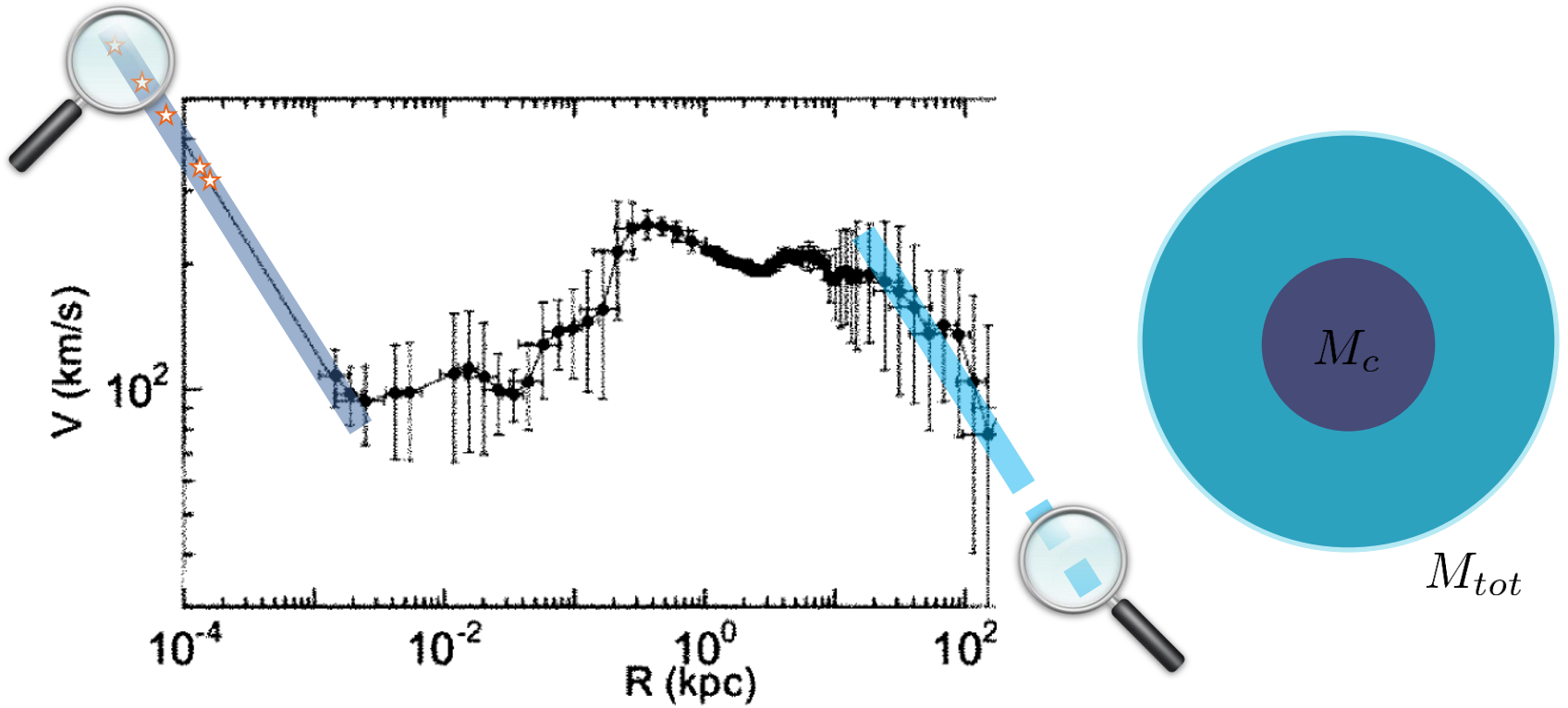
F. Zwicky



V. Rubin



Miky Way rotation curve



The outer halo

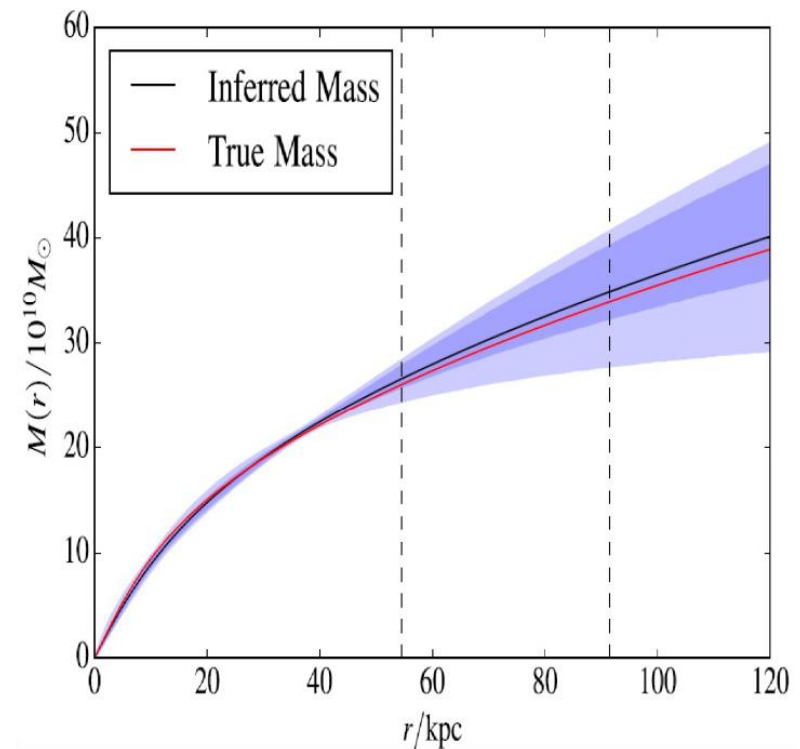
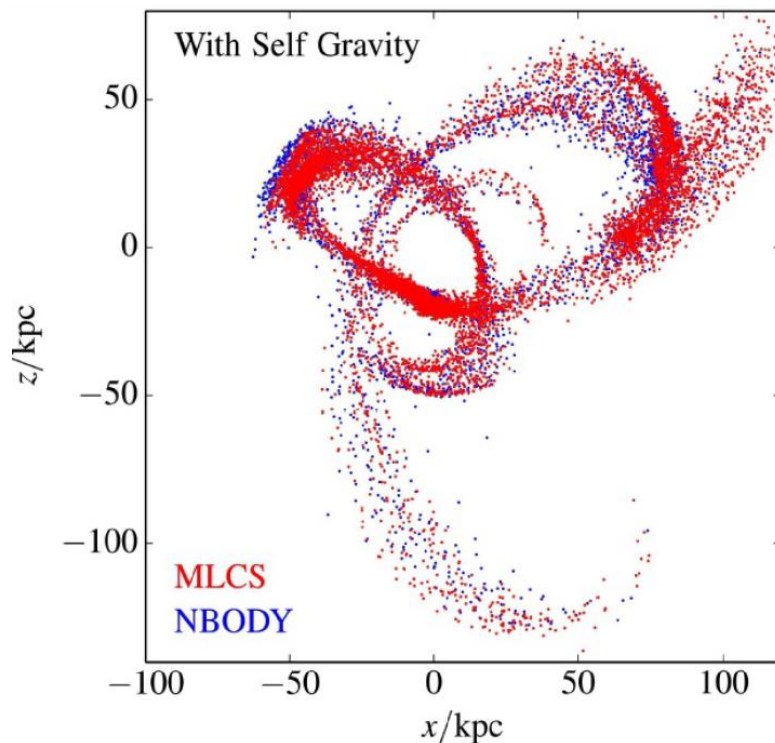


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satellite galaxies and streams

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- Outermost satellite galaxies of the MW are excellent total DM tracers
- The Sgr-dwarf satellite with its stream motion of tidally disrupted stars was well observed and well reproduced numerically



The Galactic center

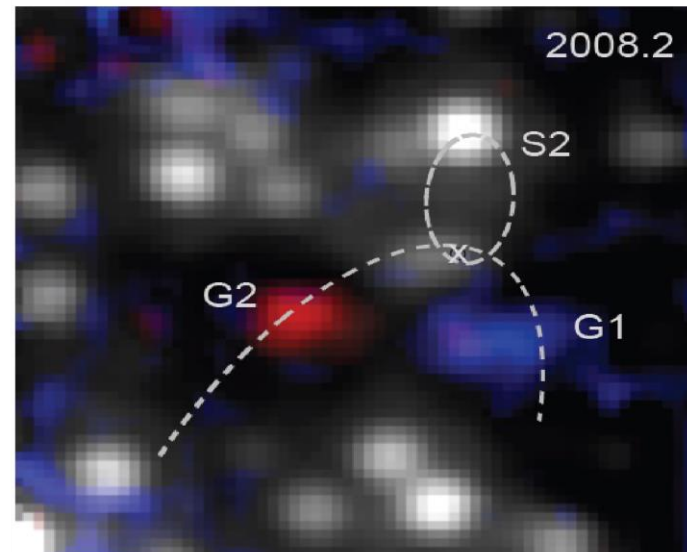
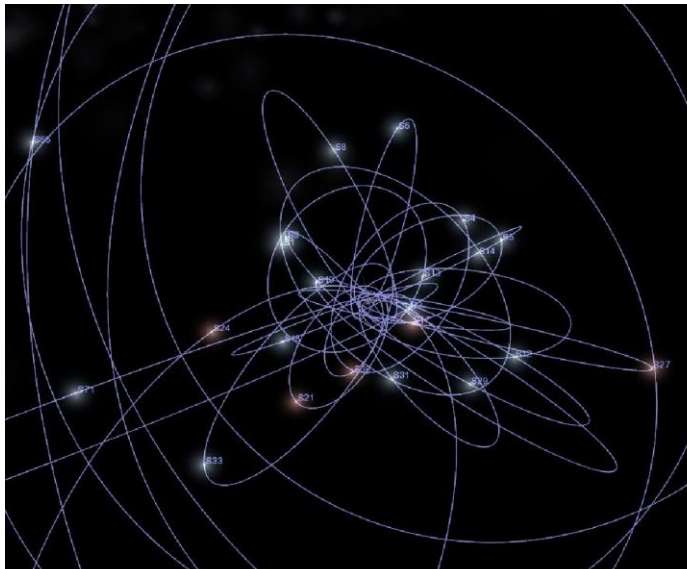


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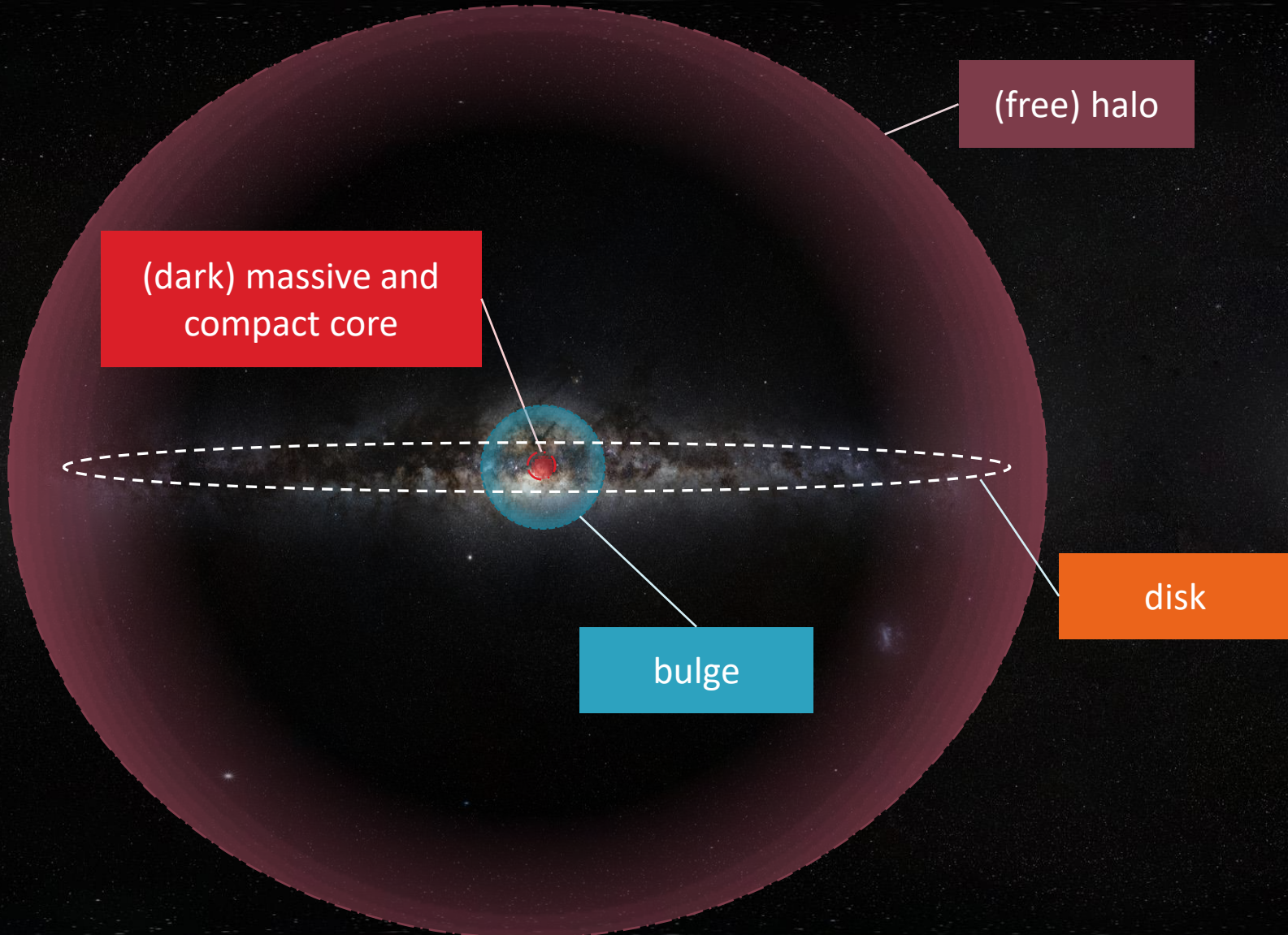
The S-stellar cluster & central gas

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- The central region ($10^{-3}\text{pc} \lesssim r \lesssim 2\text{pc}$) consist of young so-called S-stars and molecular gas
- Stars and gas dynamics are well described by Keplerian law ($v \sim r^{1/2}$)
- A dark compact object ($M \approx 4.2 \times 10^6 M_{\odot}$) is inferred within the smallest pericenter of $r_p(S_2) \approx 6 \times 10^{-4}\text{pc}$



Milky Way illustration



Basic Questions



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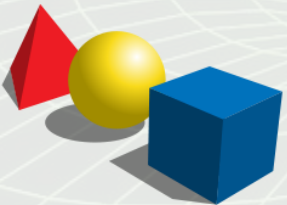
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How much of DM exists (in Universe)?
And **where** do we find it?



How does DM affect *structures*?



What is DM?

Lambda-CDM

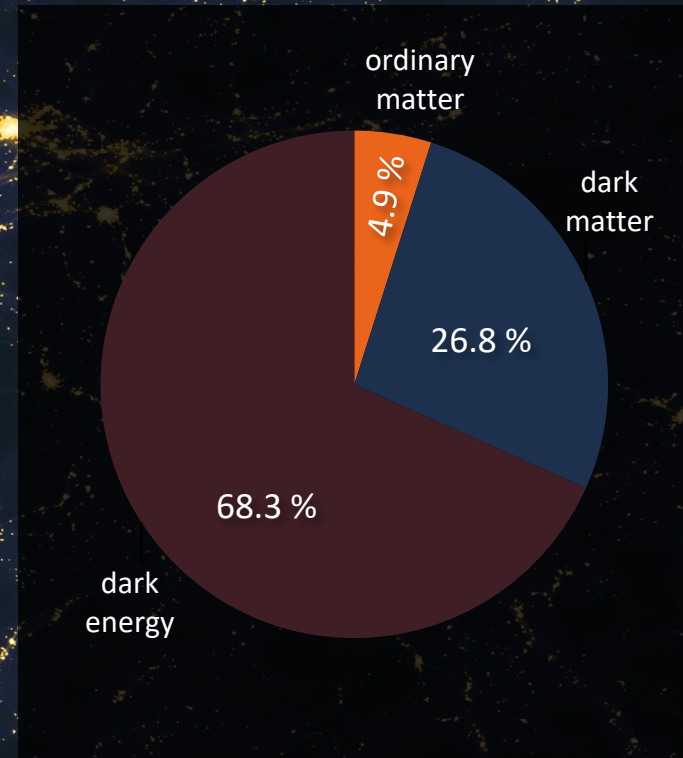


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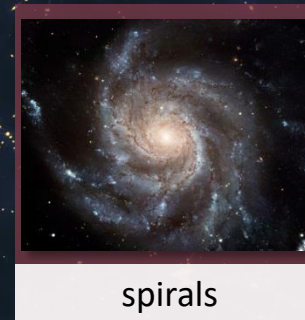
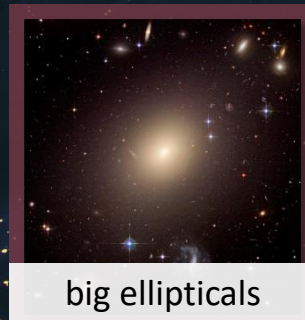
The Large Structures of the Universe

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- Λ CDM reproduces very well the Large Structure of Universe
- But few problems on smaller scales remain



> 10 Mpc



Dark Matter candidates



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particle candidate beyond SM

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↑ particle mass
 GeV/c^2
 MeV/c^2
 keV/c^2
 eV/c^2
 meV/c^2

CDM

- WIMPs
weakly interacting massive particles



WDM

- Sterile Neutrino
most promising candidate



HDM

- Neutrinos
ruled out in 80s



Fermionic DM with Cutoff



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Self-gravitating system of massive fermions in spherical symmetry

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collisionless relaxation
& escape of particles

phase space density

$$f(r, \epsilon) = \frac{1 - e^{[\epsilon - \epsilon(r)]/\beta(r)}}{e^{[\epsilon - \alpha(r)]/\beta(r)} + 1}$$

with $1 \leq \epsilon \leq \epsilon(r)$

energy $\epsilon^2 = 1 + \frac{p^2}{mc^2}$

temperature parameter

$$\beta(r) = \frac{k_B T(r)}{mc^2}$$

chemical potential (relativistic)

$$\alpha(r) = 1 + \beta(r)\theta(r)$$

escape energy (relativistic)

$$\epsilon(r) = 1 + \beta(r)W(r)$$

degeneracy parameter $\theta(r) = \frac{\mu(r)}{k_B T(r)}$

cutoff parameter $W(r) = \frac{E_c(r)}{k_B T(r)}$

Ruffini R., Stella L., 1983, A&A, 119, 35

Merafina M., Ruffini R., 1989, A&A, 221, 4

Ruffini R., Argüelles C. R., Rueda J. A., 2015, MNRAS, 451, 622

Perfect fluid in equilibrium



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Self-gravitating system of massive fermions in spherical symmetry

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$$g^{\mu\nu} = \text{diag}(e^{\nu(r)}, -e^{\lambda(r)}, -r^2, -r^2 \sin^2 \vartheta)$$

spherically
symmetric

GR

metric
potential

$$\frac{\partial \nu}{\partial r/R} = \frac{R^2}{r^2} \left[\frac{M(r)}{M} + \frac{r^3}{R^3} \frac{P(r)}{\rho c^2} \right] \left[1 - \frac{R}{r} \frac{M(r)}{M} \right]^{-1}$$

mass

$$\frac{\partial}{\partial r/R} \frac{M(r)}{M} = \frac{r^2}{R^2} \frac{\rho(r)}{\rho}$$

TOV approach

Statistics

mass
density

$$\frac{\rho(r)}{\rho} = \frac{4}{\sqrt{\pi}} \int \epsilon^2 \sqrt{\epsilon^2 - 1} f(r, \epsilon) d\epsilon$$

pressure

$$\frac{P(r)}{\rho c^2} = \frac{4}{3\sqrt{\pi}} \int (\epsilon^2 - 1)^{3/2} f(r, \epsilon) d\epsilon$$

EOS needed!

Perfect fluid in equilibrium



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Self-gravitating system of massive fermions in spherical symmetry

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EOS needed!

Perfect fluid in equilibrium



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Self-gravitating system of massive fermions in spherical symmetry

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TOV approach

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density ▶

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EOS needed!

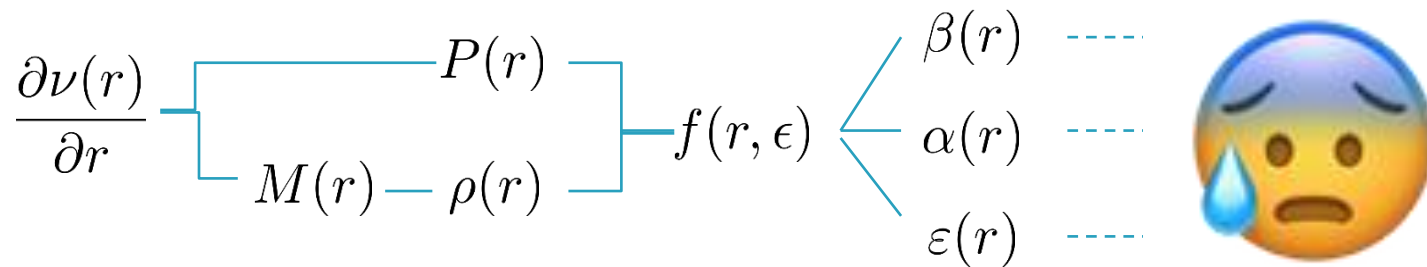
Thermodynamic equilibrium



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EOS

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conservation of energy along a geodesic

Tolman & Ehrenfest (1930)

$$\beta(r)e^{\nu(r)/2} = \text{const}$$

Klein (1949)

$$\alpha(r)e^{\nu(r)/2} = \text{const}$$

Merafina, Ruffini (1989)

$$\epsilon(r)e^{\nu(r)/2} = \text{const}$$

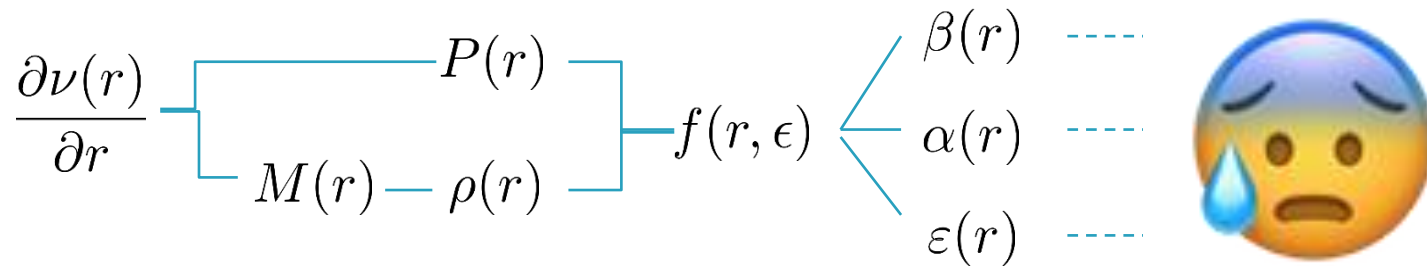
Thermodynamic equilibrium



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EOS

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dim-less initial conditions



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unit system and model parameter

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boundary condition

focus on the galactic center without BH!

$$\left. \frac{\partial M}{\partial r} \right|_{r=0} = 0 \rightarrow M_0 = 0$$

$$\left. \frac{\partial \nu}{\partial r} \right|_{r=0} = 0 \rightarrow \{\beta_0, \theta_0, W_0\}$$

unit system

eliminate constant particle mass from equations

$$\rho = \frac{gm^4}{h^3} [\pi c]^{3/2} \sim m^4$$

$$R = \left[\frac{c^2}{8\pi G\rho} \right]^{1/2} \sim m^{-2}$$

$$M = 4\pi R^3 \rho \sim m^{-2}$$

$$\{m, \beta_0, \theta_0, W_0\}$$

4 parameters

Milky Way fit



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constraints form core, halo and entire galaxy

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- boundary condition: r_h , $M_h = M(r_h)$, $M_{tot} = M(r_b)$ and $M_c = M(r_c)$ where r_c is core radius (of degenerate core), r_h is halo radius and r_b is boundary radius

observation

core

$$M_c = 4.2 \times 10^6 M_\odot$$

halo

$$M_h = 9 \times 10^{10} M_\odot$$

$$r_h = 20 \text{kpc}$$

total

$$M_{tot} = 2.4 \times 10^{11} M_\odot$$

Milky Way fit



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constraints form core, halo and entire galaxy

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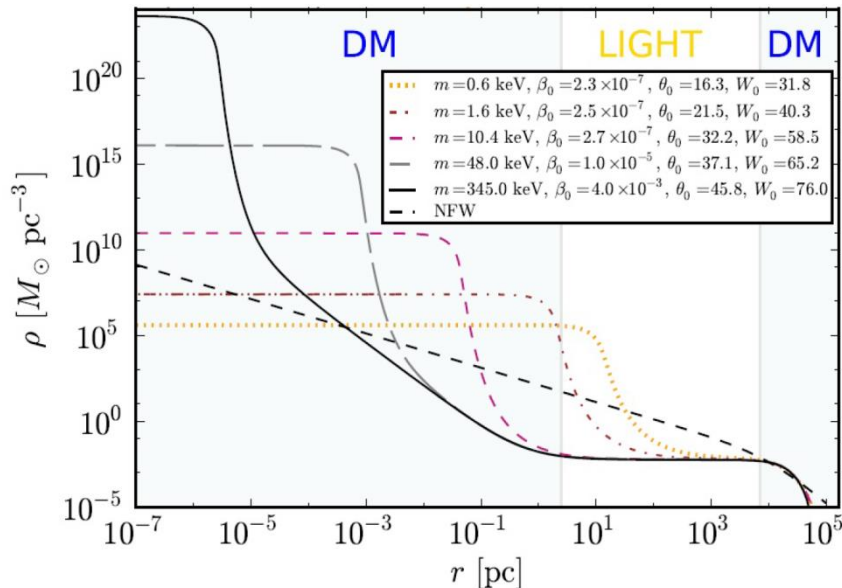
halo

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Milky Way fit



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constraints form core, halo and entire galaxy

21/32

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observation

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$$M_c = 4.2 \times 10^6 M_\odot$$

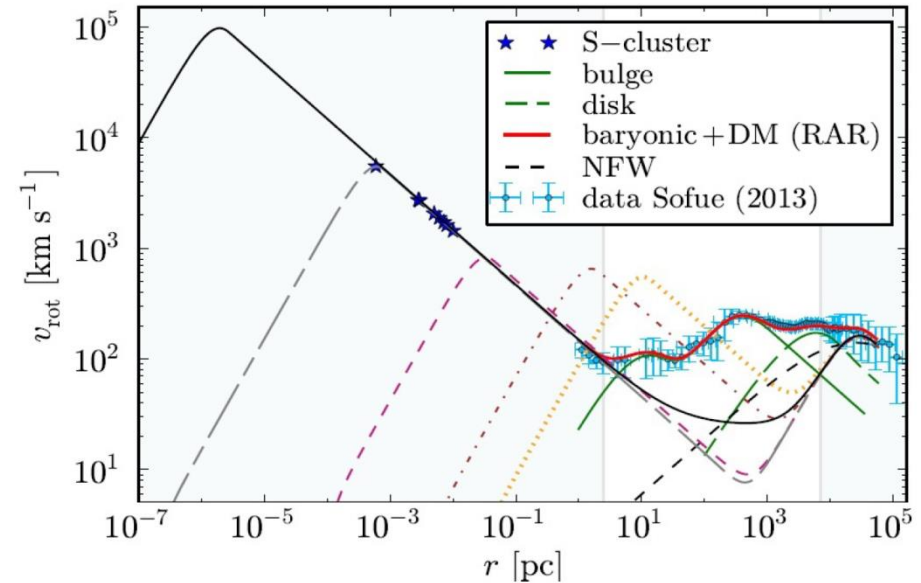
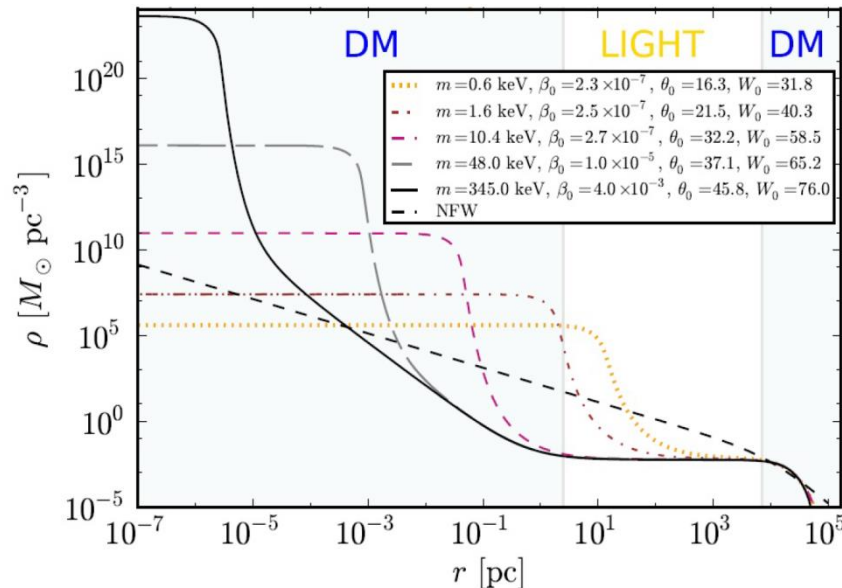
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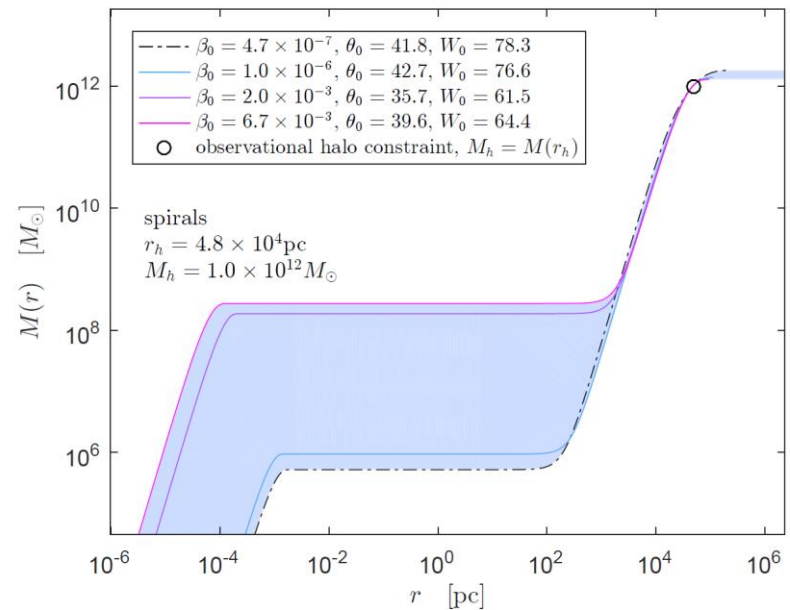
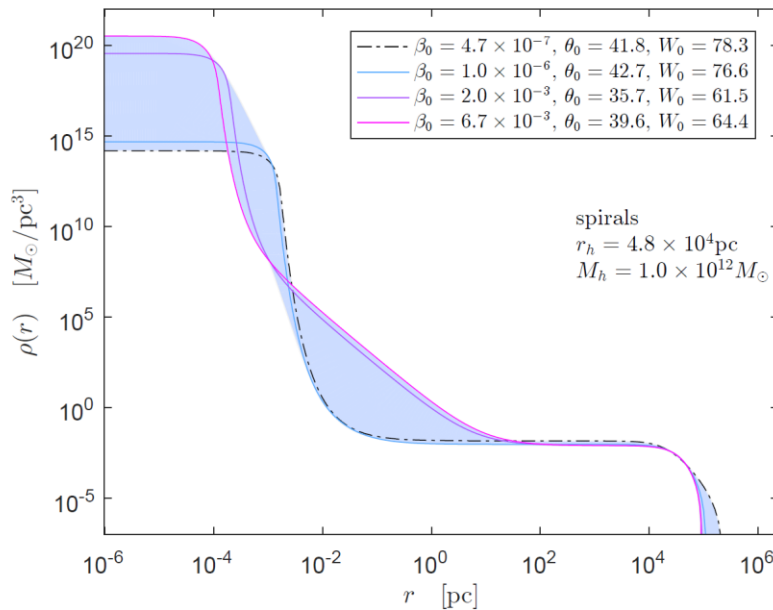
$$M_{tot} = 2.4 \times 10^{11} M_\odot$$



Typical Seyfert-like spirals



- The case of $m = 48$ keV fermionic dark matter
- Observationally-given boundary condition: $r_h = 48$ kpc and $M_h = 1 \times 10^{12} M_\odot$



core

$$M_c \in (4 \times 10^5, 2 \times 10^8) M_\odot$$

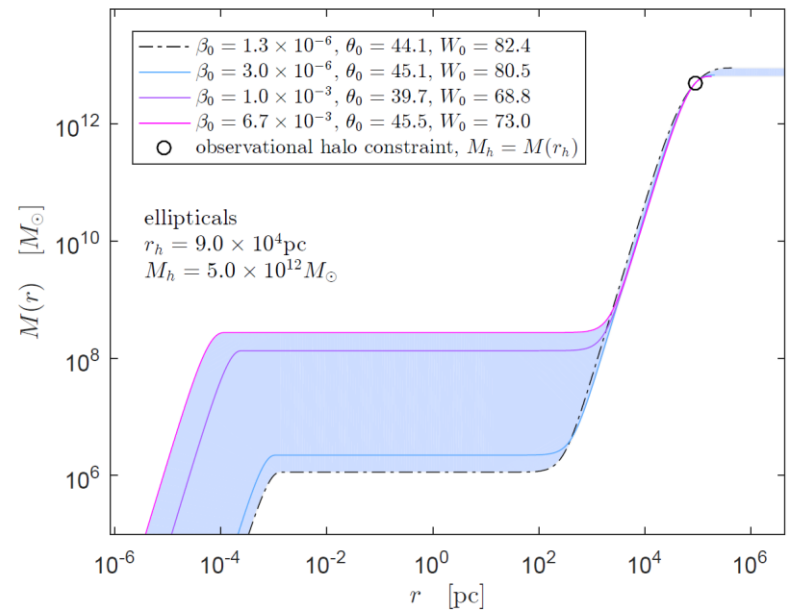
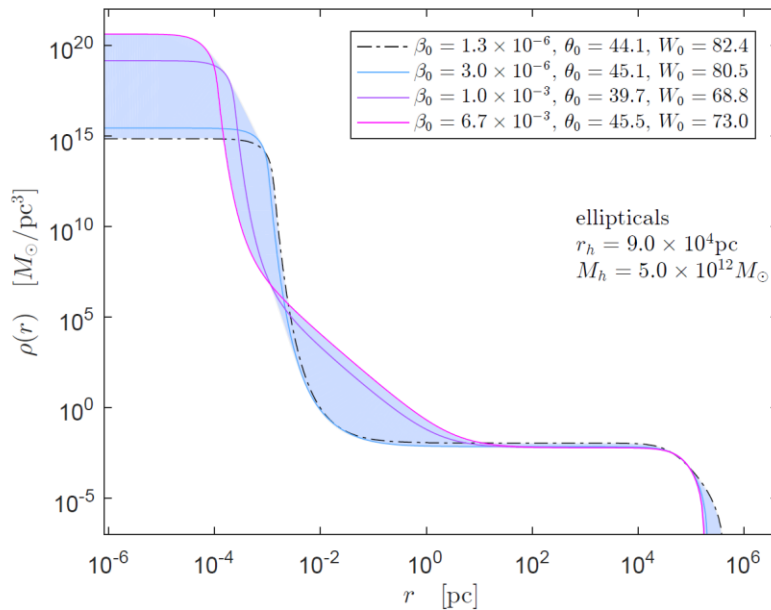
total

$$M_{tot} \in (1, 2) \times 10^{12} M_\odot$$

Typical normal ellipticals



- The case of $m = 48$ keV fermionic dark matter
- Observationally-given boundary condition: $r_h = 90$ kpc and $M_h = 5 \times 10^{12} M_\odot$



core

$$M_c \in (1 \times 10^6, 2 \times 10^8) M_\odot$$

total

$$M_{tot} \in (6, 9) \times 10^{12} M_\odot$$

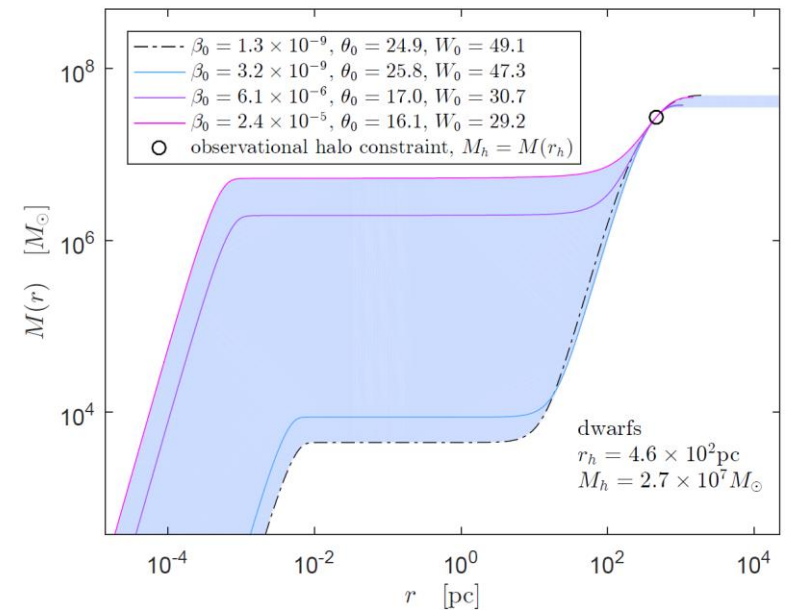
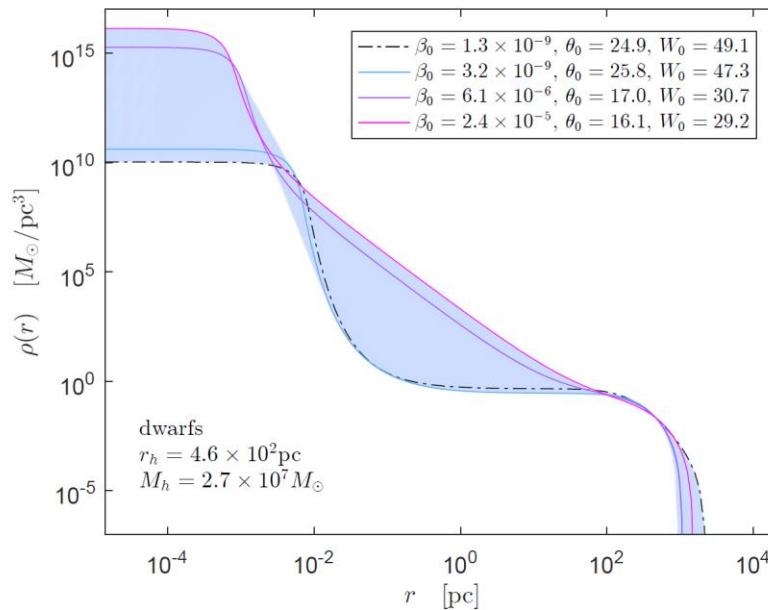
Typical dwarf spheroidals



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- The case of $m = 48$ keV fermionic dark matter
- Observationally-given boundary condition: $r_h = 0.46$ kpc and $M_h = 2.7 \times 10^7 M_\odot$



prediction

core

$$M_c \in (3 \times 10^3, 4 \times 10^6) M_\odot$$

total

$$M_{tot} \in (3, 5) \times 10^7 M_\odot$$

Prediction confirmed!



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Detection of Supermassive Black Holes in Two Virgo Ultracompact Dwarf Galaxies

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Abstract

We present the detection of supermassive black holes (BHs) in two Virgo ultracompact dwarf galaxies (UCDs), VUCD3 and M59cO. We use adaptive optics assisted data from the Gemini/NIFS instrument to derive radial velocity dispersion profiles for both objects. Mass models for the two UCDs are created using multi-band *Hubble Space Telescope* imaging, including the modeling of mild color gradients seen in both objects. We then find a best-fit stellar mass-to-light ratio (M/L) and BH mass by combining the kinematic data and the deprojected stellar mass profile using Jeans Anisotropic Models. Assuming axisymmetric isotropic Jeans models, we detect BHs in both objects with masses of $4.4_{-3.0}^{+2.5} \times 10^6 M_{\odot}$ in VUCD3 and $5.8_{-2.8}^{+2.5} \times 10^6 M_{\odot}$ in M59cO (3σ uncertainties). The BH mass is degenerate with the anisotropy parameter, β_z ; for the data to be consistent with no BH requires $\beta_z = 0.4$ and $\beta_z = 0.6$ for VUCD3 and M59cO, respectively. Comparing these values with nuclear star clusters shows that, while it is possible that these UCDs are highly radially anisotropic, it seems unlikely. These detections constitute the second and third UCDs known to host supermassive BHs. They both have a high fraction of their total mass in their BH; $\sim 13\%$ for VUCD3 and $\sim 18\%$ for M59cO. They also have low best-fit stellar M/L s, supporting the proposed scenario that most massive UCDs host high-mass fraction BHs. The properties of the BHs and UCDs are consistent with both objects being the tidally stripped remnants of $\sim 10^9 M_{\odot}$ galaxies.

Galaxy universal correlations



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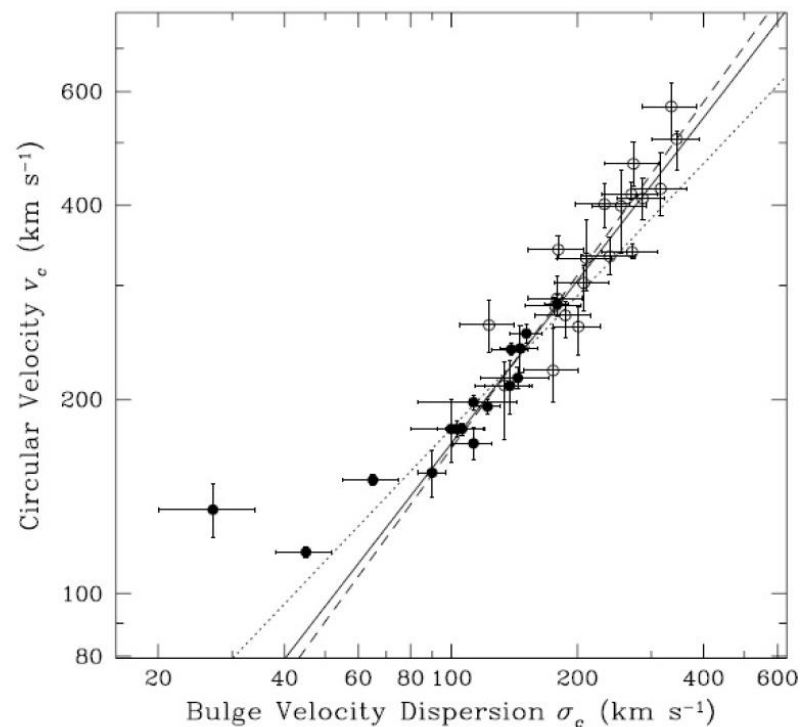
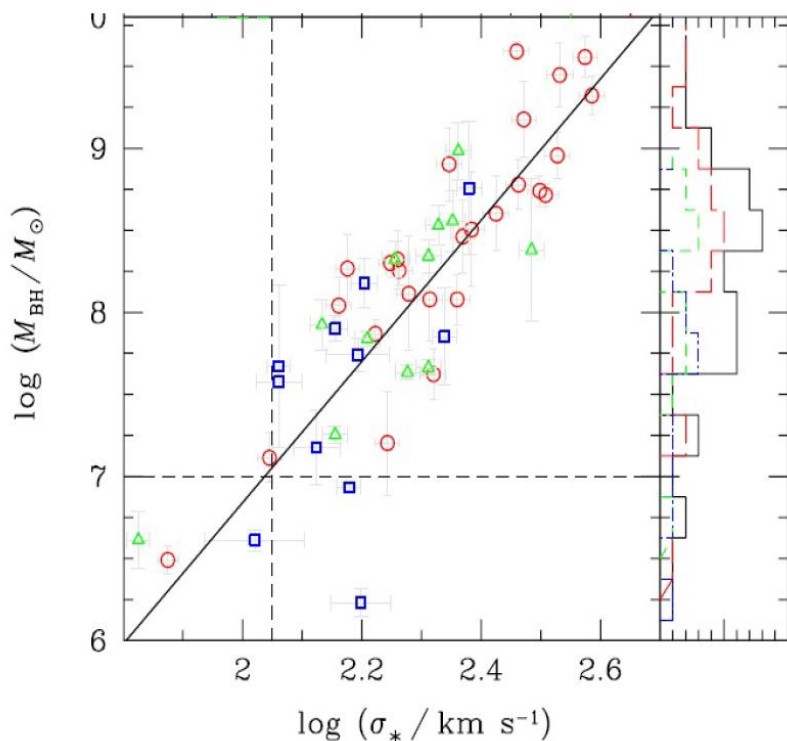
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- $\log(M_{\text{BH}}/M_{\odot}) \approx 8.12 + 4.24 \log(\sigma_*/[200\text{km/s}])$

K. Gültekin et al. (Apj) 2009

- $\log(v_{\text{circ}}) \approx 0.84 \log(\sigma_*) + 0.55$

L. Ferrarese (Apj) 2002; Kormendy et al. (Nature) 2011



The core-halo correlation

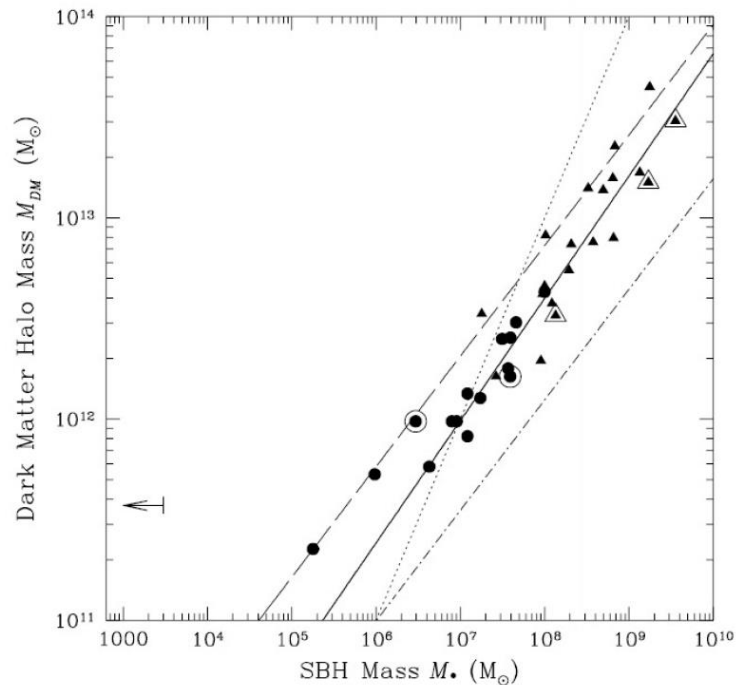


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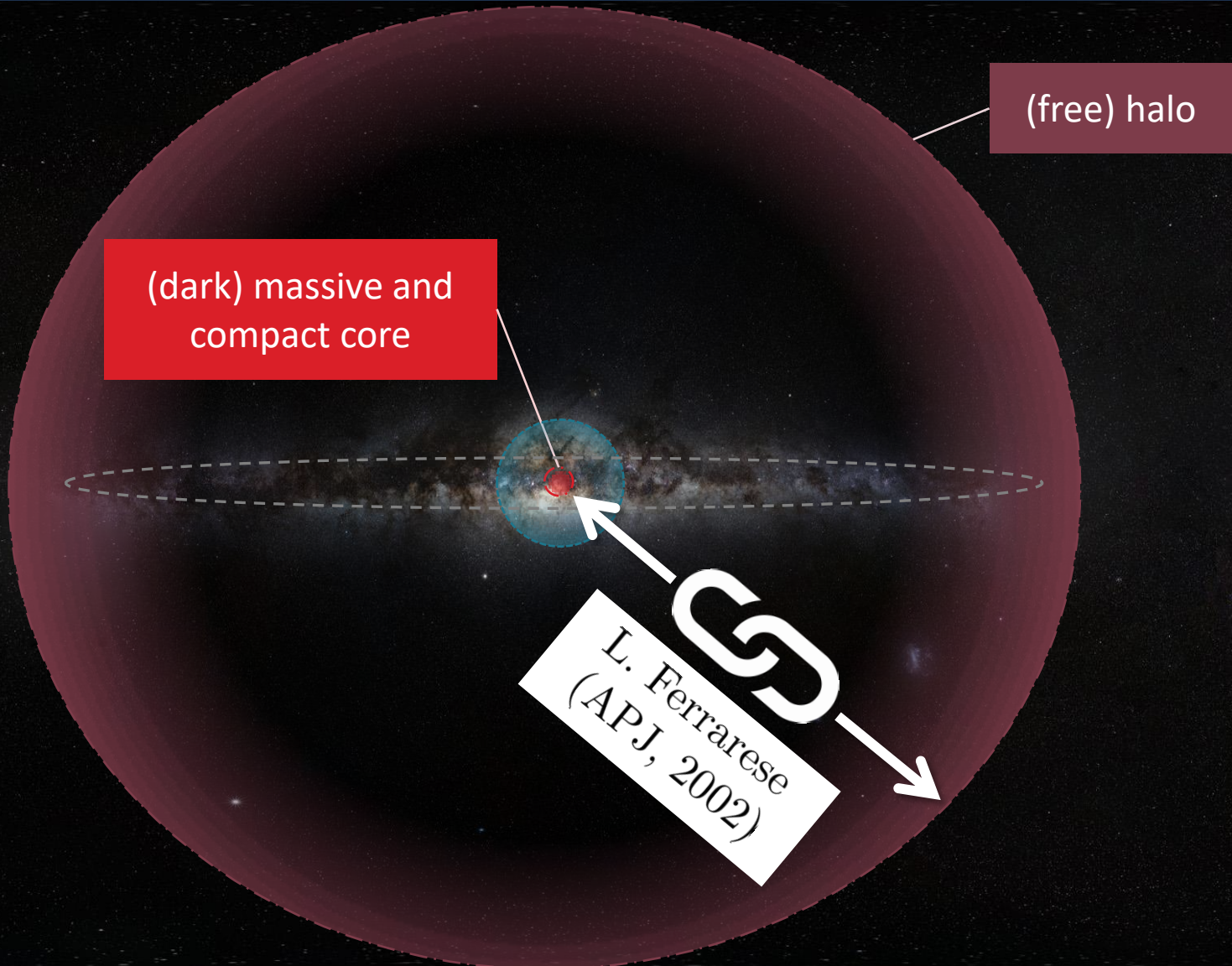
and RAR predictions

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- By transitivity between the $M_{\text{BH}}-\sigma_*$ and σ_*-v_{circ} correlations by Ferrarese '02, she found in 2002 the $M_{\text{BH}}-M_{\text{DM}}$ correlation L. Ferrarese (APJ, 2002)



Core-halo connection



The core-halo correlation

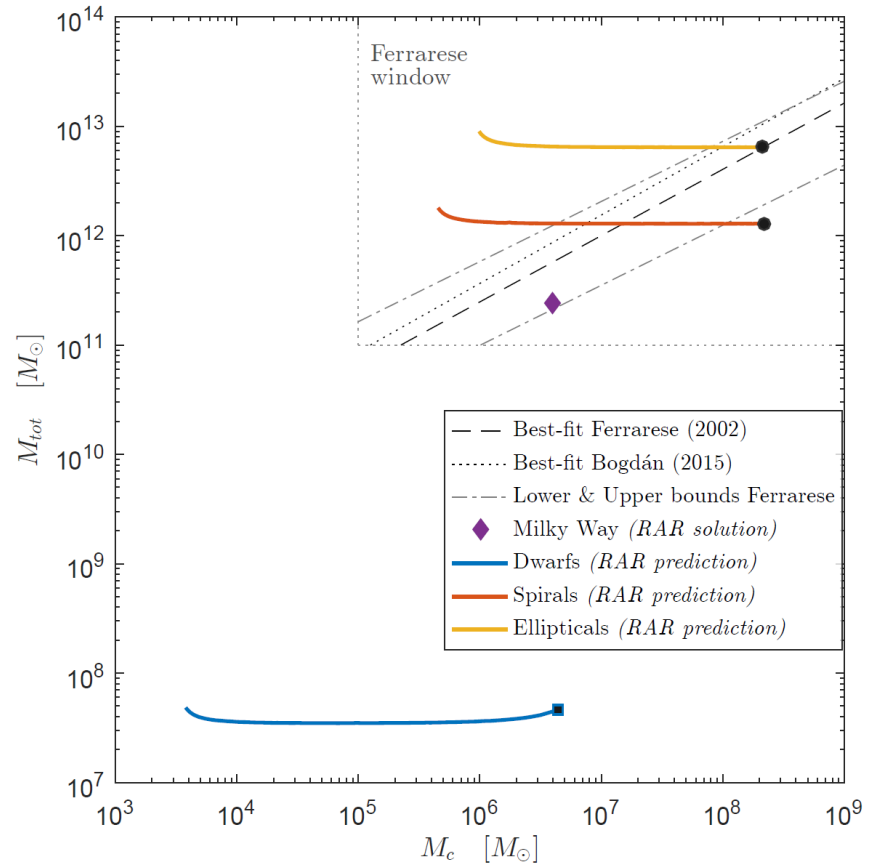
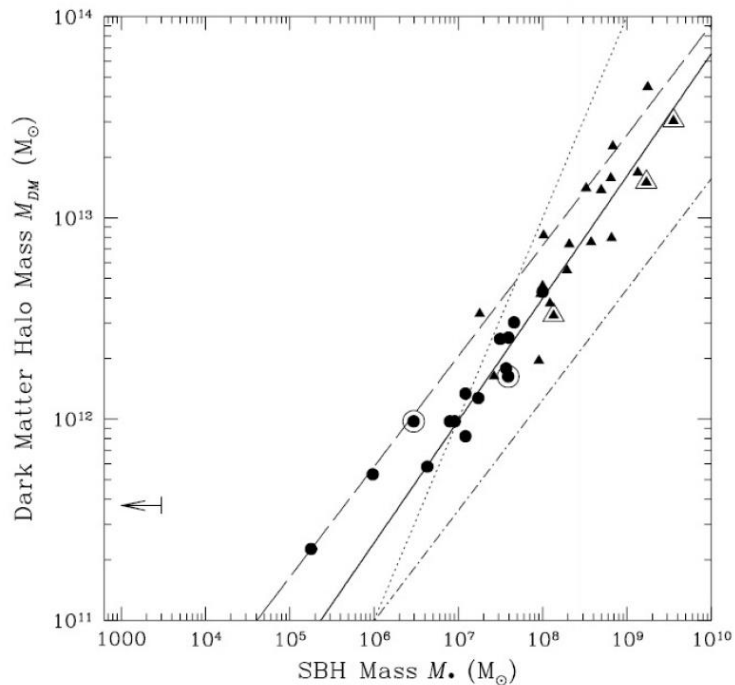


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and RAR predictions

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Normal and active galaxies

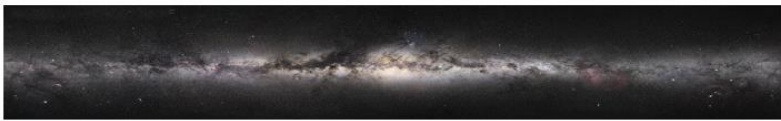


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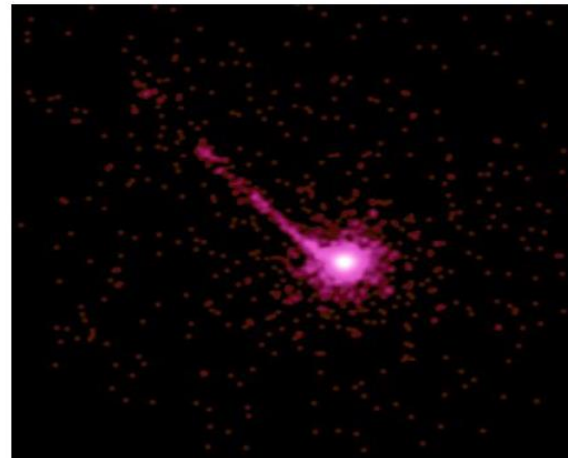
normal galaxies

no active nuclei nor jets ($M_c \sim 10^{6-7} M_\odot$)



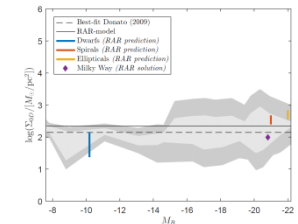
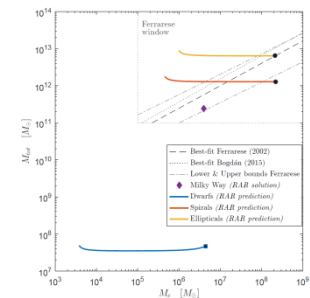
active galaxies

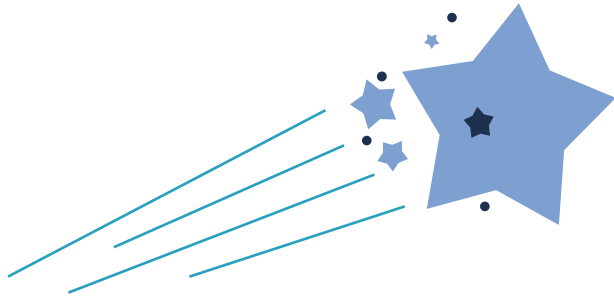
yes, active nuclei and jets ($M_c \sim 10^{9-10} M_\odot$)



Summary

- A continuous distribution of $\sim 10^1 - 10^2 \text{keV}$ fermions can be an **alternative to the black hole** scenario in SgrA* and being at the same time in agreement with the Milky Way DM halo without spoiling the baryonic (bulge and disk) components
- The RAR model is able to fulfill the observed properties of galaxies, such as the **$M_{\text{BH}}-M_{\text{DM}}$ relation** and the **$\Sigma_{0D} \approx \text{constant universal law}$** , for a unique DM fermionic mass
- For $m \sim 50 \text{keV}/c^2$ with massive fermionic quantum cores reaching a M_c^{cr} mass, it exists the interesting possibility that the less massive (not-yet collapsed) cores of $M_c \sim 10^{6-8} M_{\odot}$ **explain the less-active galaxies**; while the more massive cores of M_c^{cr} plus subsequent more violent baryonic accretion would **explain the most massive SMBHs** with associated jets \rightarrow the latter requires more work to be done
- High precision strong lensing measurements, as the expected EHT observations, may unveil the **nature of SgrA*** in the Galactic center and distinguish between a BH and a dense DM core \rightarrow first results expected in early 2018...





Thank you