# SUSY enhancement from T-branes 

Raffaele Savelli<br>University of Rome<br>"Tor Vergata"

Based on work with
F. Carta, S. Giacomelli
arXiv: I809.04906

## SUSY enhancement

- Remarkable phenomenon along certain RG flows of 4d SFT's. [Maruyoshi, Song `16]; [Agarwal, Maruyoshi, Song `I6, 'I8]
- 4d $\mathrm{N}=2$ SCFT w/ non-Abelian flavor $\mathrm{G}_{\mathrm{F}}+$ free chiral field $M \in \operatorname{Adj}\left(\mathrm{G}_{\mathrm{F}}\right)$.
- Deform with $\delta W=\operatorname{Tr}(\mu M): \mathscr{N}=2 \rightarrow \mathcal{N}=1 \quad \boldsymbol{\mu}$ is the flavor moment map.
- Give vev $\langle M\rangle=\rho\left(\sigma_{+}\right) \quad \rho: s u(2) \rightarrow G_{F}$ is a nilpotent orbit.
- Depending on $\rho, \delta \mathrm{W}$ may trigger RG flow $\mathrm{w} / \mathcal{N}=2$ fixed point!


## SUSY enhancement

- Very interesting, but also particularly useful:
- IR is non-Lagrangian, but UV may have Lagrangian.
- Use it to compute RG-invariant quantities for IR theory!


## SUSY enhancement

- Very interesting, but also particularly useful:
- IR is non-Lagrangian, but UV may have Lagrangian.
- Use it to compute RG-invariant quantities for IR theory!
- How was it discovered? a-maximization:
- $\langle M\rangle: \operatorname{Adj}\left(G_{F}\right) \rightarrow \oplus_{j} V_{j} \quad V_{j} \quad$ spin-j irrep. of su(2).
- Only $M \mathrm{w} / \mathrm{spin}\left(\mathrm{j}, \mathrm{j}_{3}=-\mathrm{j}\right)$ stay coupled, i.e. $\left[M, \rho\left(\sigma_{-}\right)\right]=0$.
- Maximize $32 \mathrm{a}=9 \mathrm{Tr}^{3} 3-3 \operatorname{Tr} \mathrm{R}$ w.r.t. trial R -charge $\mathrm{R}=\mathrm{R}_{0}+\epsilon \mathrm{F}$.
- Re-iterate until all fields have $\mathrm{D}_{*}=3 / 2 \mathrm{R}_{*}>\mathrm{I}$.


## SUSY enhancement

- Checks of enhancement: found rational IR central charges \& superconformal index in agreement w/ known $\mathcal{N}=2$ theories.
- Involved process, hiding the origin of the phenomenon.


## SUSY enhancement

- Checks of enhancement: found rational IR central charges \& superconformal index in agreement w/ known $\mathcal{N}=2$ theories.
- Involved process, hiding the origin of the phenomenon.
- Our purpose: show why it occurs by explaining it geometrically.
- Engineer rank-I SCFT's by probing F-theory singularities.
- Interpret deformations as T-brane backgrounds.
- Get simple algebraic criterion to exclude enhancement.


## SUSY enhancement

- Checks of enhancement: found rational IR central charges \& superconformal index in agreement w/ known $\mathcal{N}=2$ theories.
- Involved process, hiding the origin of the phenomenon.
- Our purpose: show why it occurs by explaining it geometrically.
- Engineer rank-I SCFT's by probing F-theory singularities.
- Interpret deformations as T-brane backgrounds.
- Get simple algebraic criterion to exclude enhancement.
- Bonus: in case of enhancement, IR conformal dimensions are derived algebraically!


## Warm-up: 3d SQED

- Easier setting, but key geometric features already apparent.
- Single D2 probing a stack of N D6-branes in flat space:
$\Rightarrow 3 \mathrm{~d} \mathscr{N}=4 \mathrm{U}(\mathrm{I})$ gauge theory $w /$ flavors in fund of $G_{F}=U(N)$.
$\Rightarrow$ IR dynamics described by an $M 2$ probing $\operatorname{ALE} \mathbb{C}^{2} / \mathbb{Z}_{N}$.


## Warm-up: 3d SQED

- Easier setting, but key geometric features already apparent.
- Single D2 probing a stack of N D6-branes in flat space:
$\Rightarrow$ 3d $\mathscr{N}=4 \mathrm{U}(\mathrm{I})$ gauge theory $w /$ flavors in fund of $G_{F}=U(N)$.
$\Rightarrow$ IR dynamics described by an M2 probing ALE $\mathbb{C}^{2} / \mathbb{Z}_{N}$.



## Deformations

- Add the IR relevant coupling $\delta W=\operatorname{Tr}(\mu M)$ :
- $\mu_{i}^{j} \equiv Q_{i} \tilde{Q}^{j}$ is the meson matrix.
- $M=\langle\boldsymbol{\Phi}\rangle$ is vev of scalar controlling D6's position in 8,9 plane.
- $M$ constant \& $\left[M, M^{\dagger}\right]=0 \quad$ N$=4$ preserving cplx masses.
- M 3d field


## Deformations

- Add the IR relevant coupling $\delta \mathrm{W}=\operatorname{Tr}(\mu M)$ :
- $\mu_{i}^{j} \equiv Q_{i} \tilde{Q}^{j}$ is the meson matrix.
- $M=\langle\Phi\rangle$ is vev of scalar controlling D6's position in 8,9 plane.
- $M$ constant \& $\left[M, M^{\dagger}\right]=0 \quad$ lilt $\mathcal{N}=4$ preserving cplx masses.
- M 3d field Int D6's coordinate dependence $<\Phi>\left(s_{1}, S_{2}\right)$.
- Give $M$ nilpotent vev ${ }^{\mathrm{In}} \quad\left[M_{1}, M^{\dagger}\right] \neq 0 \quad \& \quad \mathcal{N}=4 \rightarrow \mathcal{N}=2$ :
- Probing a T6-brane background.
[Cecotti, Cordova, Heckman, Vafa `IO]
- 2 fields among $M_{(\mathrm{j},-\mathrm{j})} \leftrightarrow s_{1}, s_{2} \quad$ take the 2 highest spins.
- CB gets deformed:


## Deformations

- Add the IR relevant coupling $\delta \mathrm{W}=\operatorname{Tr}(\mu M)$ :
- $\mu_{i}^{j} \equiv Q_{i} \tilde{Q}^{j}$ is the meson matrix.
- $M=\langle\Phi\rangle$ is vev of scalar controlling D6's position in 8,9 plane.
- $M$ constant \& $\left[M, M^{\dagger}\right]=0 \quad$ lilt $\mathcal{N}=4$ preserving cplx masses.
- M 3d field Int D6's coordinate dependence $<\Phi>\left(s_{1}, S_{2}\right)$.
- Give $M$ nilpotent vev ${ }^{\mathrm{In}} \quad\left[M_{1}, M^{\dagger}\right] \neq 0 \quad \& \quad \mathcal{N}=4 \rightarrow \mathcal{N}=2$ :
- Probing a T6-brane background.
[Cecotti, Cordova, Heckman, Vafa `IO]
- 2 fields among $M_{(\mathrm{j}, \mathrm{j})} \Leftrightarrow s_{1}, s_{2} \quad$ take the 2 highest spins.
- CB gets deformed: $V_{+} V_{-}=\operatorname{det}\left[\phi \mathbb{1}_{N}-M\left(s_{1}, s_{2}\right)\right]$


## su(3) principal orbit

- Consider $\quad M=\underbrace{\left(\begin{array}{lll}0 & m & 0 \\ 0 & 0 & m \\ 0 & 0 & 0\end{array}\right)}_{\langle M\rangle}+\underbrace{\left(\begin{array}{lll}0 & 0 & 0 \\ s_{2} & 0 & 0 \\ s_{1} & s_{2} & 0\end{array}\right)}_{\delta M}$

Then $W=\phi \sum_{i=1}^{3} Q_{i} \tilde{Q}^{i}+m\left(Q_{2} \tilde{Q}^{1}+Q_{3} \tilde{Q}^{2}\right)+s_{2}\left(Q_{1} \tilde{Q}^{2}+Q_{2} \tilde{Q}^{3}\right)+s_{1} Q_{1} \tilde{Q}^{3}$

## su(3) principal orbit

- Consider

$$
M=\underbrace{\left(\begin{array}{lll}
0 & m & 0 \\
0 & 0 & m \\
0 & 0 & 0
\end{array}\right)}_{\langle M\rangle}+\underbrace{\left(\begin{array}{lll}
0 & 0 & 0 \\
s_{2} & 0 & 0 \\
s_{1} & s_{2} & 0
\end{array}\right)}_{\delta M}
$$

Then $W=\phi \sum_{i=1}^{3} Q_{i} \tilde{Q}^{i}+m\left(Q_{2} \tilde{Q}^{1}+Q_{3} \tilde{Q}^{2}\right)+s_{2}\left(Q_{1} \tilde{Q}^{2}+Q_{2} \tilde{Q}^{3}\right)+s_{1} Q_{1} \tilde{Q}^{3}$

- Integrating out


## su(3) principal orbit

- Consider

$$
M=\underbrace{\left(\begin{array}{lll}
0 & m & 0 \\
0 & 0 & m \\
0 & 0 & 0
\end{array}\right)}_{\langle M\rangle}+\underbrace{\left(\begin{array}{ccc}
0 & 0 & 0 \\
s_{2} & 0 & 0 \\
s_{1} & s_{2} & 0
\end{array}\right)}_{\delta M}
$$

Then $W=\phi \sum_{i=1}^{3} Q_{i} \tilde{Q}^{i}+m\left(Q_{2} \tilde{Q}^{1}+Q_{3} \tilde{Q}^{2}\right)+s_{2}\left(Q_{1} \tilde{Q}^{2}+Q_{2} \tilde{Q}^{3}\right)+s_{1} Q_{1} \tilde{Q}^{3}$

- Integrating out $W^{\text {eff }}=\left(\frac{\phi^{3}}{m^{2}}-\frac{2 \phi s_{2}}{m}+s_{1}\right) Q_{1} \tilde{Q}^{3}$


## su(3) principal orbit

- Consider

$$
M=\underbrace{\left(\begin{array}{lll}
0 & m & 0 \\
0 & 0 & m \\
0 & 0 & m
\end{array}\right)}_{\langle M\rangle}+\underbrace{\left(\begin{array}{ccc}
0 & 0 & 0 \\
s_{2} & 0 & 0 \\
s_{1} & s_{2} & 0
\end{array}\right)}_{\delta M}
$$

Then $W=\phi \sum_{i=1}^{3} Q_{i} \tilde{Q}^{i}+m\left(Q_{2} \tilde{Q}^{1}+Q_{3} \tilde{Q}^{2}\right)+s_{2}\left(Q_{1} \tilde{Q}^{2}+Q_{2} \tilde{Q}^{3}\right)+s_{1} Q_{1} \tilde{Q}^{3}$


## su(3) principal orbit

- Consider $\quad M=\underbrace{\left(\begin{array}{ccc}0 & m & 0 \\ 0 & 0 & m \\ 0 & 0 & 0\end{array}\right)}_{\langle M\rangle}+\underbrace{\left(\begin{array}{lll}0 & 0 & 0 \\ s_{2} & 0 & 0 \\ s_{1} & s_{2} & 0\end{array}\right)}_{\delta M}$
- Then $W=\phi \sum_{i=1}^{3} Q_{i} \tilde{Q}^{i}+m\left(Q_{2} \tilde{Q}^{1}+Q_{3} \tilde{Q}^{2}\right)+s_{2}\left(Q_{1} \tilde{Q}^{2}+Q_{2} \tilde{Q}^{3}\right)+s_{1} Q_{1} \tilde{Q}^{3}$
- Integrating out $W^{\text {eff }}=\left(\frac{\phi^{2}}{m^{2}} \frac{2 \phi s_{2}}{w^{2}}+s_{1}\right) Q_{1} \tilde{Q}^{3} \leadsto \mathbb{R}^{\square} \sim W^{\mathrm{IR}}=s_{1} Q_{1} \tilde{Q}^{3}$

IR irrelevant!

- $\phi$ decouples \& $s_{1}$ is the new CB operator back to $\mathcal{N}=4$ !


## su(3) principal orbit

- Consider $\quad M=\underbrace{\left(\begin{array}{ccc}0 & m & 0 \\ 0 & 0 & m \\ 0 & 0 & 0\end{array}\right)}_{\langle M\rangle}+\underbrace{\left(\begin{array}{lll}0 & 0 & 0 \\ s_{2} & 0 & 0 \\ s_{1} & s_{2} & 0\end{array}\right)}_{\delta M}$
- Then $W=\phi \sum_{i=1}^{3} Q_{i} \tilde{Q}^{i}+m\left(Q_{2} \tilde{Q}^{1}+Q_{3} \tilde{Q}^{2}\right)+s_{2}\left(Q_{1} \tilde{Q}^{2}+Q_{2} \tilde{Q}^{3}\right)+s_{1} Q_{1} \tilde{Q}^{3}$
- Integrating out $W^{\text {eff }}=\left(\frac{\phi^{2}}{m^{2}} \frac{2 \phi s_{2}}{m_{2}}+s_{1}\right) Q_{1} \tilde{Q}^{3} \leadsto \mathbb{R}^{\mathrm{R}} \leadsto W^{\mathrm{IR}}=s_{1} Q_{1} \tilde{Q}^{3}$

IR irrelevant!

- $\phi$ decouples \& $s_{1}$ is the new CB operator back to $\mathcal{N}=4$ !
- This is $U(I) w /$ I flavor $\Longleftrightarrow$ theory of I free hyper.


## su(3) principal orbit

- Geometrically: M2 probing the fourfold $\left(C Y_{4}\right)$.

$$
V_{+} V_{-}=\phi^{3}-2 m \phi s_{2}+m^{2} s_{1}
$$

field-dependent deformation of $\mathbb{C}^{2} / \mathbb{Z}_{3}$

## su(3) principal orbit

- Geometrically: M2 probing the fourfold $\left(\mathrm{CY}_{4}\right)$.

$$
V_{+} V_{-}=\phi^{3}-2 m \phi s_{2}+m^{2} s_{1} \quad \begin{aligned}
& \text { field-dependent } \\
& \text { deformation of } \mathbb{C}^{2} / \mathbb{Z}_{3}
\end{aligned}
$$

- In the IR, M2 only "sees" tiny neighbor of the origin.
$\Rightarrow \mathrm{CY}_{4} \leadsto \leadsto\left\{V_{+} V_{-}=s_{1}\right\} \times \mathbb{C}^{2} \Rightarrow \mathscr{N}=2 \rightarrow \mathcal{N}=4$.


## su(3) principal orbit

- Geometrically: M2 probing the fourfold $\left(\mathrm{CY}_{4}\right)$.

$$
V_{+} V_{-}=\phi^{3}-2 m \phi s_{2}+m^{2} s_{1} \quad \begin{aligned}
& \text { field-dependent } \\
& \text { deformation of } \mathbb{C}^{2} / \mathbb{Z}_{3}
\end{aligned}
$$

- In the IR, M2 only "sees" tiny neighbor of the origin.
$\Rightarrow \mathrm{CY}_{4} \leftrightarrow\left\{V_{+} V_{-}=s_{1}\right\} \times \mathbb{C}^{2} \Rightarrow \mathscr{N}=2 \rightarrow \mathcal{N}=4$.
- Physically: D2 hasn't enough energy to "feel" D6's curvature.


## su(3) principal orbit

- Geometrically: M2 probing the fourfold $\left(\mathrm{CY}_{4}\right)$.

$$
V_{+} V_{-}=\phi^{3}-2 m \phi s_{2}+m^{2} s_{1} \quad \begin{aligned}
& \text { field-dependent } \\
& \text { deformation of } \mathbb{C}^{2} / \mathbb{Z}_{3}
\end{aligned}
$$

- In the IR, M2 only "sees" tiny neighbor of the origin.
$\Rightarrow \mathrm{CY}_{4} \leadsto\left\{V_{+} V_{-}=s_{1}\right\} \times \mathbb{C}^{2} \Rightarrow \mathscr{N}=2 \rightarrow \mathcal{N}=4$.
- Physically: D2 hasn't enough energy to "feel" D6's curvature.



## su(3) minimal orbit

- Consider instead $M=\underbrace{\left(\begin{array}{lll}0 & m & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0\end{array}\right)}_{\langle M\rangle}+\underbrace{\left(\begin{array}{lll}0 & 0 & 0 \\ s_{1} & 0 & 0 \\ s_{2} & 0 & 0\end{array}\right)}_{\delta M}$
- $\mathbb{C}^{2} / \mathbb{Z}_{3}$ gets deformed to $V_{+} V_{-}=\phi^{3}-2 m \phi s_{1}$
- In the $\operatorname{IR}$, this is equivalent to a conifold! $\quad V_{+} V_{-}=\phi s_{1}$
- No factorized twofold no no enhancement.
- D2 sees two intersecting D6-branes:



## Rank-1 4d SCFT's \& F-theory

- 3d case boring: only maximal orbit enhances \& IR theory is free.
- Single D3 probing a stack of 7-branes (possibly non-perturb.):
$\Rightarrow$ 4d $\mathscr{N}=2$ rank-I theory (possibly non-Lagrangian).
$\Rightarrow$ IR physics given by D3 prob. loc. elliptic K3 w/ Kodaira sing. [Banks, Douglas, Seiberg '96]


## Rank-I 4d SCFT's \& F-theory

- 3d case boring: only maximal orbit enhances \& IR theory is free.
- Single D3 probing a stack of 7-branes (possibly non-perturb.): $\Rightarrow 4 \mathrm{~d} \mathscr{N}=2$ rank-I theory (possibly non-Lagrangian).
$\Rightarrow$ IR physics given by D3 prob. loc. elliptic K3 w/ Kodaira sing. [Banks, Douglas, Seiberg `96]


K3 is NOT a moduli space

singularity $\leftrightarrow$ flavor structure

## Deformations

- Add the $\mathscr{N}=2 \rightarrow \mathscr{N}=\mathrm{I}$ coupling $\delta W=\operatorname{Tr}(\mu M)$ :
- $M=\langle\boldsymbol{\Phi}\rangle$ is vev of 7-brane worldvolume scalar.
- $M=\rho\left(\sigma_{+}\right)+\sum_{j=s_{1}, s_{2}} \delta M_{(j,-j)}$
- $\left[M, M^{\dagger}\right] \neq 0$ Int D3 probes a T7-brane background.
- $\mathrm{D}^{\mathrm{UV}}\left(s_{i}\right)=\operatorname{spin}\left(s_{i}\right)+1$.
- $\boldsymbol{\delta} W \Rightarrow$ field-dependent deformations of K3:
- Casimir invariants of $M \leadsto \leadsto$ Versal deform'ns of singularity.


## Enhancement criterion

- For any RG flow \& any Energy scale, F-theory geometry is:
$y^{2}=x^{3}+f\left(\phi, s_{1}, s_{2}\right) x+g\left(\phi, s_{1}, s_{2}\right) \quad$ Elliptic $\mathrm{CY}_{4}$ in Weierstrass form
- $f \& g$ depend on UV theory \& choice of nilpotent orbit.
- Original K3 retrieved by $s_{1} \equiv s_{2} \equiv 0$.
- If SUSY enhances in IR, then:

$$
\mathrm{K} 3^{\mathrm{UV}} \xrightarrow{\mathrm{Def}} \mathrm{CY}_{4} \xrightarrow{\mathrm{IR}} \quad \mathrm{~K}_{3}{ }^{\mathrm{IR}} \times \mathbb{C}^{2}
$$

- K3UV $\equiv \mathrm{K} 3 \mathrm{R}^{\mathrm{R}} \Longleftrightarrow$ trivial orbit.
- Educated guess: $\mathrm{K} 3^{\mathrm{IR}}=\left.\mathrm{CY}_{4}\right|_{\phi \equiv s_{2} \equiv 0}$ (1N+ $s_{1}$ new CB in IR.


## Enhancement criterion

- Conformal dim. of operators $\Longleftrightarrow \mathbb{C}^{*}$-action on coordinates:
- Promote affine variables to projective ones.
- Relative scalings of operators are RG-flow invariants:
- Impose homogeneity of $\mathrm{CY}_{4}$ polynomial.
- Ansatz for $\operatorname{IR}$ geom. consistent $\Longleftrightarrow \mathrm{D}^{\mathbb{R}}(\phi) \leq 1 \quad \& \quad \mathrm{D}^{\mathbb{R}}\left(s_{2}\right) \leq 1$.


## Enhancement criterion

- Conformal dim. of operators $\Longleftrightarrow \mathbb{C}^{*}$-action on coordinates:
- Promote affine variables to projective ones.
- Relative scalings of operators are RG-flow invariants:
- Impose homogeneity of $\mathrm{CY}_{4}$ polynomial.
- Ansatz for $\mathbb{R}$ geom. consistent $\Longleftrightarrow \mathrm{D}^{\mathrm{R}}(\phi) \leq 1 \quad \& \quad \mathrm{D}^{\mathrm{R}}\left(s_{2}\right) \leq 1$.
- Assign $\mathbb{C}^{*}$-actions: Fiber K3 adiabatic. \& force total space be CY.
$\Rightarrow \mathrm{D}^{\mathbb{R}}(x)-\mathrm{D}^{\mathbb{R}}(y)+\mathrm{D}^{\mathbb{R}}\left(s_{1}\right)=1$.
- Equivalent to $\mathrm{D}\left(\lambda^{s W}\right)=1$


## Enhancement criterion

- Conformal dim. of operators $\Longleftrightarrow \mathbb{C}^{*}$-action on coordinates:
- Promote affine variables to projective ones.
- Relative scalings of operators are RG-flow invariants:
- Impose homogeneity of $\mathrm{CY}_{4}$ polynomial.
- Ansatz for $\mathbb{R}$ geom. consistent $\Longleftrightarrow \mathrm{D}^{\mathrm{R}}(\phi) \leq 1 \quad \& \quad \mathrm{D}^{\mathrm{R}}\left(s_{2}\right) \leq 1$.
- Assign $\mathbb{C}^{*}$-actions: Fiber K3 adiabatic. \& force total space be CY.
$\Rightarrow \mathrm{D}^{\operatorname{R}}(x)-\mathrm{D}^{\mathbb{R}}(y)+\mathrm{D}^{\operatorname{R}}\left(s_{1}\right)=1$.
- Equivalent to $\mathrm{D}\left(\lambda^{\mathrm{sW}}\right)=1$

$$
\partial \lambda^{\mathrm{SW}} / \partial \mathcal{O}_{\mathrm{CB}}=\Omega_{\mathrm{SW}}^{1,0}=\mathrm{d} x / y
$$

## Example: $\mathrm{D}_{4} \leadsto \mathscr{T}_{2}$

- $\mathrm{D}^{\mathrm{UV}}\left(s_{1}\right) \leq \mathrm{D}^{\mathrm{UV}}(\phi)$ nime no enhancement!
- Highest spin multiply populated
- Consider the Lagrangian th. $\mathrm{SU}(2) \mathrm{w} / 4$ flavors $\mathrm{G}_{\mathrm{F}}=\mathrm{SO}(8)$.
- From D3 probing K3UV: $y^{2}=x^{3}+\tau \phi^{2} x+\phi^{3}$.


## Example: $\mathrm{D}_{4} \leadsto \mathscr{T}_{2}$

- $\mathrm{D}^{\mathrm{UV}}\left(s_{1}\right) \leq \mathrm{D}^{\mathrm{UV}}(\phi)$ nime no enhancement!
- Highest spin multiply populated
- Consider the Lagrangian th. $\mathrm{SU}(2) \mathrm{w} / 4$ flavors $\mathrm{G}_{\mathrm{F}}=\mathrm{SO}(8)$.
- From D3 probing K3UV: $y^{2}=x^{3}+\tau \phi^{2} x+\phi^{3}$.
- Deform:


## Example: $\mathrm{D}_{4} \leadsto \mathscr{T}_{2}$

- $\mathrm{D}^{\mathrm{UV}}\left(s_{1}\right) \leq \mathrm{D}^{\mathrm{UV}}(\phi)$ nime no enhancement!
- Highest spin multiply populated mm no enhancement!
- Consider the Lagrangian th. $\operatorname{SU}(2) \mathrm{w} / 4$ flavors $\mathrm{m} \| \mathrm{GF}=\mathrm{SO}(8)$.
- From D3 probing K3UV: $y^{2}=x^{3}+\tau \phi^{2} x+\phi^{3}$.
- Deform: $\left.M=\frac{\left(\begin{array}{cc}A & 0 \\ 0 & -A^{T}\end{array}\right)}{\left(\begin{array}{cccc}0 & \sqrt{2} & 0 & 0 \\ s_{2} & 0 & \sqrt{2} & 0 \\ s_{1} & s_{2} & 0 & 0 \\ 0 & 0 & 0 & 0\end{array}\right)} \quad \begin{array}{l}P_{M}(z)=z^{8}-4 \sqrt{2} s_{2} z^{6}+8 s_{2}^{2} z^{4}-4 s_{1}^{2} z^{2}\end{array}\right\}$


## Example: $\mathrm{D}_{4} \leadsto \mathscr{T}_{2}$

- $\mathrm{D}^{\mathrm{UV}}\left(s_{1}\right) \leq \mathrm{D}^{\mathrm{UV}}(\phi)$ nime no enhancement!
- Highest spin multiply populated mm no enhancement!
- Consider the Lagrangian th. $\operatorname{SU}(2) \mathrm{w} / 4$ flavors $\mathrm{m} \| \mathrm{GF}=\mathrm{SO}(8)$.
- From D3 probing K3UV: $y^{2}=x^{3}+\tau \phi^{2} x+\phi^{3}$.
- Deform: $M=\frac{\left(\begin{array}{cc}A & 0 \\ 0 & -A^{T}\end{array}\right)}{\left(\begin{array}{cccc}0 & \sqrt{2} & 0 & 0 \\ s_{2} & 0 & \sqrt{2} & 0 \\ s_{1} & s_{2} & 0 & 0 \\ 0 & 0 & 0 & 0\end{array}\right)} \quad \begin{gathered}P_{M}(z)=z^{8}-4 \sqrt{2} s_{2} z^{6}+8 s_{2}^{2} z^{4}-4 s_{1}^{2} z^{2}\end{gathered}$

$$
y^{2}=x^{3}+\tau \phi^{2} x+2 s_{2} \phi x+s_{2}^{2} x+\phi^{3}+s_{1}^{2}
$$

## Example: $\mathrm{D}_{4} \leadsto \mathscr{T}_{2}$

- $\mathrm{D}^{\mathrm{UV}}\left(s_{1}\right) \leq \mathrm{D}^{\mathrm{UV}}(\phi)$ NIL no enhancement!
- Highest spin multiply populated mm no enhancement!
- Consider the Lagrangian th. $\operatorname{SU}(2) \mathrm{w} / 4$ flavors $\ln \mathrm{G}_{\mathrm{F}}=\mathrm{SO}(8)$.
- From D3 probing K3Uv: $y^{2}=x^{3}+\tau \phi^{2} x+\phi^{3}$.
- Deform: $M=\begin{aligned} & \left(\begin{array}{cc}A & 0 \\ 0 & -A^{T}\end{array}\right) \\ & \underline{\left[3^{2}, 1^{1}\right]} \longrightarrow\left(\begin{array}{cccc}0 & \sqrt{2} & 0 & 0 \\ s_{2} & 0 & \sqrt{2} & 0 \\ s_{1} & s_{2} & 0 & 0 \\ 0 & 0 & 0 & 0\end{array}\right) \\ & 2 V_{0} \oplus 7 V_{1} \oplus V_{2}\end{aligned} \quad \begin{gathered}P_{M}(z)=z^{8}-4 \sqrt{2} s_{2} z^{6}+8 s_{2}^{2} z^{4}-4 s_{1}^{2} z^{2} \\ C_{(2)}=-4 \sqrt{2} s_{2} ; \\ \quad C_{(6)}=8 s_{2}^{2} ; \quad C_{(6)}=-4 s_{1}^{2}\end{gathered}$

$$
y^{2}=x^{3}+\tau \phi^{2} x+2 s_{2} \phi x+s_{2}^{2} x+\phi^{3}+s_{1}^{2}
$$

- $\mathrm{D}^{\mathbb{R}}(y)=\mathrm{D}^{\mathbb{R}}\left(s_{1}\right)=3 / 2 ; \mathrm{D}^{\mathbb{R}}(x)=\mathrm{D}^{\mathbb{R}}\left(s_{2}\right)=\mathrm{D}^{\mathbb{R}}(\phi)=1$.


## Example: $\mathrm{D}_{4} \leadsto \mathscr{T}_{2}$

- $\mathrm{D}^{\mathrm{UV}}\left(s_{1}\right) \leq \mathrm{D}^{\mathrm{UV}}(\phi)$ nime no enhancement!
- Highest spin multiply populated mm no enhancement!
- Consider the Lagrangian th. $\operatorname{SU}(2) \mathrm{w} / 4$ flavors $\ln \mathrm{G}_{\mathrm{F}}=\mathrm{SO}(8)$.
- From D3 probing K3 ${ }^{U v}: y^{2}=x^{3}+\tau \phi^{2} x+\phi^{3}$.


$$
y^{2}=x^{3}+\epsilon^{2} x+2 s_{2} \phi x+s_{2}^{2} x+\phi^{3}+s_{1}{ }^{2} \Rightarrow \mathscr{T}(2
$$

- $\mathrm{D}^{\mathbb{R}}(y)=\mathrm{D}^{\mathbb{R}}\left(s_{1}\right)=3 / 2 ; \mathrm{D}^{\mathbb{R}}(x)=\mathrm{D}^{\mathbb{R}}\left(s_{2}\right)=\mathrm{D}^{\mathbb{R}}(\phi)=1$.


## Example: No enhancement

- Consider E8 Minahan-Nemeschanski.
- From D3 probing K3uv: $y^{2}=x^{3}+\phi^{5}$.
- Deform by $\mathrm{E}_{8}\left(a_{2}\right)$ orbit: $V_{1} \oplus V_{3} \oplus V_{5} \oplus V_{7} \oplus V_{8} \oplus V_{9} \oplus 2 V_{11} \oplus V_{13} \oplus V_{14} \oplus V_{17} \oplus V_{19}$
$\Rightarrow s_{1}$ : 20th-order Casimir ; $s_{2}$ : 18th-order Casimir.
$\Rightarrow y^{2}=x^{3}+\phi^{2} S_{2}+s_{1} x+\phi^{5}$.
- $\mathrm{D}^{\operatorname{Rr}}(y)=1 ; \mathrm{D}^{\operatorname{R}}(x)=2 / 3 ; \mathrm{D}^{\operatorname{R}}\left(s_{1}\right)=4 / 3$.
- $\mathrm{D}^{\operatorname{IR}}(\phi)=2 / 5 \quad$ InIt $\quad$ decouples!
- BUT $\mathrm{D}^{\text {IR }}\left(S_{2}\right)=6 / 5 \quad \mathrm{ln}$ " prevents enhancement.
- Incorrect conformal dimensions... a-maximization needed!


## Lessons \& Outlook

$\checkmark$ Geometry of IR SUSY enhancement involves local structure of some algebraic space around (singular) point.
$\checkmark$ Branch of moduli space for 3d case, but auxiliary space for 4d case (rank-I
$\checkmark$ Enhancement only arises when non-trivial factorization occurs nut holonomy reduction in M/F-theory.

- Higher-rank theories ? In class-S, factorization IIt hyperkähler structure of moduli space of sol'ns of Hitchin system. [in progress with Carta, Giacomelli, Hayashi]
- Geometry behind factorization still obscure... RG-flow changes singularity structure, reminiscent of simultaneous resolutions!

