

Inflation, Primordial Black Holes, Gravitational Waves



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Motivations

- Slow-roll inflation for the first 7 e-foldings ($61 > N > 54$)

Planck CMB + LSS + SNe + ...

- Binary black holes detected by aLIGO/VIRGO may be primordial black holes produced during inflation ($N \sim 40$) ?

PBHs may be dark matter and aLIGO has seen dark matter!

Then...

- While we will learn more details about the slow-roll inflation (r , n_s , running n_s ,...), there may be an opportunity to probe beyond the inflationary standard model
- Formation of PBHs, associated with gravitational waves in inflation

Primordial Black Holes

- Formed at high-density contrasts over a wide range of scales or masses in the radiation-dominated Universe
- There have been stringent astrophysical and cosmological constraints on M_{PBH}
- $10M_{\odot}$ PBHs could be the binary BHs observed by aLIGO gravity-wave detectors

Bird et al. 16., Clesse et al. 16, Sasaki et al. 16

- PBHs behave like cold dark matter

García-Bellido, Linde, Wands 96

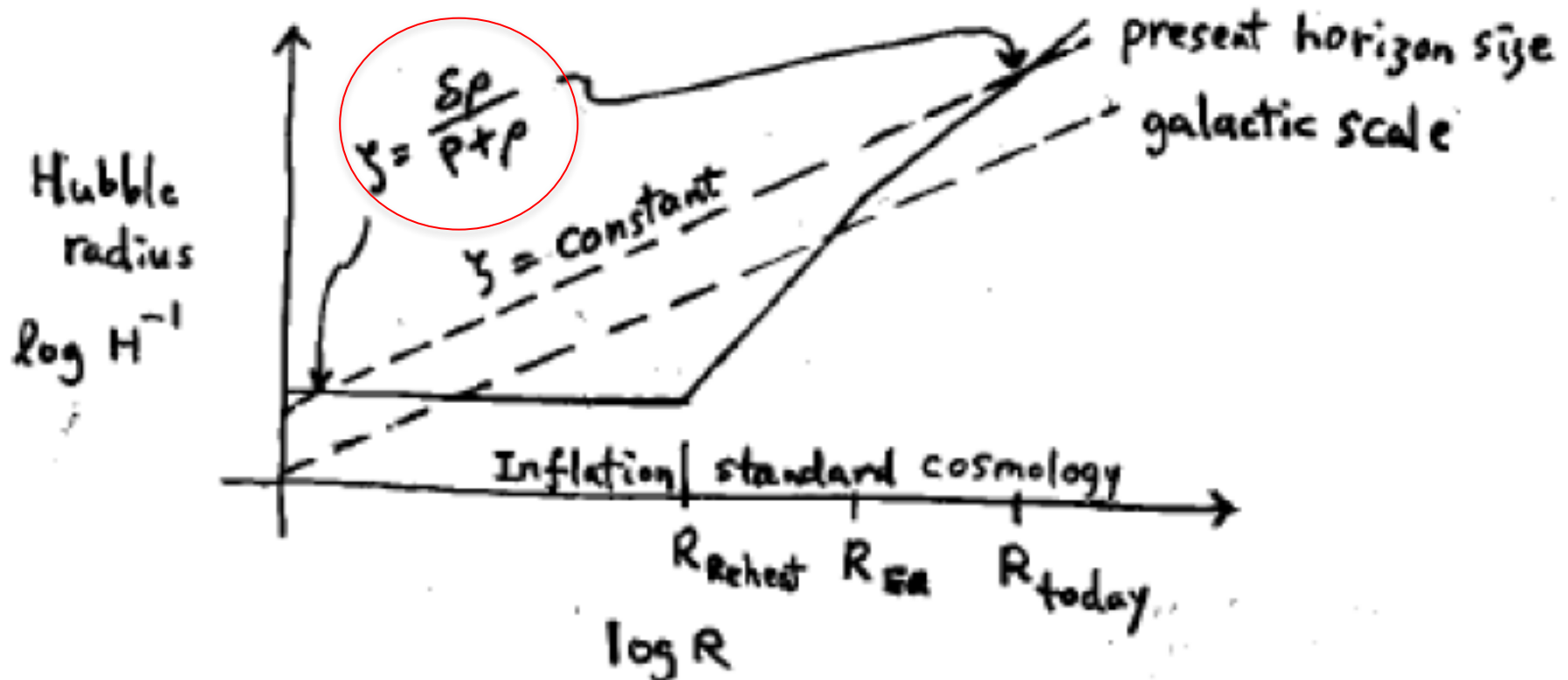
- They, although being of baryonic origin, do not participate in big-bang nucleosynthesis

PBH Formation

- spontaneously formed in phase transitions, e.g., 1st order electroweak, quark-hadron phase transition $M_{\text{BH}} \sim M_{\odot}$
- arisen from the collapse of horizon-sized large matter inhomogeneity seeded by quantum fluctuations during inflation that re-enter the horizon in the subsequent radiation-dominated Universe, with M_{BH} spans a wide range of mass

Production of PBHs in inflation

Evolution of cosmological density perturbation



$$M_{\text{BH}} = \frac{4\pi M_p^2}{H} e^{2N} = 2.74 \times 10^{-38} e^{2N} \left(\frac{M_p}{H} \right) M_{\odot}$$

N is e-foldings
before the end
of inflation

PBH Production in Inflation

- Single-field slow-roll inflation models, matter density perturbation ($\delta\rho/\rho \sim 10^{-5}$) too small
- Modified inflation potential to achieve blue-tilted matter power spectra or running spectral indices, leading to large density perturbation at the end of inflation, but mostly $M_{\text{PBH}} \ll M_{\odot}$

García-Bellido, Linde, Wands 96

- To boost M_{PBH} , hybrid inflation, double inflation, curvaton models by inflating small-scale density perturbation to the size of a stellar-mass to supermassive PBH.

Kawasaki, Kohri, Yokoyama, Yanagida....

Interacting inflaton

Quantum environment Wu et al. JCAP 07

$$\mathcal{L} = \frac{1}{2}g^{\mu\nu}\partial_\mu\phi\partial_\nu\phi + \frac{1}{2}g^{\mu\nu}\partial_\mu\sigma\partial_\nu\sigma - V(\phi) - \frac{m_\sigma^2}{2}\sigma^2 - \frac{g^2}{2}\phi^2\sigma^2$$

Trapped Inflation Kofman, Linde et al. 04, Green et al.09

Inflaton coupled to
instantaneously
massless fields

$$\mathcal{L} = \frac{1}{2}\partial_\mu\phi\partial^\mu\phi - V(\phi) + \frac{1}{2}\sum_i(\partial_\mu\chi_i\partial^\mu\chi_i - g^2(\phi - \phi_i)^2\chi_i^2) + \dots,$$

Natural Inflation with axion-like couplings Sorbo, Kaloper, Peloso,...

$$\mathcal{S} = \int d^4x \sqrt{-g} \left[\frac{M_p^2}{2} R - \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi - V(\varphi) - \frac{1}{4} F^{\mu\nu} F_{\mu\nu} - \frac{\alpha}{4f} \varphi \tilde{F}^{\mu\nu} F_{\mu\nu} \right]$$

Trapped axion Inflation

We consider a version of the trapped inflation driven by a pseudoscalar φ that couples to a $U(1)$ gauge field A_μ :

$$\mathcal{S} = \int d^4x \sqrt{-g} \left[\frac{M_p^2}{2} R - \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi - V(\varphi) - \frac{1}{4} F^{\mu\nu} F_{\mu\nu} - \frac{\alpha}{4f} \varphi \tilde{F}^{\mu\nu} F_{\mu\nu} \right], \quad (3)$$

$$\varphi = \phi(\eta) + \delta\varphi(\eta, \vec{x})$$

Under the temporal gauge, $A_\mu = (0, \vec{A})$, we decompose $\vec{A}(\eta, \vec{x})$ into its right and left circularly polarized Fourier modes, $A_\pm(\eta, \vec{k})$, whose equation of motion is then given by

$$k/(aH) < 2|\xi|$$

Spinoidal
instability

$$\left[\frac{d^2}{d\eta^2} + k^2 \mp 2aHk\xi \right] A_\pm(\eta, k) = 0, \quad \xi \equiv \frac{\alpha}{2fH} \frac{d\phi}{dt}. \quad (5)$$

Background

$$\frac{d^2\phi}{dt^2} + 3H \frac{d\phi}{dt} + \frac{dV}{d\phi} = \frac{\alpha}{f} \langle \vec{E} \cdot \vec{B} \rangle,$$

$$3H^2 = \frac{1}{M_p^2} \left[\frac{1}{2} \left(\frac{d\phi}{dt} \right)^2 + V(\phi) + \frac{1}{2} \langle \vec{E}^2 + \vec{B}^2 \rangle \right]$$

$$\langle \vec{E} \cdot \vec{B} \rangle \simeq -2.4 \cdot 10^{-4} \frac{H^4}{\xi^4} e^{2\pi\xi},$$

$$\left\langle \frac{\vec{E}^2 + \vec{B}^2}{2} \right\rangle \simeq 1.4 \cdot 10^{-4} \frac{H^4}{\xi^3} e^{2\pi\xi}.$$

$$\frac{1}{2} \langle \vec{E}^2 + \vec{B}^2 \rangle = \int \frac{dk k^2}{4\pi^2 a^4} \sum_{\lambda=\pm} \left(\left| \frac{dA_\lambda}{d\eta} \right|^2 + k^2 |A_\lambda|^2 \right),$$

$$\langle \vec{E} \cdot \vec{B} \rangle = - \int \frac{dk k^3}{4\pi^2 a^4} \frac{d}{d\eta} (|A_+|^2 - |A_-|^2).$$

Perturbation

$$\left[\frac{\partial^2}{\partial t^2} + 3\beta H \frac{\partial}{\partial t} - \frac{\vec{\nabla}^2}{a^2} + \frac{d^2V}{d\phi^2} \right] \delta\varphi(t, \vec{x}) = \frac{\alpha}{f} \left(\vec{E} \cdot \vec{B} - \langle \vec{E} \cdot \vec{B} \rangle \right)$$

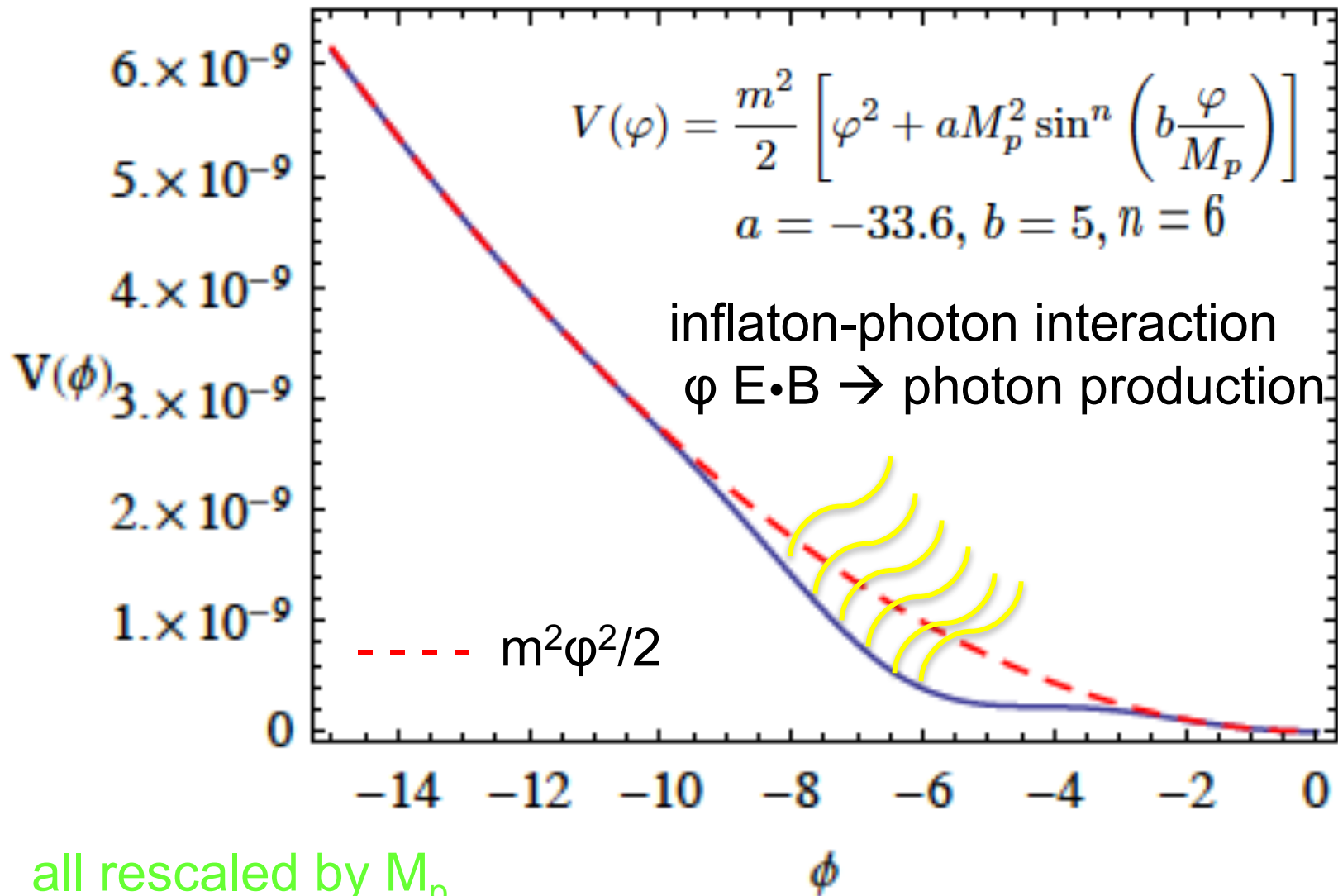
$$\delta\varphi = \frac{\alpha}{3\beta f H^2} \left(\vec{E} \cdot \vec{B} - \langle \vec{E} \cdot \vec{B} \rangle \right)$$

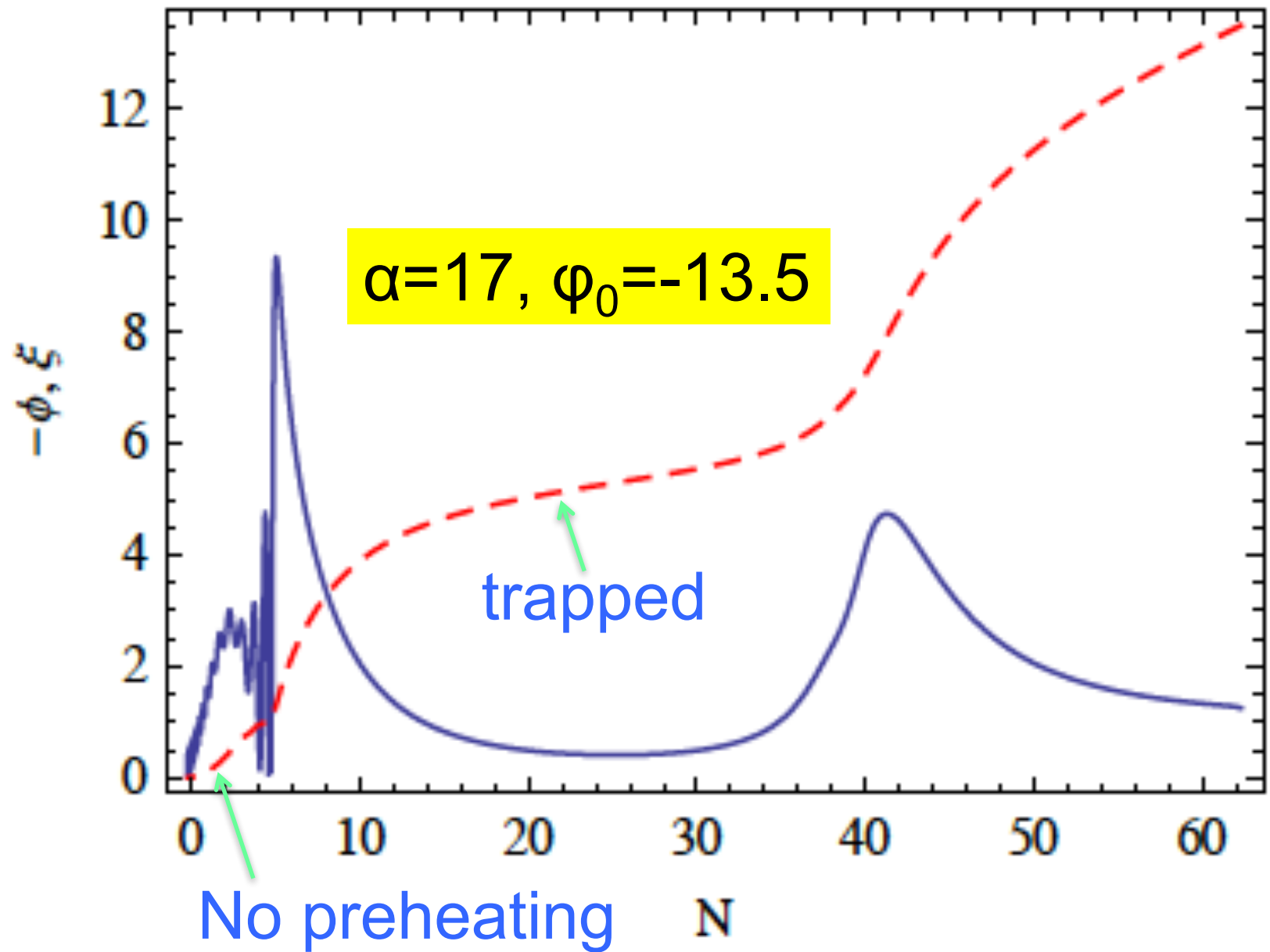
$$\beta \equiv 1 - 2\pi\xi \frac{\alpha}{f} \frac{\langle \vec{E} \cdot \vec{B} \rangle}{3H(d\phi/dt)}$$

$$\Delta_\zeta^2(k) = \langle \zeta(x)^2 \rangle = \frac{H^2 \langle \delta\varphi^2 \rangle}{(d\phi/dt)^2} = \left[\frac{\alpha \langle \vec{E} \cdot \vec{B} \rangle}{3\beta f H (d\phi/dt)} \right]^2$$

e.g. Trapped axion inflation with a steep potential

Cheng, Lee, Ng 16

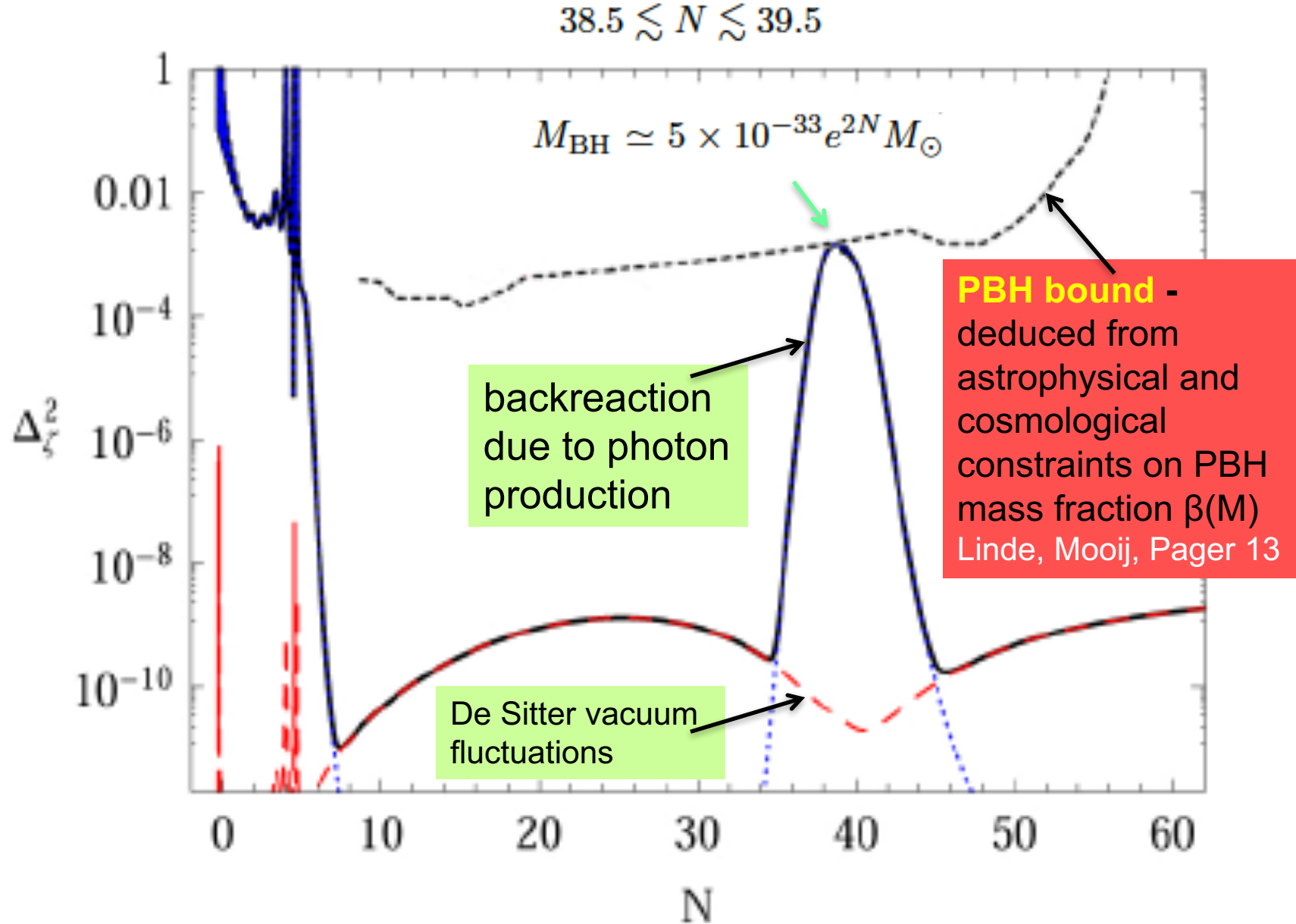




Note that inflation starts at $N=61$

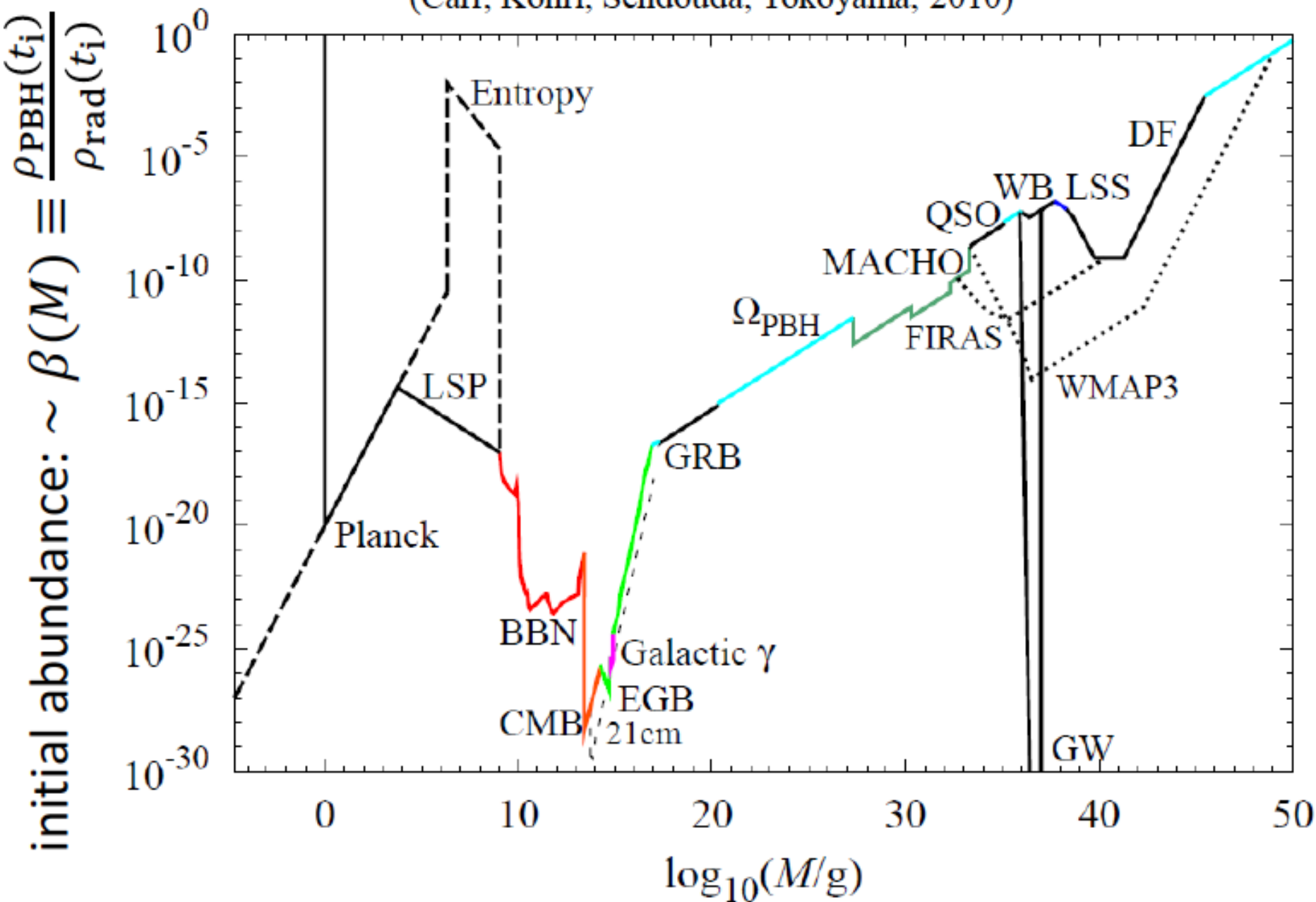
Production of $10\text{-}100M_{\odot}$ PBHs

Density perturbation $\Delta_{\zeta}^2 = (\delta\rho/\rho)^2$



Observational constraint on PBH abundance

(Carr, Kohri, Sendouda, Yokoyama, 2010)



PBH Associated Gravitational Waves in Inflation

$$\left[\frac{\partial^2}{\partial \eta^2} + \frac{2}{a} \frac{da}{d\eta} \frac{\partial}{\partial \eta} - \vec{\nabla}^2 \right] h_{ij} = 0 \quad \text{Free gravitational wave equation}$$

De Sitter vacuum fluctuations during inflation lead to almost scale-invariant primordial gravitational waves

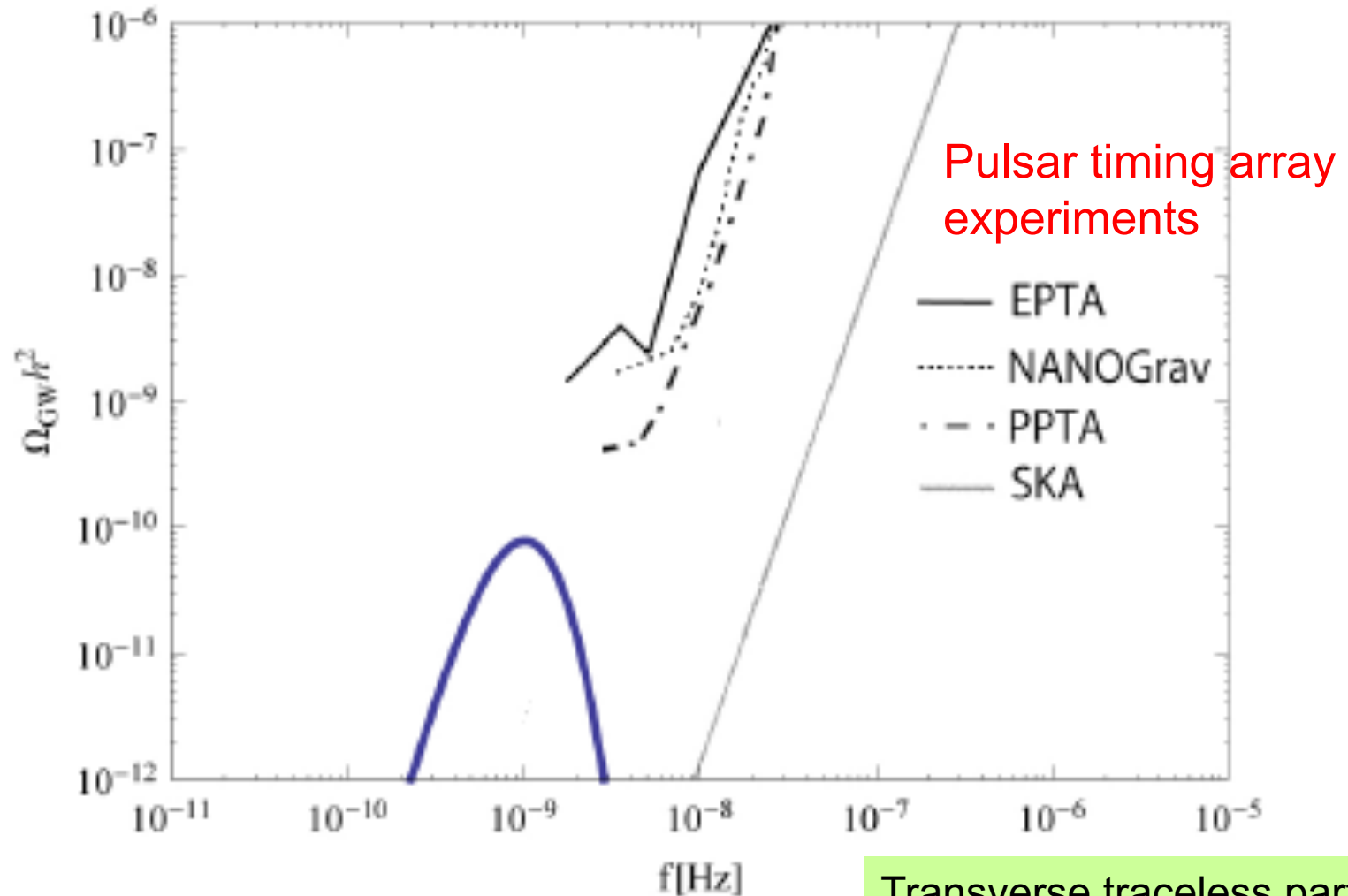
$$\Delta_\zeta^2 = \langle \zeta \zeta \rangle = (\delta \rho / \rho)^2 \sim 2 \times 10^{-9}$$

$$\left[\frac{\partial^2}{\partial \eta^2} + \frac{2}{a} \frac{da}{d\eta} \frac{\partial}{\partial \eta} - \vec{\nabla}^2 \right] h_{ij} = \text{GT}_{\mu\nu} \quad \text{Sources due to transverse traceless part of second-order density perturbation } \zeta\zeta$$

$$\Delta_\zeta^2 = \langle \zeta \zeta \rangle = (\delta \rho / \rho)^2 \sim 10^{-3}$$

Associated Gravitational Waves in Trapped Inflation

Spectral energy density

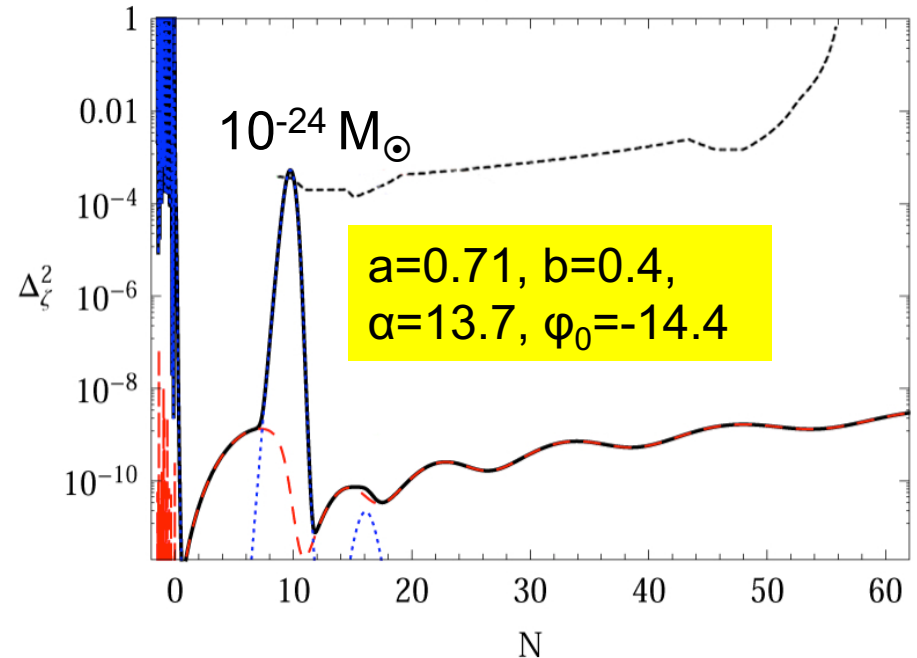
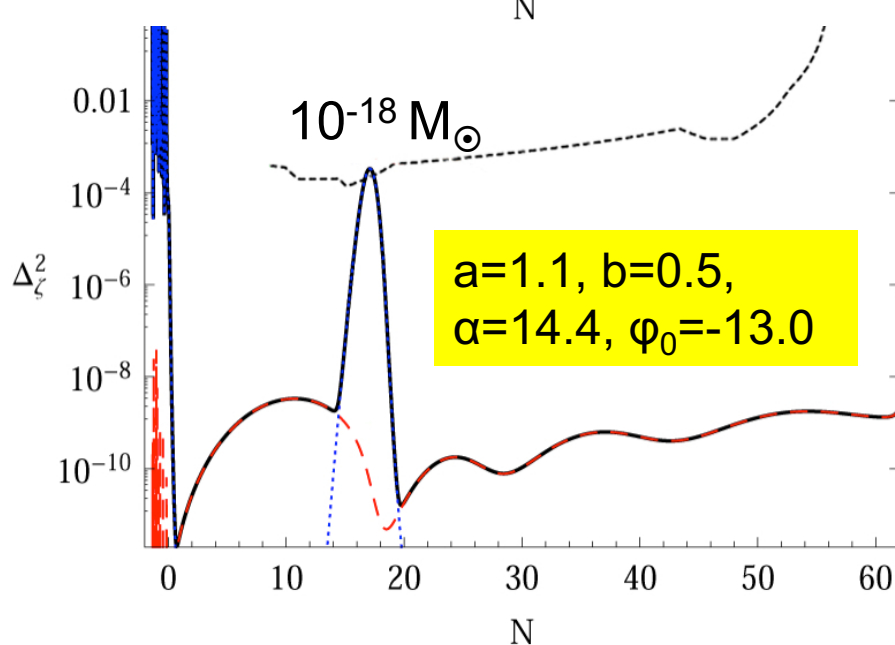
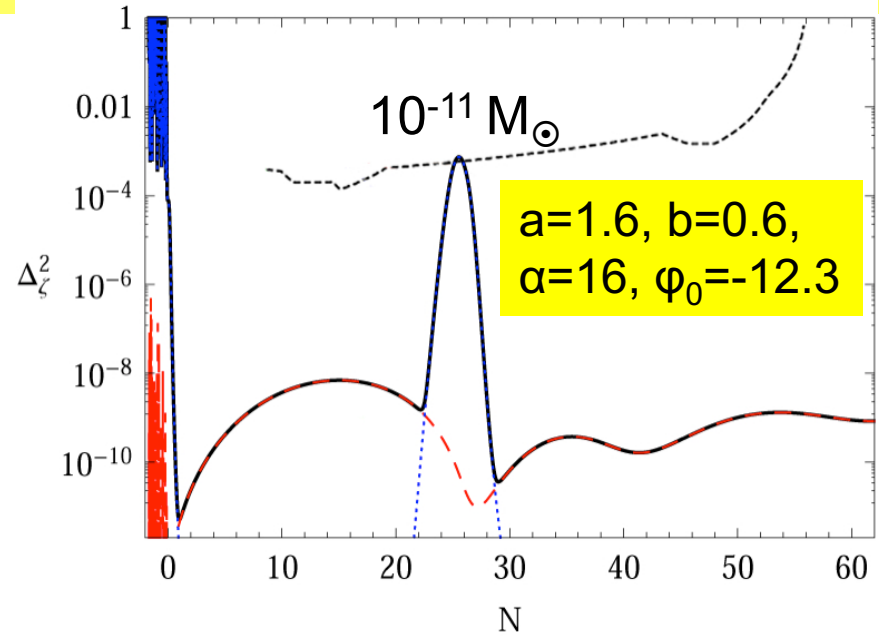
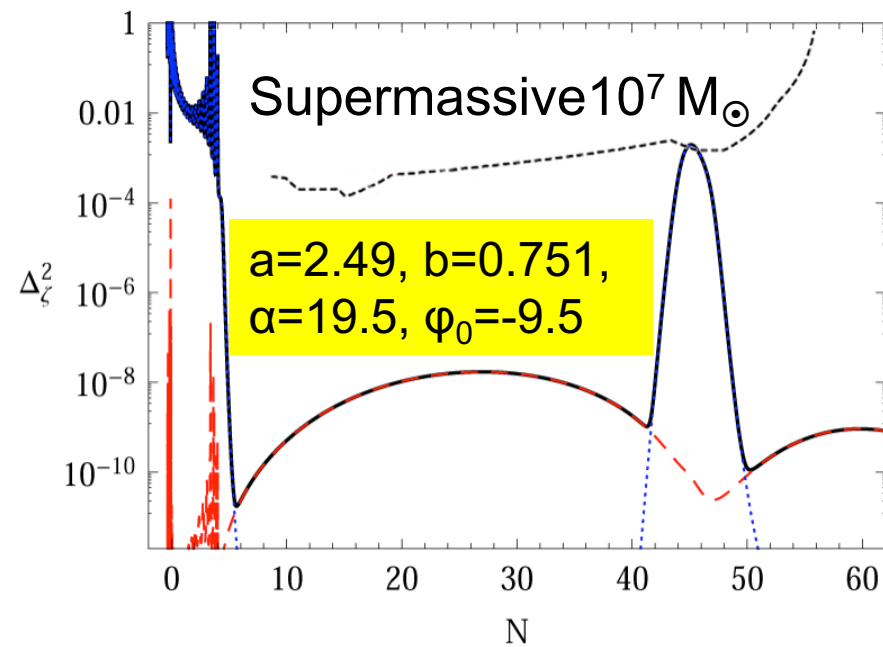


$$\left[\frac{\partial^2}{\partial \eta^2} + \frac{2}{a} \frac{da}{d\eta} \frac{\partial}{\partial \eta} - \vec{\nabla}^2 \right] h_{ij} = \frac{2a^2}{M_p^2} (-E_i E_j - B_i B_j)^{TT}$$

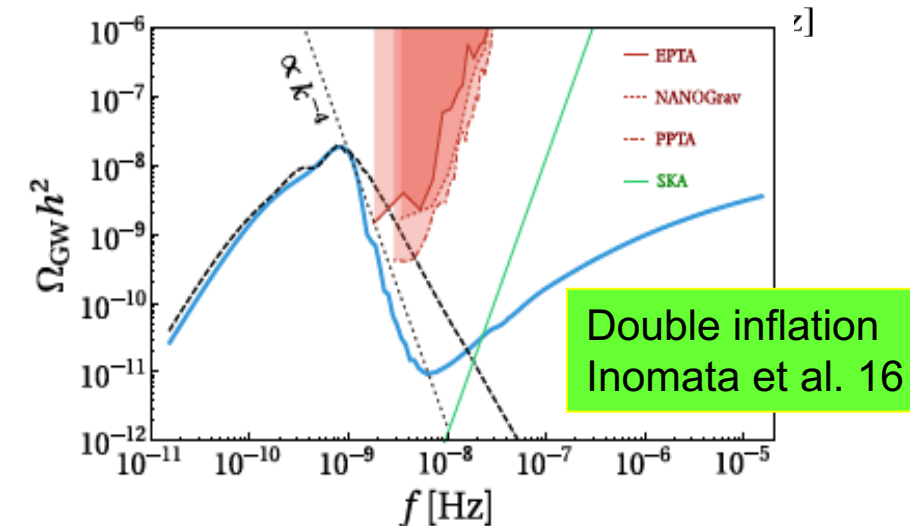
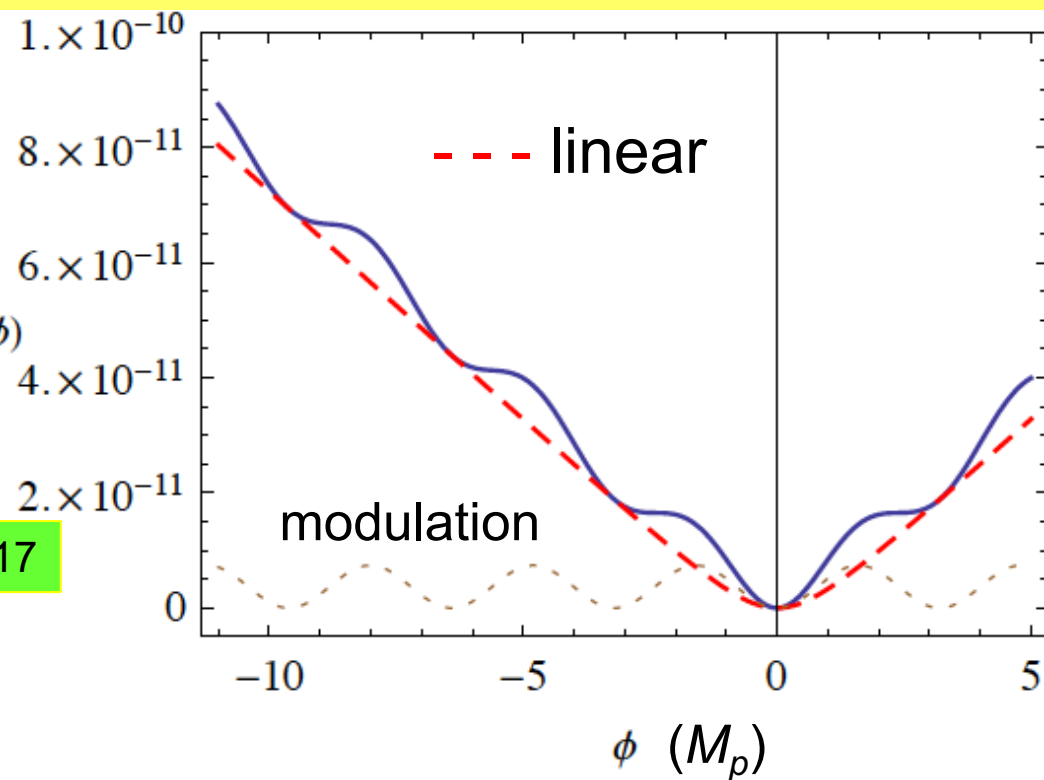
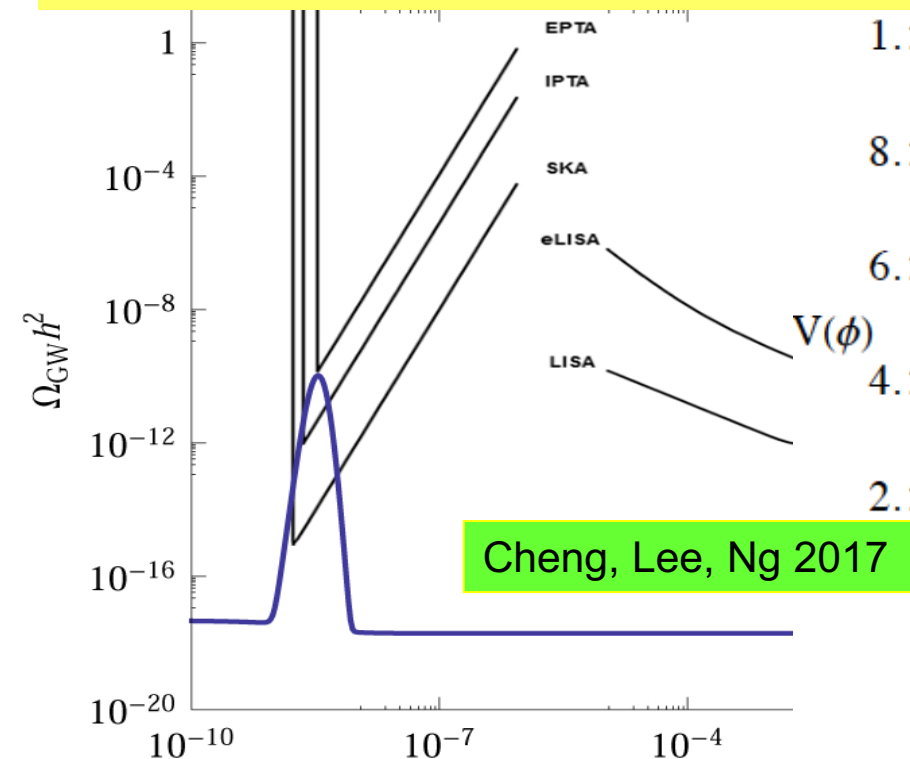
Transverse traceless part
of photon energy-momentum

$$\langle \mathbf{E} \cdot \mathbf{E} \rangle \sim \langle \mathbf{E} \cdot \mathbf{B} \rangle \sim \zeta$$

Producing PBHs with a wide range of masses



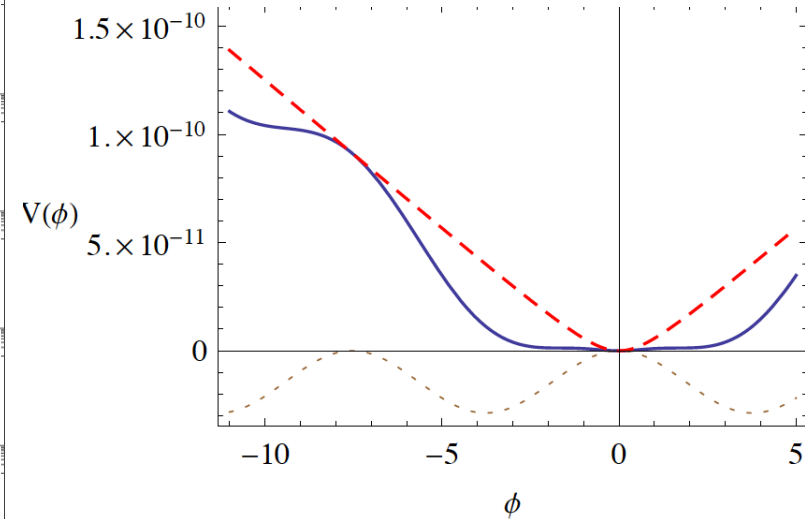
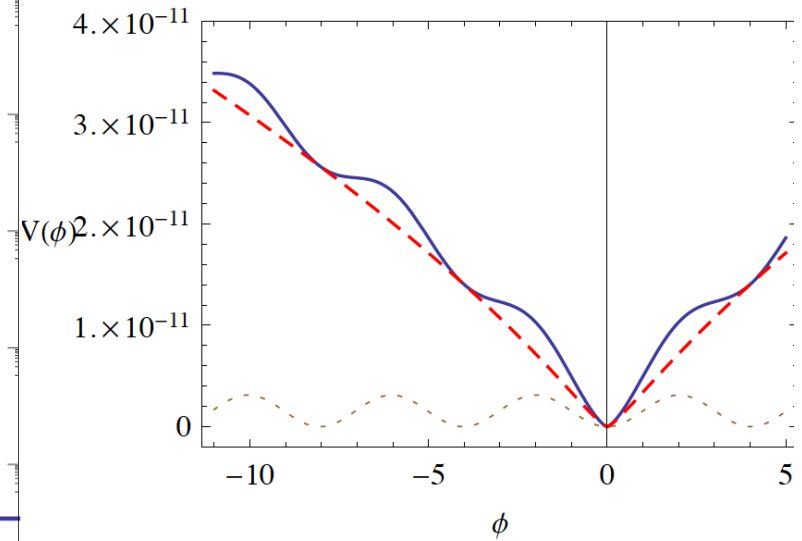
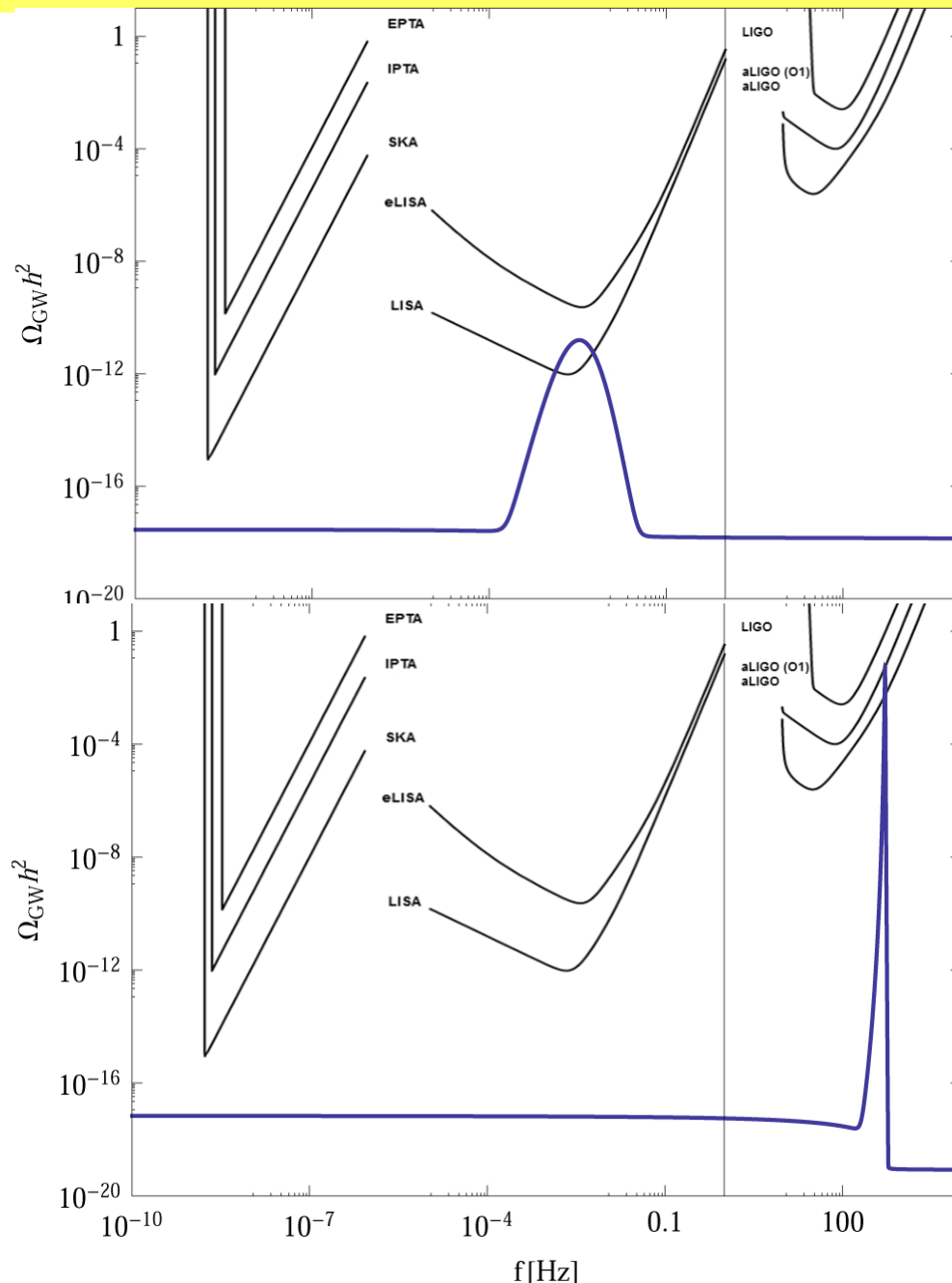
GWs associated with $10M_\odot$ PBHs in modulated axion inflation



$$V(\varphi) = M_p \mu^3 \sqrt{1 + \left(\frac{\varphi}{M_p}\right)^2} - M_p \mu^3 + \Lambda^4 \left[1 + \cos \left(\frac{\varphi}{f} + \pi \right) \right]$$

M_p reduced Planck mass
 f axion decay constant

GWs associated with smaller PBHs in modulated axion inflation



Higgs Inflation

Bezrukov+Shaposhnikov

SM Higgs potential

Jordan
frame

$$S_J = \int d^4x \sqrt{-g} \left\{ -\frac{M^2 + \xi h^2}{2} R + \frac{\partial_\mu h \partial^\mu h}{2} - \frac{\lambda}{4} (h^2 - v^2)^2 \right\}$$

Conformal transformation

$$\hat{g}_{\mu\nu} = \Omega^2 g_{\mu\nu}, \quad \Omega^2 = 1 + \frac{\xi h^2}{M_P^2}$$

Einstein
frame

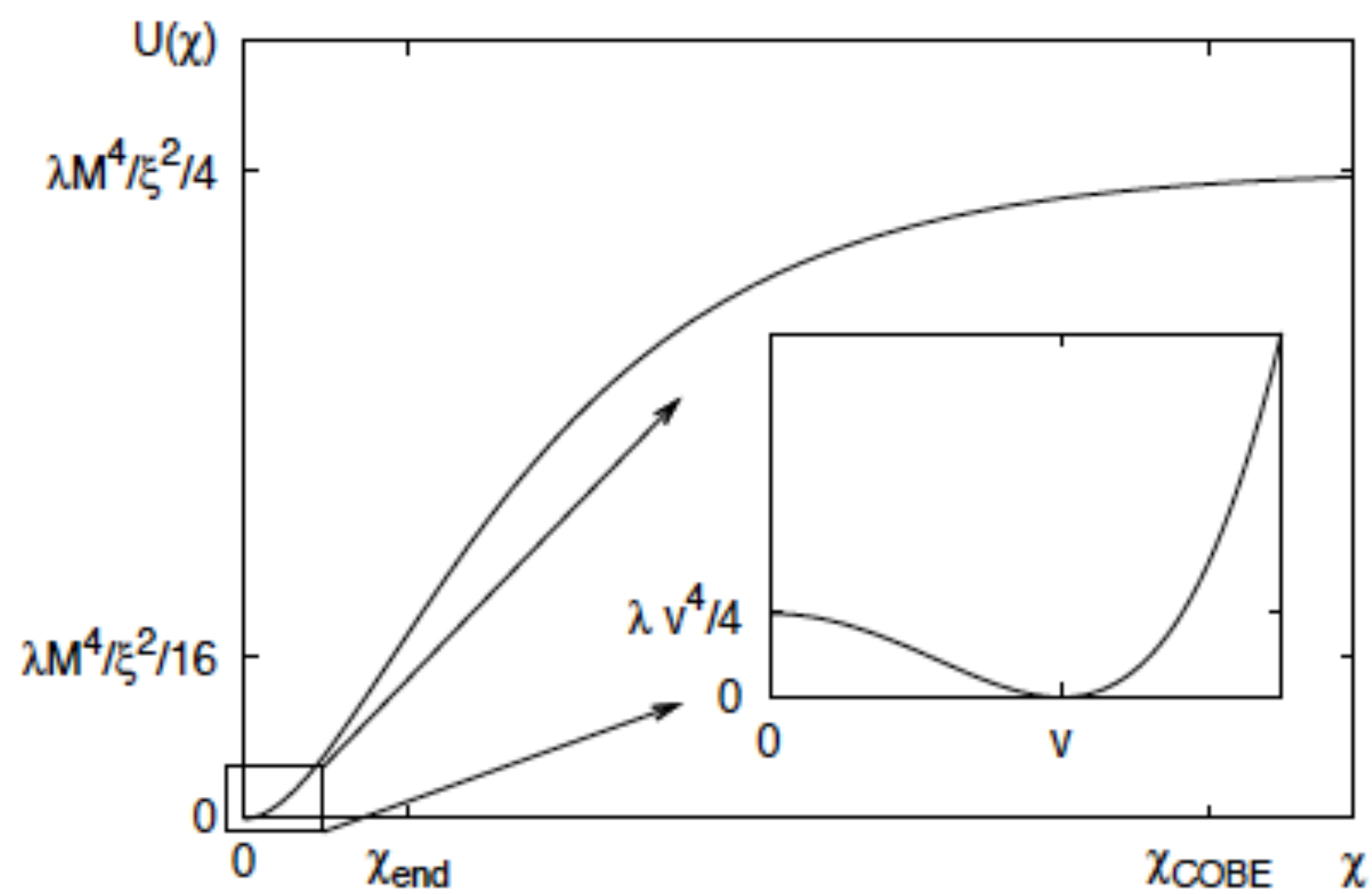
$$S_E = \int d^4x \sqrt{-\hat{g}} \left\{ -\frac{M_P^2}{2} \hat{R} + \frac{\partial_\mu \chi \partial^\mu \chi}{2} - U(\chi) \right\}$$

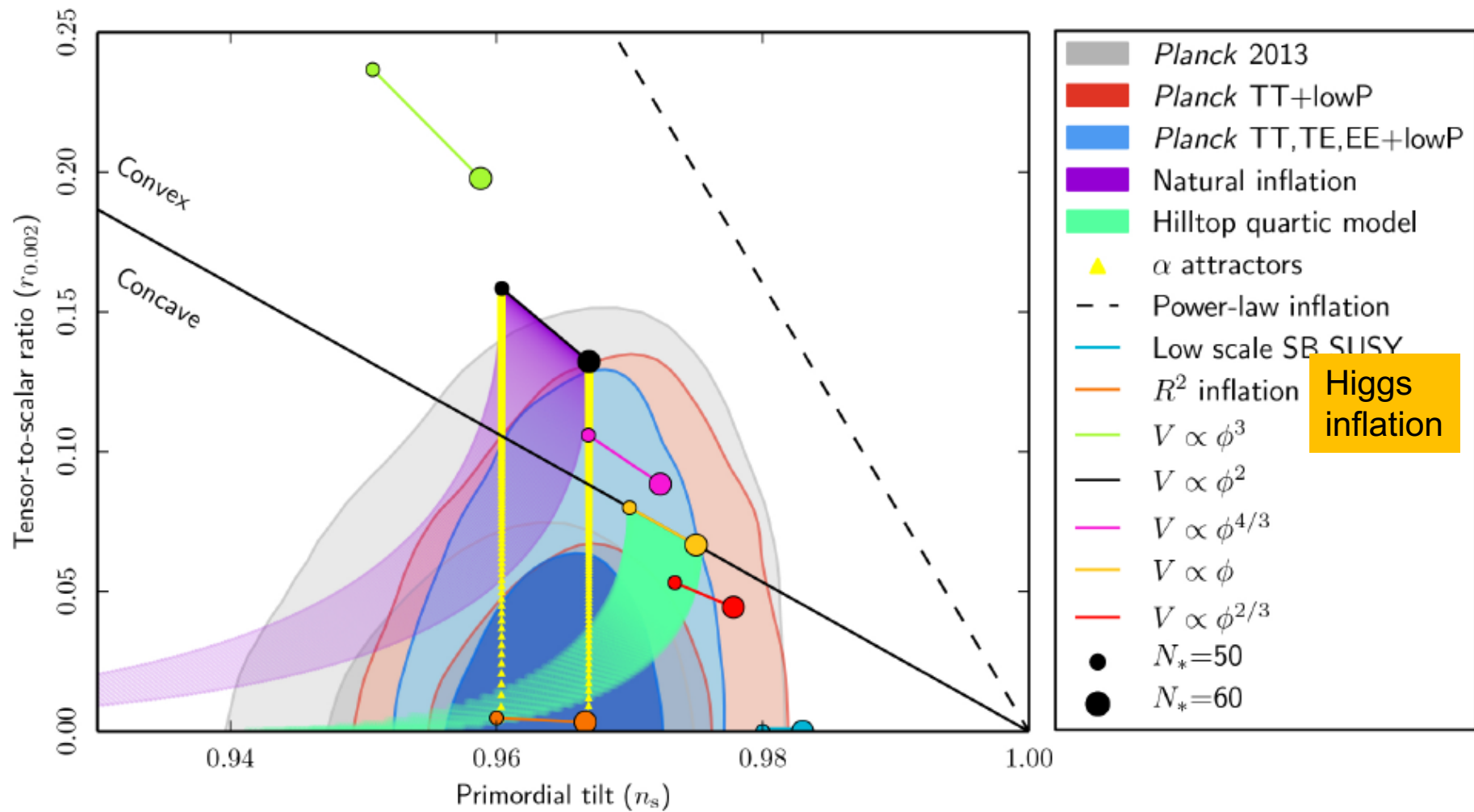
$$\frac{d\chi}{dh} = \sqrt{\frac{\Omega^2 + 6\xi^2 h^2 / M_P^2}{\Omega^4}}$$

Canonical scalar field χ as inflaton

$$U(\chi) = \frac{1}{\Omega(\chi)^4} \frac{\lambda}{4} (h(\chi)^2 - v^2)^2$$

Inflation potential





$$\xi \sim 10^4 m_H / v \text{ to give } \Delta_\zeta^2 \sim 2 \times 10^{-9}$$

Higgs + photon

magnetogenesis

$$S_E = \int d^4x \sqrt{-\hat{g}} \left\{ -\frac{M_P^2}{2} \hat{R} + \frac{\partial_\mu \chi \partial^\mu \chi}{2} - U(\chi) - \frac{1}{4} I(\chi) F^{\mu\nu} F_{\mu\nu} \right\}$$

$$\frac{\partial^2 \vec{A}}{\partial \tau^2} - \vec{\nabla}^2 \vec{A} + \frac{I'}{I} \frac{\partial \vec{A}}{\partial \tau} = 0$$

$$\vec{E} = -\frac{1}{a^2} \frac{\partial \vec{A}}{\partial \tau}, \quad \vec{B} = \frac{1}{a^2} \vec{\nabla} \times \vec{A}.$$

$$\vec{A}(\tau, \vec{x}) = \sum_{\lambda=\pm} \int \frac{d^3k}{(2\pi)^{3/2}} \left[\vec{\epsilon}_\lambda(\vec{k}) a_\lambda(\vec{k}) A_\lambda(\tau, \vec{k}) e^{i\vec{k} \cdot \vec{x}} + \text{h.c.} \right]$$

$$\left[\frac{\partial^2}{\partial \tau^2} + \frac{I'}{I} \frac{\partial}{\partial \tau} + k^2 \right] A_\pm(\tau, k) = 0$$

$$[a_\lambda(\vec{k}), a_{\lambda'}^\dagger(\vec{k}')] = \delta_{\lambda\lambda'} \delta(\vec{k} - \vec{k}')$$

Background

$$\ddot{\phi} + (3H + \Gamma)\dot{\phi} + \frac{dV}{d\phi} = \frac{1}{2} \frac{dI}{d\phi} \langle \vec{E}^2 - \vec{B}^2 \rangle \quad \chi = \phi + \delta\phi$$

$$3H^3 = \frac{1}{M_p^2} \left[\frac{1}{2} \dot{\phi}^2 + V(\phi) + \rho_{EB} + \rho_\gamma \right] \quad \Gamma \text{ perturbative reheating}$$

$$\dot{\rho}_\gamma + 4H\rho_\gamma = \Gamma\dot{\phi}^2$$

$$\rho_{\text{EM}} \equiv \frac{I}{2} \langle \vec{E}^2 + \vec{B}^2 \rangle = \frac{I}{4\pi^2 a^4} \int dk k^2 \sum_{\lambda=\pm} (|A'_\lambda|^2 + k^2 |A_\lambda|^2)$$

Conclusion

- The binary BHs observed by aLIGO/VIRGO gravity-wave detectors may be **primordial** BHs
- PBHs could be dark matter
- aLIGO saw indirect signals (GWs) of DM !?
- How to distinguish PBHs from astrophysical BHs?

PBHs have less spins – need more data and statistics

- Production of PBHs with a wide range of masses in trapped inflation models
- From a toy trapped inflation model to axion monodromy inflation with **modulations**
- PBHs need axion decay constant $f \sim M_{Planck}$

Oscillations in CMB from **modulations** with
axion decay constant $f \ll M_{Planck}$ [Flauger et al. 2010](#), [Choi+Kim 2016](#)

- Associated primordial gravitational waves in the frequency ranges of PTAs, LISA, aLIGO/VIRGO
polarized (+ mode / - mode) GWs in axion inflation